## THE ROLE OF CONSTITUENT QUARK MASS IN THE STANDARD MODEL

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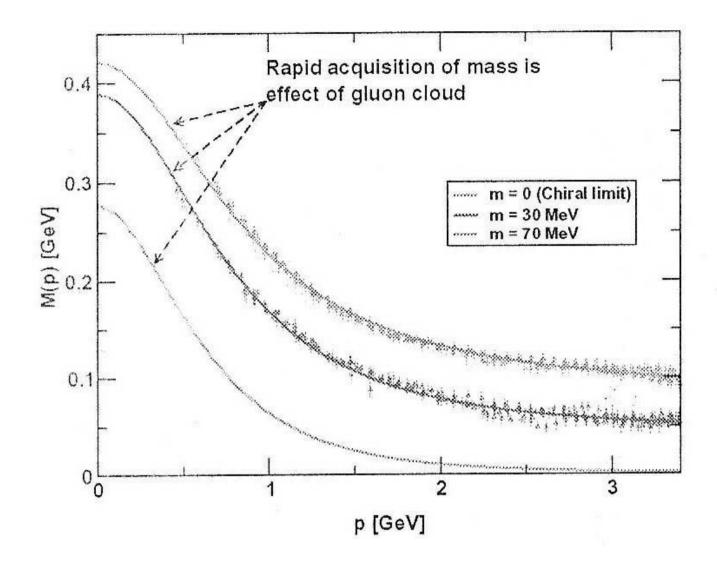


Fig.1. QCD gluon-quark-dressing effect calculated with Dyson-Schwinger Equation [9,13-17], initial masses m=0, 30 and 70 MeV; the constituent-quark mass arises from a cloud of low-momentum gluons attaching themselves to the current-quark; this is dynamical chiral symmetry breaking: an nonperturbative effect that generates a quark mass from nothing even at chiral limit m=0, bottom curve) [9].

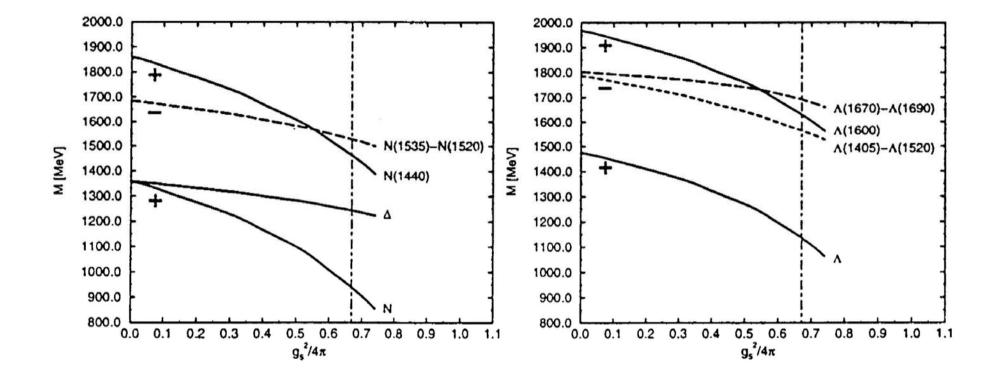


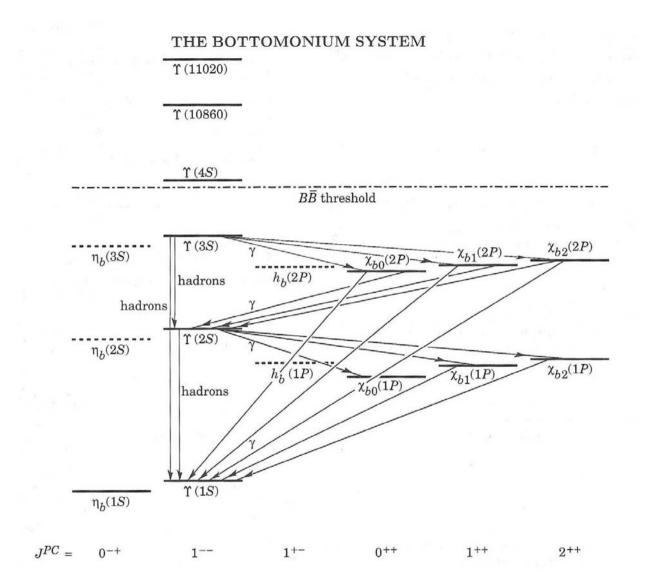
Fig.2. Calculation of nonstrange baryon masses (left) and  $\Lambda$ -hyperon masses as a function of interaction strength within GBECQM – Goldstone Boson Exchange interaction Constituent Quark Model [19]; initial baryon mass 1350 MeV=3×450 MeV=3 $M_q$  is near the bottom "+" on the left vertical axis.

**Table 1.** Discussed in the literature closeness of masses or mass differences to the integer numbers (k) of the pion mass value  $m_{\pi}^{\pm}=139.5$  MeV or to  $2m_{\pi}^{\pm}+m_{\pi}^{\circ}=\kappa_{T}=409$  MeV.

Particle	Λ	Ω	(bb)(2S-1S)	(bb)(4S-2S)	$\Delta E_B$
Mass or $\Delta M$ (MeV)	1115.683(6) $1672.45(29)$		10023-9460	10579-10023	408.9
			$=\!563$	=556	
$km_{\pi} \text{ or } 2m_{\pi}^{\pm} + m_{\pi}^{\circ} = \kappa_T$	1116 k=8	1672 k=12	558 $k=4$	558 k=4	409 $\kappa_T$
difference, reference	0, [1,23-26]	0, [1, 25, 26]	-5, [1,26]	-2, [1,27]	0, [28]

**Table 2.** Discussed in the literature closeness of particle masses or mass differences to the integer numbers (k) of the nucleon  $\Delta$ -excitation parameter  $\Delta_{\Delta}=147$  MeV [23].

Particle	(cc) (2S-1S)	(cc) (1S)	Ξ	$\omega_3$ - $\omega$	$K_3^*-K^*(892)$	$\Delta E_B$
Mass or	3686.1-3096.9	3096.92(1)	1321.3	1667(4) - 782.7	1776(7)-891.7	441.5 (Fig.4)
$\Delta M (MeV)$	=589.2 k $=4$	3087 k=21	k=9	884(4) k=6	=884(7) k=6	k=4
$k\Delta M_{\Delta}$	588	588	1323	882	882	441
-diff., ref.	1, [1, 25, 29]	1, [1, 25, 29]	-2, [??]	2(4), [1, 25, 29]	2(7), [1, 25, 29]	0,[28]



**Fig.3.** Top Mesons of (bb)-structure with radial excitations n=1 and n=1-3 close to  $4m_{\pi}$  [1,27]; Bottom Mesons of (cc)-structure with radial excitation (n=1) close to  $4\Delta M_{\Delta}=4\times 147$  MeV=588 MeV.

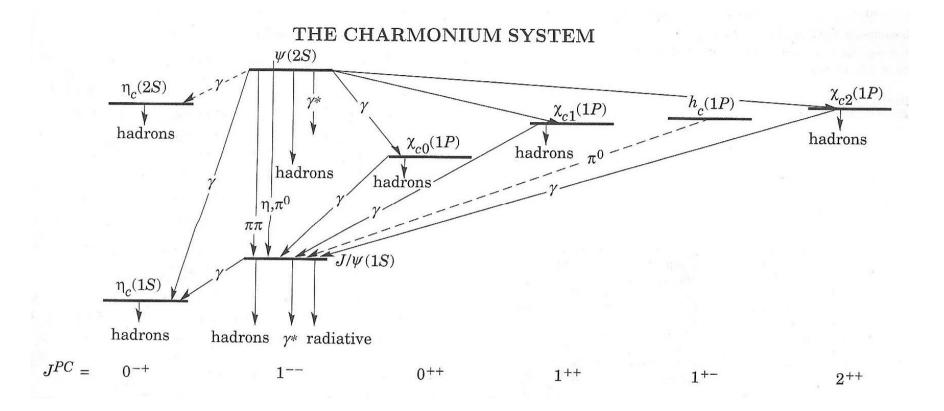


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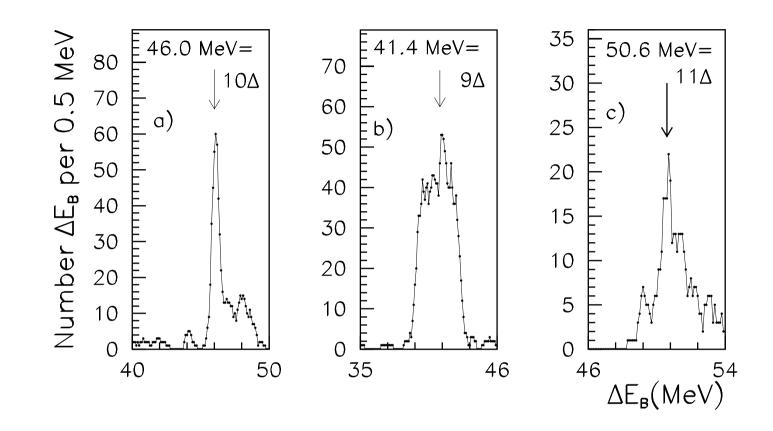


Fig. 4, *Top:* Distribution of  $\Delta E_B$  in nuclei with  $\Delta Z=2$ ,  $\Delta N=4$ , Z=50-58, 64-82 and  $Z\leq28$  [26] *Bottom:* The same in all even-even and nuclei [26], in even-even nuclei with  $Z\leq58$  [29].

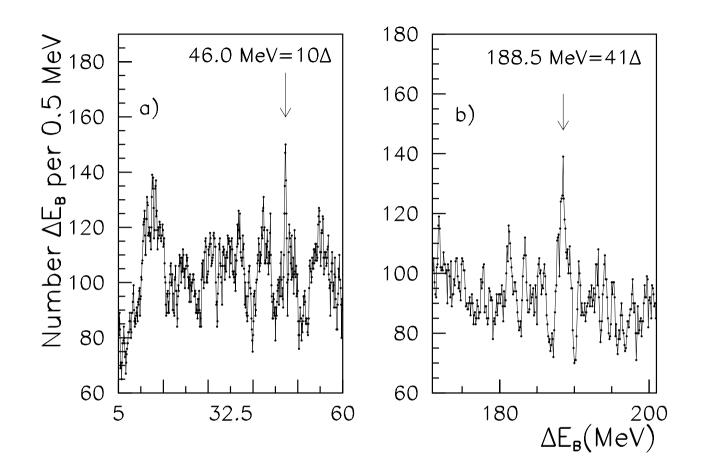


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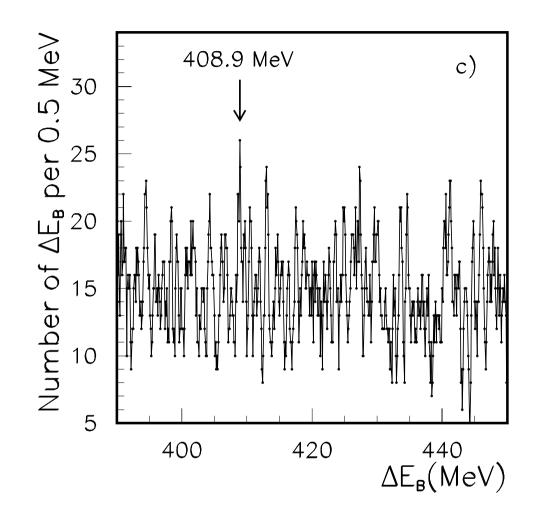
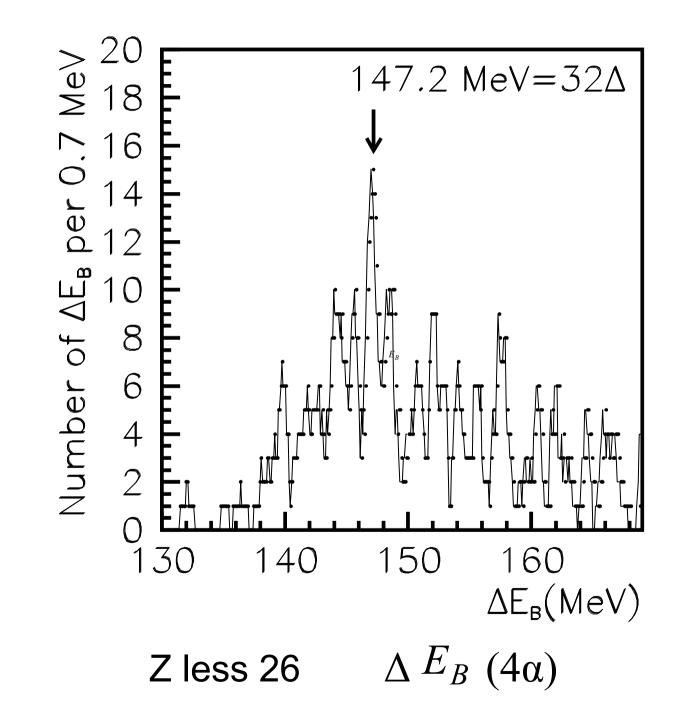
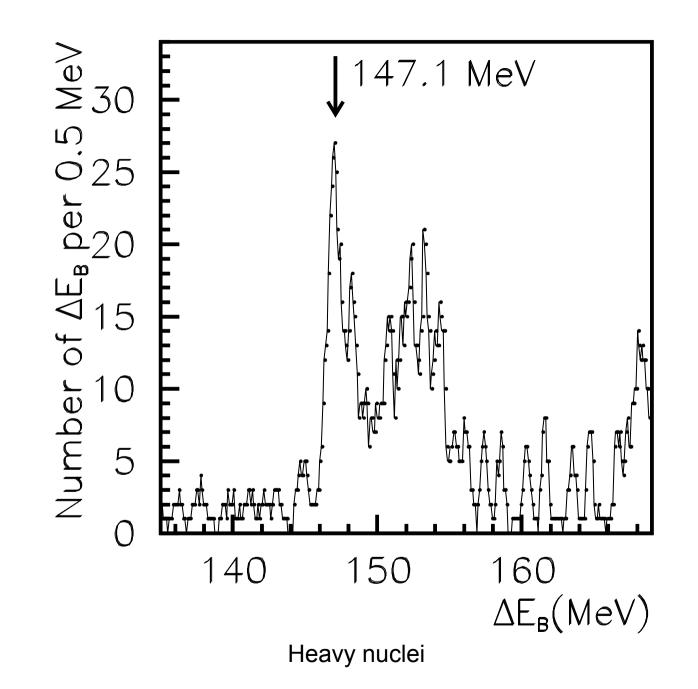


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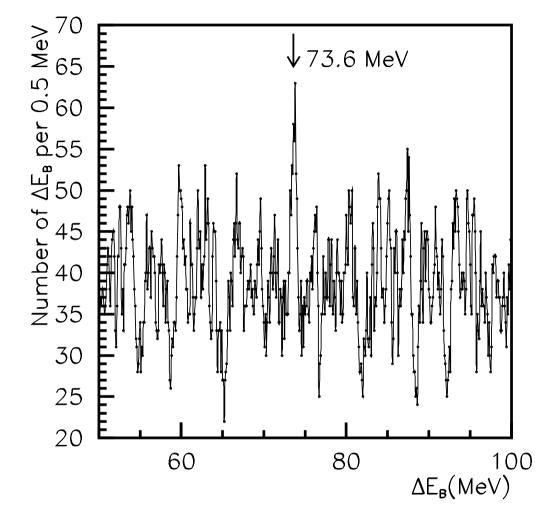


Fig. 5, *top:* Distribution of ΔEB in nuclei with ΔZ≤26; 4α- 2α-config. and all nuclei [6]); *center*. Distribution of adjacent intervals ΔE<sub>B-AIM</sub> in nuclei with Z≤26 for x=147.2 and 73.6 MeV [6]); *bottom*: Distribution of ΔE<sub>B</sub> in nuclei with ΔZ=8, ΔN=14 (Z=50-82); Distribution of ΔE<sub>B-AIM</sub> in nuclei with ΔZ=65-81 for x=147.1 MeV [6]; Distribution of ΔE<sub>B</sub> in all odd-odd nuclei [26].

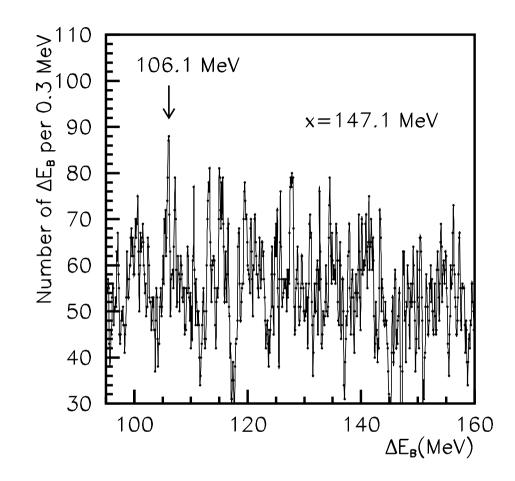


Fig. 5, *top:* Distribution of  $\Delta$ EB in nuclei with  $\Delta$ Z≤26; 4α- 2α-config. and all nuclei [6]); *center*. Distribution of adjacent intervals  $\Delta$ EB-AIM in nuclei with Z≤26 for x=147.2 and 73.6 MeV [6]); *bottom*: Distribution of  $\Delta$ EB in nuclei with  $\Delta$ Z=8,  $\Delta$ N=14 (Z=50-82); Distribution of  $\Delta$ EB-AIM in nuclei with  $\Delta$ Z=65-81 for x=147.1 MeV [6]; Distribution of  $\Delta$ EB in all odd-odd nuclei [26].

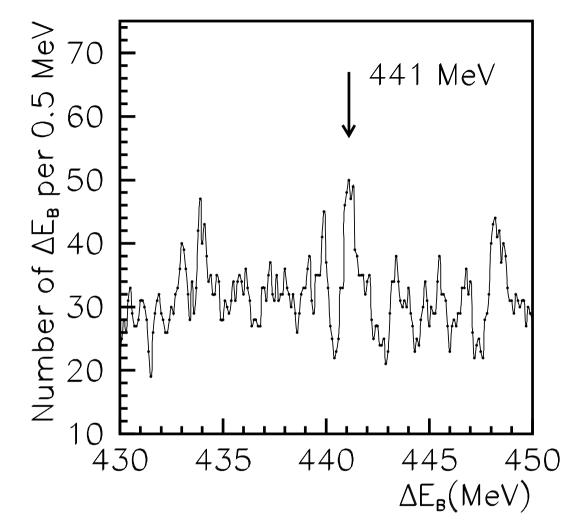


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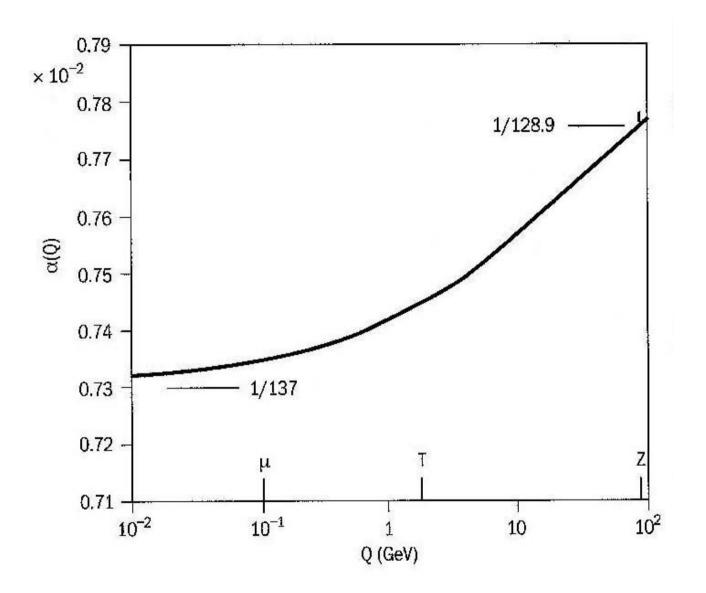


Fig.6. Momentum transfer evolution of QED effective electron charge squired. The momotonically rising theoretical curve is confronted with precise measurements at the Z mass at CERN LEP collider

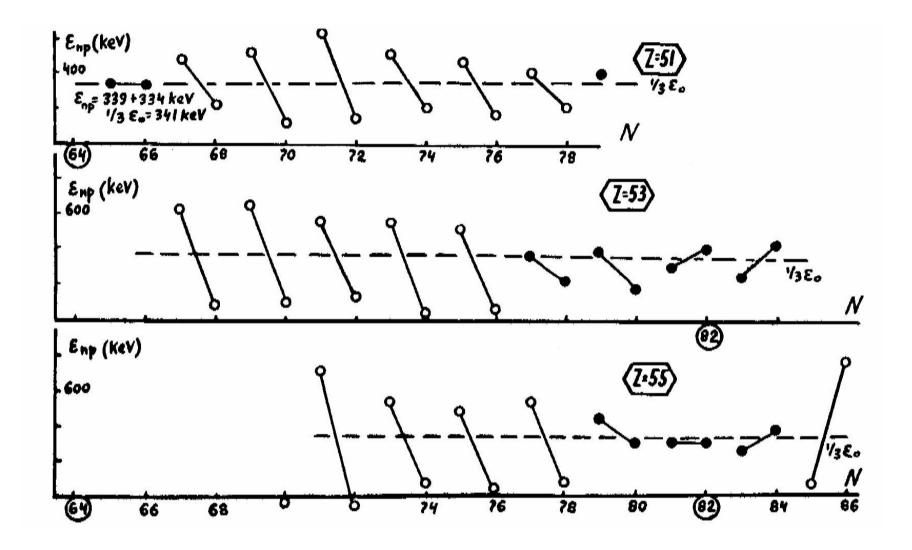
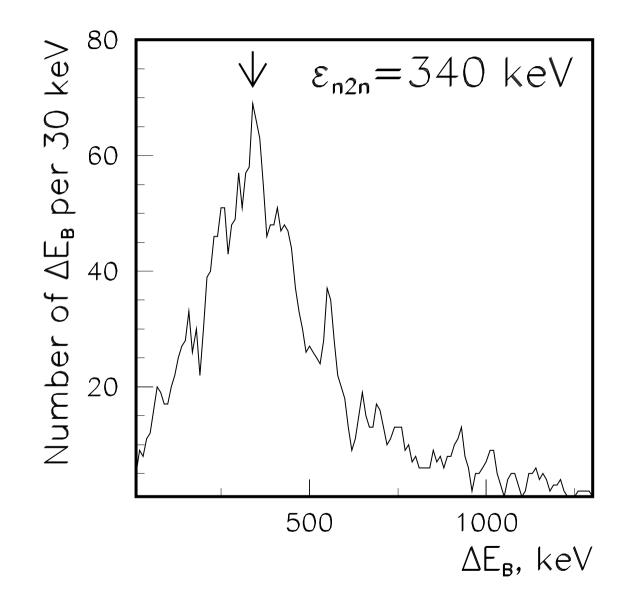
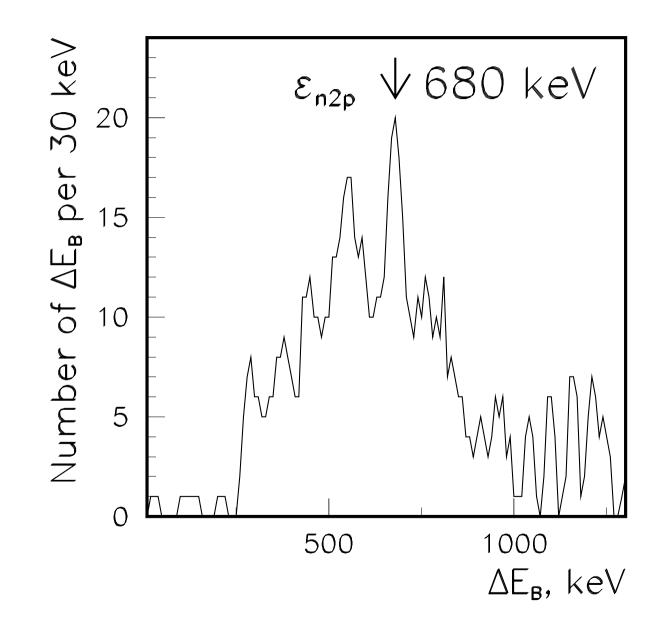
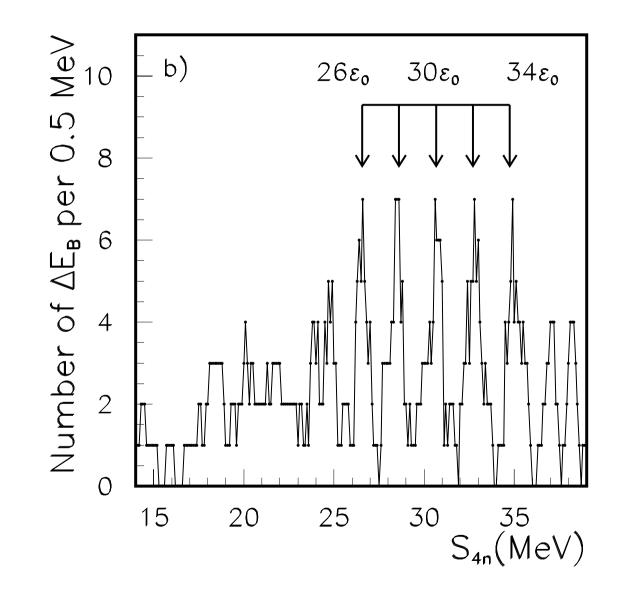
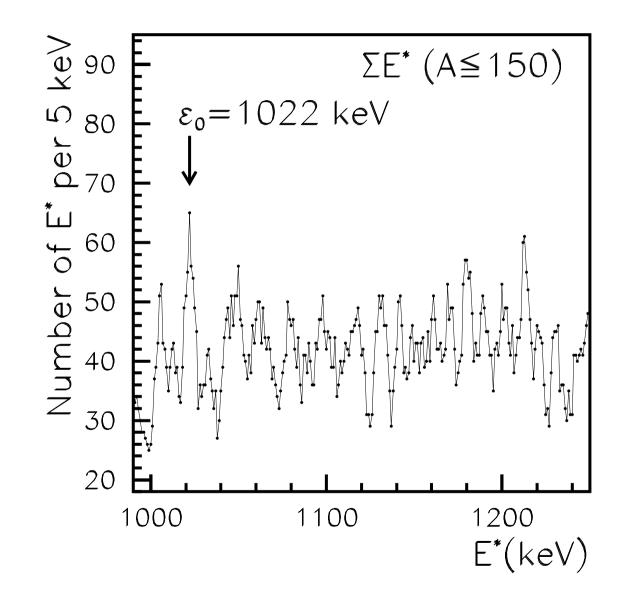


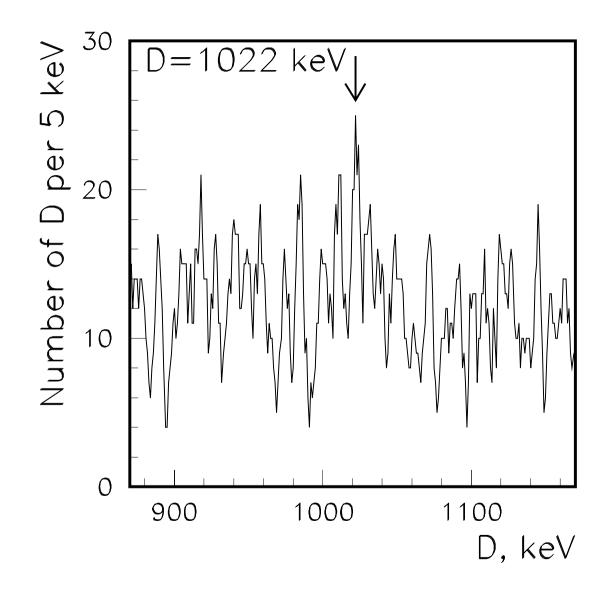
Fig.7. Parameters of the residual interaction  $\varepsilon_{np}$  from differences of  $S_p$  in Z-odd nuclei ( $\Delta N = 1$ ). The value  $\varepsilon_0/3=341$  keV is given by horizontal line

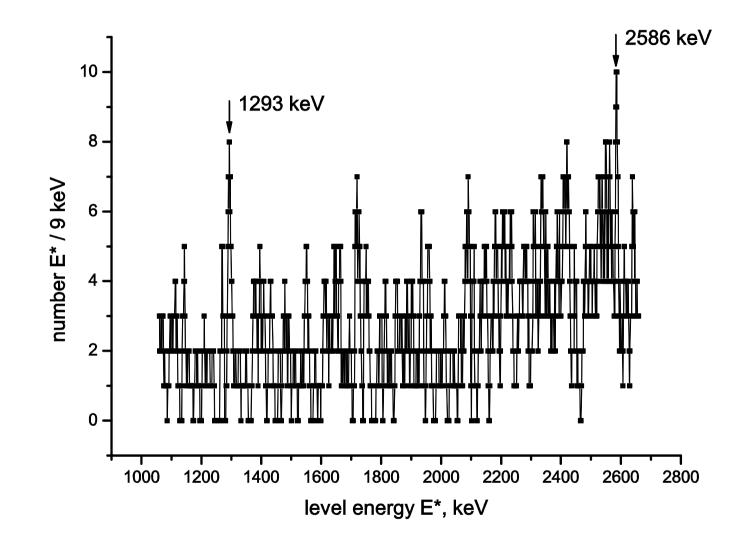


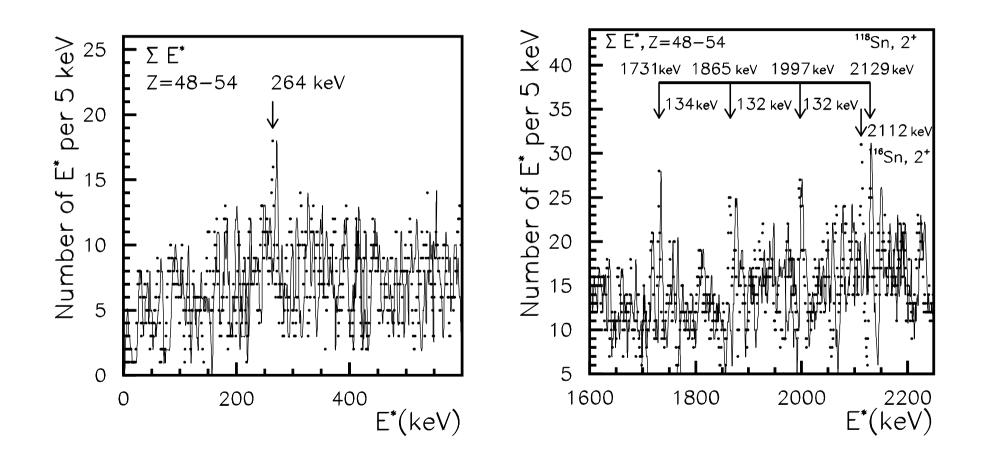


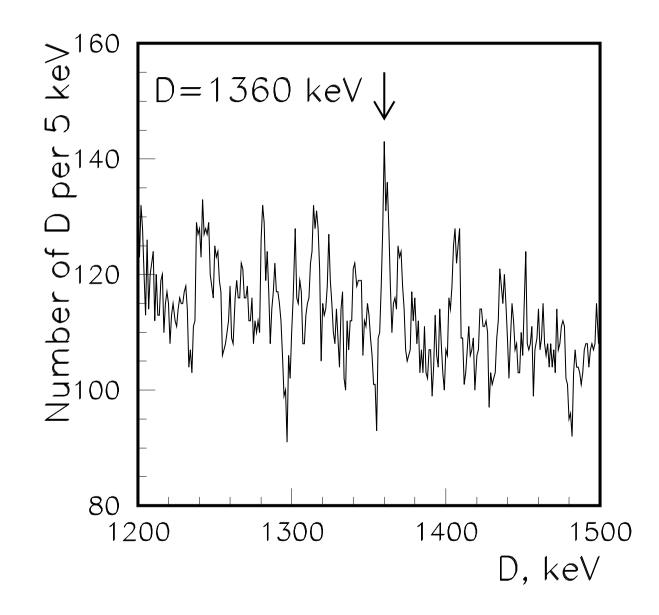


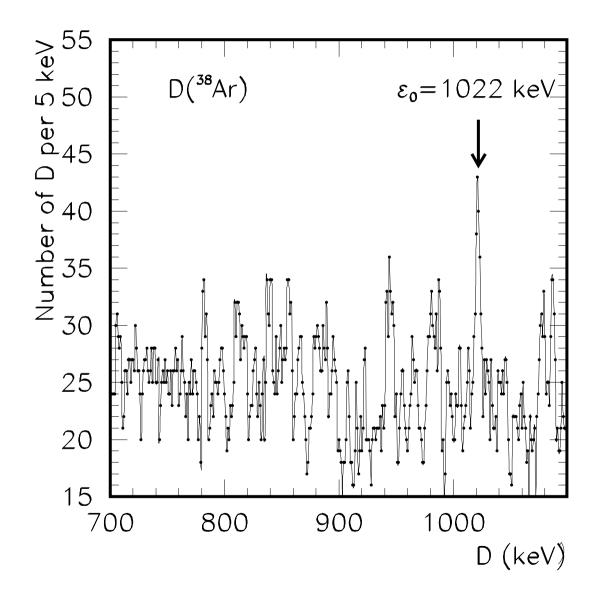


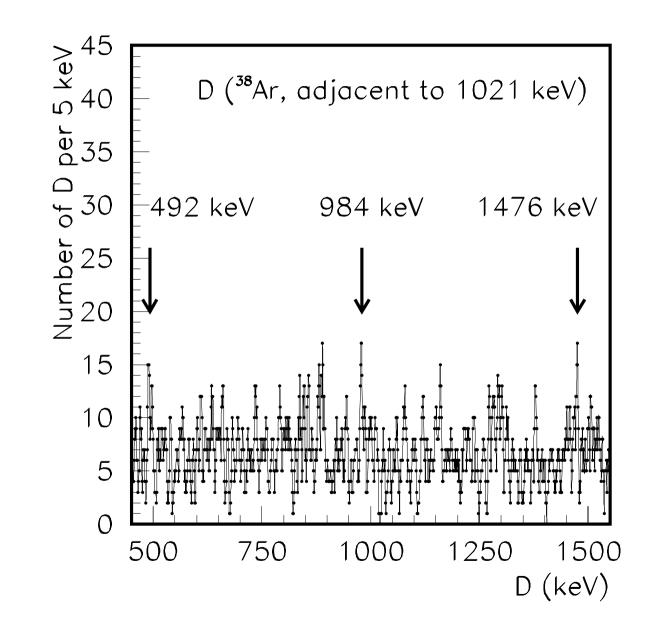


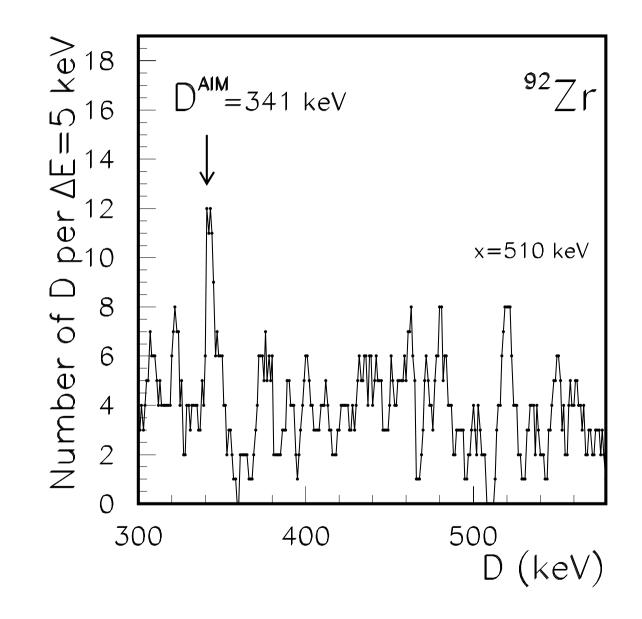












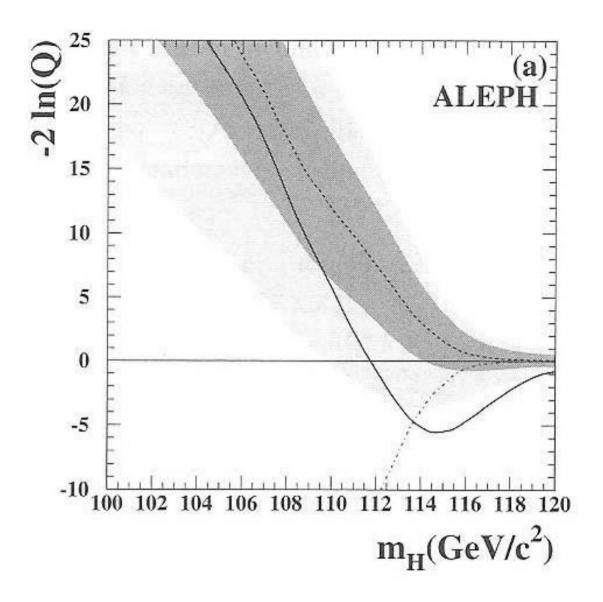


Fig.8. ALEPH results with about 3 standard deviation at mass 115 GeV; observed (solid line) and the expected behaviors of the test statistic (dark region) are presented and discussed in [1].

**Table 3.** Representation of parameters of tuning effects in particle masses (upper part) and in nuclear data by the expression  $(n \times m_e (\alpha/2\pi)^x) \times m$  with  $\alpha = 137^{-1}$  [5]. Asterisk marks stable intervals observed in low-energy excitations and neutron resonances;  $\varepsilon_{np} = 340$  keV is discussed in the text, Fig.7

Х	m	n=1/8	n=1	n=13	n=16	n=17	n=18
-1	1			$M_Z = 91.188$	$M_{H} = 115$		
${\rm GeV}$	3				$2m_t = 348$		
0	1	$2m_e$	$16m_e$	$m_{\mu} = 105.658$	· · ·	$m_{\pi}$ - $m_e$	$147=\Delta M_{\Delta}$
MeV	1	$\varepsilon_o$	δ	$106.4 = \Delta E_B$	$130 = \Delta E_B$	$140 = \Delta E_B$	$147 = \Delta E_B$
	3				$M_q''{=}m_\rho/2\approx m_\omega/2$	$M'_q = 420$	$M_q$ =441= $\Delta E_B$
1	1	$1.2^{*}$	$9.48 = \delta'^{*}$	123*	$152^{*}$	161*	$170^{*}$
$\mathrm{keV}$	2			246*	$303^{*}$	$321^{*}$	$\varepsilon_{np}=340=\varepsilon_o/3$
	3			368*	$455^{*}$	481*	$511 = \varepsilon_o/2$
	8	$9.5^{*}$	$76^{*}$	984, Fig.8	1212*	$1293 = D_o$	1360, Fig.8
2	1		11*	$143^{*}$	176*	187*	D in neutron
eV	4	$5.5^{*}$	44*	$572^{*}$		$750 - 1500^*$	resonances

## Lepton ratio as the distinguished parameter

Earlier, as a realization of Nambu's suggestion to search for empirical mass relations needed for SM-development, it was noticed in [5,6] that

- 1) the well-known lepton ratio  $L=m\mu/me=206.77$  becomes the integer 207=9\*23=13\*16-1 after a small QED radiative correction applied to me (it becomes  $m\mu/me(1-\alpha/2\pi)=207.01$ )
- 2) the same ratio L=207 exists between masses of vector bosons Mz=91.188(2) GeV and Mw=80.40(3) GeV and two above discussed estimates of baryon/meson constituent quark masses Mq=441 MeV=mΞ<sub>i</sub>/3=(3/2)(m△-mN) and M" q =mp/2=775.5(4) MeV/2=387.8(2) MeV
- [1] (*MZ*/441 MeV=206.8;  $Mw/(m_{\rho}/2)$ =207.3 [5,6]). The origin of these effects should be considered in the complex analysis of tunung effects in particle masses and in nuclear data [5,6].

## **Conclusions**

The QCD-based estimates of the constituent quark masses (M0 q=420 MeV, Mq=441 MeV, M'' q) could play important role in the description of Standard Model dynamics if the observed now empirical relations in particle masses (and value MH) would be confirmed in the experiment.

Nuclear data can provide some important additional information on fundamental properties of strong nucleon interactions and nuclear matter as well as general properties of fermion systems.