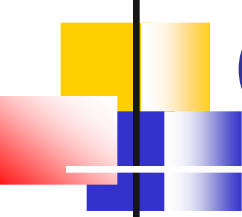


Non-Renormalization of Axial Anomaly in Perturbation Theory and Beyond

Advances of QFT (DIK-60)

October 5 2011

Oleg Teryaev
JINR, Dubna



Outline

Recollection on 3 loop “renormalization” of
Anomaly and its interpretation (80’s)

Anomaly sum rules (example – anomaly for
virtual photons and transition formfactors)

in collaboration with

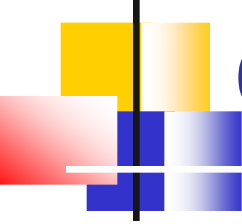
Yaroslav Klopot (JINR), Armen Oganesian (ITEP&JINR)

Phys.Lett.B695:130-135,2011.

e-Print: **arXiv:1009.1120** [hep-ph]

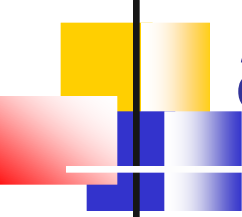
Phys. Rev. D 84, 051901(R) (2011).

e-Print: **arXiv:1106.3855** [hep-ph]



Symmetries and conserved operators

- (Global) Symmetry $\rightarrow$ conserved current ($\partial^\mu J_\mu = 0$)
- Exact:
- U(1) symmetry – charge conservation - electromagnetic (vector) current
- Translational symmetry – energy momentum tensor $\partial^\mu T_{\mu\nu} = 0$



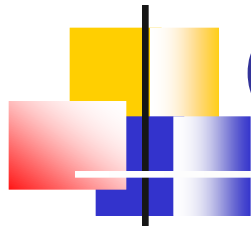
Massless fermions (quarks) – approximate symmetries

- Chiral symmetry (mass flips the helicity)

$$\partial^\mu J^5_\mu = 0$$

- Dilatational invariance (mass introduce dimensional scale – c.f. energy-momentum tensor of electromagnetic radiation)

$$T_{\mu\mu} = 0$$



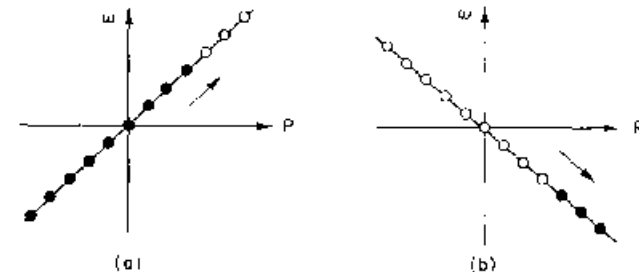
Quantum theory

- Currents $\rightarrow$ operators
- Not all the classical symmetries can be preserved $\rightarrow$ anomalies
- Enter in pairs (triples?...)
- Vector current conservation $\leftrightarrow$ chiral invariance
- Translational invariance $\leftrightarrow$ dilatational invariance

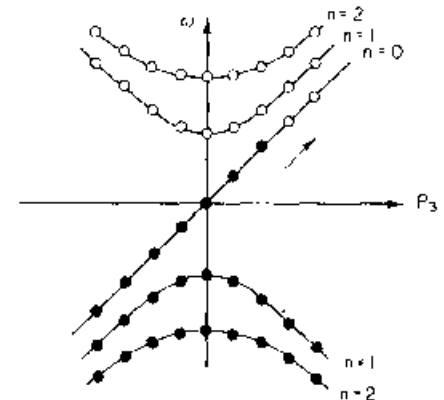
Calculation of anomalies

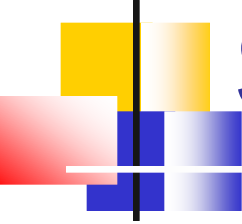
- Many various ways
- All lead to the same operator equation

$$\partial^\mu j_{5\mu}^{(0)} = 2i \sum_q m_q \bar{q} \gamma_5 q - \left(\frac{N_f \alpha_s}{4\pi} \right) G_{\mu\nu}^a \tilde{G}^{\mu\nu,a}$$



- UV vs IR languages-
understood in physical
picture (Gribov, Feynman,
Nielsen and Ninomiya)
of Landau levels flow (E||H)





Some history

- **Radiative Corrections To The Axial Anomaly.**
[A.A. Anselm](#), ([Birkbeck Coll.](#)) , [A.A. Johansen](#), ([St. Petersburg, INP](#)) .
Print-89-0622 (BIRKBECK), Jul 24, 1989. (Published Jul 1989). 16pp.
Published in **JETP Lett.49:214-218,1989, Sov.Phys.JETP 69:670-682,1989, Pisma Zh.Eksp.Teor.Fiz.49:185-189,1989.**
- Cited 38 times
- **In the course of our work on anomaly for nucleon spin structure -**
- **On Renormalization Of Axial Anomaly.**
[A.V. Efremov](#), [O.V. Teryaev](#), ([Dubna, JINR](#)) . JINR-E2-89-392, Jun 1989. 4pp.
Published in **Sov.J.Nucl.Phys.51:943-944,1990, Yad.Fiz.51:1492-1494,1990.**
- Cited 6 times
- My seminar at BLTP (chaired by DIK) in the role of devil's attorney...
- Elucidating discussions with Dima

What (non-)renormalization means?

- Extra power of coupling in K
- Non-conservation -> MULTIPLICATIVE (K is gauge-dependent) renormalization of J
- Various ways to calculate the same anomalous dimension

$$\partial \mu(J-K) \mu = 0$$

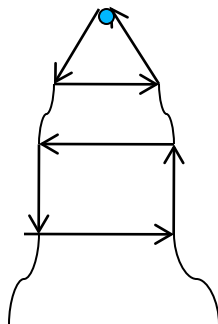
$$\dot{J} = \gamma J$$

$$\dot{J} = \gamma \partial J - \langle e | \dots | e' \rangle - 2 \text{loops (Adler'69)}$$

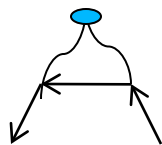
$$\dot{J} = \gamma \partial K - \langle \gamma | \dots | \gamma' \rangle - 3 \text{loops (AJ'89)}$$

$$\partial K = \gamma \partial J - \langle e | \dots | e' \rangle - 1 \text{loop (DGLAP'70s)}$$

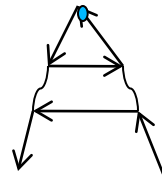
$$\partial K = \gamma \partial K - \langle \gamma | \dots | \gamma' \rangle - 2 \text{loops (nobody?)}$$



■ '89



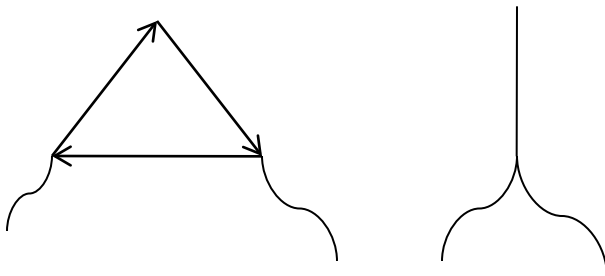
70's



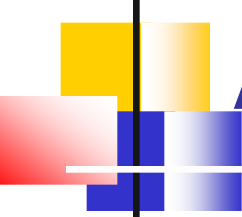
'69

NP stability of anomaly – t'Hooft consistency principle

- Zero quark mass pole
- Baryons are massive ->
- Massless pion (complementary to NG – stable against small corrections)
- Virtual photon – no pole in quark triangle, but pion pole persists – t'Hooft principle questioned (Achasov 90's)



- Principle is valid, but no pole (Horejsi, OT '95)!



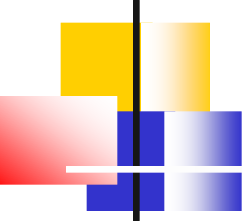
Anomaly and virtual photons

- Often assumed that only manifested in real photon amplitudes
- Not true – appears at any Q^2
- Natural way – dispersive approach to anomaly (Dolgov, Zakharov'70) - anomaly sum rules
- One real and one virtual photon – Horejsi, OT'95

- where
$$\int_{4m^2}^{\infty} A_3(t; q^2, m^2) dt = \frac{1}{2\pi}$$

$$F_j(p^2) = \frac{1}{\pi} \int_{4m^2}^{\infty} \frac{A_j(t)}{t - p^2} dt, \quad j = 3, 4$$

$$T_{\alpha\mu\nu}(k, q) = F_1 \varepsilon_{\alpha\mu\nu\rho} k^\rho + F_2 \varepsilon_{\alpha\mu\nu\rho} q^\rho + F_3 q_\nu \varepsilon_{\alpha\mu\rho\sigma} k^\rho q^\sigma + F_4 q_\nu \varepsilon_{\alpha\mu\rho\sigma} k^\rho q^\sigma + F_5 k_\mu \varepsilon_{\alpha\nu\rho\sigma} k^\rho q^\sigma + F_6 q_\mu \varepsilon_{\alpha\nu\rho\sigma} k^\rho q^\sigma$$



Dispersive derivation

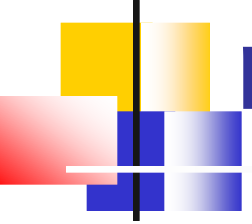
- Axial WI $F_2 - F_1 = 2mG + \frac{1}{2\pi^2}$
- GI $F_2 - F_1 = (q^2 - p^2)F_3 - q^2 F_4$
- No anomaly for imaginary parts

$$(q^2 - t)A_3(t) - q^2 A_4(t) = 2mB(t) \qquad F_j(p^2) = \frac{1}{\pi} \int_{4m^2}^{\infty} \frac{A_j(t)}{t - p^2} dt, \quad j = 3, 4$$

- Anomaly as a finite subtraction

$$F_2 - F_1 - 2mG = \frac{1}{\pi} \int_{4m^2}^{\infty} A_3(t) dt \qquad \int_{4m^2}^{\infty} A_3(t; q^2, m^2) dt = \frac{1}{2\pi}$$

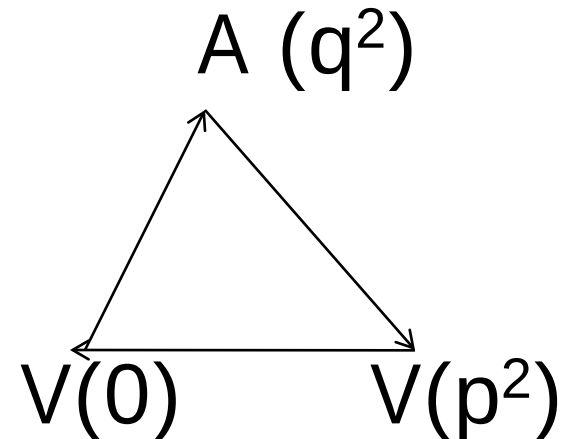
Properties of anomaly sum rules



- Valid for any Q^2 (and quark mass)
- No perturbative QCD corrections (Adler-Bardeen theorem)
- No non-perturbative QCD corrections (t'Hooft consistency principle)
- Exact – powerful tool
- For $Q^2=0$ – pole (no pole required for extra dimensionful parameter – temperature, chemical potential?)

Anomaly in axial and vector channels

- Dispersion relations in q^2 or p^2



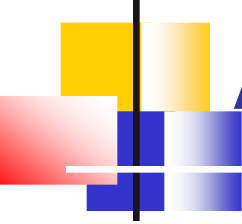
- Results coincide
- 2 ASRs – V-A global duality
- Local duality – Vainshtein relation ('02) between L and T (not related to anomaly!)
V-A correlators in external field
- Corrections... AdS/QCD - Son, Yamamoto

Mesons contributions

(Klopot, Oganesian, OT)

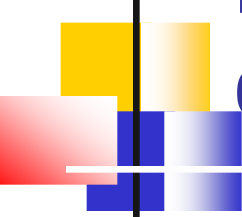
Phys.Lett.B695:130-135,2011 (1009.1120) and 1106.3855

- Pion – saturates sum rule for real photons $ImF_3 = \sqrt{2}f_\pi\pi F_{\pi\gamma\gamma^*}(Q^2)\delta(s - m_\pi^2)$ $F_{\pi\gamma^*\gamma}(0) = \frac{1}{2\sqrt{2}\pi^2 f_\pi}$
- For virtual photons – pion contribution is rapidly decreasing $F_{\pi\gamma\gamma^*}^{asympt}(Q^2) = \frac{\sqrt{2}f_\pi}{Q^2} + \mathcal{O}(1/Q^4)$
- This is also true also for axial and higher spin mesons (longitudinal components are dominant)
- Heavy PS decouple in a chiral limit



Anomaly as a collective effect

- One can never get constant summing finite number of decreasing function
- Anomaly at finite Q^2 is a **collective** effect of meson spectrum
- **General** situation –occurs for any scale parameter (playing the role of **regulator** for massless pole)
- For quantitative analysis – quark-hadron duality



Mesons contributions within quark hadron duality – transition FFs

- Pion:
$$F_{\pi\gamma\gamma^*}(Q^2) = \frac{1}{2\sqrt{2}\pi^2 f_\pi} \frac{s_0}{s_0 + Q^2}$$
- Cf Brodsky&Lepage, Radyushkin – comes now from anomaly!
- Axial mesons contribution to ASR

$$\int_0^\infty A_3(s; Q^2) ds = \frac{1}{2\pi} = I_\pi + I_{a_1} + I_{cont}. \quad I_{a_1} = \frac{1}{2\pi} Q^2 \frac{s_1 - s_0}{(s_1 + Q^2)(s_0 + Q^2)}$$

Content of Anomaly Sum Rule ("triple point")

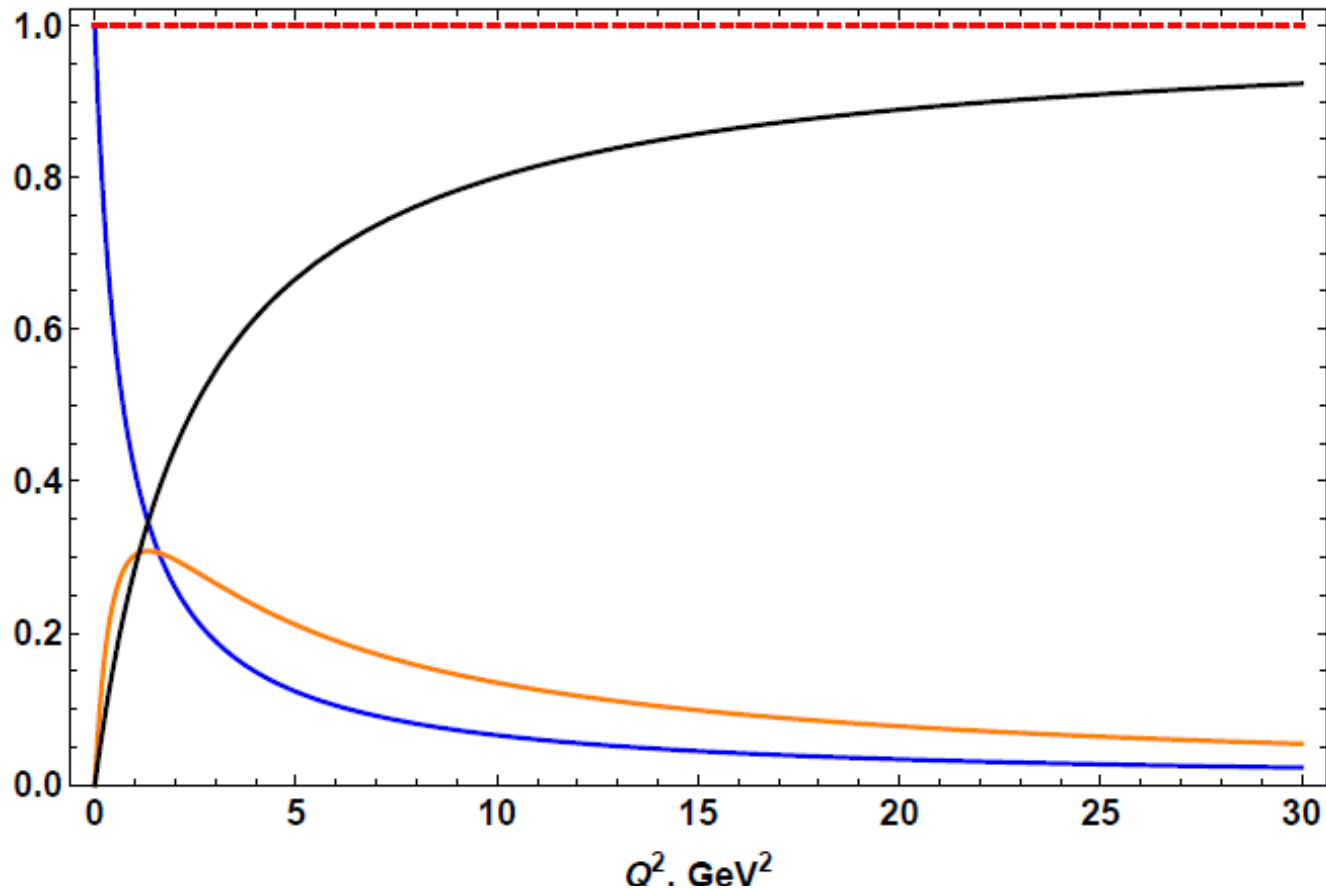
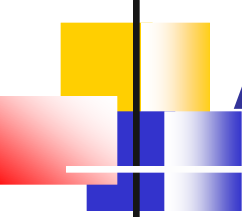


Figure 1: Relative contributions of π (blue line) and a_1 (orange line) mesons, intervals of duality are $s_0 = 0.7 \text{ GeV}^2$ and $s_1 - s_0 = 1.8 \text{ GeV}^2$ respectively, and continuum (black line), continuum threshold is $s_1 = 2.5 \text{ GeV}^2$



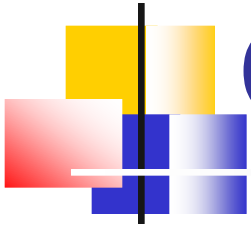
ASR and BaBar data

- In the BaBar(2009) region – main contribution comes from the continuum
- Small relative correction to continuum –due to exactness of ASR **must** be compensated by large relative contributions to lower states!

- Amplification of corrections

$$\frac{\delta I_{cont}/I_{cont}^0}{\delta I_{\pi}/I_{\pi}^0} = \frac{s_0}{Q^2} \simeq \frac{1}{30} \quad Q^2 = 20 \text{ GeV}^2, s_0 = 0.7 \text{ GeV}^2$$

- Smaller for eta because of larger duality interval (supported by BaBar)



Corrections to Continuum

- Perturbative – zero at 2 loops level (massive-Pasechnik&OT – however cf Melnikov; massless-Jegerlehner&Tarasov)
- Non-perturbative (e.g. instantons, short strings)
- The general properties of ASR require decrease at asymptotically large Q^2 (and $Q^2=0$)

- Corresponds to logarithmic contribution (cf Radyushkin, Polyakov, Dorokhov).

$$I_{cont} = \frac{1}{2\pi} \frac{Q^2}{s_0 + Q^2} - cs_0 \frac{\ln(Q^2/s_0) + b}{Q^2},$$

$$I_{\pi} = \frac{1}{2\pi} \frac{s_0}{s_0 + Q^2} + cs_0 \frac{\ln(Q^2/s_0) + b}{Q^2}.$$

Modelling of corrections

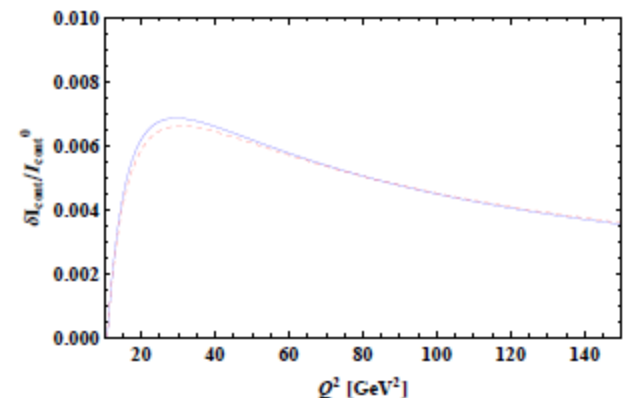
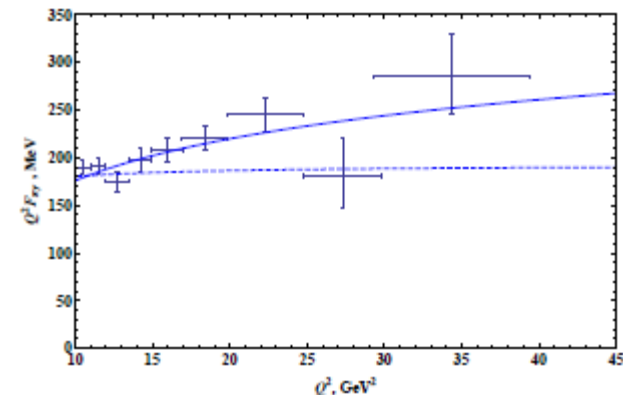
- Continuum vs pion

- Fit $b = -2.74, \quad c = 0.045.$

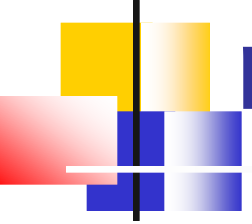
- Continuum contribution similar for Radyushkin's approach

$$I_{cont} = \frac{1}{2\pi} \frac{Q^2}{s_0 + Q^2} - cs_0 \frac{\ln(Q^2/s_0) + b}{Q^2},$$

$$I_{\pi} = \frac{1}{2\pi} \frac{s_0}{s_0 + Q^2} + cs_0 \frac{\ln(Q^2/s_0) + b}{Q^2}.$$

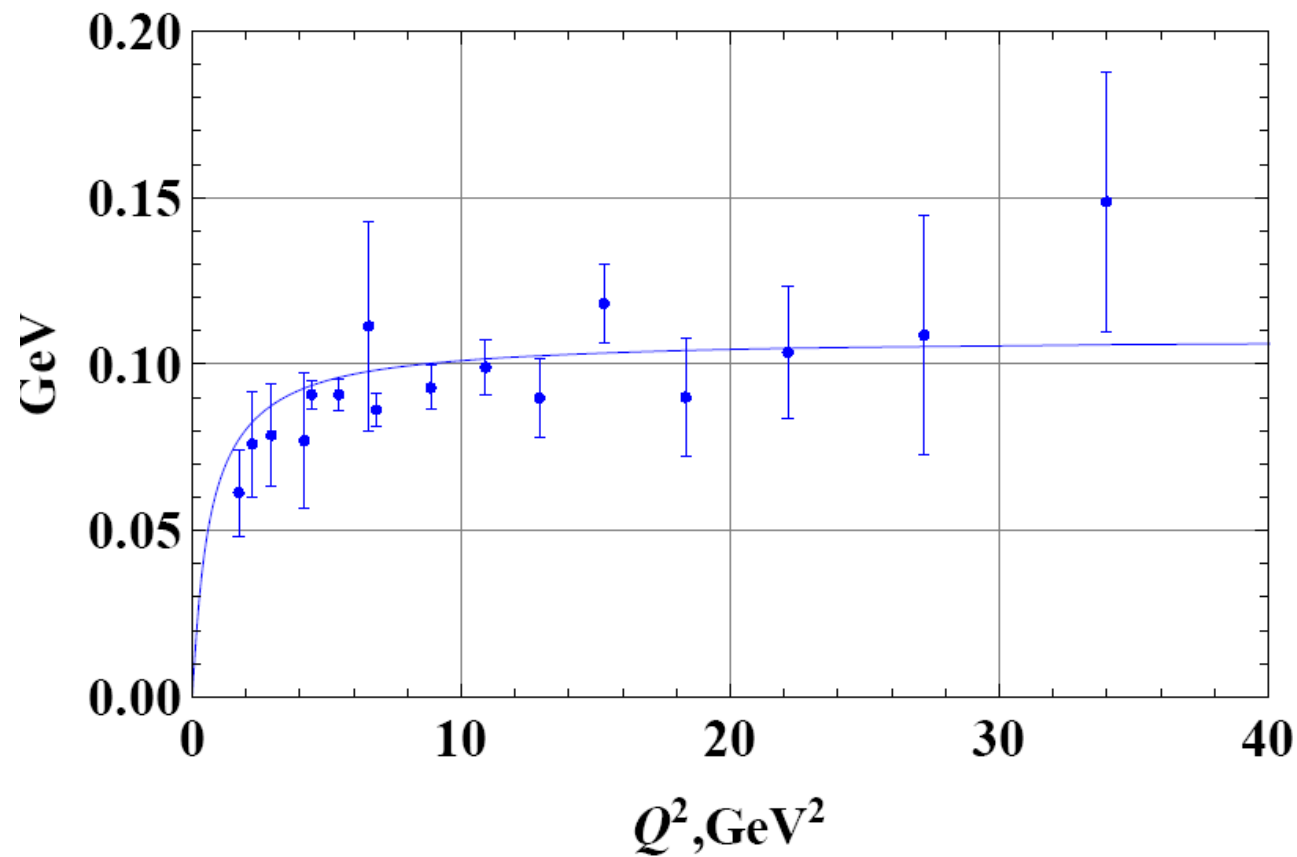


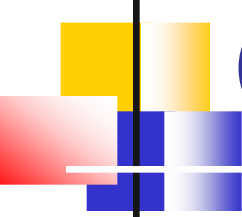
Interplay of pion with lower resonances

- 
-
- Small (NP) corrections to continuum – interplay of pion with higher states
 - A1 – decouples for real photons
 - Relation between transition FF's of pion and A1 (testable!)

Generalization for $\eta(\prime)$

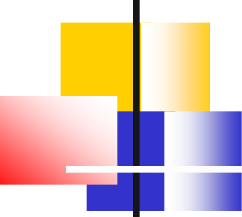
- Octet channel sum rule (gluon anomaly free)





Conclusions/Discussion-I

- New manifestation of Axial Anomaly - Anomaly Sum Rule – exact NPQCD tool- do not require QCD factorization
- Anomaly for virtual photons – collective effect (with fast excitation of collective mode)
- Similar collective effect is expected for finite temperature and/or chemical potential
- Exactness of ASR – very unusual situation when small pion contribution can be studied on the top of large continuum – amplification of corrections to continuum
- BaBar data – small negative correction to continuum
- If continuum is precisely described by Born term–interplay with A_1 (TO BE STUDIED THEORETICALLY AND EXPERIMENTALLY)



HAPPY

BIRTHDAY!