

# Imaginary Acceleration

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# Explanation of the title of the talk

Based on:

“On axial current in rotating and accelerating medium”

G. Prokhorov, O. Teryaev, V. Zakharov,

arXiv:1805.12029 [hep-th] published in

“Rapid Communications” PRD

one evaluates matrix element  $\langle \vec{J}_5 \rangle_{\text{statistical}}$

Result can be reproduced through the substitution:

$$\mu \rightarrow \mu + \frac{\Omega}{2} - i\frac{\mathbf{a}}{2}$$

where  $\mu$  is the chemical potential,  $\Omega$  - angular velocity,  $\mathbf{a}$  - acceleration (Signs are rather  $\pm$ )

# Alexander Vilenkin, PRD 21 (1980) 2260

The field is started by A. Vilenkin, in case of **rotation**:

$$\langle \mathbf{J}^\mu(\vec{x}) \rangle = \text{Tr}(\rho \mathbf{J}^\mu(\vec{x}, t))$$

where  $\mathbf{J}^\mu = \frac{1}{2}[\bar{\Psi}, \gamma^\mu \Psi]$  is the current density operator and

$$\rho = \mathbf{C} \exp\left(-\beta(H - \vec{M} \cdot \vec{\Omega} - \sum_i \mu_i \mathbf{N}_i)\right)$$

$\beta = T^{-1}$ ,  $\rho$  is the statistical operator (see Landau&Lifshitz)  
 $\vec{M}$  is the angular momentum,  $\vec{\Omega}$  is the angular velocity  
 $\mu_i$  is chemical potential,  $\mathbf{N}_i$  is number of charged particles  
 $\rho$  is built on conserved operators

# Vilenkin, cnt'd

One-loop effect in finite-temperature QFT

$$\langle \mathbf{J}^\mu(\vec{\mathbf{x}}) \rangle = -\text{Tr} \gamma^\mu \mathbf{S}(\vec{\mathbf{x}}, \tau; \vec{\mathbf{x}}, \tau + \epsilon)_{\epsilon \rightarrow 0}$$

reduces to one of the so called Sommerfeld integrals  
(M.Stone (2018))

$$\langle \mathbf{J}_\Omega \rangle = \frac{1}{4\pi^2} \int_{-\infty}^{+\infty} \epsilon^2 d\epsilon \cdot \left( \frac{1}{1 + e^{\beta(\epsilon - (\mu + \Omega/2))}} - \frac{1}{1 + e^{\beta(\epsilon - (\mu - \Omega/2))}} \right)$$

Finally, (for a single Weyl fermion of unit charge):

$$\langle \vec{\mathbf{J}}(\mathbf{0}) \rangle = -\vec{\Omega} \left( \mu^2/4\pi^2 + T^2/12 + \Omega^2/(\pi^2 48) \right)$$

so called **Chiral Vortical Effect**

## More recent generalizations

Basic points on  $\rho$  from Landau&Lifshitz:

$$\ln \rho_{12} = \ln \rho_1 + \ln \rho_2 \quad (\text{additivity})$$

$$\partial\rho/\partial t = i[H, \rho], \text{ in equilibrium } [H, \rho] = 0\dots$$

Recently:

- Relativistic formulation,  $\beta \rightarrow \beta_\nu$ ,  $\beta_\nu \equiv \mathbf{u}_\nu/T_0$
- Possible “conflict” between massless and massive cases (generally speaking, non-trivial in view of the anomaly)
- Phenomenology (of heavy-ion collisions) suggests that the medium is accelerated (STAR Collaboration)

# Francesco Becattini + coauthors, $\sim$ 2017

As a response to the challenge, these authors suggested

$$\rho = \frac{1}{Z} \exp[-\hat{H}/T_0 + \mathbf{a}_z \hat{K}_z/T_0]$$

where  $\mathbf{a}_z$  is acceleration and the operator  $\hat{K}_z$  is the boost (both in  $\mathbf{z}$ -direction). [A basic generalization of L&L.](#)

There exists prehistory, mostly for scalar particles

Not without problems, e.g.:

- explicit dependence on time is introduced
- should we factorize motion of the center of mass?

[A strong case in favor of the extension of  \$\hat{\rho}\$  has been made](#)

## Prokhorov et al. (2018)

It is in this framework (and Prokhorov+Teryaev (2017)) that imaginary acceleration has been introduced, see above  
A more precise formulation of Prokhorov et al. result:

$$2 \langle J_{\Omega}^5 \rangle = \int \frac{d^3p}{(2\pi)^3} \left( n_F(E_p - \mu - \Omega/2 + ia/2) - n_F(E_p - \mu + \Omega/2 - ia/2) + n_F(E_p + \mu - \Omega/2 + ia/2) - n_F(E_p + \mu + \Omega/2 - ia/2) + \text{c.c.} \right)$$

where  $n_F$  is the corresponding Fermi distribution (mass of the fermion is **not** necessarily vanishing)

# Imaginary acceleration as a sign of instability

Because of the factor  $i$  the second order in  $ia$  is **negative**

As a result, there emerges **instability** at temperature  $T < T_{Unruh}$ :

- The axial current oscillates like mad at  $T < a/(2\pi)$
- If the average energy-momentum is evaluated along similar lines, density of energy is negative at temperatures below the Unruh temperature (Becattini)
- The axial current is somewhat **stabilized** with increasing angular velocity (Prokhorov et al.(2018))

At least qualitatively, the picture is similar to the Hwking/Unruh effect



# Conclusions to Part I (“Phenomenology”)

- Following the phenomenology path we apparently arrive at another incarnation of the Hawking radiation, or Unruh effect
- It is better to become “more theoretical” at this point. Let’s go back to FT

## Part II: Is there place for $ia$ in Field Theory?

The answer is definite although sounds unexpected:  
“Yes, one could have readily dug out the imaginary acceleration long time ago”

The most straightforward way is to start with the **chiral limit** ( L&L, Becattini... are rather non-relativistic)

Namely, it is a textbook statement that generators of the Lorentz transformation are realized on the fundamental (massless) fermion representation as:

$$\hat{J}_z = \frac{\hat{\sigma}_z}{2}, \quad \hat{K}_z = \frac{i\hat{\sigma}_z}{2}$$

and we immediately come to  $\mu \rightarrow \mu + \Omega/2 + ia/2$

# Imaginary acceleration as signal of instability

In other words, we do not notice that  $\hat{K}_z$  is anti-hermitian as far as work with Lagrangians and use  $\psi^\dagger \gamma_0$  instead of  $\psi^\dagger$ . Statistics (or systems with finite densities) make  $\langle \hat{K}_z \rangle$  observable and reveal that boost is realized as

$$(\hat{K}_z)^\dagger = -\hat{K}_z$$

Imaginary acceleration  $ia$ , entering along with real  $\mu$  looks as imaginary energy and, apparently, signals **instability**

What kind of instability?—**Field-theoretic** instability of finite-density, accelerated matter with  $T = 0$ .

# Lorenzian signature vs Unruh effect

Back to text-books: one chooses  $K_z = i\sigma_z/2$  to imitate commutators of the Hermitian boost operators.

Thus, in the second order we observe the minus-sign inherent to the commutators rather than  $\sqrt{-1} \mathbf{a}$ .

As a result, the Unruh-instability gets related to the Lorenzian signature  
(amazing!)

# From FT to QFT : from Facts to Speculations

**A fact:** Quantizing allows to introduce  $\psi(\vec{p})$ , (and the corresponding  $a^\dagger(\vec{p})$ ,  $a(\vec{p})$ ) a kind of “infinite-component representation”.

Boosts can transform density of fermions with momentum  $\vec{p}$  into density with boosted value of the momentum.

Restoring hermitian  $\hat{K}_z$ .

Apparently, this trick should work for integration over virtual momenta as well.

**Except for the anomalous cases** when shifts of momenta are not allowed. Hence, speculation:

in QFT the instability can come back with anomaly

# Support for re-discovering the Unruh effect

- Proceeding to evaluate energy-density of the equilibrium, accelerated state one finds out (Becattini (2018)) that for temperatures below  $T_{Unruh}$  (better to say, below or order of) **energy-density is negative**
- Similarly (Prokhorov et al. (2018)), the axial current at temperatures below  $T_{Unruh}$  **oscillates like mad**
- The axial current is somewhat **stabilized** with increasing angular velocity (Prokhorov et al.(2018))

All these observations (qualitatively) agree with the Unruh effect/instability.

# Quantum Field Theory: role of anomalies

Independently, the role of anomalies in generating thermal chiral effects was clarified by M Stone (2018), who generalized approach due to Wilzcek (2005)

$$\nabla_{\mu} T^{\mu\nu} = F^{\nu A} J_{N,A} - \frac{1}{384\pi^2} \frac{\epsilon^{\rho\sigma\alpha\beta}}{\sqrt{-g}} \nabla_{\mu} [F_{\rho\sigma} R^{\nu\mu}_{\alpha\beta}]$$

$$\nabla_{\mu} J_N^{\mu} = -\frac{1}{32\pi^2} \frac{\epsilon^{\mu\nu\rho\sigma}}{\sqrt{-g}} F_{\rho\nu} F_{\mu\sigma} - \frac{1}{768\pi^2} \frac{\epsilon^{\mu\nu\rho\sigma}}{\sqrt{-g}} R^{\alpha}_{\beta\mu\nu} R^{\beta}_{\alpha\rho\sigma}$$

where  $J_N$  is the number-current, corresponds to a single right-handed Weyl spinor, and other notations are standard

## Anomalies, cont'd

In all the cases the metric is factorized into a 2d black hole,

$$ds^2 = -f(r)dt^2 + \frac{1}{f(r)}dr^2, \quad f(r = r_H) = 0$$

and transverse coordinates which are occupied by a uniform magnetic ( $\vec{B} \neq 0$ ) or rotational ( $\vec{\Omega} \neq 0$ ) fields

The **2d** gravitational anomaly is quite readily integrated out in time-independent (equilibrium) case

$$\sqrt{|g|} \nabla_\mu (T^{\mu\nu} \eta_\nu) = \frac{c}{96\pi} \epsilon^{\nu\sigma} \eta_\nu \partial_\sigma R$$

to produce a flux of particles far off from the horizon (path to the Hawking radiation found by Wilzcek)



# Anomalies, cnt'd

Reproduced thermal effects, with  $T = T_{\text{Black hole}}$

Energy flux:

$$\vec{J}_{\text{energy}} = \vec{B} \left( \frac{\mu^2}{8\pi i^2} + \frac{1}{24} T^2 \right)$$

Number-of-particles flux:

$$\vec{J}_N = \vec{\Omega} \left( \frac{\mu^2}{4\pi^2} + \frac{|\vec{\Omega}|^2}{48\pi^2} + \frac{1}{12} T^2 \right)$$

The whole of thermal CVE given exactly by the anomalies.  
No a priori mention of the temperature. Pure field theory.  
Temperature is nothing else but acceleration (on the horizon)

## Conclusions to Part II

QFT, through anomalies, clearly signals that matter occupying part of space— because of the horizon —is unstable against particle production. Quantitatively, Hawking radiation from a black hole is reproduced (F. Wilzcek, M. Stone....)

Hint on the instability, contained in classical (in Grassmann numbers) theory of fermions is promoted to a fully quantitative framework by QFT

# Questions not answered

Many more questions to answer. E.g.:

- In a way, we have not started yet. Would like to consider **imaginary part** So far, quadratic in acceleration terms.
- What is analogy to

$$eA_\alpha \rightarrow eA_\alpha + \mu \cdot u_\alpha$$

- Possible dependence on IR, and so on
- Critique (?)

## Part III: $T \leftrightarrow a$ Duality as Guiding Principle

Absolutely unfinished part.

To reiterate, duality between what and what:

- On one hand, one can apply thermal field theory in flat space and evaluate the thermal part of the chiral vortical effect.  $\vec{j} = (T^2/12) \cdot \vec{\Omega}$
- On the other hand, one can evaluate flow of charge from the horizon in terms of the gravitational anomaly (no explicit temperature). The input is metric near the horizon,  $g_{00} \rightarrow 0$
- Results coincide for (massless) spin-1/2 constituents, provided  $T = a/(2\pi)$

## Further tests of the duality

Chiral vortical effect for photons (in grav. field). The basic new ingredient is the bosonic chiral anomaly:

$$\nabla^\alpha K_\alpha = (\text{const}) R\tilde{R}$$

where  $K_\alpha = \epsilon_{\alpha\beta\gamma\delta} A^\beta \partial^\gamma A^\delta$ ,  $R\tilde{R}$ - product of Riemann tensors

The two ways of evaluating the CVE for photons seem to disagree with each other by a factor of two.

Should we modify the field-theoretic calculation and impose the duality?

Also, composite particles are treated differently in the two approaches

# Overall conclusion

Equations with imaginary acceleration seem to be true and shed new light on the Hawking radiation in FT

The crucial question is whether there are novel applications