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MFD Models

2-Fluid Models 3-Fluid Models 3FD Friction Freeze-out

Phys. Input

Results

Baryon Stopping Part. Production Hadron Ratios pt spectra Flow

Summary

Heavy-Ion Collisions within Multi-Fluid Simulations: Scenarios with and without Deconfinement Transition

Multi-Fluid Dynamics

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Outline

3FD model

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2-Fluid Mode 3-Fluid Mode 3FD Friction

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Summary

Multi-Fluid Models:

minimal extension of hydrodynamics required by heavy-ion dynamics

- 2-Fluid Models
- 3-Fluid Models
- 3-Fluid Dynamics (3FD), present version
- Friction
- Freeze-out
- Physical Input of the Model
- Results
 - Baryon Stopping
 - Part. Production
 - Hadron Ratios
 - opt spectra
 - Flow



Exploring Nuclear Phase Diagram

3FD model





3FD Friction Freeze-out

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Summary



However, nonequilibrium prevents direct application of Hydrodynamics



Hydrodynamics versus Kinetics

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Summary

Why we are not satisfied with kinetics?

• In practice, kinetics \Rightarrow only binary collisions mean free path $\lambda \approx 1/(n_B \sigma)$ if $\sigma \approx 4 \text{ fm}^2$ and $n_B \approx 5n_0 \Rightarrow \lambda \approx 0.3 \text{ fm} \sim \text{nucleon core}$ $(n_0 = 0.15 \text{ fm}^{-3} = \text{normal nuclear density})$

Approximation of binary collisions is bad!

• Phase transition into QGP is inaccessible in kinetics as a rule

Two exceptions based on simple combinatorics of quarks: A Multi-Phase Transport (AMPT) model [Lin, Ko and Pal, PRL 89, 152301 (2002)] Parton-Hadron-String Dynamics [Cassing, Bratkovskaya, arXiv:0907.5331 (2009)]

Hydrodynamics

- takes into account any multi-particle interactions
- directly addresses Equation of State (EoS)!
- Phase transition in QGP is accessible through EoS
- However, there are certain problems



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2-Fluid Models

Summarv



- Distributions are separated in momentum space ⇒ different fluids
- Leading particles carry baryon charge
 - \Rightarrow 2 baryon-rich fluids: projectile-like and target-like
- Los Alamos 1978–1986: Amsden, Harlow, Nix, Clare, Strottman 2-fluid hydrodynamics
- Kurchatov Inst. 1988–1991: Mishustin, Russkikh, and Satarov 2-fluid hydro with free-streaming radiation of pions
- GSI 1991–1997: Iv., Russkikh, Nörenberg 2-fluid mean-field hydro with hadrochemistry for SIS energies



From 2 Fluids to 3 Fluids

3FD model

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- Summary

- Distributions are separated in momentum space
 - ⇒ different fluids
- Leading particles carry baryon charge
 - \Rightarrow 2 baryon-rich fluids: projectile-like and target-like
- At high incident energies ($E_{lab} \gtrsim 10A \, {\rm GeV}$)
- Produced particles populate mid-rapidity
 ⇒ fireball fluid



momentum along beam



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MFD Models

2-Fluid Models 3-Fluid Models 3FD Friction

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Results

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Summary

 Kurchatov Inst. 1988–1991:
 2-fluid hydro with free-streaming radiation of pions Mishustin, Russkikh, and Satarov

 Frankfurt University 1993–2000:
 3-fluid hydrodynamics with instant formation of fireball Brachmann, Katscher, Dumitru, Rischke, Maruhn, Stöcker, Greiner, Mishustin, Satarov, et al.

 GSI 2003–now:
 3-fluid hydrodynamics with delayed formation of fireball Iv., Russkikh, Toneev

3-Fluid Dynamics, present version



Total energy-momentum conservation:

 $\partial_{\mu}(T^{\mu\nu}_{\rho}+T^{\mu\nu}_{t}+T^{\mu\nu}_{f})=0$



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Baryon current:

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$J^{\mu}_{\alpha} = n_{\alpha} u^{\mu}_{\alpha}$

 n_{α} = baryon density of α -fluid

 u^{μ}_{α} = 4-velocity of α -fluid

Energy-momentum tensor:

 $T^{\mu\nu}_{\alpha} = (\varepsilon_{\alpha} + P_{\alpha})u^{\mu}_{\alpha}u^{\nu}_{\alpha} - g_{\mu\nu}P_{\alpha}$ ε_{α} = energy density P_{α} = pressure

Part. Production

Summary

+ Equation of state:

 $P = P(n,\varepsilon)$

Final Aim: To find a proper EoS, which reproduces all data



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$$F_{\alpha}^{\nu} = \rho_{p}^{\xi} \rho_{t}^{\xi} \left[\left(u_{\alpha}^{\nu} - u_{\bar{\alpha}}^{\nu} \right) D_{P} + \left(u_{p}^{\nu} + u_{t}^{\nu} \right) D_{E} \right],$$

 $s_{\rm pt} = m_{\rm AI}^2 (u_{\rm p}^{\nu} + u_{\rm t}^{\nu})^2$, $\alpha = p \text{ or t}, \bar{p} = t \text{ and } \bar{t} = p, \text{ and } \rho_{\alpha}^{\xi}$ is

$$\rho_{\alpha}^{\xi}(\boldsymbol{s}_{pt}) = \left(\rho_{\alpha}^{bar.} + \frac{2}{3}\rho_{\alpha}^{mes.}\right) \boldsymbol{\xi}_{h}(\boldsymbol{s}_{pt}) + \frac{1}{3} \left(\rho_{\alpha}^{q} + \rho_{\alpha}^{g}\right) \boldsymbol{\xi}_{q}(\boldsymbol{s}_{pt}),$$

Friction

Part. Production

Summarv

 $\rho_{\alpha}^{\text{bar.}}, \rho_{\alpha}^{\text{mes.}}, \rho_{\alpha}^{q}$ and ρ_{α}^{g} are scalar densities of all baryons, all mesons, guarks and gluons, respectively,

$$\rho_{\alpha}^{a}(x) = m_{a} \int \frac{d^{3}p}{p_{0}} f_{\alpha}^{a}(x,p)$$

in terms of where equilibrium distribution function f_{α}^{a} . Factors like 2/3 and 1/3 are from naive valence-quark count.

D_P and D_F were estimated based on proton-proton cross sections [Satarov, 1990]

In view of uncertainties of the estimated friction, tuning factors are introduced for hadronic and quark-gluon phases: ξ_h and ξ_q , respectively.



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When system becomes dilute, we have to stop hydro

Freeze-out

Spectra of observed particles

$$E\frac{dN}{d^{3}\rho} = \sum_{\alpha} \int dV \left(u^{\mu}_{\alpha} \rho^{\mu} \right) \frac{\text{degeneracy factor}}{\exp[\left(u^{\mu}_{\alpha} \rho^{\mu} - \mu_{\alpha} \right) T_{\alpha}] \pm 1}$$

where $u^{\mu}_{\alpha}, \mu_{\alpha}, \mathcal{T}_{\alpha}$ are taken at freeze-out instant

In-plain evolution of energy density





Physical Input I

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Summary

I. Equation of State

 Hadronic EoS Galitsky&Mishustin (1979)

- 1st-order transition to QGP (2-phase EoS*)
- crossover EoS*

*[Khvorostukhin, Skokov, Redlich, Toneev, (2006)]



Phase transition \Longrightarrow EoS softening



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II. Friction was fitted to reproduce the baryon stoppingHadronic EoS

$$\xi_h^2(s) = 1 + \left[\ln \left(\frac{s^{1/2}}{2m_N} \right) \right]^{1/4}, \quad \tau = 2 \text{ fm/c}, \quad \varepsilon_{\text{frz}} = 0.4 \text{ GeV/fm}^3.$$

Friction is enhanced to reproduce the baryon stopping at $\sqrt{s_{NN}} > 5$ GeV. Though the enhancement looks too high: $\xi_h^2 = 2.2$ at $\sqrt{s_{NN}} = 17.3$ GeV. • 2-phase EoS

$$\xi_h^2(s) = 1, \quad \xi_q^2(s) = 60 \; rac{4m_N^2}{s}, \quad au = 0.17 \; {
m fm/c}, \quad arepsilon_{
m frz} = 0.4 \; {
m GeV/fm^3}.$$

crossover EoS

 $\xi_h^2(s) = 1, \quad \xi_q^2(s) = 200 \; rac{4m_N^2}{s}, \quad au = 0.17 \; {
m fm/c}, \quad arepsilon_{
m frz} = 0.4 \; {
m GeV/fm^3}.$

 $\xi_h^2(s) = 1$ is advantage of deconfinement scenarios





Crossover transition by Toneev et al. is too smooth

Lattice QCD predicts a fast crossover.

Therefore, a true EoS is somewhere in between the "Toneev"-crossover and "Toneev"-2-phase EoS's.

Onset of deconfinement happens at top-AGS-low-SPS energies.



3ED

Net-Proton Rapidity distributions



 Predictions of different scenarios differ to the largest extent in the energy region 8A GeV $< E_{lab} <$ 40A GeV.

"peak-dip-peak-dip" irregularity at midrapidity



Reduced Curvature at Midrapidity

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Summary

To quantify the "peak-dip-peak-dip" irregularity, net-proton rapidity distributions are fitted by $(a, y_s \text{ and } w_s \text{ are parameters of the fit})$

$$\frac{dN}{dy} = a \left(\exp \left\{ -(1/w_s) \cosh(y - y_{cm} - y_s) \right\} + \exp \left\{ -(1/w_s) \cosh(y - y_{cm} + y_s) \right\} \right)$$

A reduced curvature of the spectrum at midrapidity





Updated experimental results at energies 20A and 30A GeV are badly needed

Physical Origin of the Irregularity

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Summary

The irregularity is a combined effect of

- softest point of a EoS

 (a minimum of the sound speed)
 - spherical fireball
 - ⇒ essentially 3D expansion
 - ⇒ a peak at midrapidity
 - strongly deformed fireball \Rightarrow approximately 1D expansion
 - \Rightarrow a dip at midrapidity

the more softer matter

- \Rightarrow the more deformed fireball
- ⇒ a dip at midrapidity



from E. G. Nikonov, A. A. Shanenko and V. D. Toneev, Heavy Ion Phys. 8, 89 (1998)

• a change in the nonequilibrium regime from hadronic to partonic one

a change in cross sections \Rightarrow an irregularity in baryon stopping

This irregularity is a signal from hot and dense stage of nuclear collision



Particle Production



Summary

3ED

- Hadr. scenario considerably overestimates experimental antibaryons already at lower SPS energies, i.e. ≈20A GeV.
- At $\sqrt{s_{NN}}$ > 17.4 GeV, hadronic scenario overestimates available RHIC data on all species.
- Problem with ϕ -meson yields at lower SPS energies

hadronic scenario fails at high incident energies





Summary

• "Horn" anomaly in the K^+/π^+ ratio is not reproduced in any scenario.

?????????



Pt Spectra

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Flow

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Summary
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hadronic scenario again fails for antiprotons



Inverse Slopes and Mean Transverse Masses

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Summary



$$\langle m_T \rangle = \frac{\int d^2 p_T \ m_T \left(\frac{d^2 N}{m_T \ dm_T \ dy} \right)}{\int d^2 p_T \ \left(\frac{d^2 N}{m_T \ dm_T \ dy} \right)}$$



"Step" is not a signal of deconfinement







Deconfinement scenarios look preferable everywhere except for protons at 158A GeV

No comments as yet





Summary

Afterburner (Hirano, et al., nucl-th/0701075) or viscosity reduces v_2 achieved at hydro stage. Therefore, v_2 at 158A GeV is reasonable. However, v_2 at 40A GeV is too high.

Critical point? E.Shuryak, arXiv:hep-ph/0504048



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Summary

- Deconfinement scenarios look preferable at $\sqrt{s_{NN}}$ > 5 GeV
- net-proton rapidity distributions:
 Irregularity signaling deconfinement onset (not yet seen)
- particle production: hadronic scenario overestimates
 - exp. antibaryons yields already at $\sqrt{s_{NN}} > 5$ GeV, and at $\sqrt{s_{NN}} > 17.4$ GeV yields of all species
- m_T -spectra: "Step" is not a signal of deconfinement
- Directed and elliptic flow:
 - reasonably well reproduced so far, analysis is still in progress
- (i) Disappearance of v_2 at 40A GeV (semicentral collisions) and (ii) maximum in K^+/π^+ ratio at 20A – 30A GeV (central collisions) cannot be reproduced by any scenario
- Analysis is still in progress