

Beam energy scan using a viscous hydro+cascade model

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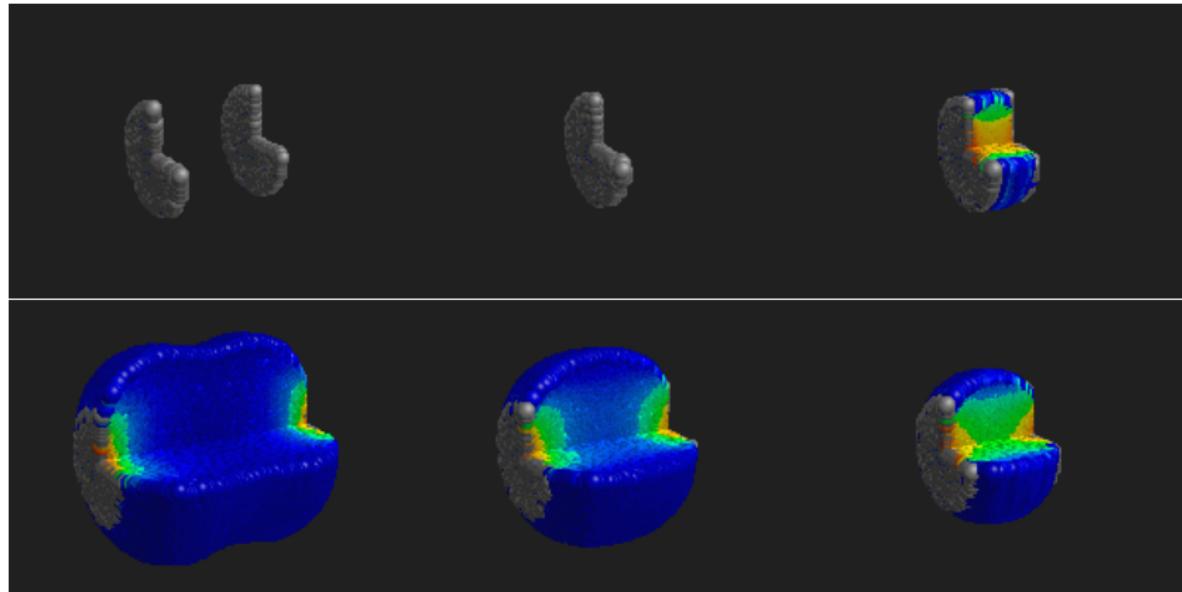
BLTP JINR Dubna, Apr 16, 2014

In collaboration with M. Bleicher, P. Huovinen, H. Petersen

arXiv:1310.0702 arXiv:1311.0133 arXiv:1312.4160 [nucl-th]



Introduction: heavy ion collision in pictures¹



Typical size
 $10 \text{ fm} \approx 10^{-14} \text{ m}$

Typical lifetime
 $10 \text{ fm/c} \approx 10^{-23} \text{ s}$

10^{-8} sec after the collision: hadrons are detected

¹https://www.jyu.fi/fysiikka/tutkimus/suurenenergia/urhic/anim1.gif/image_view_fullscreen

//www.jyu.fi/fysiikka/tutkimus/suurenenergia/urhic/anim1.gif/image_view_fullscreen

"Stages of Heavy Ion Collision"

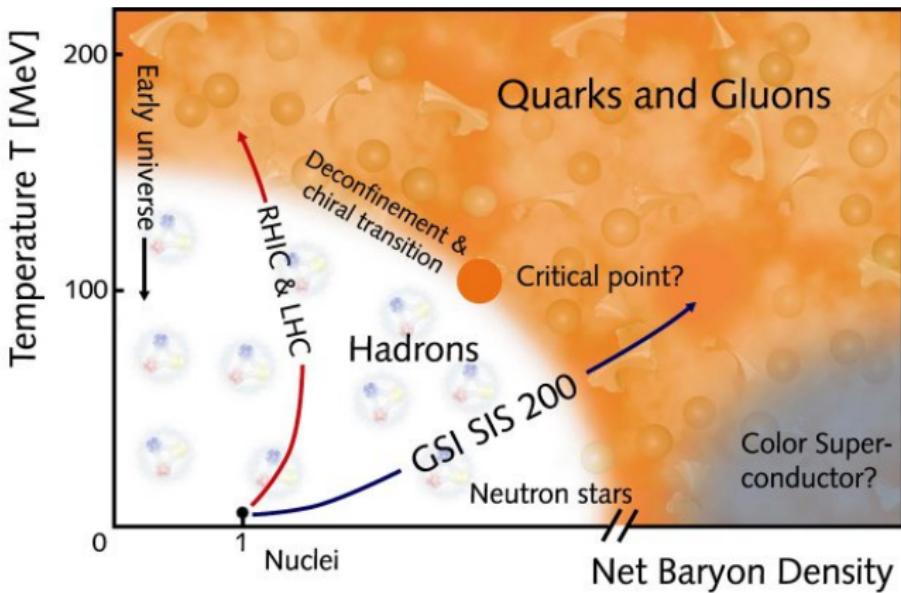
- ① Initial(pre-thermal) stage
 - ▶ Thermalization
- ② Hydrodynamic expansion
 - ▶ Quark-gluon plasma phase
 - ▶ Phase transition
 - ▶ Hadron Gas phase
 - ▶ Chemical freeze-out
 - ▶ End of hydrodynamic regime
- ③ Kinetic stage
 - Kinetic freeze-out
 - ↓
 - Free streaming, then hadrons are detected



1. Ingredients of hydro+cascade model:

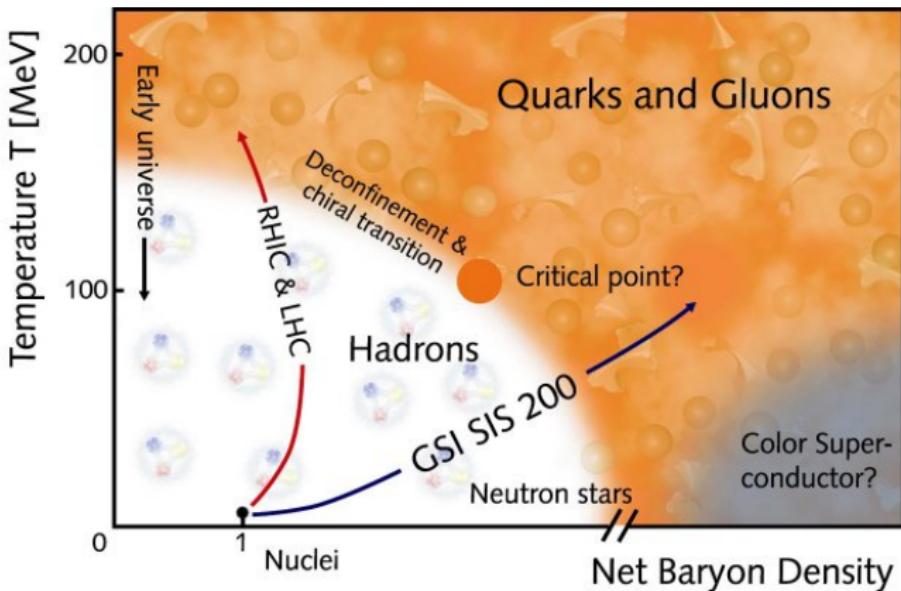
- ① Initial stage model
Enforced thermalization
- ② Hydrodynamic solution
 - ▶ Equation of state for hydrodynamics
 - ▶ transport coefficients
- ③ Particilization and switching to a cascade

Where do we want to apply our model

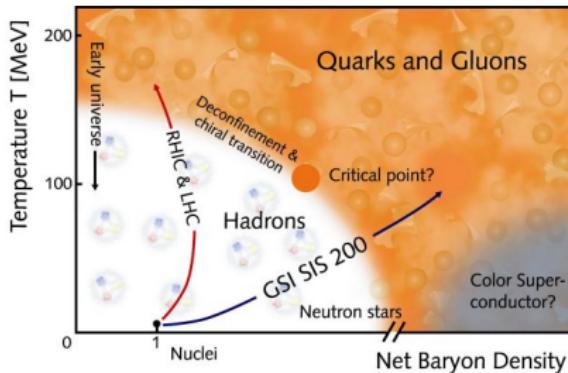


- small net baryon density: hydro(+cascade) model is well established
arXiv: "hydrodynamic" + "RHIC" = 44 manuscripts. Existing codes (by author):
Kolb, Song, Hirano, Nonaka, Chaudhuri, Mota, Luzum, Holopainen, Schenke, Bozek,
Molnar, Del Zanna, IuK,...

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- large net baryon density:
arXiv: "hydrodynamic" + "SPS" = 8 manuscripts
arXiv: "hydrodynamic" + "FAIR" = 3 manuscripts



Ingredients essential for beam energy scan studies are marked red.

EoS reference: J. Steinheimer,
S. Schramm and H. Stocker,
J. Phys. G 38, 035001 (2011).

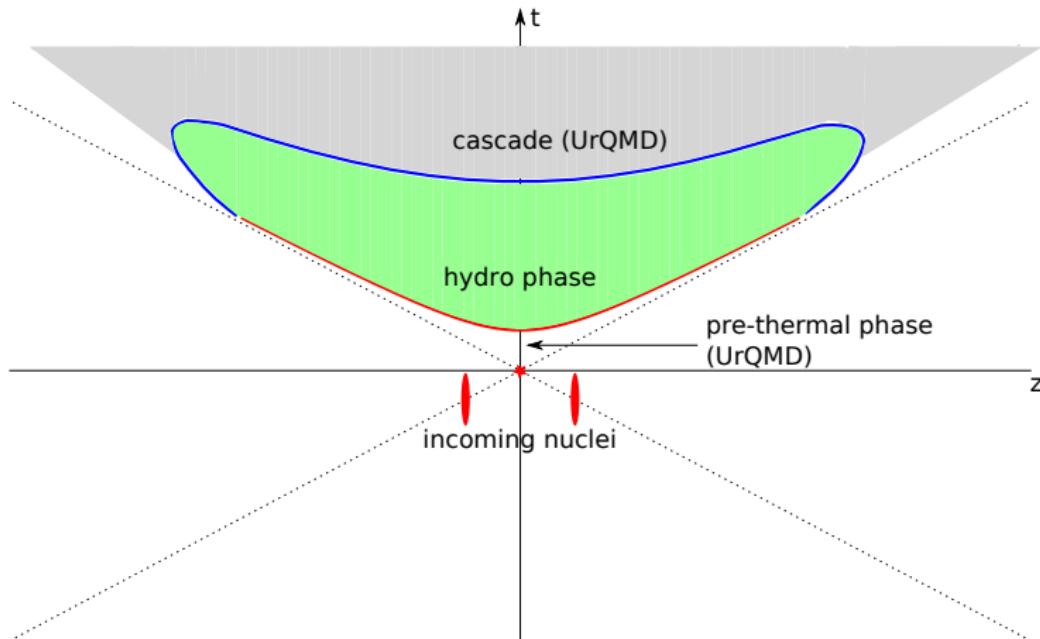
1. Ingredients of the model:

- ➊ Initial stage:
UrQMD
- ➋ Hydrodynamic solution
 - ▶ Equation of state for hydrodynamics:
Chiral model coupled to Polyakov loop to include the deconfinement phase transition
 - ★ good agreement with lattice QCD data at $\mu_B = 0$
 - ★ Applicable also at finite baryon densities
 - ▶ transport coefficients
- ➌ Particilization and switching back to cascade (UrQMD)

Initial conditions for hydrodynamic evolution

Time to switch from UrQMD to fluid description: $\tau = \frac{2R}{\gamma v_z} = \frac{2R}{\sqrt{(\sqrt{s}/2m_N)^2 - 1}}$

where $\tau = \sqrt{t^2 - z^2}$. Switching surface is the red curve



$\{T^{0\mu}, N_b^0, N_q^0\}$ of fluid = averaged $\{T^{0\mu}, N_b^0, N_q^0\}$ of particles

Hydrodynamic phase

The hydrodynamic equations in arbitrary coordinate system:

$$\partial_{;\nu} T^{\mu\nu} = \partial_\nu T^{\mu\nu} + \Gamma_{\nu\lambda}^\mu T^{\nu\lambda} + \Gamma_{\nu\lambda}^\nu T^{\mu\lambda} = 0, \quad \partial_{;\nu} N^\nu = 0 \quad (1)$$

where (we choose Landau definition of velocity)

$$T^{\mu\nu} = \varepsilon u^\mu u^\nu - (p + \Pi)(g^{\mu\nu} - u^\mu u^\nu) + \pi^{\mu\nu} \quad (2)$$

and $\Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$

Evolutionary equations for shear/bulk, coming from **Israel-Stewart** formalism:

$$\langle u^\gamma \partial_{;\gamma} \pi^{\mu\nu} \rangle = -\frac{\pi^{\mu\nu} - \pi_{\text{NS}}^{\mu\nu}}{\tau_\pi} - \frac{4}{3} \pi^{\mu\nu} \partial_{;\gamma} u^\gamma \quad (3a)$$

where

$$\langle A^{\mu\nu} \rangle = \left(\frac{1}{2} \Delta_\alpha^\mu \Delta_\beta^\nu + \frac{1}{2} \Delta_\alpha^\nu \Delta_\beta^\mu - \frac{1}{3} \Delta^{\mu\nu} \Delta_{\alpha\beta} \right) A^{\alpha\beta}$$

* Bulk viscosity $\zeta = 0$, charge diffusion=0

Coordinate transformations (hydro phase)

Milne coordinates

The coordinate system is defined as follows:

$$0) \tau = \sqrt{t^2 - z^2}$$

$$g^{\mu\nu} = \text{diag}(1, -1, -1, -1/\tau^2)$$

$$1) x = x$$

Nonzero Christoffel symbols are:

$$2) y = y$$

$$\Gamma_{\tau\eta}^\eta = \Gamma_{\eta\tau}^\eta = 1/\tau, \quad \Gamma_{\eta\eta}^\tau = \tau$$

$$3) \eta = \frac{1}{2} \ln \frac{t+z}{t-z}$$

$$T^{\mu\nu} = (\varepsilon + p) u^\mu u^\nu - p \cdot g^{\mu\nu}, \text{ where}$$

$$u^\mu = \{\cosh(\eta_f - \eta) \cosh \eta_T, \sinh \eta_T \cos \phi, \sinh \eta_T \sin \phi, \frac{1}{\tau} \sinh(\eta_f - \eta) \cosh \eta_T\}$$

$$(\text{cf. } u_{\text{Cart}}^i = \{\cosh(\eta_f) \cosh \eta_T, \sinh \eta_T \cos \phi, \sinh \eta_T \sin \phi, \sinh(\eta_f) \cosh \eta_T\})$$

Additional transformations:

EM conservation equations are

$$\partial_{;\nu} T^{\mu\nu} = 0$$

or

$$\mu = 0 : \partial_v T^{\tau v} + \tau T^{\eta\eta} + \frac{1}{\tau} T^{\tau\tau} = 0$$

$$\mu = 1 : \partial_v T^{xv} + \frac{1}{\tau} T^{x\tau} = 0$$

$$\mu = 2 : \partial_v T^{yv} + \frac{1}{\tau} T^{y\tau} = 0$$

$$\mu = 3 : \partial_v T^{\eta v} + \frac{3}{\tau} T^{\eta\tau} = 0$$

$$T^{\mu\eta} \rightarrow T^{\mu\eta}/\tau, \mu \neq \eta,$$
$$T^{\eta\eta} \rightarrow T^{\eta\eta}/\tau^2$$



$$\partial_v(\tau T^{\tau v}) + \frac{1}{\tau}(\tau T^{\eta\eta}) = 0$$

$$\partial_v(\tau T^{xv}) = 0$$

$$\partial_v(\tau T^{yv}) = 0$$

$$\partial_v(\tau T^{\eta v}) + \frac{1}{\tau}\tau T^{\eta\tau} = 0$$

Conservative variables are
 $Q^\mu = \tau \cdot T^{\mu\eta}$

Closer to numerics:

$$\partial_\mu (T_{\text{id}}^{\mu\nu} + \delta T^{\mu\nu}) = S^\nu, \quad \text{S=geometrical source terms}$$

$$\underbrace{\partial_\tau (T_{\text{id}}^{\tau i} + \delta T^{\tau i})}_{Q_i} + \underbrace{\partial_j (T^{ji})}_{\text{id.flux}} + \underbrace{\partial_j (\delta T^{ji})}_{\text{visc.flux}} = \underbrace{S_{\text{id}}^\nu + \delta S^\nu}_{\text{source terms}}$$

Finite-volume realization:

$$\frac{1}{\Delta\tau} (Q_{\text{id}}^{n+1} + \delta Q^{n+1} - Q_{\text{id}}^n - \delta Q^n) + \frac{1}{\Delta x} (\Delta F_{\text{id}}^{n+1/2} + \Delta \delta F^{n+1/2}) = S_{\text{id}}^{n+1/2} + \delta S^{n+1/2}$$

now, a small trick:

$$\frac{1}{\Delta\tau} (Q_{\text{id}}^{n+1} + \delta Q^{n+1} - \underbrace{Q_{\text{id}}^{*n+1} + Q_{\text{id}}^{*n+1}}_{=0} - Q_{\text{id}}^n - \delta Q^n) + \frac{1}{\Delta x} (\Delta F_{\text{id}} + \Delta \delta F) = S_{\text{id}} + \delta S$$

Then, split the equation into two parts²:

$$\frac{1}{\Delta t} (Q_{\text{id}}^{*n+1} - Q_{\text{id}}^n) + \frac{1}{\Delta x} \Delta F_{\text{id}} = S_{\text{id}} \quad (\text{using finite volume, HLLE approx}) \quad (4)$$

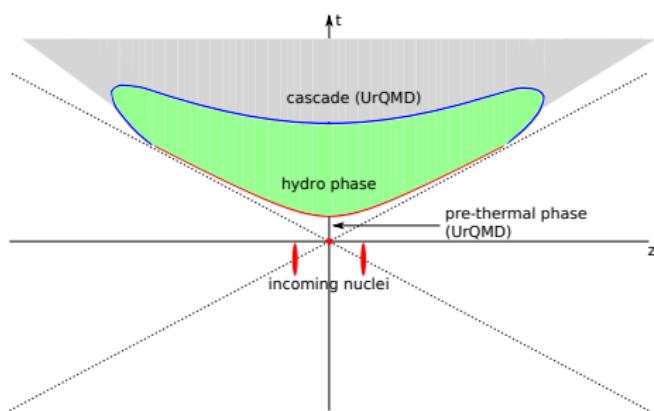
$$\frac{1}{\Delta t} (Q_{\text{id}}^{n+1} + \delta Q^{n+1} - Q_{\text{id}}^{*n+1} - \delta Q^n) + \frac{1}{\Delta x} \Delta \delta F = \delta S \quad (\text{upwind/Lax-Wendroff}) \quad (5)$$

²Makoto Takamoto, Shu-ichiro Inutsuka, J.Comput.Phys. 230 (2011), 7002

Fluid→particle transition

$\varepsilon = \varepsilon_{sw} = 0.5 \text{ GeV/fm}^3$ (blue curve):

$\{T^{0\mu}, N_b^0, N_q^0\}$ of hadron-resonance gas = $\{T^{0\mu}, N_b^0, N_q^0\}$ of fluid



▷ Space and momentum distribution from Cooper-Frye prescription:

$$p^0 \frac{d^3 n_i}{d^3 p} = \int (f_{i,\text{eq.}}(x, p) + \delta f(x, p)) p^\mu d\sigma_\mu$$

- ▷ Cornelius subroutine* is used to compute $\Delta\sigma_i$ on transition hypersurface.
- ▷ UrQMD cascade is employed after particlization surface.

*Huovinen P and Petersen H 2012, *Eur.Phys.J. A* **48** 171

Particle sampling

For each surface element:

$$\Delta N_i = \Delta\sigma_\mu u^\mu n_{i,\text{th}} = \Delta\sigma_0^* n_{i,\text{th}}$$

Momentum distribution in a fluid rest frame³:

$$\frac{d^3 N_i}{dp^* d(\cos\theta) d\phi} = \underbrace{\frac{\Delta\sigma_\mu^* p^{*\mu}}{p^{*0}}}_{W_{\text{residual}}} \underbrace{p^{*2} f_{\text{eq}}(p^{*0}; T, \mu_i)}_{\text{isotropic}} \underbrace{\left[1 + (1 \mp f_{\text{eq}}) \frac{p_\mu^* p_v^* \pi^{*\mu\nu}}{2T^2(\varepsilon + p)} \right]}_{W_{\text{visc}}} \quad (6)$$

Ideal case: $\max\left(\frac{\Delta\sigma_\mu^* p^{*\mu}}{p^{*0}}\right) = \max(\Delta\sigma_0^* + |\Delta\vec{\sigma}_i^*|)$

Momentum generation procedure:

- momentum sampling according to isotropic part of DF
- correction according to W_{residual} or $W_{\text{residual}} \cdot W_{\text{visc}}$
- Lorentz boost to global frame

³N.S. Amelin et al, Phys.Rev.C74:064901, 2006 (FASTMC event generator)

Finally, generated particle are fed to UrQMD cascade.

Model validation at $\sqrt{s} = 200$ GeV RHIC energy

Setup: smooth 3D initial conditions:

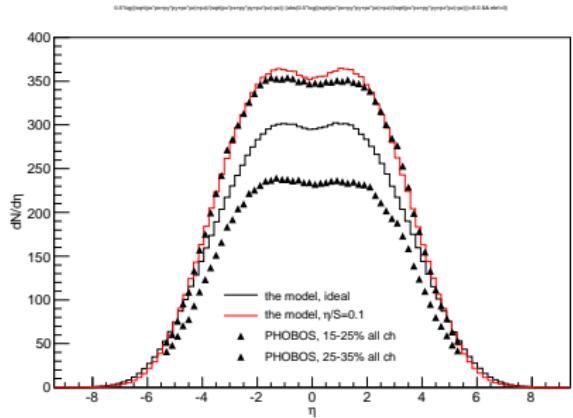
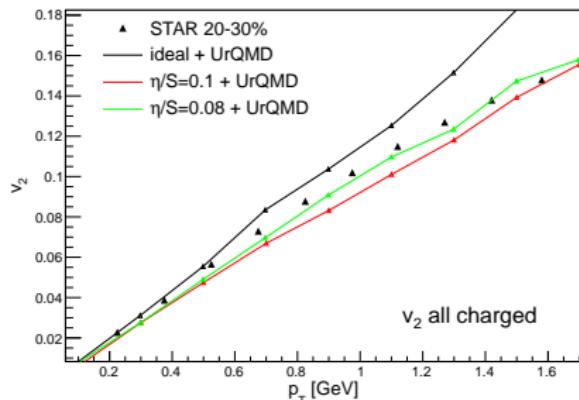
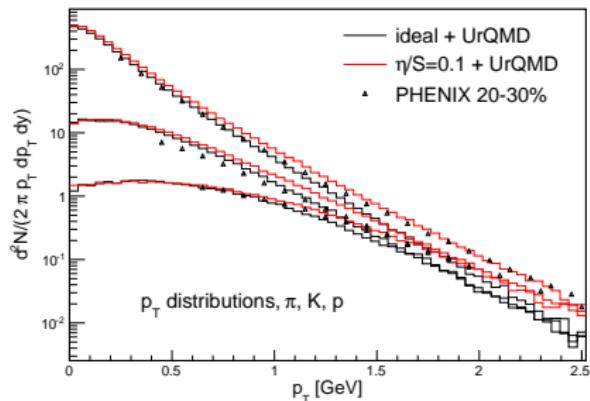
2D Monte Carlo Glauber + parametrized rapidity dependence

no UrQMD for initial state

$$\varepsilon(\tau_0, \vec{r}_T, \eta) = \varepsilon_{\text{MCG}}(\vec{r}_T) \cdot \theta(Y_b - |\eta|) \exp \left[-\theta(|\eta| - \Delta\eta) \frac{(|\eta| - \Delta\eta)^2}{\sigma_\eta^2} \right]$$

Y_b is beam rapidity, parameters: $\Delta\eta = 1.3$, $\sigma_\eta = 2.1$
(chosen from the fit to PHOBOS $dN_{\text{ch}}/d\eta$)

Model validation at $\sqrt{s} = 200$ GeV RHIC energy



experimental data (points), 15-25% and 25-35% central

ideal hydro+cascade (black curve), 20-30% central
viscous hydro+cascade (red curve), 20-30% central

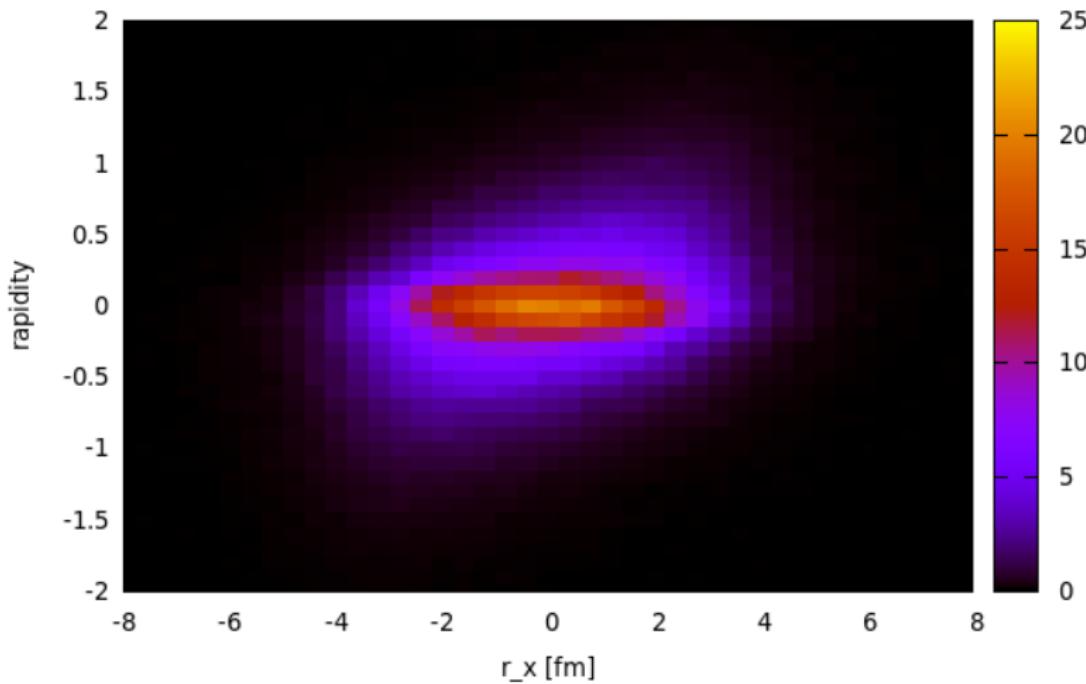
Beam energy scan

First round of simulations:

- single-shot hydro (1 hydro simulation for a given energy and centrality)
- smooth initial conditions taken as an average from many UrQMD initializations

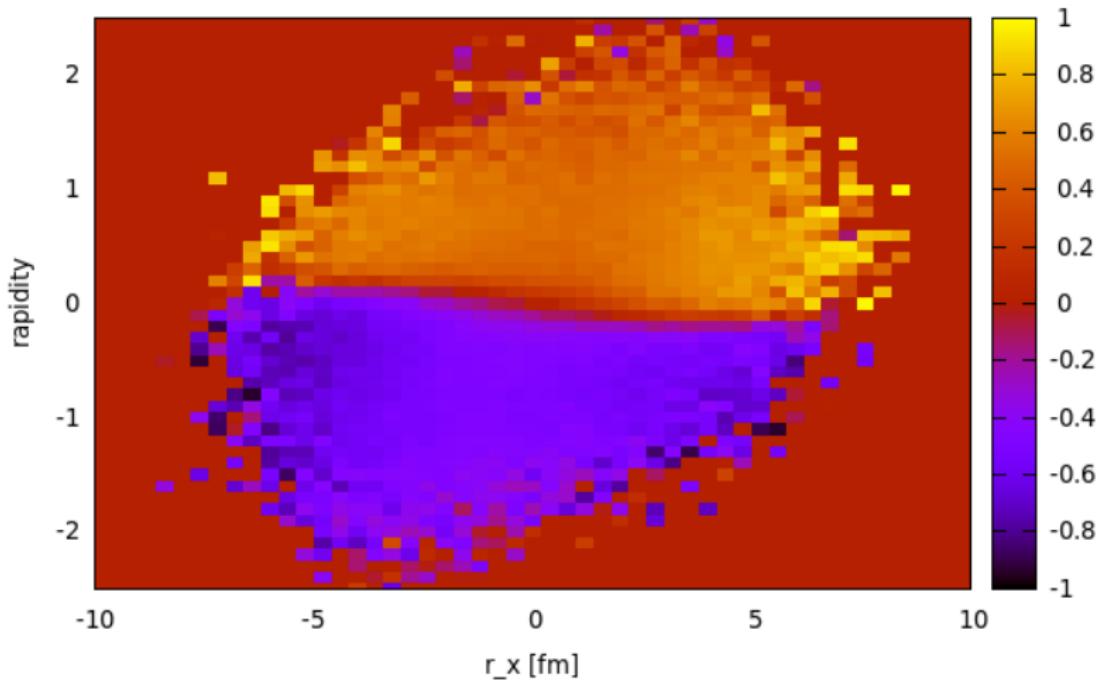
Typical smooth (event-averaged) initial condition for $E_{\text{lab}} = 168$ A GeV midcentral SPS collisions.

energy density [GeV/fm³] distribution:



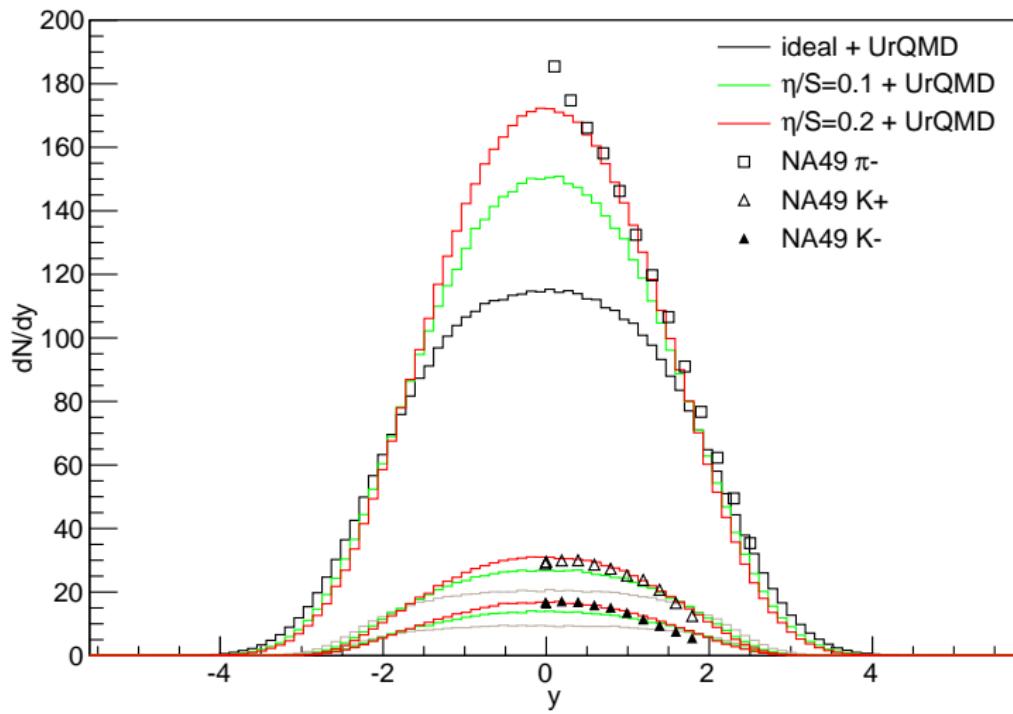
Typical smooth (event-averaged) initial condition for $E_{\text{lab}} = 168$ A GeV midcentral SPS collisions.

v_η distribution (notice nonzero angular momentum!):



Results: $E_{\text{lab}} = 158 \text{ A GeV Pb-Pb (SPS)}$

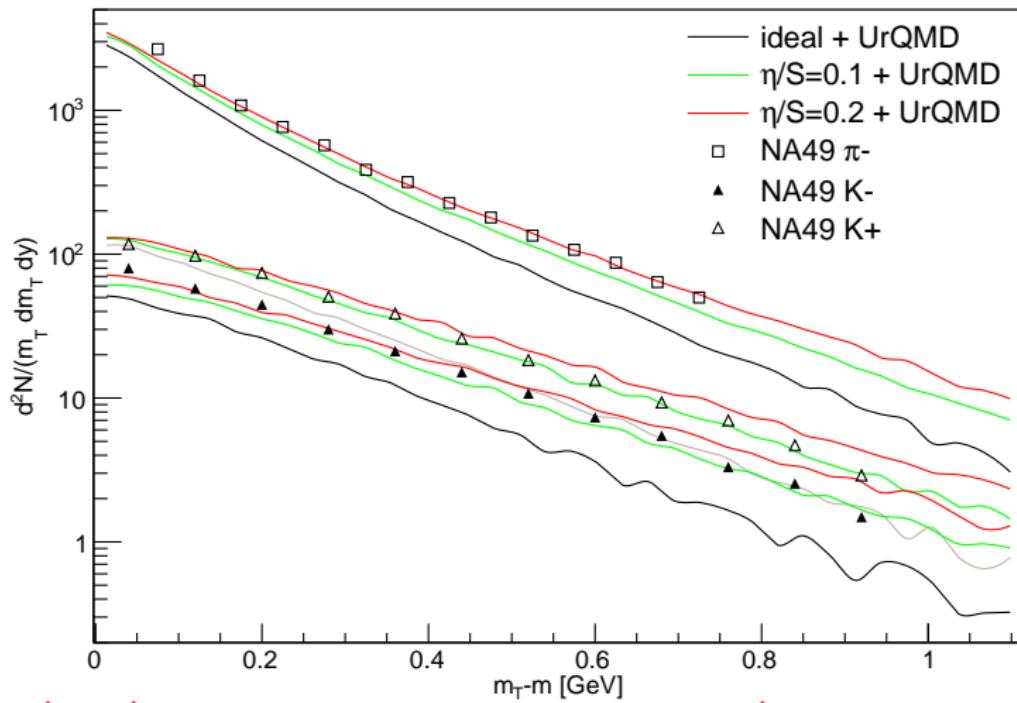
$\sqrt{s_{NN}} = 17.3 \text{ GeV}$, 0-5% central collisions ($b = 0 \dots 3.4 \text{ fm}$)



→ strong viscous entropy production

Results: $E_{\text{lab}} = 158 \text{ A GeV Pb-Pb (SPS)}$

$\sqrt{s_{NN}} = 17.3 \text{ GeV}, 0\text{-}5\% \text{ central collisions } (b = 0 \dots 3.4 \text{ fm})$

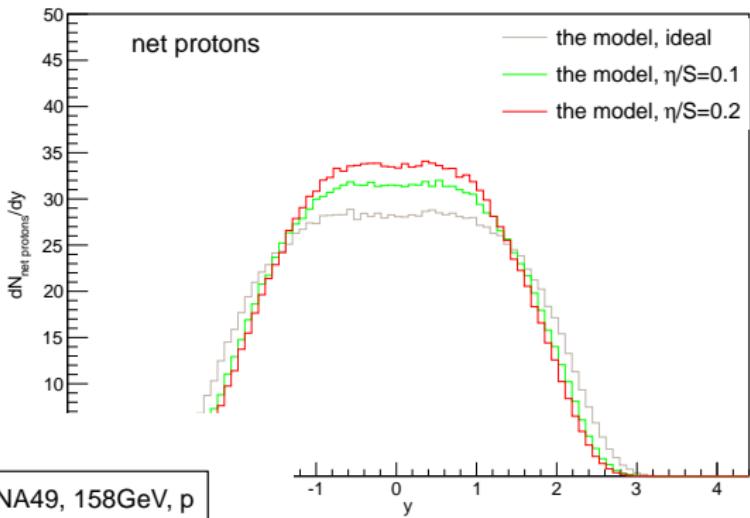
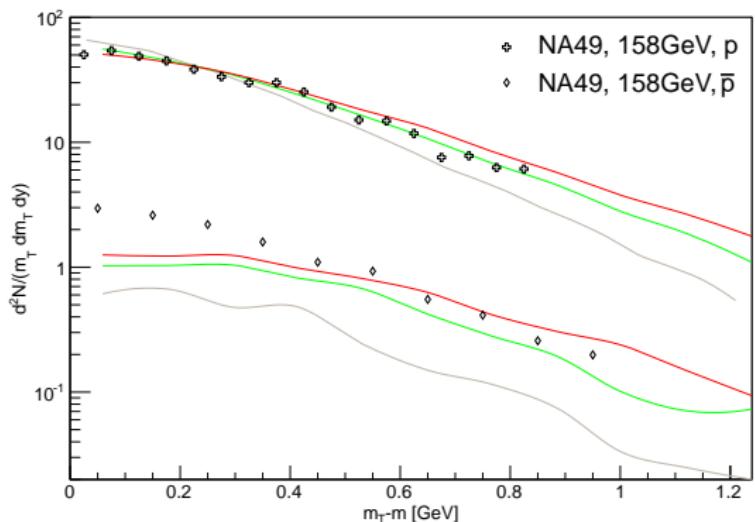


→ viscosity causes stronger transverse expansion

Results: 158 GeV SPS

protons & antiprotons

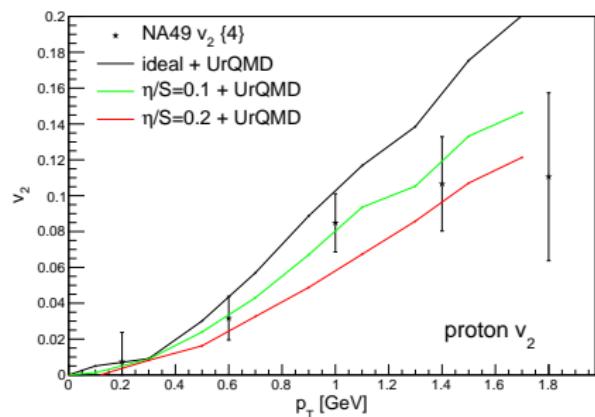
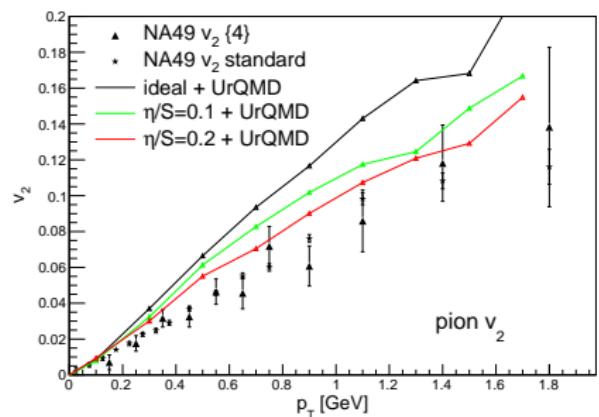
most central events
($b = 0..3.4$ fm)



Hydrodynamic
 $\tau_{\text{start}} = 1.42$ fm/c

Results: $E_{\text{lab}} = 158 \text{ A GeV Pb-Pb (SPS)}$

Mid-central events as defined by NA49 ($c = 12.5 - 33.5\%$)

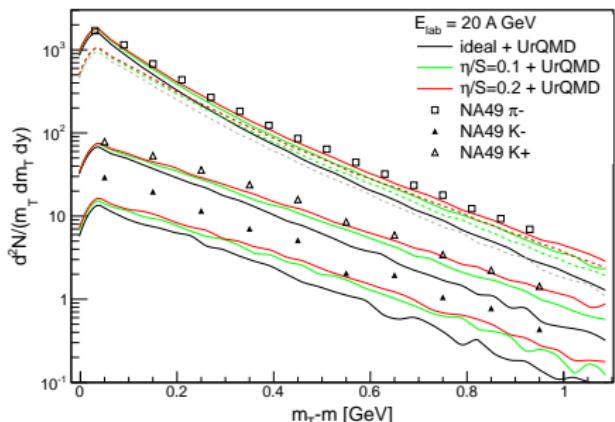
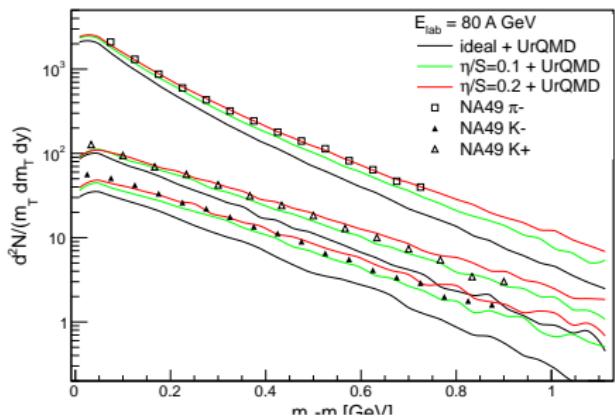
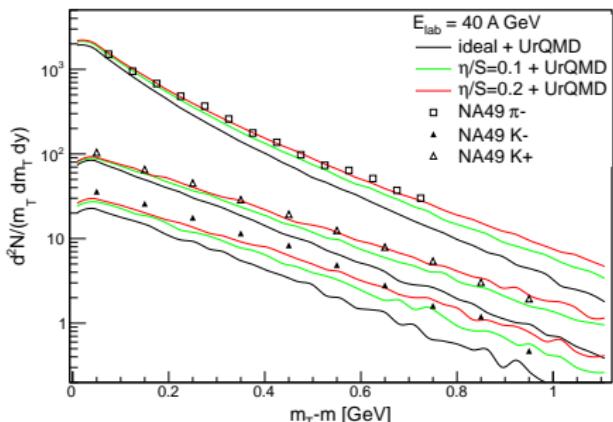


Results: $E_{\text{lab}} = 80, 40, 20 \text{ A GeV Pb-Pb (SPS)}$

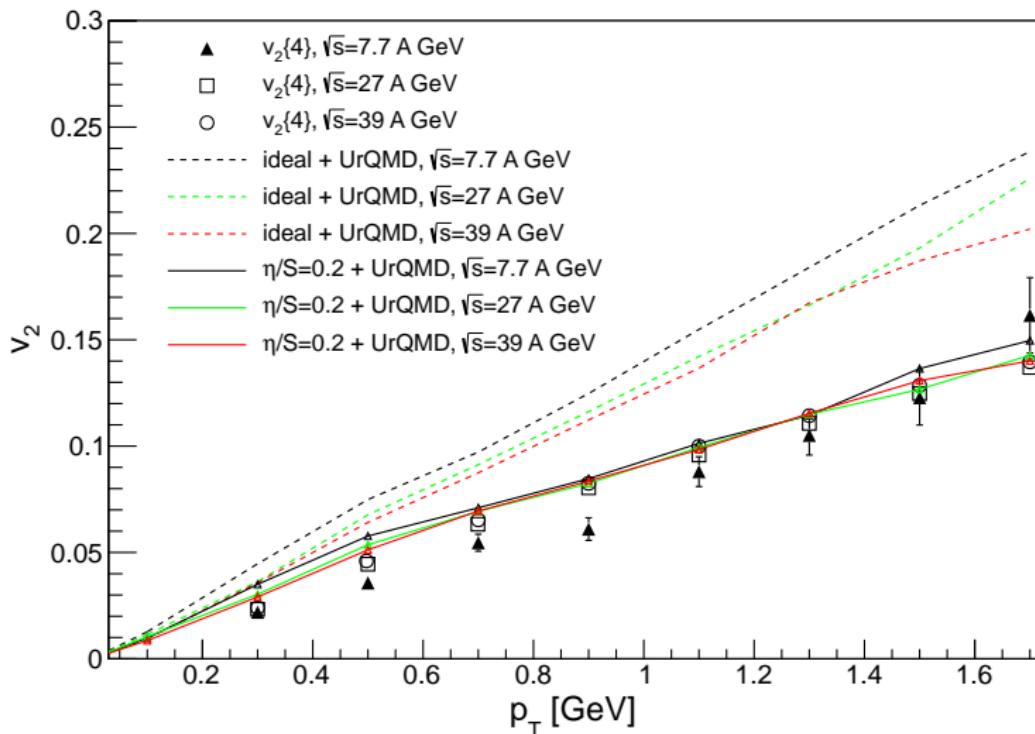
Corresp. $\sqrt{s_{NN}} = 12.3, 8.8, 6.3 \text{ GeV}$

Pion & kaon pt-distributions for most central events ($c = 0 - 5\%$, $b = 0 \dots 3.4 \text{ fm}$)

Overall good description with $\eta/S = 0.2$ except for K^- for lowest energies



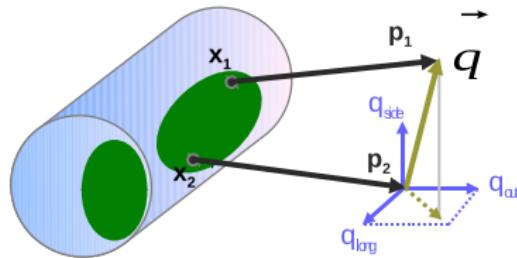
v_2 for BES at RHIC ($\sqrt{s_{NN}} = 7.7, 27, 39$ GeV Au-Au)



(20-30% central) $\eta/S \geq 0.2$ is required in hydro phase for all BES energies.

HBT(interferometry) measurements

The only tool for space-time measurements at the scales of 10^{-15}m , 10^{-23}s



$$\vec{q} = \vec{p}_2 - \vec{p}_1$$

$$\vec{k} = \frac{1}{2}(\vec{p}_1 + \vec{p}_2)$$

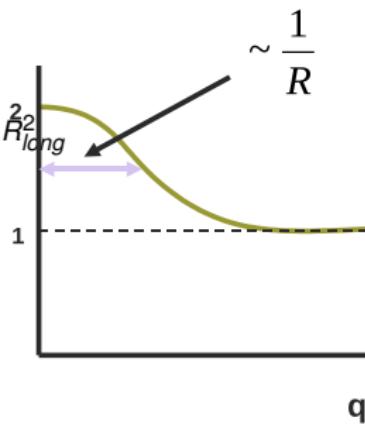
$$C(p_1, p_2) = \frac{P(p_1, p_2)}{P(p_1)P(p_2)} = \frac{\text{real event pairs}}{\text{mixed event pairs}}$$

Gaussian approximation of CFs ($q \rightarrow 0$):

$$C(\vec{k}, \vec{q}) = 1 + \lambda(k) e^{-q_{out}^2 R_{out}^2 - q_{side}^2 R_{side}^2 - q_{long}^2 R_{long}^2}$$

$R_{out}, R_{side}, R_{long}$ (HBT radii) correspond to *homogeneity lengths*, which reflect the space-time scales of emission process

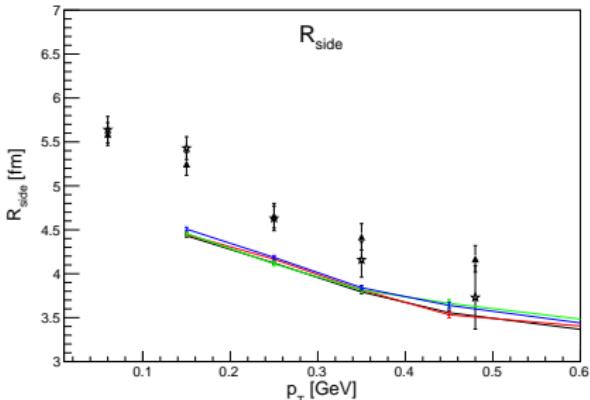
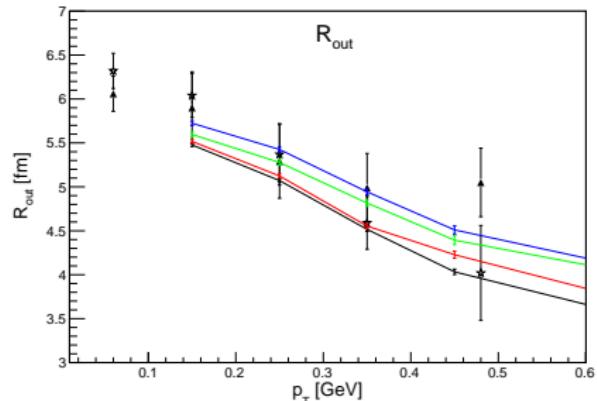
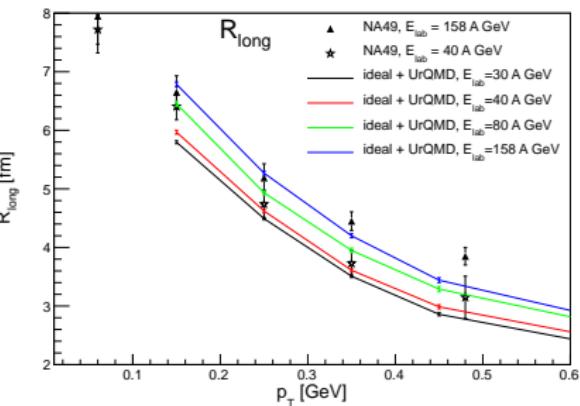
In an event generator, BE/FD two-particle amplitude (anti)symmetrization must be introduced



Femtoscopy at SPS energies

Corresponding $\sqrt{s_{NN}} = 12.3, 8.8, 6.3 \text{ GeV}$,
NA49, most central collisions ($c = 0 - 5\%$)

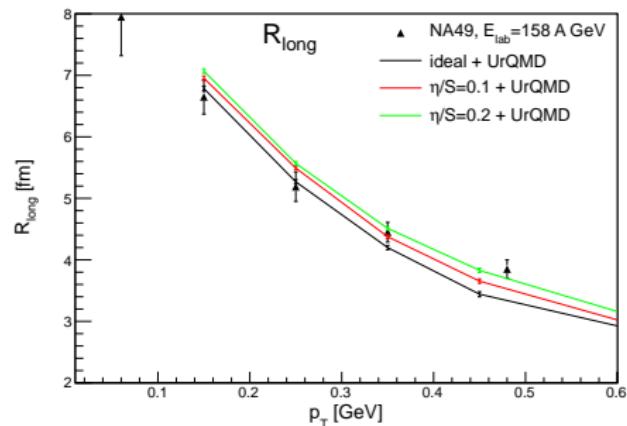
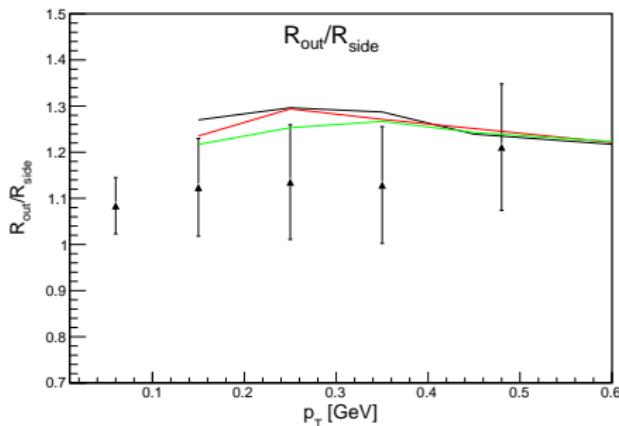
Femtoscopic radii for $\pi^-\pi^-$ pairs:
 $R_{\text{long}}, R_{\text{out}}$ consistent with NA49 data,
 R_{side} underestimated.



Femtoscopy at top SPS energy

$E_{\text{lab}} = 158 \text{ A GeV SPS}$ ($\sqrt{s_{NN}} = 17.3 \text{ GeV}$)

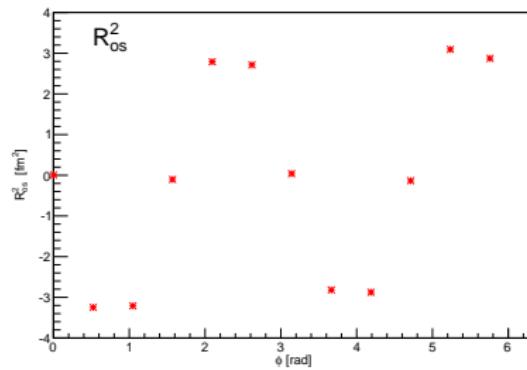
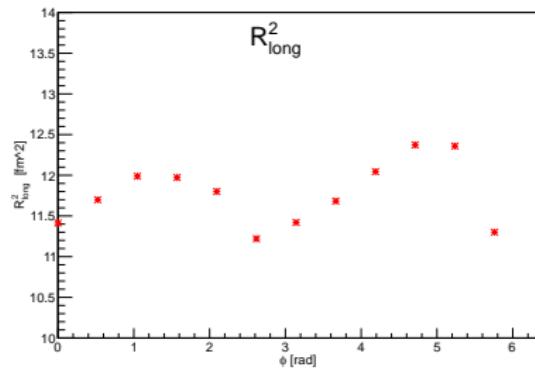
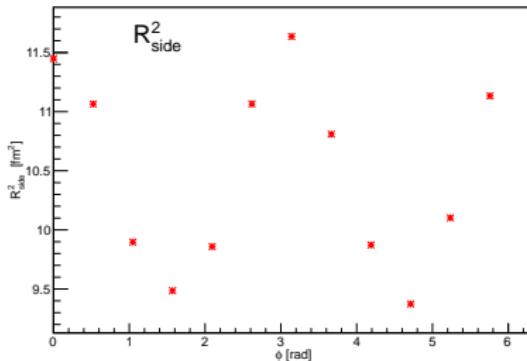
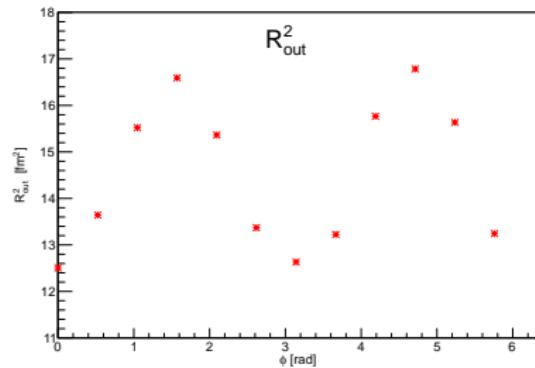
Dependence on η/S



R_{long} is increased and $R_{\text{out}}/R_{\text{side}}$ is slightly improved by viscosity

Azimuthally-sensitive femtoscopy

$\sqrt{s_{NN}} = 7.7 \text{ GeV}$, 10-30% central AuAu; $p_T = 0.15 \dots 0.6 \text{ GeV}$; $\phi = \psi_{\text{pair}} - \Psi_{\text{RP}}$



Azimuthally-sensitive femtoscopy

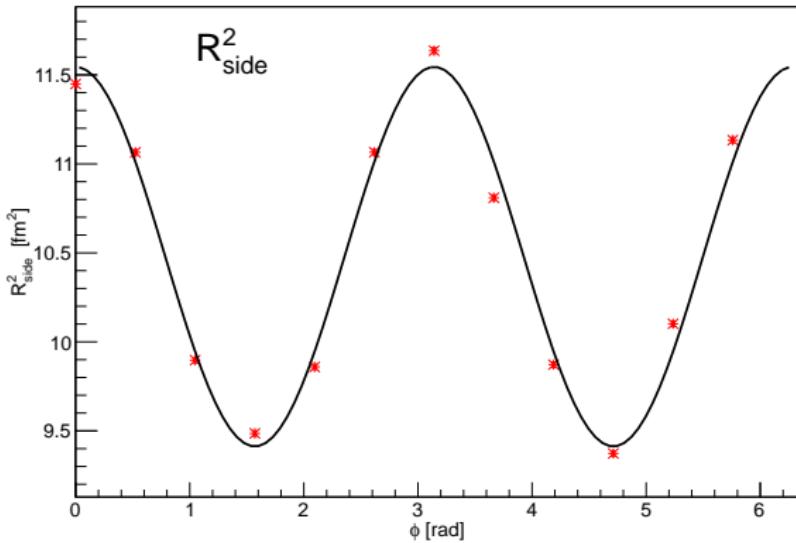
$$R_i^2(\phi) = R_{i,0}^2 + 2 \sum_{n=2,4,6,\dots} R_{i,n}^2 \cos(n\phi), \quad i = \text{out, side, long}$$

$$R_i^2(\phi) = 2 \sum_{n=2,4,6,\dots} R_{i,n}^2 \sin(n\phi), \quad i = \text{os}$$

solid curve:

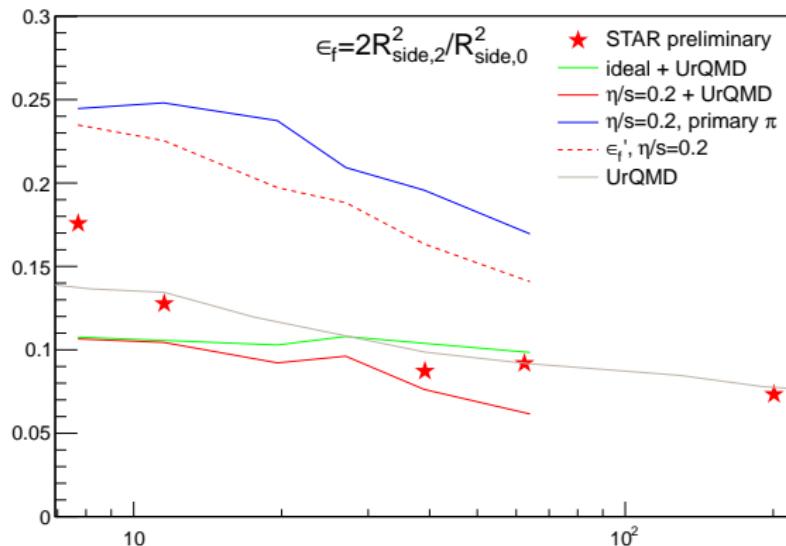
$$R_{s,0}^2 + 2R_{s,2}^2 \cos(2\phi) \Rightarrow$$

$$\varepsilon_f = 2 \frac{R_{\text{side},2}^2}{R_{\text{side},0}^2}$$



F. Retiere and M. Lisa, Phys.Rev. C70:044907, 2004

Azimuthally-sensitive femtoscopy



STAR: C. Anson,
J.Phys. G38:124148,2011

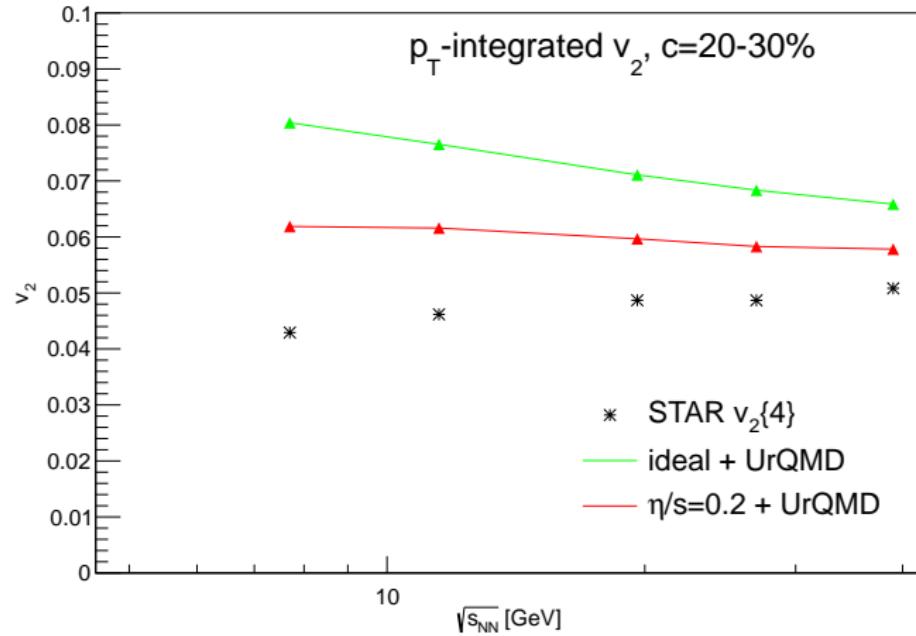
10-30% central AuAu,
 $p_T = 0.15\dots 0.6 \text{ GeV}$

$$\epsilon' = \frac{\int (y^2 - x^2) u^\mu d\sigma_\mu}{\int (y^2 + x^2) u^\mu d\sigma_\mu} - 1$$

- ¹ C. Shen, U. Heinz, Phys.Rev. C 85, 054902 (2012)
² UrQMD: M.A. Lisa, et al., New J.Phys.13:065006,2011

Rescatterings and
resonance decays
decrease the
eccentricity

...and p_T -integrated elliptic flow



Large v_2 for $\eta/s = 0$ in hydro phase: feature of ICs used?

Second round of simulations:

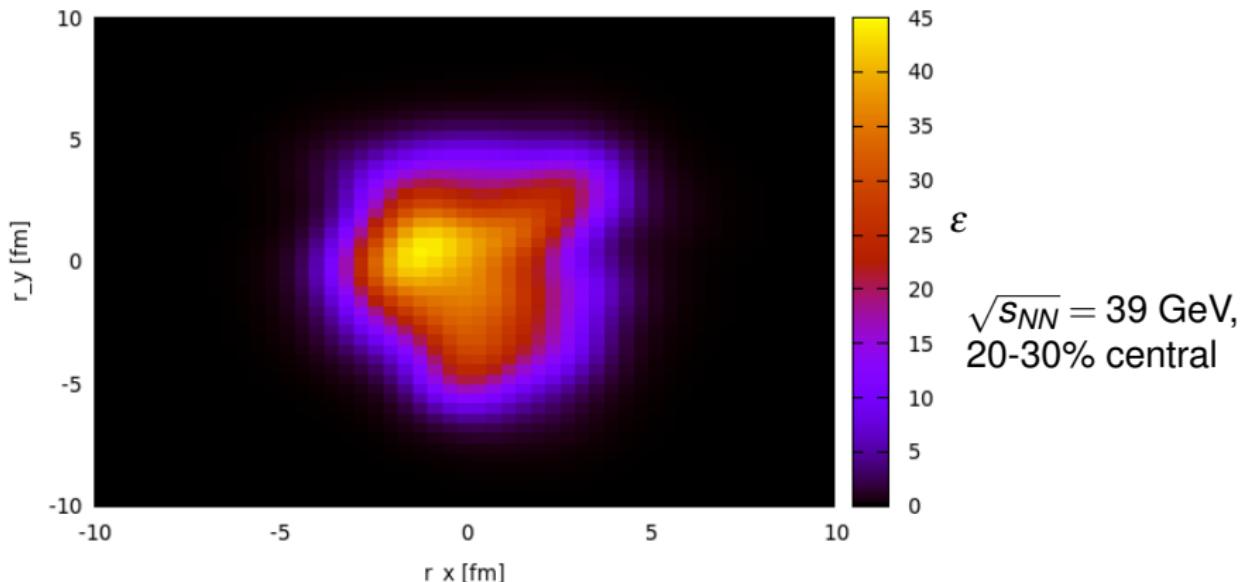
- event-by-event hydrodynamic evolution
- fluctuating initial conditions taken from single UrQMD initialization each

Fluctuating initial state

Fluctuating, but smoothed initial state:

$$E \propto \exp\left(-\frac{(x-x_{part})^2 + (y-y_{part})^2 + \gamma_z^2(z-z_{part})^2}{2R^2}\right), \text{ where } R = 1 \text{ fm}$$

see e.g. H. Petersen et al., Phys. Rev. C78 (2008) 044901



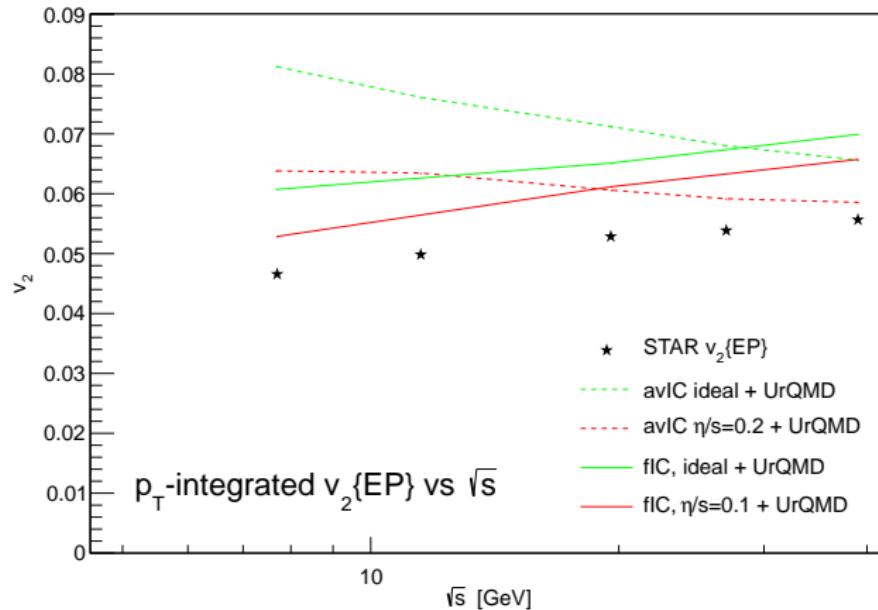
23% larger $dS/dy(y=0)$ than averaged IC.

Fluctuating vs. averaged initial state

Fluctuating, but smoothed initial state:

$$E \propto \exp\left(-\frac{(x-x_{part})^2 + (y-y_{part})^2 + \gamma_z^2(z-z_{part})^2}{2R^2}\right), \text{ where } R = 1 \text{ fm}$$

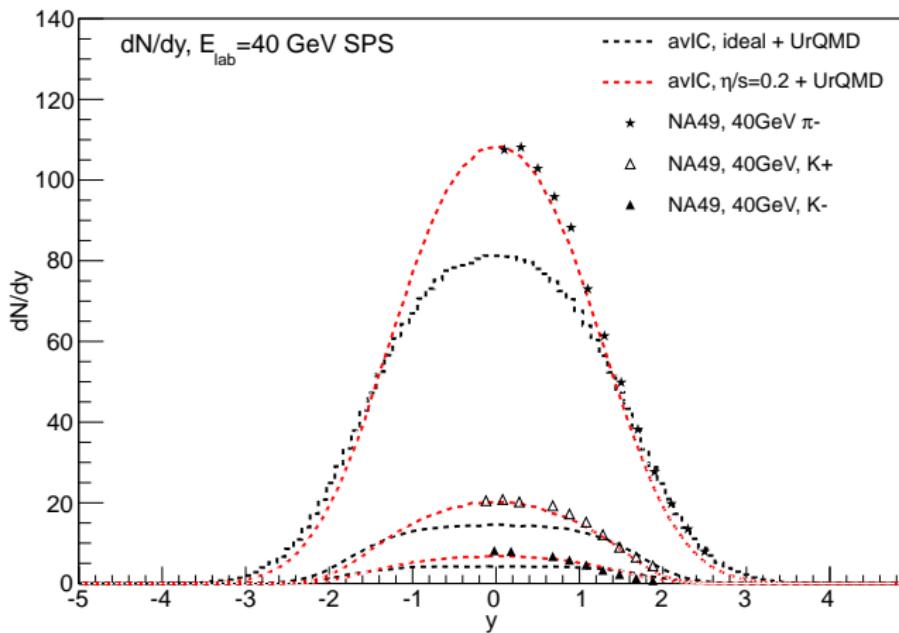
see e.g. H. Petersen et al., Phys.Rev. C78 (2008) 044901



$v_2(\sqrt{s})$ depends on how the initial state is constructed.

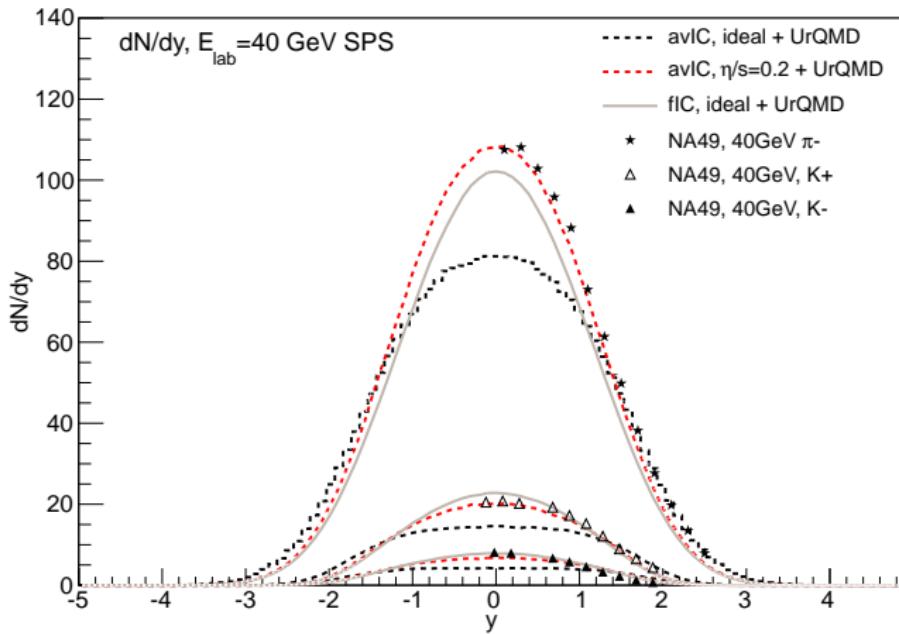
Influence on rapidity distribution

Results from single-shot hydro runs



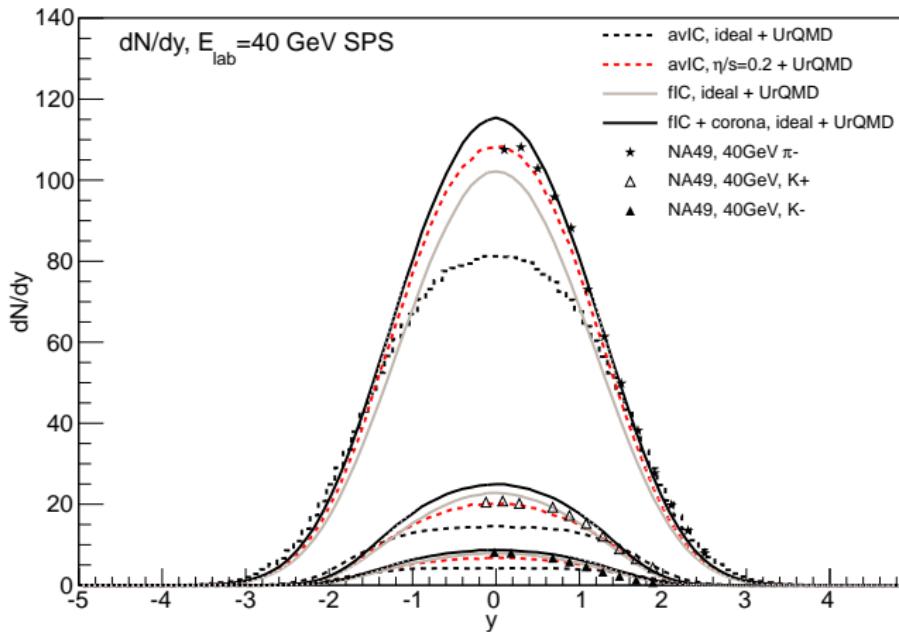
Influence on rapidity distribution

Single-shot hydro vs EbE ideal hydro



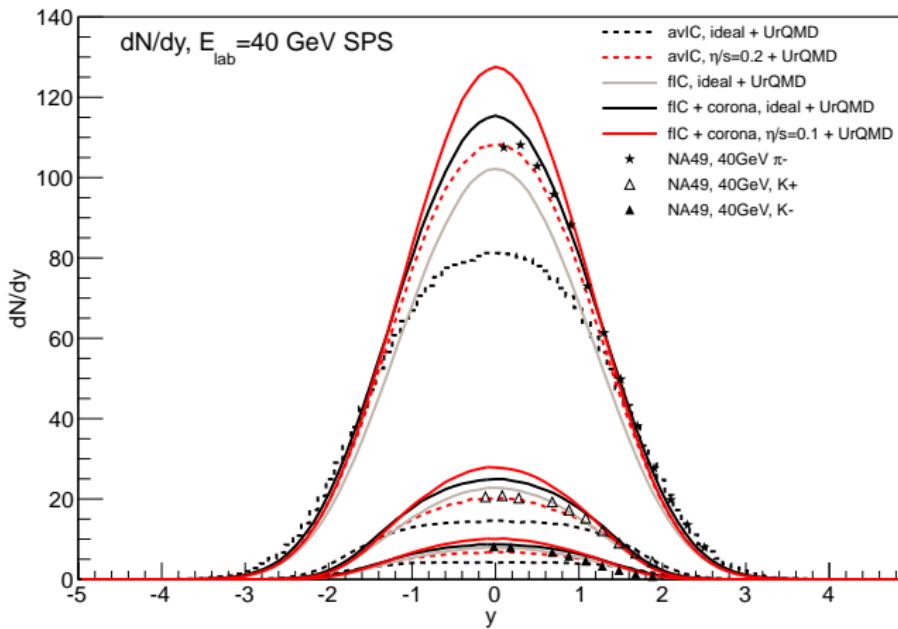
Influence on rapidity distribution

Single-shot hydro vs EbE id. hydro vs EbE id. hydro + corona



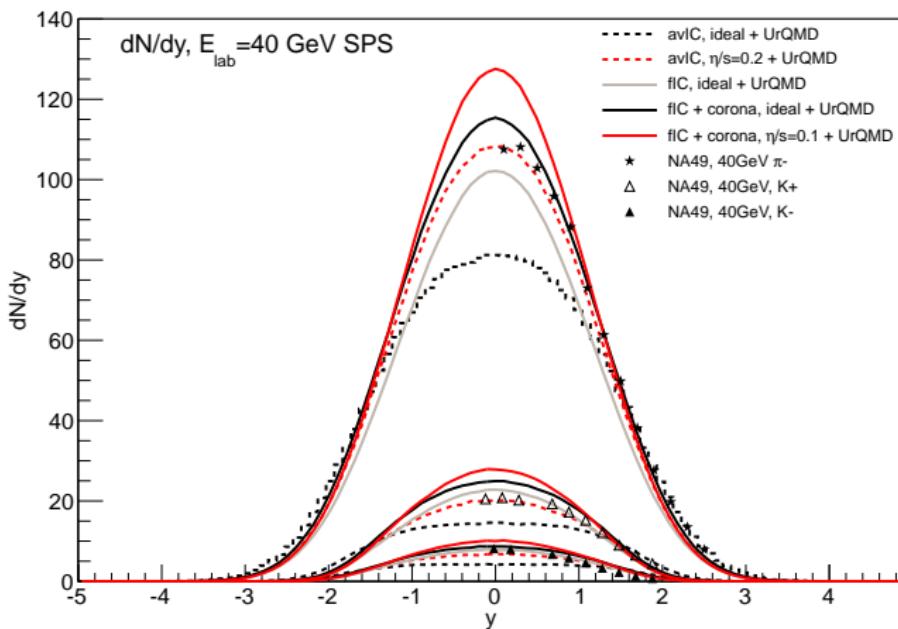
Influence on rapidity distribution

All that vs EoS viscous hydro + corona



Influence on rapidity distribution

All that vs EoS viscous hydro + corona



Too large initial entropy to accommodate viscous hydro phase



This can be regulated by decreasing Gaussian radius R .

Summary

UrQMD + 3D viscous hydro + UrQMD model including EoS at finite μ_B

Conclusions:

- model validated at top RHIC energy, and applied for Beam Energy Scan.
- averaged IC: and single-shot hydro
shear viscosity in hydro phase improves description of
 - ▶ p_T -spectra
 - ▶ dN/dy
 - ▶ elliptic flow
 - ▶ femtoscopic radii
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 - ▶ p_T -spectra
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- v_2 from RHIC BES suggests $\eta/S \geq 0.2$
- EbE hydro with fluctuating IC corrects $v_2(\sqrt{s})$
- in EbE case the a simultaneous fit to dN/dy and v_2 should allow to fix both initial state granularity and shear viscosity in hydro phase (provided that the initial state is from UrQMD).

Outlook:

- Effects of EoS with 1st order PT?

Thank you for your attention!

Extra slides

Viscous hydrodynamic equations

The hydrodynamic equations in arbitrary coordinate system:

$$\partial_{;\nu} T^{\mu\nu} = \partial_\nu T^{\mu\nu} + \Gamma^\mu_{\nu\lambda} T^{\nu\lambda} + \Gamma^\nu_{\nu\lambda} T^{\mu\lambda} = 0 \quad (7)$$

where (we choose Landau definition of velocity)

$$T^{\mu\nu} = \epsilon u^\mu u^\nu - (p + \Pi)(g^{\mu\nu} - u^\mu u^\nu) + \pi^{\mu\nu} \quad (8)$$

and $\Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$

Evolutionary equations for shear/bulk, coming from **Israel-Stewart** formalism:

$$\langle u^\gamma \partial_{;\gamma} \pi^{\mu\nu} \rangle = -\frac{\pi^{\mu\nu} - \pi_{\text{NS}}^{\mu\nu}}{\tau_\pi} - \frac{4}{3} \pi^{\mu\nu} \partial_{;\gamma} u^\gamma \quad (9a)$$

$$u^\gamma \partial_{;\gamma} \Pi = -\frac{\Pi - \Pi_{\text{NS}}}{\tau_\Pi} - \frac{4}{3} \Pi \partial_{;\gamma} u^\gamma \quad (9b)$$

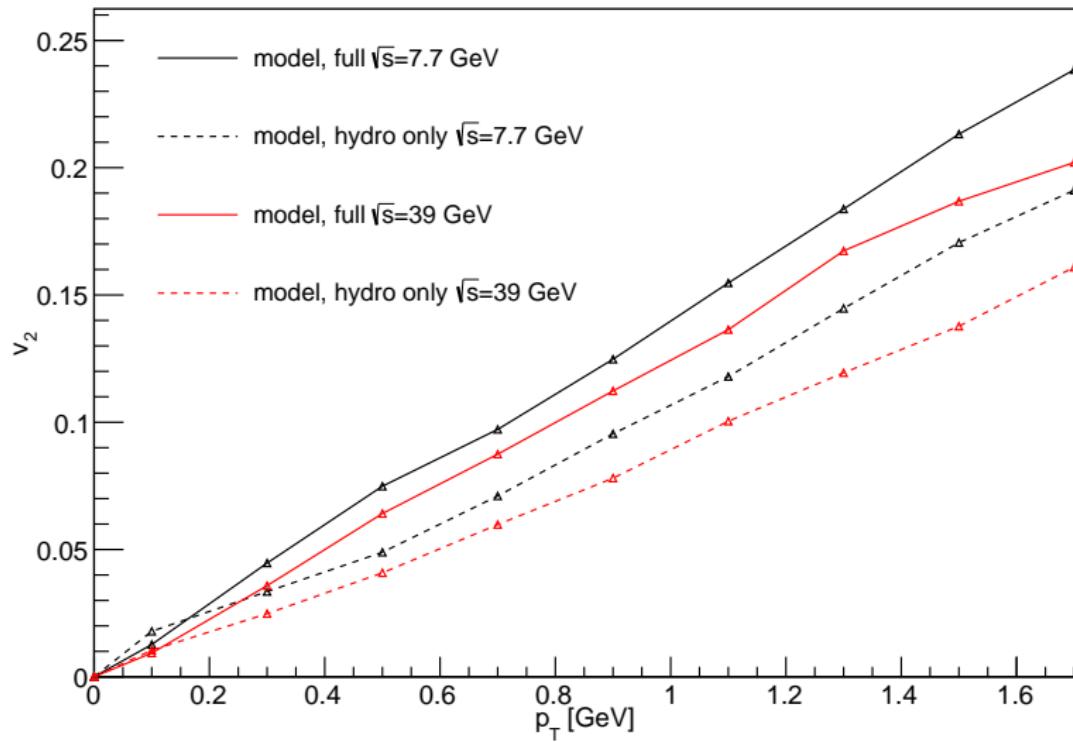
where

$$\langle A^{\mu\nu} \rangle = \left(\frac{1}{2} \Delta_\alpha^\mu \Delta_\beta^\nu + \frac{1}{2} \Delta_\alpha^\nu \Delta_\beta^\mu - \frac{1}{3} \Delta^{\mu\nu} \Delta_{\alpha\beta} \right) A^{\alpha\beta}$$

v_2 before and after the cascade

$\eta/S = 0$

full vs hydro_only



Transition surfaces

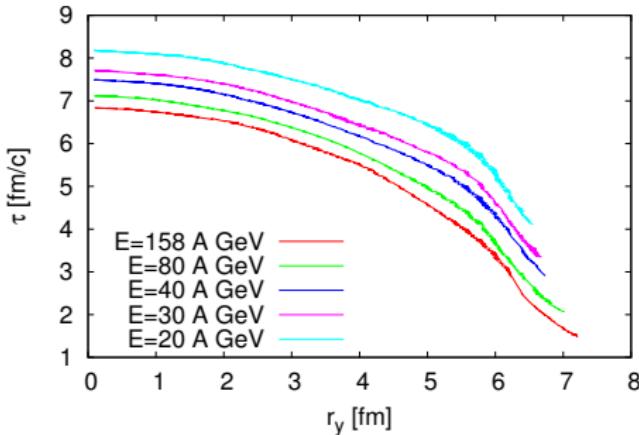
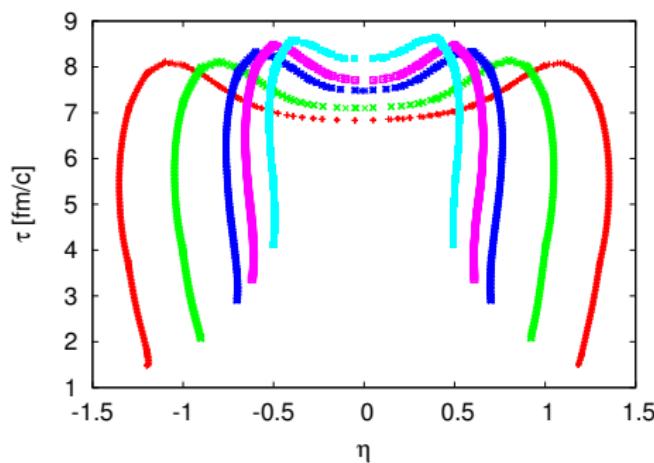
hydro→cascade transition

Most central collisions,

$E_{\text{lab}} = 20 \text{ GeV}$ (cyan) ... 158 GeV (red)

$\sqrt{s_{\text{NN}}} = 6.27 \dots 17.3 \text{ GeV}$

Transition criterion: $\varepsilon = \varepsilon_{\text{crit}} = 0.5 \text{ GeV/fm}^3$,
same for all energies



System squeezes in rapidity with decreasing collision energy,
hydro phase still lasts about 4.5 fm/c at lowest SPS energy.

Thermodynamics on transition surface

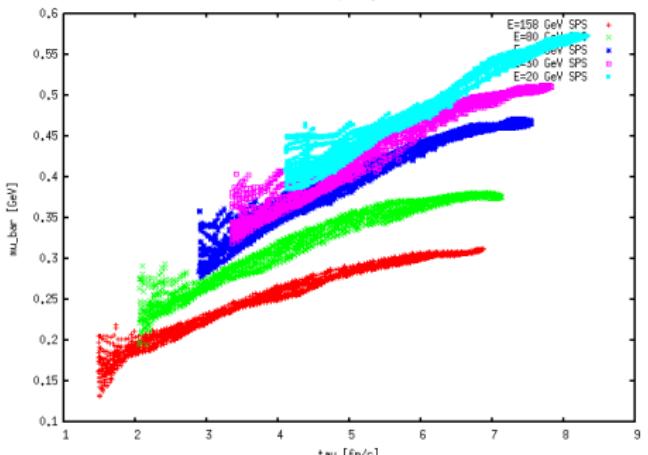
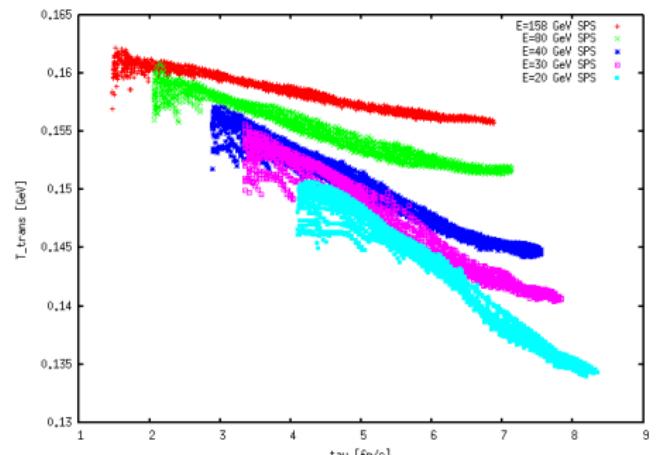
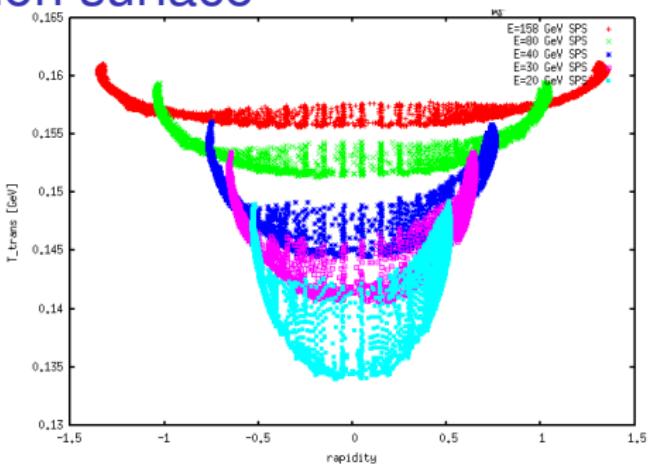
Procedure (for each surface element):

$$\{\varepsilon = \varepsilon_{\text{crit}}, n_B, n_Q\} \xrightarrow{\text{EoS}} \{T, \mu_B, \mu_Q, \mu_S\}$$

Most central collisions,

$E_{\text{lab}} = 20 \text{ GeV}$ (cyan) ... 158 GeV (red)

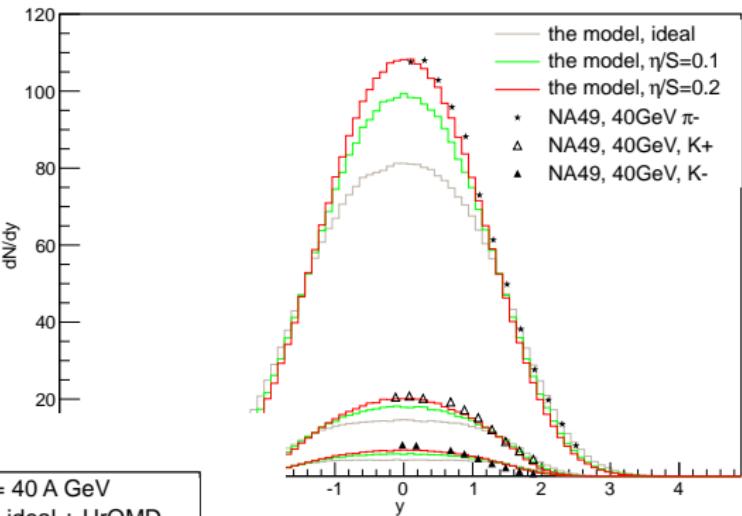
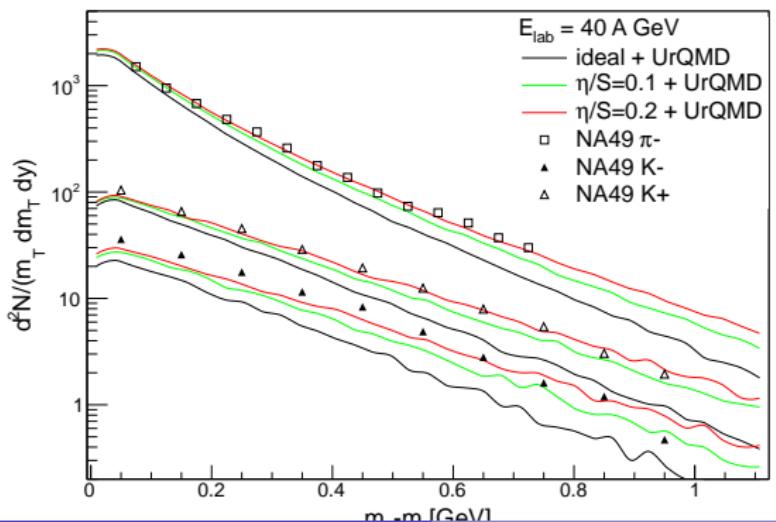
T (rapidity) (top), $T(\tau)$ (bottom left),
 $\mu_B(\tau)$ (bottom right)



Results: 40 GeV SPS

pions & kaons

most central events
($b = 0..3.4$ fm)

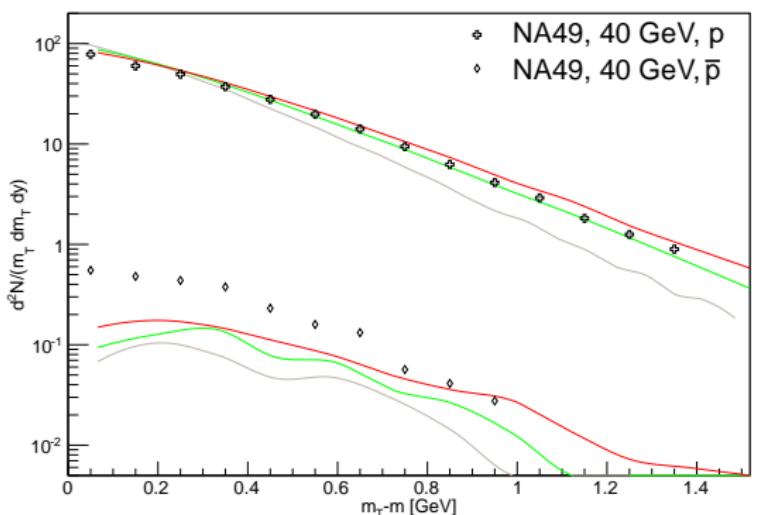
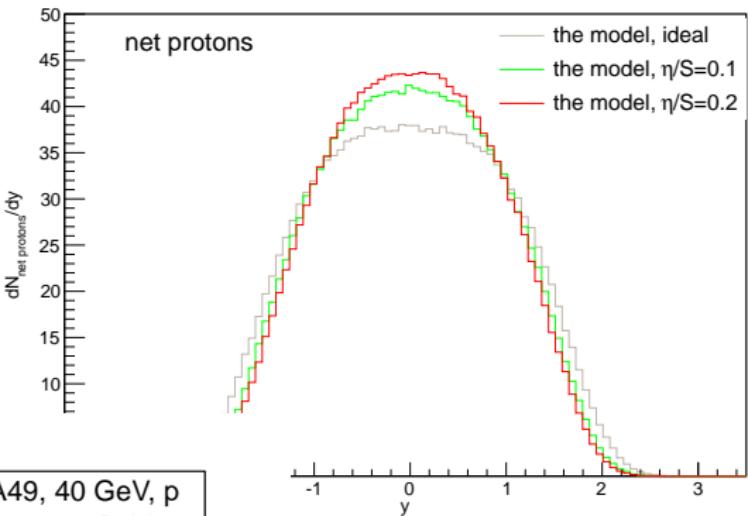


Hydrodynamic
 $\tau_{start} = 2.83$ fm/c

Results: 40 GeV SPS

protons & antiprotons

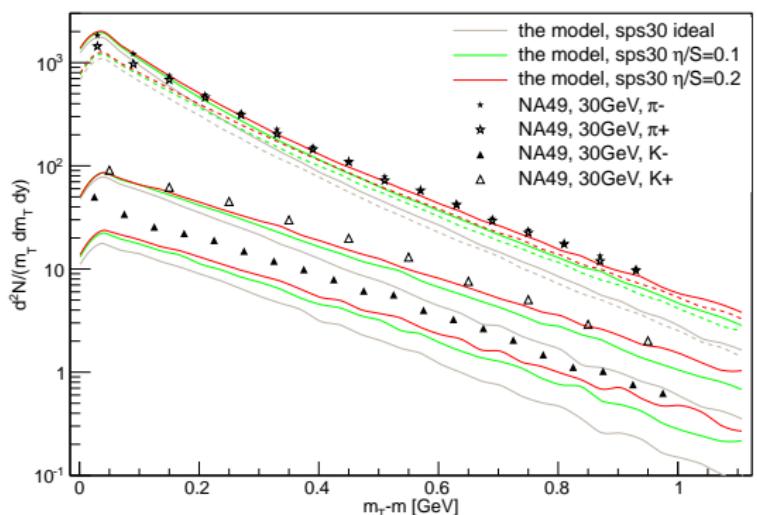
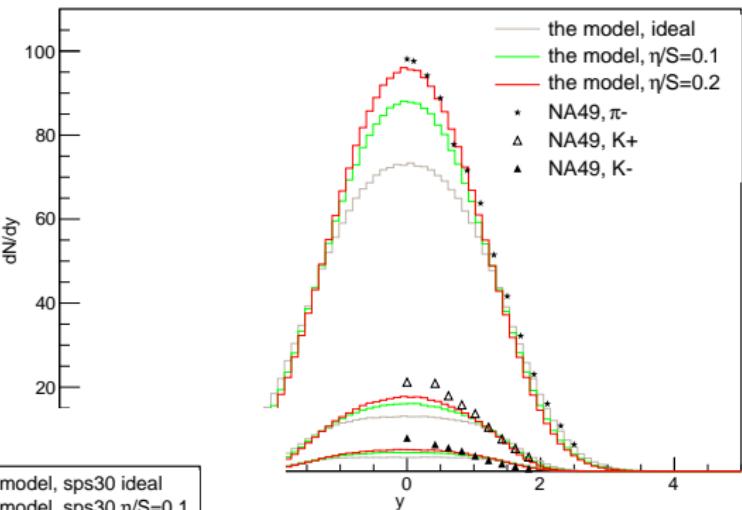
most central events
 $(b = 0..3.4 \text{ fm})$



Results: 30 GeV SPS

pions & kaons

most central events
($b = 0..3.4$ fm)

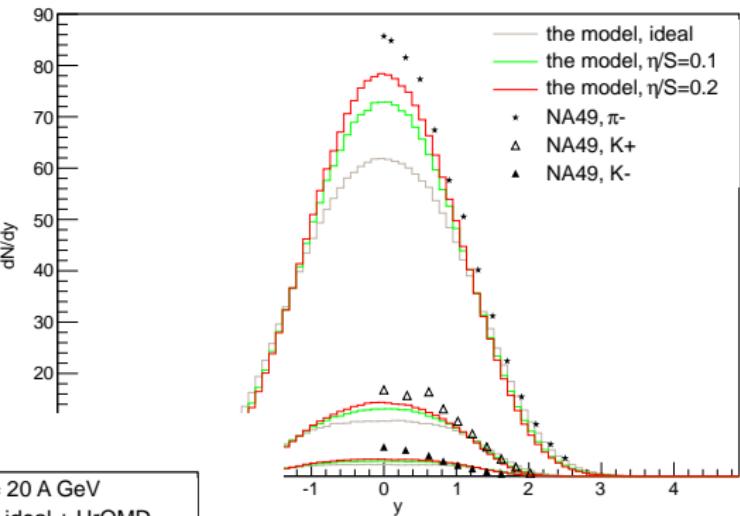
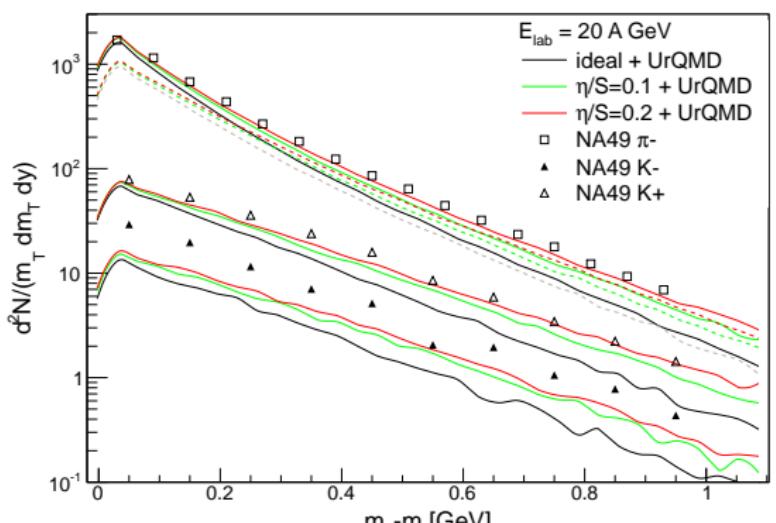


Hydrodynamic
 $\tau_{\text{start}} = 3.28$ fm/c

Results: 20 GeV SPS

pions & kaons

most central events
($b = 0..3.4$ fm)



Hydrodynamic
 $\tau_{\text{start}} = 4.05$ fm/c