# A new RMF based quark-nuclear matter EoS for applications in astrophysics and heavy-ion collisions

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# Comparition to current EOS<sup>1 2</sup>



 <sup>&</sup>lt;sup>1</sup>A.S. Khvorostukhin, V.V. Skokov, V.D. Toneev, K. Redlich, Eur.Phys.J.C48:531-543,2006
 <sup>2</sup>Y. B. Ivanov, V. N. Russkikh and V. D. Toneev, Phys. Rev. C 73 (2006) 044904

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- orientation on astrophysical constraints (e.g. two solar mass neutron stars)
- improvements in the low temperature / high density regime

### density functional approach with tdyn consistency

• start with approach for grand canonical potential density

$$\omega = -U - \sum_{i} g_i \int \frac{\mathrm{d}^3 p}{(2\pi)^3} \left\{ T \ln[1 + \mathrm{e}^{(\tilde{E}_i^+ - \mu_i)/T}] + T \ln[1 + \mathrm{e}^{(\tilde{E}_i^- + \mu_i)/T}] \right\}$$

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one can introduce an effective mass  $M_i = m - S$  and an effective chemical potential  $\tilde{\mu}_i = \mu_i - V_i$  with self-energies

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• rearrangement contributions  $U, m_i^{\rm R}, E_i^{\rm R}$  ensure consistency

# density functional approach with tdyn consistency derivation of rearrangement

• to preserve thermodynamical consistency the definition of the particle density

$$n_i = \frac{\partial p}{\partial \mu_i} = g_i \int \frac{\mathrm{d}^3 p}{(2\pi)^3} \frac{1}{1 + \mathrm{e}^{(\sqrt{p^2 - M_i} - \tilde{\mu}_i)/T}}$$

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• this leads us to the differential equation

$$\frac{\partial U}{\partial \mu_i} = \sum_j \left\{ n_j \frac{\partial \Delta E_j}{\partial \mu_i} + n_j \frac{\partial E_j^{\rm R}}{\partial \mu_i} - n_j^{\rm s} \frac{\partial \Delta m_j}{\partial \mu_i} - n_j^{\rm s} \frac{\partial m_j^{\rm R}}{\partial \mu_i} \right\}$$

# density functional approach with tdyn consistency $_{\mbox{\tiny derivation of rearrangement}}$

• one solution, provided by Stefan Typel (but slightly altered), is

$$E_i^{\rm R} = \sum_j n_j \frac{\partial \Delta E_j}{\partial n_i} - \sum_j n_j^{\rm s} \frac{\partial \Delta m_j}{\partial n_i}$$
$$m_i^{\rm R} = -\sum_j n_j \frac{\partial \Delta E_j}{\partial n_i^{\rm s}} + \sum_j n_j^{\rm s} \frac{\partial \Delta m_j}{\partial n_i^{\rm s}}$$
$$U = \sum_i n_i E_i^{\rm R} - \sum_i n_i^{\rm s} m_i^{\rm R}$$

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now we are coming to an concrete example

### Stringflip modell

• confinement potential with effects of pauli quenching<sup>34</sup>

$$\Delta m_i = -C^{q} \cdot (n^{s})^{1/3} - D^{q} \cdot (n^{s})^{-1/3} \quad \Delta E_i = a^{q}n + b^{q}n^3$$

with the density dependent  $D^{q} = D^{q}(n^{s})$ .

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• The rearrangement contributions

$$E_i^{\rm R} = (a^{\rm q} + 3b^{\rm q}n^2) n = E^{\rm R}$$
$$m_i^{\rm R} = \left(-\frac{C^{\rm q}}{3}(n^{\rm s})^{-2/3} + \frac{D^{\rm q}}{3}(n^{\rm s})^{-4/3} - D^{\rm q'} \cdot (n^{\rm s})^{-1/3}\right) n^{\rm s} = m^{\rm R}$$
$$U = nE^{\rm R} - n^{\rm s}m^{\rm R}$$

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#### resulting eos pressure over baryon chemical potential



# resulting eos



#### resulting eos neutron star configurations



#### resulting eos symmetric matter



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- thank you for your attention!

# Phase transition bare models



#### Phase transition including bag constant



# Phase transition over density



# Phase transition phase diagram

