

# Spin Correlations in Higgs Decays

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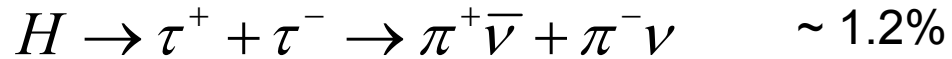
JINR, Dubna

In collaboration with R. Leitner,

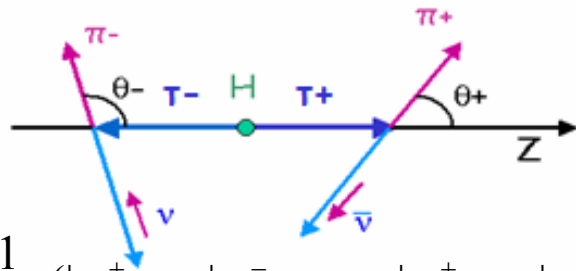
O. Teryaev

# Outline

- Spin correlation in decay products: Quantum mechanical consideration
- Spin correlation in decay products: Quantum Field Theory considerations
- Scalar vs vector particles decays
- Simulations
- Conclusions

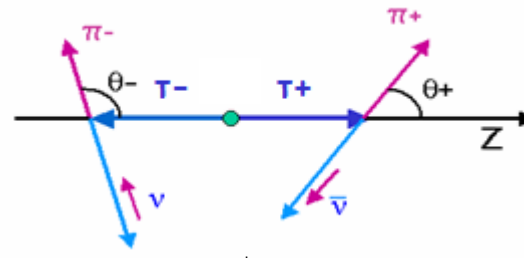


Spin correlation in decay products:  
Quantum mechanical consideration

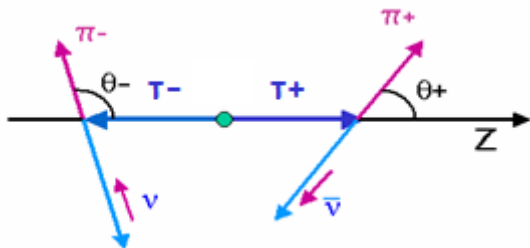


$$A_{0,0} = \frac{1}{\sqrt{2}} (|\tau^+, +\rangle |\tau^-, +\rangle - |\tau^+, -\rangle |\tau^-, -\rangle)$$

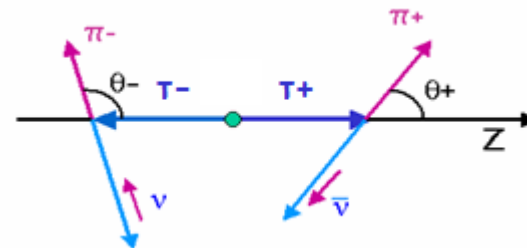
$\sin \frac{\theta_+}{2} e^{i\phi_+/2}$        $-\sin \frac{\theta_-}{2} e^{-i\phi_-/2}$        $\cos \frac{\theta_+}{2} e^{-i\phi_+/2}$        $\cos \frac{\theta_-}{2} e^{i\phi_-/2}$



$$A_{1,1} = |\tau^+, +\rangle |\tau^-, -\rangle$$



$$A_{1,0} = \frac{1}{\sqrt{2}} (|\tau^+, +\rangle |\tau^-, +\rangle + |\tau^+, -\rangle |\tau^-, -\rangle)$$



$$A_{1,-1} = |\tau^+, -\rangle |\tau^-, +\rangle$$

$$P_{J=0}(\theta^+, \theta^-, \varphi^+, \varphi^-) = |A_{0,0}|^2$$

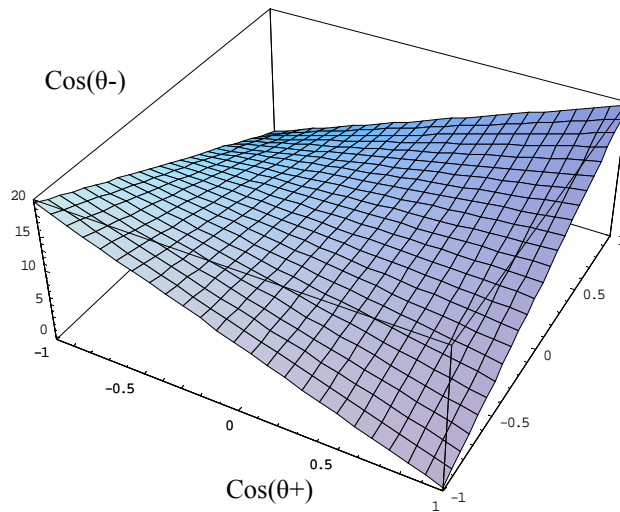
$$P_{J=0}(\theta^+, \theta^-) = \int_0^{2\pi} d\varphi^+ \int_0^{2\pi} d\varphi^- P_{J=0}(\theta^+, \theta^-, \varphi^+, \varphi^-)$$

$$P_{J=1}(\theta^+, \theta^-, \varphi^+, \varphi^-) = \frac{1}{3} (|A_{1,0}|^2 + |A_{1,-1}|^2 + |A_{1,1}|^2)$$

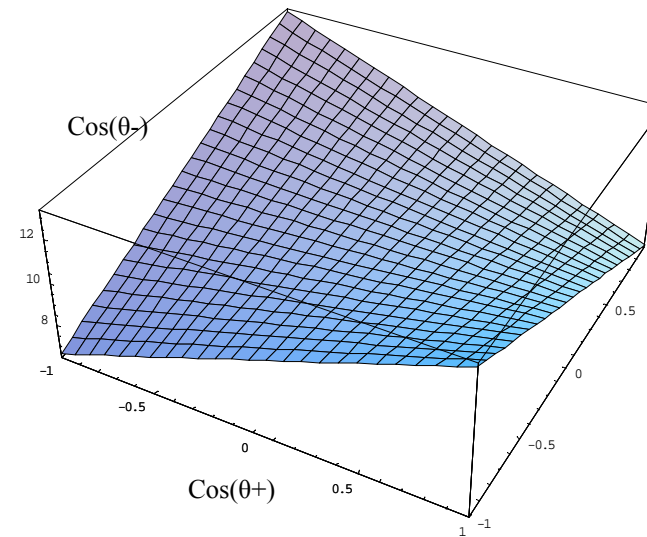
$$P_{J=1}(\theta^+, \theta^-) = \int_0^{2\pi} d\varphi^+ \int_0^{2\pi} d\varphi^- P_{J=1}(\theta^+, \theta^-, \varphi^+, \varphi^-)$$

After integration with respect to  $\varphi^+$  and  $\varphi^-$  we have probabilities dependent only on  $\theta^+$ ,  $\theta^-$ .

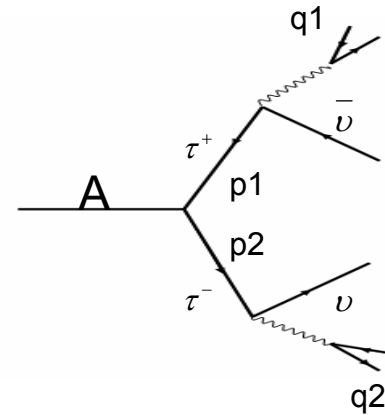
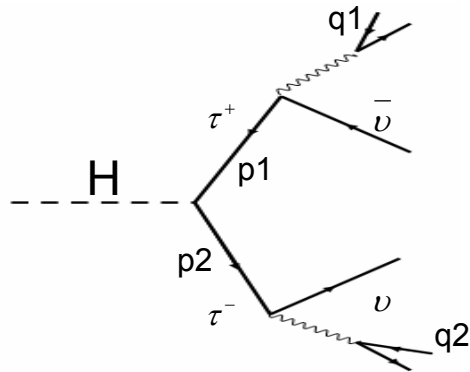
$$P_{J=0}(\theta^+, \theta^-) \sim (1 + \cos\theta^+ \cos\theta^-)$$



$$P_{J=1}(\theta^+, \theta^-) \sim (1 - \frac{1}{3} \cos\theta^+ \cos\theta^-)$$



# Spin correlation in decay products: Quantum Field Theory considerations



Sudakov variables is used for extract the angular dependence:

$$\begin{aligned}
 q_1 &= x_1 p_1 + u_1 p_2 + k_1^\perp \\
 q_2 &= x_2 p_2 + u_2 p_1 + k_2^\perp
 \end{aligned}
 \begin{array}{l}
 \longrightarrow \\
 \longleftarrow
 \end{array}
 \left\{ \begin{array}{l}
 (q_2 p_2) = x_2 (p_2 p_2) + u_2 (p_1 p_2) \\
 (q_1 p_1) = x_1 (p_1 p_1) + u_1 (p_2 p_1) \\
 (q_2 p_1) = x_2 (p_2 p_1) + u_2 (p_1 p_1) \\
 (q_1 p_2) = x_1 (p_1 p_2) + u_1 (p_2 p_2) \\
 (q_1 q_2) = (x_1 x_2 + u_1 u_2) (p_1 p_2) + \\
 + x_1 u_2 (p_1 p_1) + x_2 u_1 (p_2 p_2) + k_1^\perp k_2^\perp
 \end{array} \right.$$

$$u_{1,2} = -\frac{m^2 \cos \theta^\pm}{2\sqrt{\frac{M^4}{4} - 2m^2}}$$

$$x_{1,2} = \frac{1}{2} \left( 1 - \frac{\cos \theta^\pm}{\sqrt{1 - \frac{m^4}{(\frac{M^2}{2} - m^2)}}} \right)$$

## QFT result

$$P_{J=0}^{QFT} \sim \frac{16m^2}{\frac{M^2}{2} - m^2} (1 + \cos \theta^+ \cos \theta^-) +$$
$$+ 16(1 + \cos \theta^+ \cos \theta^-) -$$
$$- \left( \frac{64}{\frac{M^2}{2} - m^2} + \frac{64}{m^2} \right) (k_1^\perp k_2^\perp)$$

$$P_{J=1}^{QFT} \sim \frac{64m^2}{\frac{M^2}{2} - m^2} -$$
$$- 32(1 - \cos \theta^+ \cos \theta^-) -$$
$$- \frac{128}{\frac{M^2}{2} - m^2} (k_1^\perp k_2^\perp)$$

## Non-relativistic case

$$P_{J=0}^{QFT} \sim 32(1 + \cos \theta^+ \cos \theta^-)$$

$$P_{J=1}^{QFT} \sim -96\left(1 - \frac{1}{3} \cos \theta^+ \cos \theta^-\right)$$

## Relativistic case – $m \ll M$

$$P_{J=0}^{QFT} \sim 16(1 + \cos \theta^+ \cos \theta^-)$$

$$P_{J=1}^{QFT} \sim -32(1 - \cos \theta^+ \cos \theta^-)$$

## Quantum mechanical result

$$P_{J=0} \sim (1 + \cos \theta^+ \cos \theta^-)$$

$$P_{J=1} \sim \left(1 - \frac{1}{3} \cos \theta^+ \cos \theta^-\right)$$

# Comparison of methods

- J=0: For scalar particle (Higgs) decay -COINCIDE
- J=1: Quantum mechanical approach

$$P_{J=1} \sim (1 - \frac{1}{3} \cos \theta^+ \cos \theta^-)$$

- J=1: QFT approach
  - Coincide for non-relativistic tau's ( $m \sim M/2$ )

$$P_{J=1}^{QFT} \sim (1 - \frac{1}{3} \cos \theta^+ \cos \theta^-)$$

- For relativistic tau's ( $m \ll M$ ) – factor 3 difference – only longitudinal polarization of heavy vector “Higgs” survive

$$P_{J=1}^{QFT} \sim (1 - \cos \theta^+ \cos \theta^-)$$

# Experimental consideration

$$p_\tau = \frac{1}{2} \sqrt{M_H^2 - 4m^2} \approx M_H/2 = E_\tau$$

$$p_\pi^* = \frac{m^2 - m_\pi^2}{2m} \approx \frac{m}{2}$$

$$E_\pi^* = \frac{m^2 + m_\pi^2}{2m} \approx \frac{m}{2}$$

$$\beta_\tau = \frac{p_\tau}{E_\tau} \approx 1 \quad \gamma_\tau = \frac{E_\tau}{m} \approx \frac{M_H}{2m}$$

$$E_{\pi^\pm} = \gamma_\tau (E_{\pi^\pm}^* \pm \beta_\tau p_{\pi^\pm}^* \cos \theta^\pm) \approx \frac{M_H}{2m} \frac{m}{2} (1 \pm \cos \theta^\pm) = \frac{M_H}{2} (1 \pm \cos \theta^\pm)$$

## Collinear approximation

$$\sum \vec{p}_T = 0 \quad \vec{p}_T = (p_x, p_y) \quad \begin{array}{l} 6 \text{ unmeasured variables} \\ \text{two kinematic constraints} \end{array}$$

$$\vec{p}_H = c_- \vec{p}_{\pi^-} + c_+ \vec{p}_{\pi^+}$$

$$\vec{p}_{T,H} = \vec{p}_{T,jet} = c_- \vec{p}_{T,\pi^-} + c_+ \vec{p}_{T,\pi^+}$$

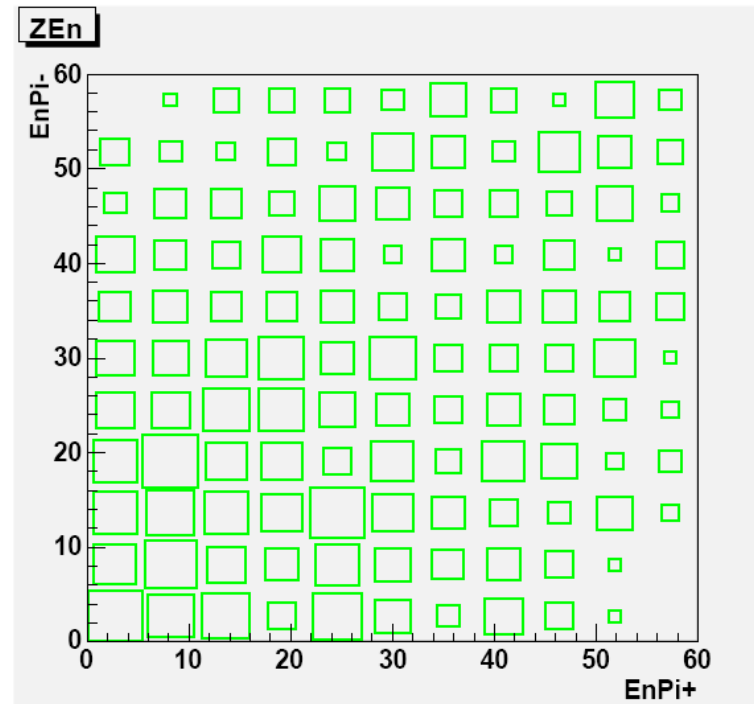
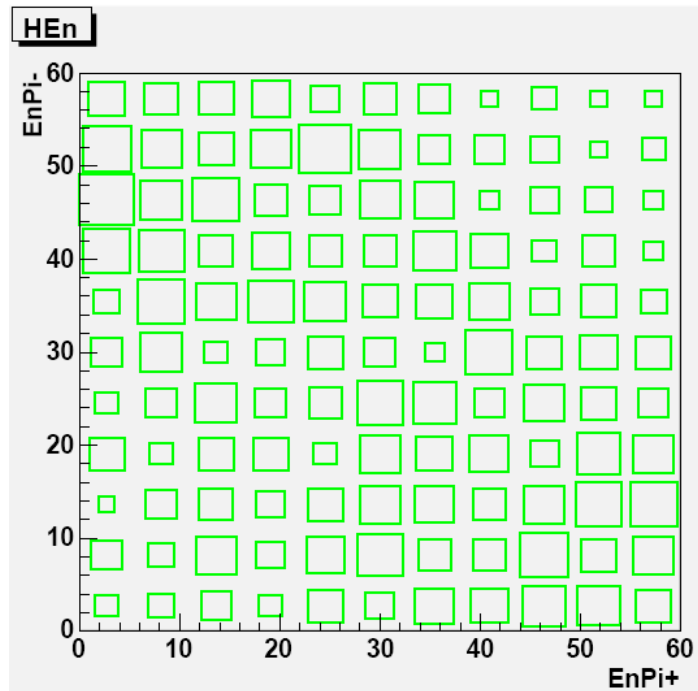
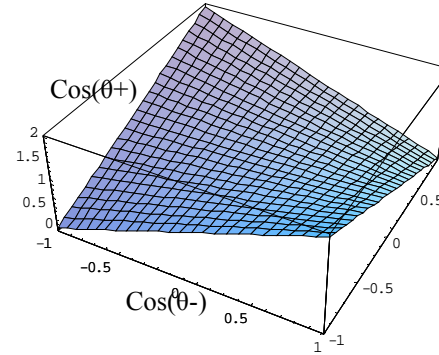
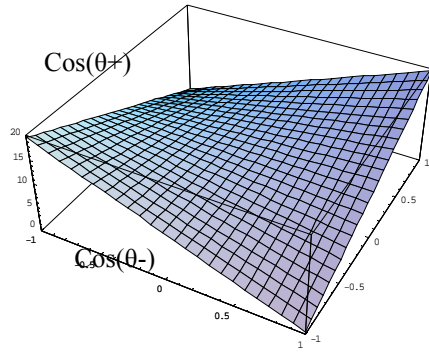
$$c_- = \frac{p_{jet,x} p_{\pi^+,y} - p_{jet,y} p_{\pi^+,x}}{p_{\pi^-,x} p_{\pi^+,y} - p_{\pi^+,x} p_{\pi^-,y}}$$

$$c_+ = \frac{p_{jet,y} p_{\pi^-,x} - p_{jet,x} p_{\pi^-,y}}{p_{\pi^-,x} p_{\pi^+,y} - p_{\pi^+,x} p_{\pi^-,y}}$$

$$M_H^2 = (c_- p_{\pi^-} + c_+ p_{\pi^+})^2 - (c_- \vec{p}_{\pi^-} + c_+ \vec{p}_{\pi^+})^2$$



# Simulations for ATLAS detector at LHC (PYTHIA with TAUOLA)



# Outlook

- Separation of signals from background ( $Z^0$ )
- Applications for other decays ( $W$  bosons,  $t$ -quarks?)
- Search of new particles, SUSY?