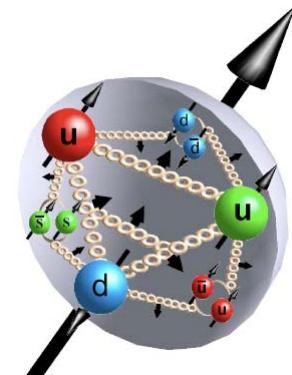


Transverse Spin program at



involving possible Collins and Sivers
effects



Mirko Planinić

University of Zagreb

Prague, July 20 - July 26, 2008



RHIC Spin Goals - I

How is the proton built from its known quark and gluon constituents?

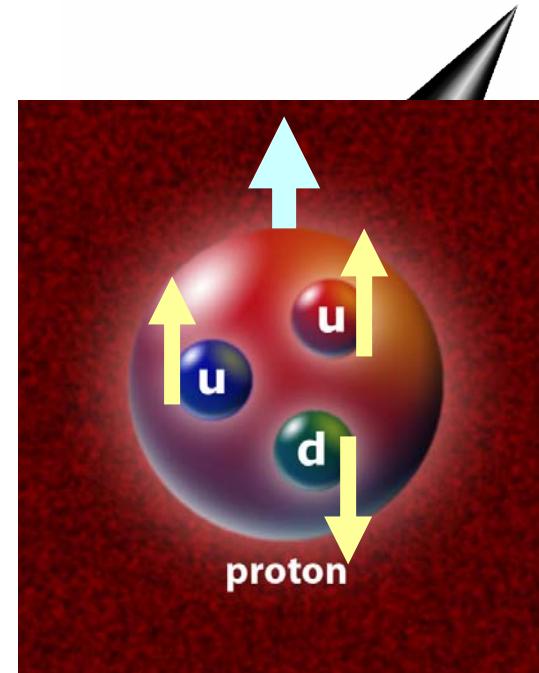
As with atomic and nuclear structure, this is an evolving understanding

In QCD: proton is not
just 3 quarks !

Recall:

simple quark model

Rich structure of quarks
anti-quarks, gluons

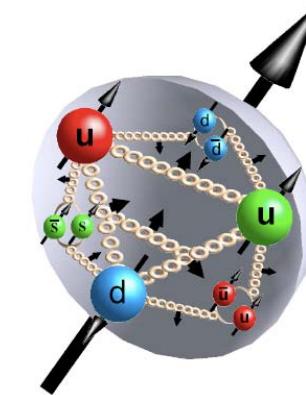


RHIC Spin Goals II



Longitudinal Spin

$$J = \frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L_z^q + L_z^g$$



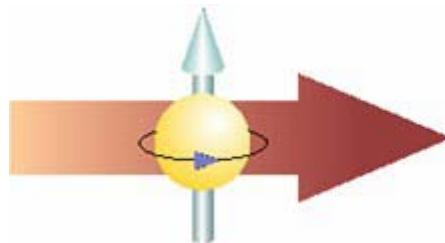
- $\Delta\Sigma = \Delta u + \Delta d + \Delta s + \dots$, Quark Contribution
- ΔG , Gluon Contribution
- L_z^q , Quarks orbital angular momentum
- L_z^g , Gluons orbital angular momentum

Frank Simon talk



RHIC Spin Goals III

Transverse Spin



PRD 70 (2004) 114001

For a proton with transverse spin vector \vec{S}_T

$$\frac{1}{2} = \frac{1}{2} \sum_{q,\bar{q}} \int dx \Delta_T q^a(x) + \sum_{q,\bar{q},G} \left\langle L_{\vec{S}_T} \right\rangle^a$$

$L_{\vec{S}_T}$ Component of partonic orbital angular momentum L along \vec{S}_T

$\Delta_T q^a(x)$ quark transverse spin distribution in the nucleon

Understanding the origin of proton spin helps to understand its structure.



RHIC Spin Goals - IV

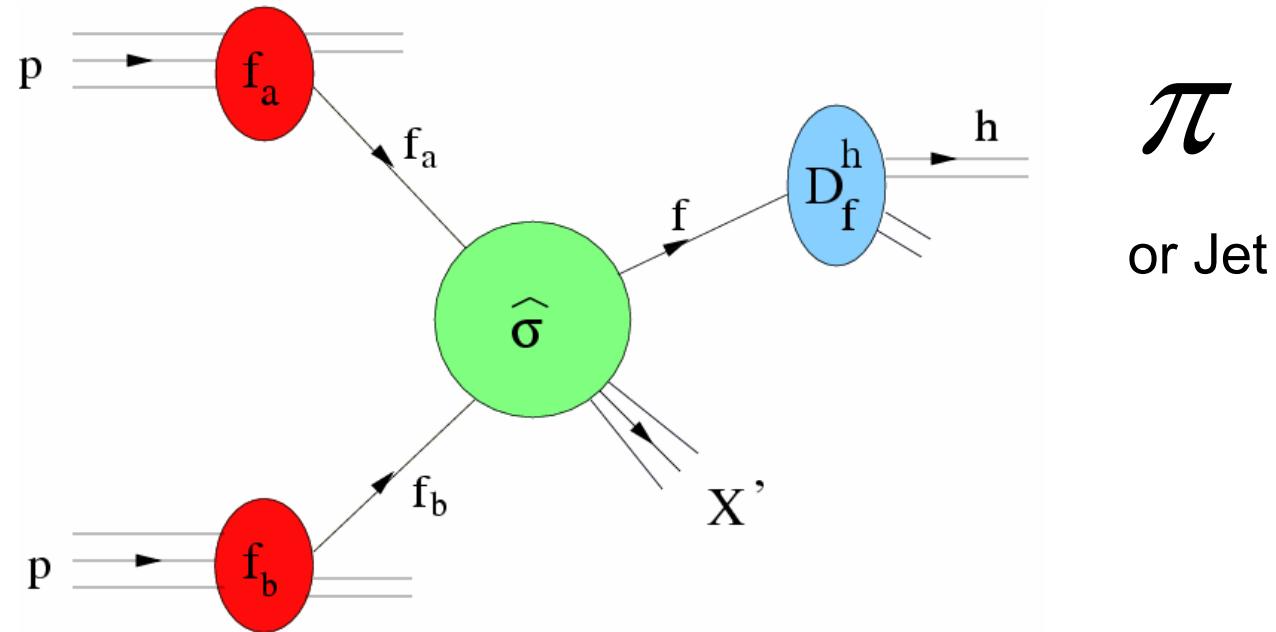
- Direct measurement of flavor identified *anti-quark polarization* using *parity violating production of W^\pm*

This talk

- *Transverse spin:* connections to partonic orbital angular momentum \mathbf{L} and transversity Δ_T



Hard scattering of protons



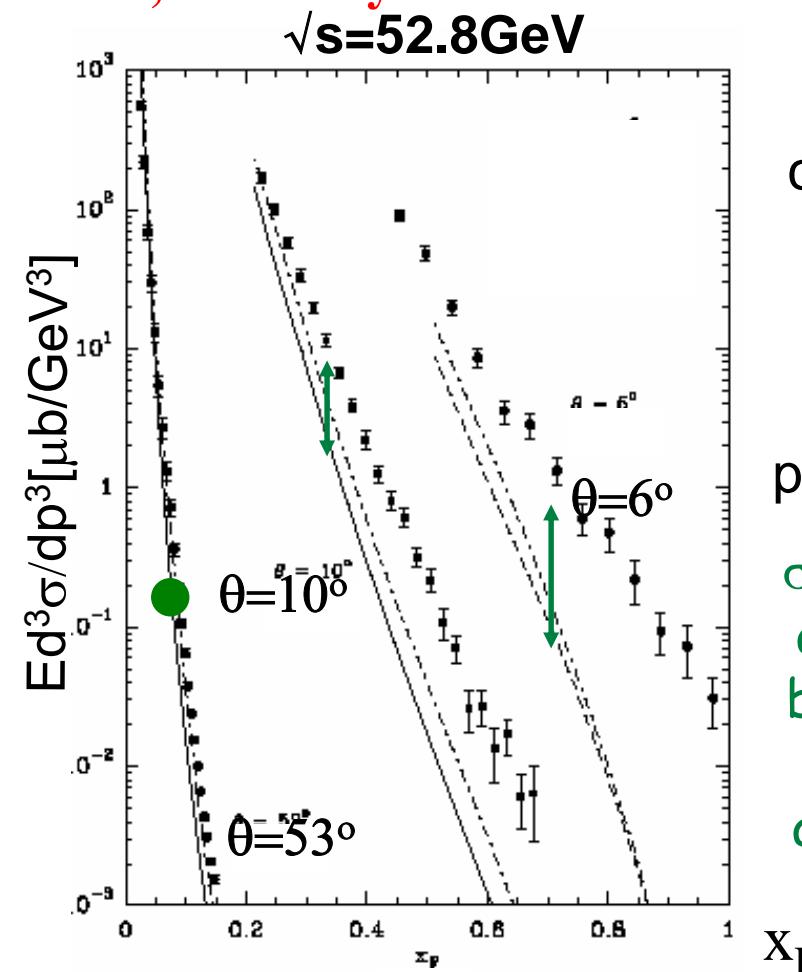
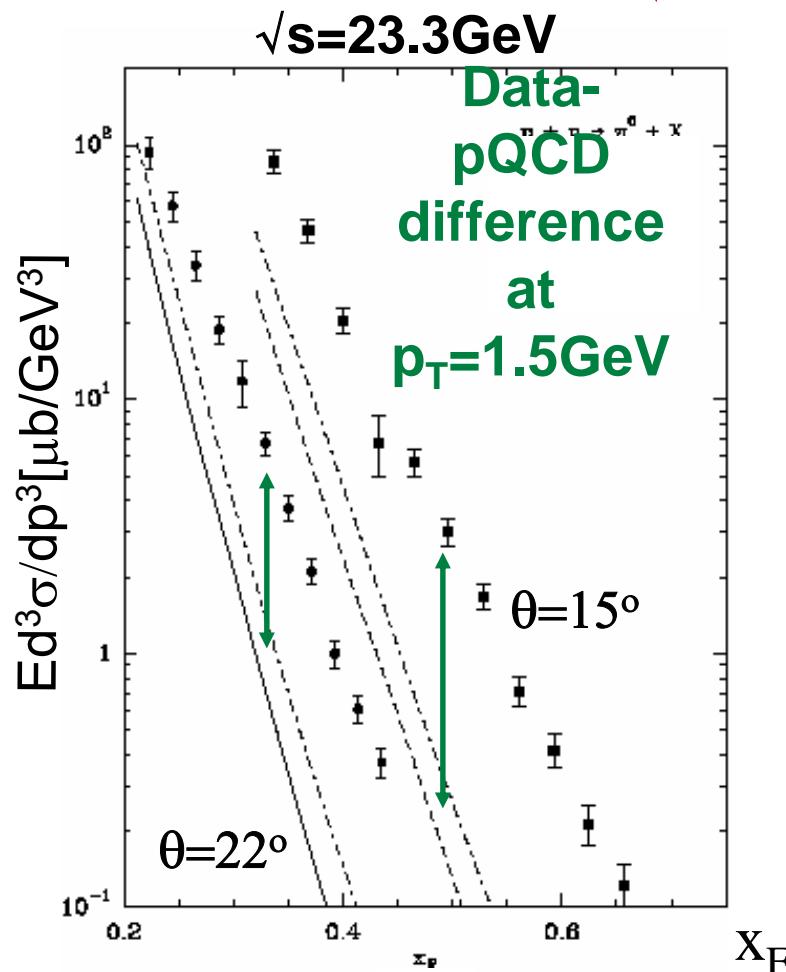
$$d\sigma = \sum_{a,b,c} \int dx_a \int dx_b \int dz_c f_a(x_a, \mu) f_b(x_b, \mu) D_c^\pi(z_c, \mu) \times d\hat{\sigma}_{ab}^c(x_a P_A, x_b P_B, P_\pi / z_c, \mu)$$

Using perturbative QCD at NLO and *universal* parton distribution functions and fragmentation functions. RHIC energy $\sqrt{s} = 200$ GeV.



But, do we understand forward π^0 production in $p + p$?

At $\sqrt{s} << 200 \text{ GeV}$, not really....

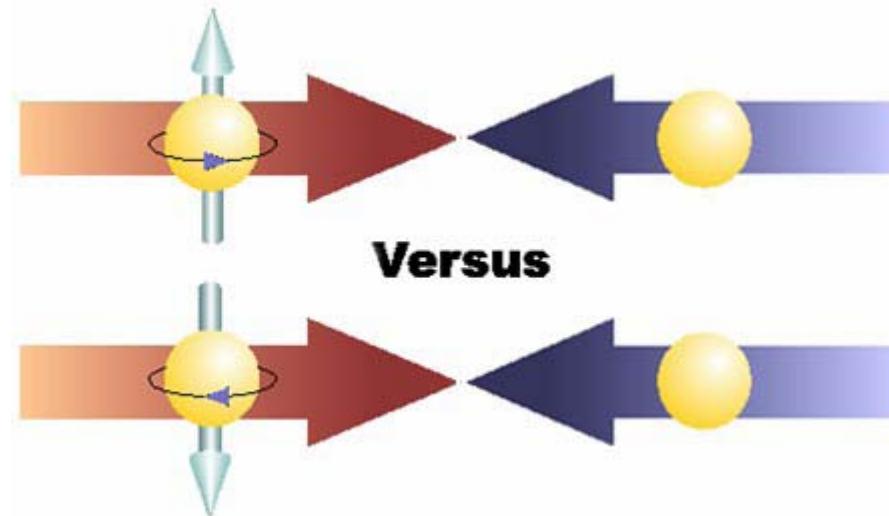


2 NLO calculation with different scale:
 p_T and $p_T/2$
 $\sigma_{\text{data}}/\sigma_{\text{pQCD}}$ appears to be function of θ, \sqrt{s} in addition to p_T

Bourrely and Soffer (Eur. Phys. J C36 (2004) 371):
NLO pQCD calculations underpredict the data at low \sqrt{s}



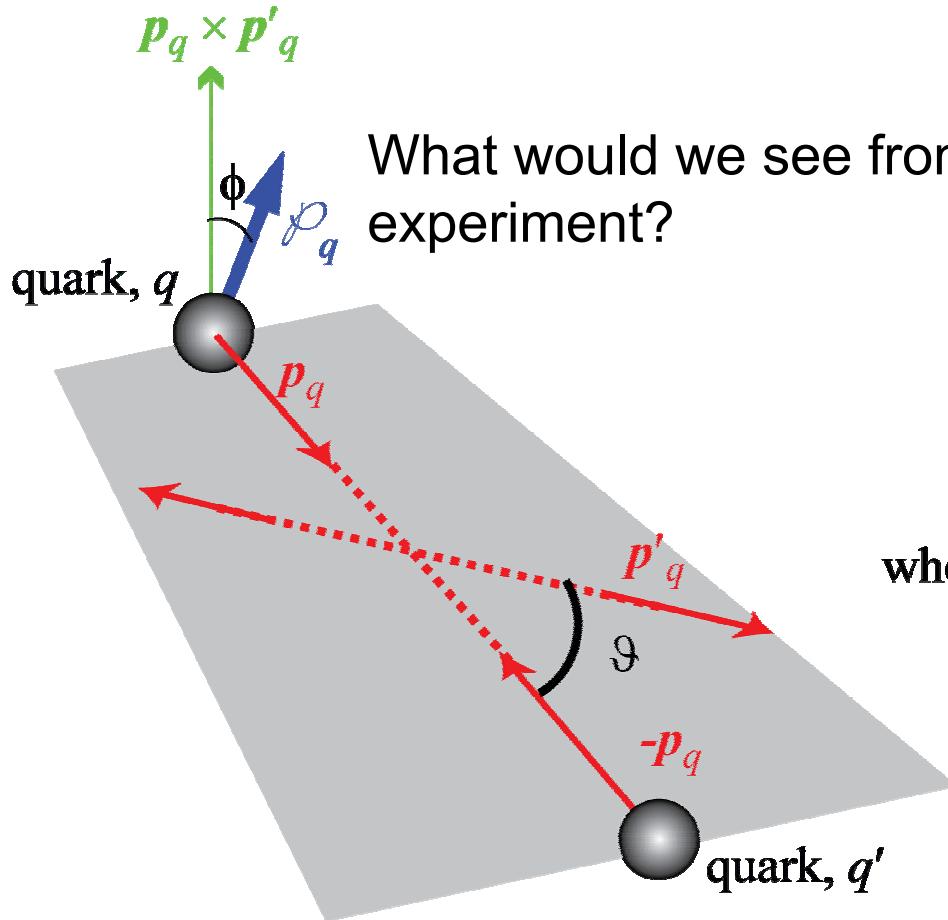
Transverse Single-Spin Asymmetries (A_N)



$$A_N = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$$



Expectations from Theory



What would we see from this gedanken experiment?

$$N(\vartheta, \phi) = N_0(\vartheta) [1 + \rho_q A_N \cos \phi]$$

$$A_N = \frac{\sigma_{\uparrow} - \sigma_{\downarrow}}{\sigma_{\uparrow} + \sigma_{\downarrow}}$$

where, ρ_q is quark transverse polarization;
 $A_N \propto \text{Im}(NF^*)$;
 N is the non-*helicity*-flip amplitude;
and F is the *helicity*-flip amplitude

$F \rightarrow 0$ as $m_q \rightarrow 0$ in vector gauge theories, so $A_N \sim m_q/p_T$

or, $A_N \sim 0.001$ for $p_T \sim 2 \text{ GeV}/c$

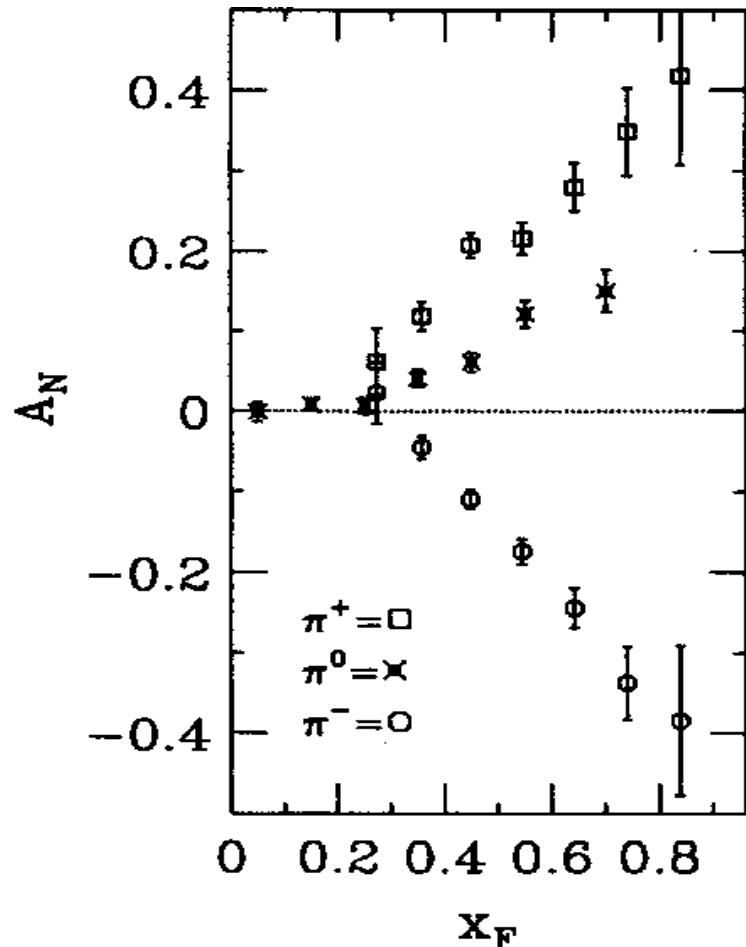
Kane, Pumplin and Repko PRL 41 (1978) 1689



History

$p_\uparrow + p \rightarrow \pi + X$

$\sqrt{s}=20 \text{ GeV}$, $p_T=0.5-2.0 \text{ GeV}/c$



π^0 – E704, PLB261 (1991) 201.

$\pi^{+/-}$ – E704, PLB264 (1991) 462.

- QCD theory expects small $A_N \sim 0.001$ transverse single spin asymmetries

- E-704 (Fermilab) – fixed target

Remember from slide 7 that:

Cross-section is **NOT** consistent
with NLO pQCD calculations

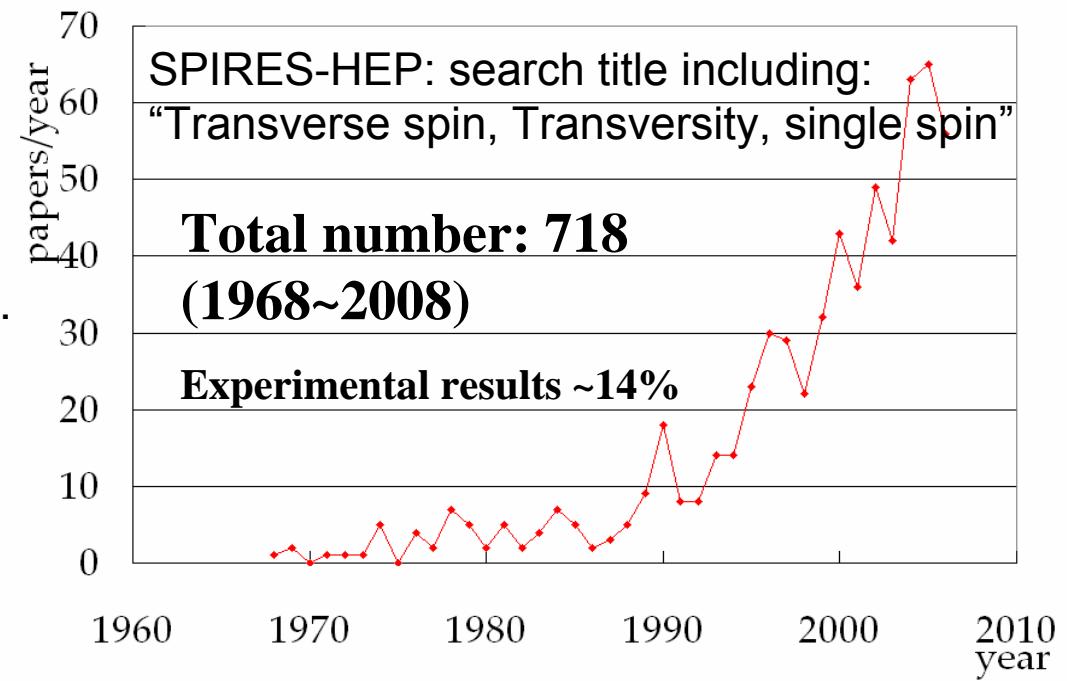


Transverse Single-Spin Asymmetries

World-wide experimental and theoretical efforts

- Transverse single-spin asymmetries are observed in semi-inclusive deep inelastic scattering with transversely polarized proton targets
 - ⇒ HERMES (e); COMPASS (μ); and planned at JLab
- Intense theory activity underway

• Four theory hires in "permanent" positions in the US within the past year.
All four were transverse spin theorists.

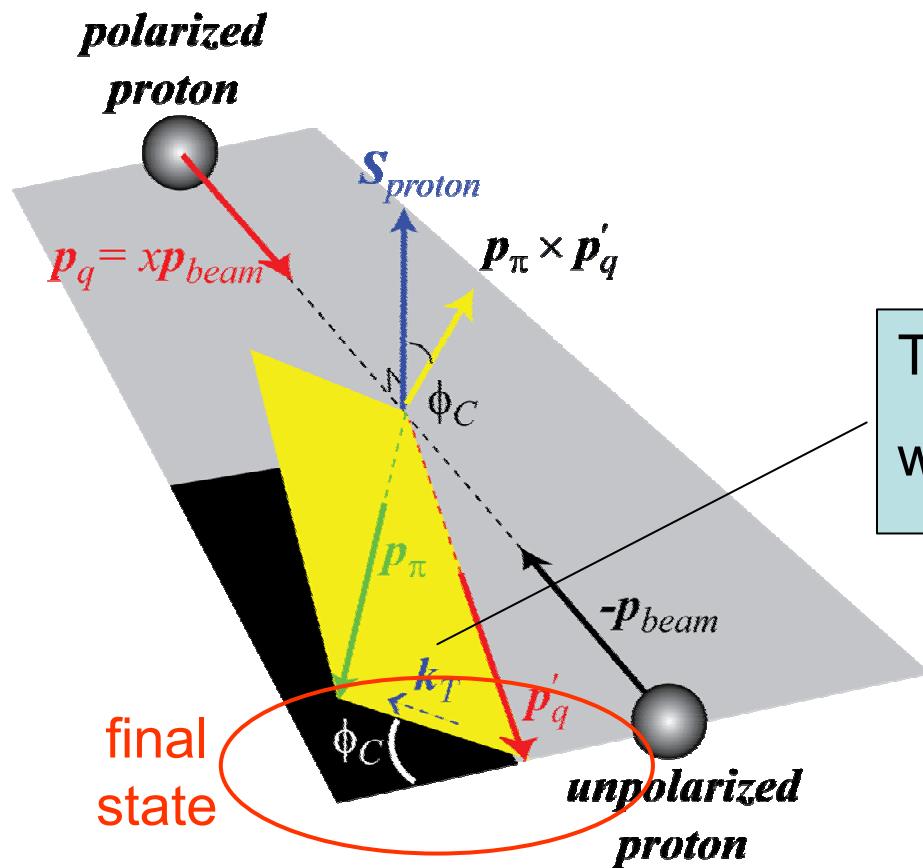


Possible explanations for Large Transverse SSA

Collins/Hepplemann mechanism

requires *transverse quark polarization* and
spin-dependent fragmentation

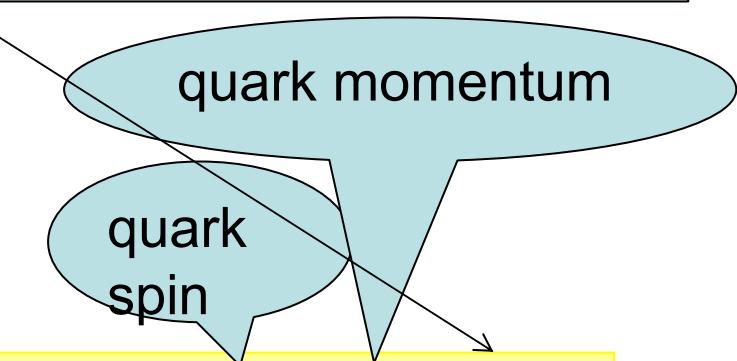
[Nucl. Phys. B396, 161 (1993)]:



asymmetry in the forward jet fragmentation

Sensitive to **transversity**

Transverse momentum of the hadron with respect to the fragmenting quark



$D_{\pi/q}(z, \mathbf{k}_\pi^\perp, \mathbf{s}_q) = \hat{D}_{\pi/q}(z, \mathbf{k}_\pi^\perp) + \frac{1}{2} \Delta^N D_{\pi/q}(z, \mathbf{k}_\pi^\perp) \frac{\mathbf{s}_q \cdot (\mathbf{p}_q \times \mathbf{k}_\pi^\perp)}{|\mathbf{p}_q \times \mathbf{k}_\pi^\perp|}$



Sivers mechanism

requires *spin-correlated transverse momentum* in the proton (orbital motion).

SSA is present for jet or γ

$$f_q(x, \mathbf{k}_q^\perp, \mathbf{S}_P) = f_q(x, k_q^\perp) + \frac{1}{2} \Delta_q^N f_q(x, k_q^\perp) \frac{\mathbf{S}_P \cdot (\mathbf{P}_P \times \mathbf{k}_q^\perp)}{|\mathbf{S}_P| |\mathbf{P}_P| |\mathbf{k}_q^\perp|}$$

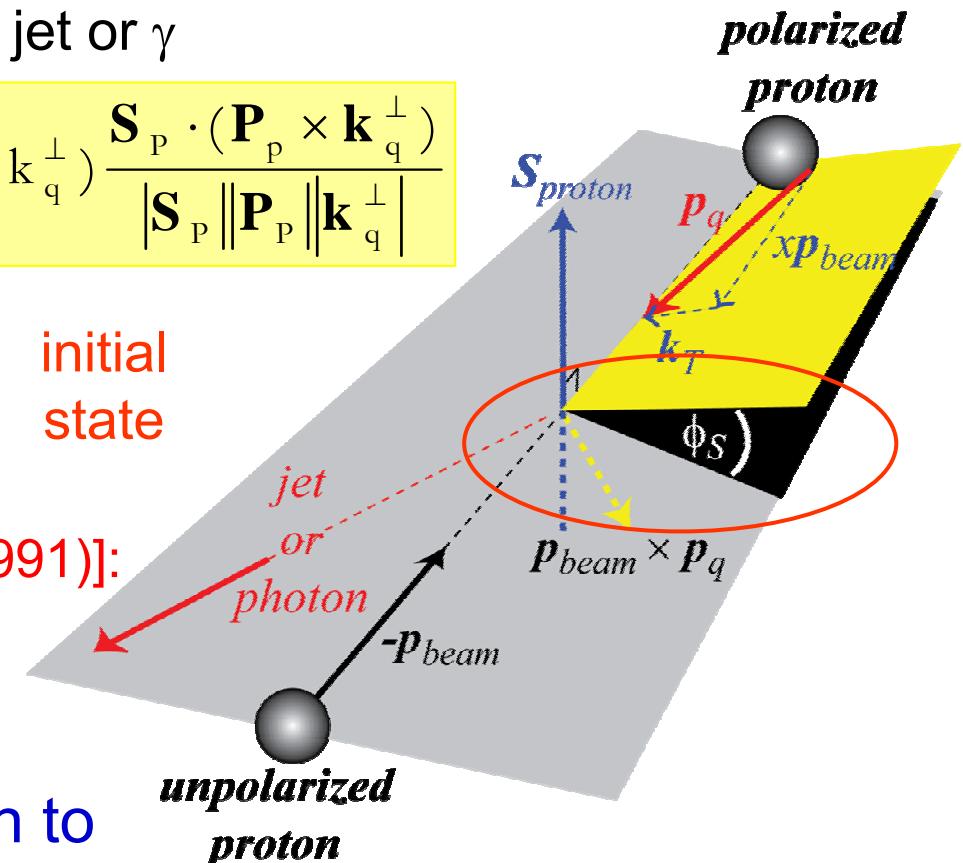
[Phys. Rev. D 41, 83 (1990); 43, 261 (1991)]:

Sensitive to **proton spin – parton transverse motion** correlations

Need to go beyond π^0 detection to

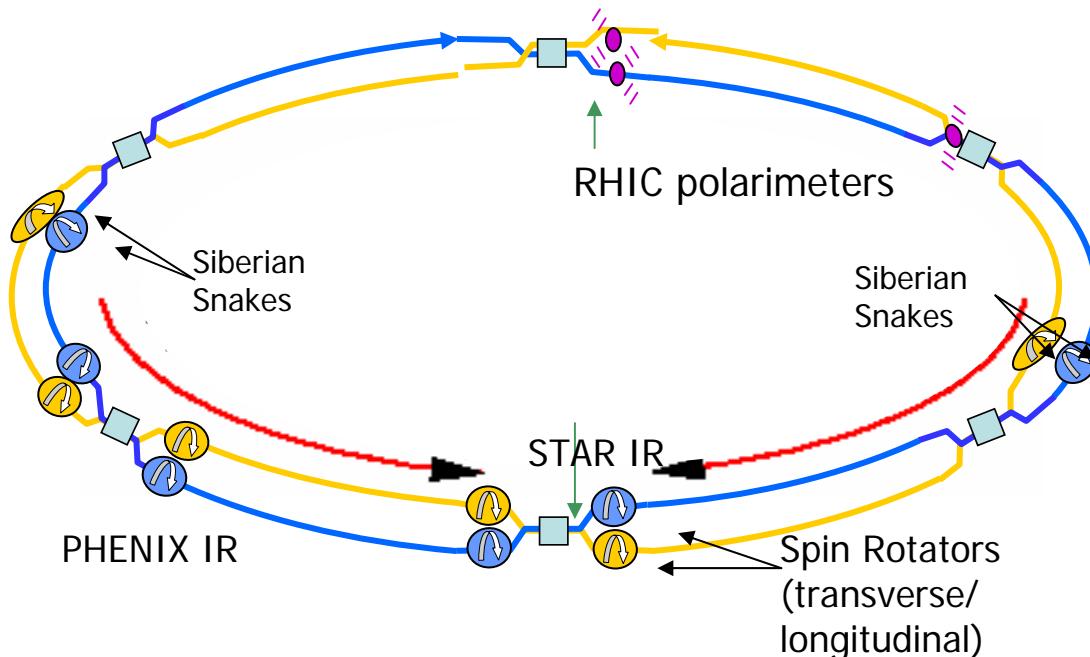
jets and direct photons

Require experimental separation of Collins and Sivers contributions



Relativistic Heavy Ion Collider

...worlds 1st $\vec{p}\vec{p}$ Collider



- 100 GeV beam proton beams
- Each bunch filled with a distinct polarization state
- Spin Rotators at STAR IR allow for transverse and longitudinal spin orientation
- Bunch Xings every 100-200ns
- CNI polarimeters + Hydrogen Jet target provide run by run & absolute polarization

pp Run Year	FOM= $P^2 L$	2002	2003	2004	2005	2006
< Polarization> %		15	30	40-45	45-50	60
L_{max} [$10^{30} \text{ s}^{-1} \text{cm}^{-2}$]		2	6	6	16	20
L_{int} [pb^{-1}] at STAR (T)		0.15	0.25	0	0.1	6.8



STAR Collaboration

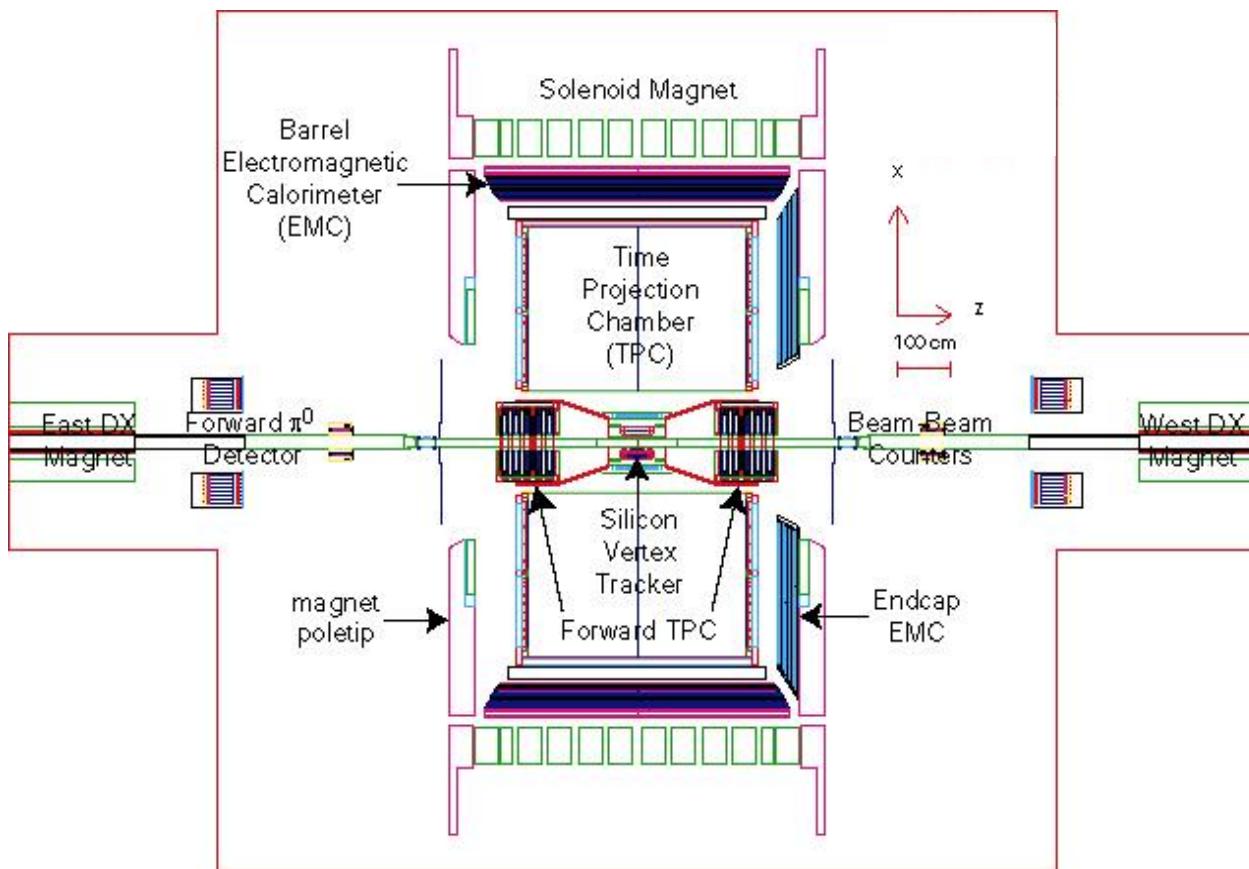
[B.I. Abelev](#), [M.M. Aggarwal](#), [Z. Ahammed](#), [B.D. Anderson](#), [D. Arkhipkin](#), [G.S. Averichev](#), [Y. Bai](#), [J. Balewski](#), [O. Barannikova](#), [I.S. Barnby](#), [J. Baudot](#), [S. Baumgart](#), [V.V. Belaga](#), [A. Bellingeri-Laurikainen](#), [R. Bellwied](#), [F. Benedosso](#), [R.R. Betts](#), [S. Bhardwaj](#), [A. Bhasin](#), [A.K. Bhati](#), [H. Bichsel](#), [J. Bielcik](#), [J. Bielcikova](#), [L.C. Bland](#), [S.-L. Blyth](#), [M. Bombara](#), [B.E. Bonner](#), [M. Bolje](#), [J. Bouchet](#), [A.V. Brandin](#), [A. Bravar](#), [T.P. Burton](#), [M. Bystersky](#), [X.Z. Cai](#), [H. Caines](#), [M. Calder](#), [M. Calderon de la Barca Sanchez](#), [J. Callner](#), [O. Catu](#), [D.A. Cebra](#), [M.C. Cervantes](#), [Z. Chajecki](#), [P. Chaloupka](#), [S. Chattopadhyay](#), [H.F. Chen](#), [J.H. Chen](#), [J.Y. Chen](#), [J. Cheng](#), [M. Cherney](#), [A. Chikianian](#), [W. Christie](#), [S.U. Chung](#), [R.F. Clarke](#), [M.J.M. Codrington](#), [J.P. Coffin](#), [T.M. Cormier](#), [M.R. Cosentino](#), [J.G. Cramer](#), [H.J. Crawford](#), [D. Das](#), [S. Dash](#), [M. Daugherty](#), [M.M. de Moura](#), [T.G. Dedovich](#), [M. DePhillips](#), [A.A. Derevschikov](#), [L. Didenko](#), [T. Dietel](#), [P. Djawotho](#), [S.M. Dogra](#), [X. Dong](#), [J.L. Drachenberg](#), [J.F. Draper](#), [F. Du](#), [V.B. Dunin](#), [J.C. Dunlop](#), [M.R. Dutta Mazumdar](#), [V. Eckardt](#), [W.R. Edwards](#), [L.G. Efimov](#), [V. Emelianov](#), [J. Engelage](#), [G. Eppley](#), [B. Erazmus](#), [M. Estienne](#), [P. Fachini](#), [R. Fatemi](#), [J. Fedorisin](#), [A. Feng](#), [P. Filip](#), [E. Finch](#), [V. Fine](#), [Y. Fisyak](#), [J. Fu](#), [C.A. Gagliardi](#), [L. Gaillard](#), [M.S. Ganti](#), [E. Garcia-Solis](#), [V. Ghazikhanian](#), [P. Ghosh](#), [Y.N. Gorbunov](#), [H. Gos](#), [O. Grebenyuk](#), [D. Grosnick](#), [B. Grube](#), [S.M. Guertin](#), [K.S.F.F. Guimaraes](#), [N. Gupta](#), [B. Haag](#), [T.J. Hallman](#), [A. 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[Argonne](#) & [Birmingham U.](#) & [Brookhaven](#) & [Caltech](#) & [UC, Berkeley](#) & [UC, Davis](#) & [UCLA](#) & [Carnegie Mellon U.](#) & [Illinois U.](#), [Chicago](#) & [Creighton U.](#) & [Rez, Nucl. Phys. Inst.](#) & [Dubna, JINR](#) & [Frankfurt U.](#) & [Bhubaneswar, Inst. Phys.](#) & [Indian Inst. Tech., Mumbai](#) & [Indiana U.](#) & [Strasbourg, IReS](#) & [Jammu U.](#) & [Kent State U.](#) & [Lanzhou, Inst. Modern Phys.](#) & [LBL, Berkeley](#) & [MIT, LNS](#) & [Munich, Max Planck Inst.](#) & [Michigan State U.](#) & [Moscow Phys. Eng. Inst.](#) & [City Coll., N.Y.](#) & [NIKHEF, Amsterdam](#) & [Utrecht U.](#) & [Ohio State U.](#) & [Panjab U.](#) & [Penn State U.](#) & [Serpukhov, IHEP](#) & [Purdue U.](#) & [Pusan Natl. U.](#) & [Rajasthan U.](#) & [Rice U.](#) & [Sao Paulo U.](#) & [Hefei, CUST](#) & [SINAP, Shanghai](#) & [SUBATECH, Nantes](#) & [Texas A-M](#) & [Texas U.](#) & [Tsinghua U., Beijing](#) & [Valparaiso U., Indiana](#) & [Calcutta, VECC](#) & [Warsaw U. of Tech.](#) & [Washington U., Seattle](#) & [Wayne State U.](#) & [Hua-Zhong Normal U.](#) & [Yale U.](#) & [Zagreb U.](#)





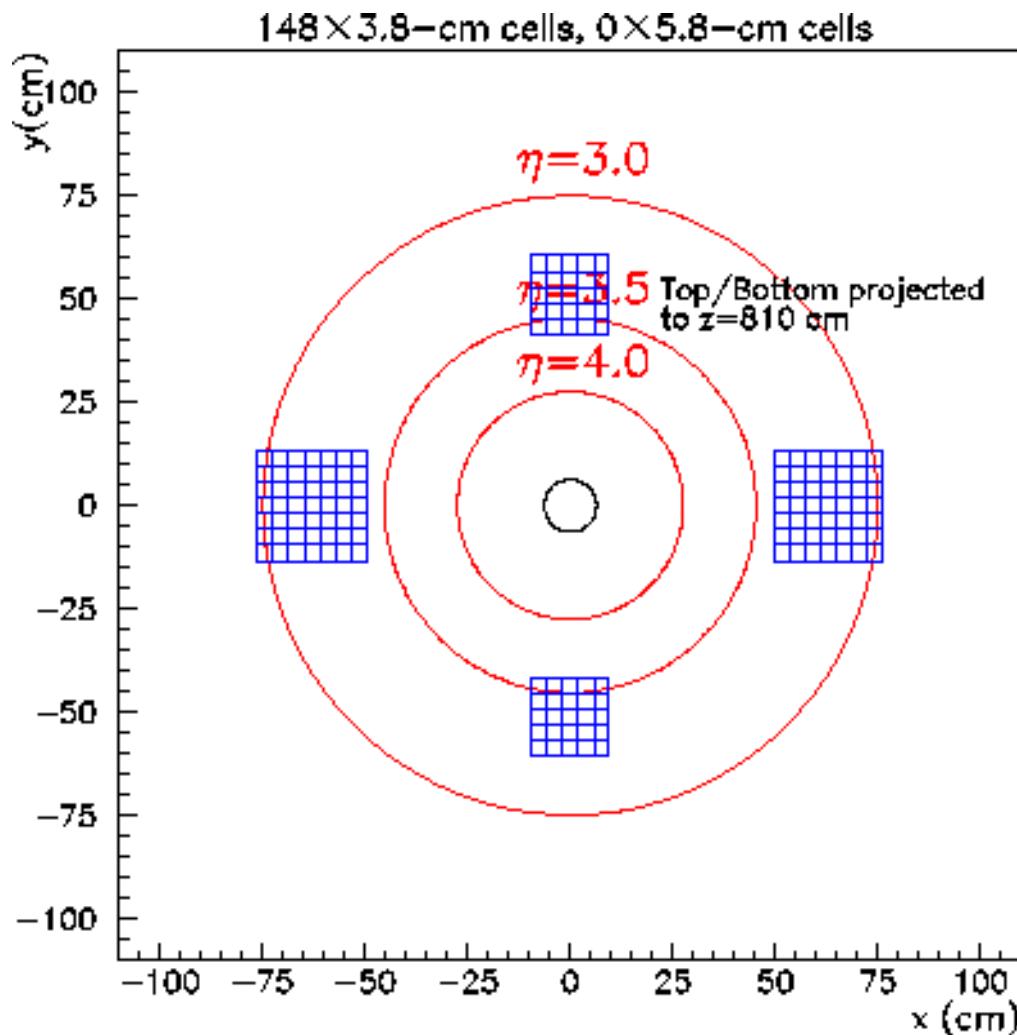
STAR detector layout



- **TPC:** $-1.0 < \eta < 1.0$
- **FTPC:** $2.8 < |\eta| < 3.8$
- **BBC :** $2.2 < |\eta| < 5.0$
- **EEMC:** $1 < \eta < 2$
- **BEMC:** $0 < \eta < 1$
- **FPD:** $|\eta| \sim 4.0 \text{ \& } \sim 3.7$

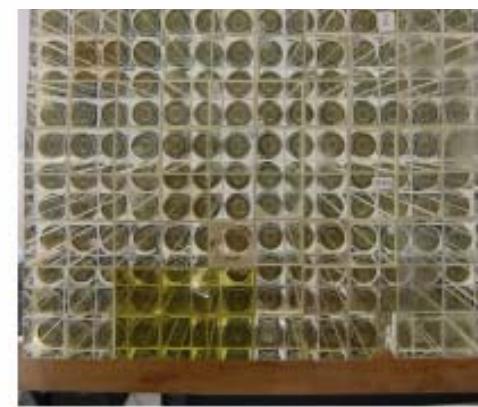


FPD-Forward Pion Detector



(3.8cm)² x 45 cm lead glass
18 radiation lengths
FEU84 + XP2972 photomultipliers

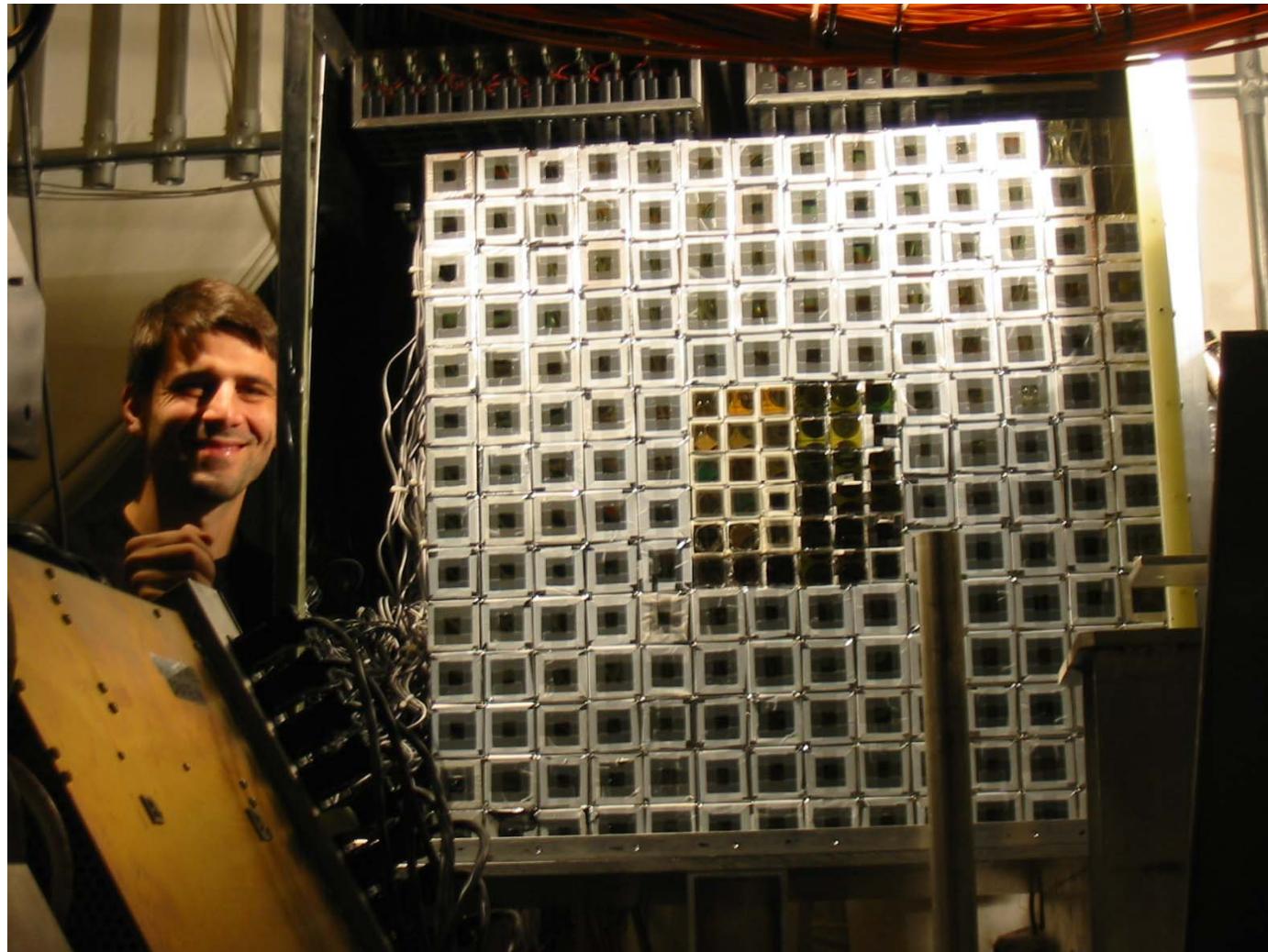
17



170 small cells prior to wrapping



Run 6 – FPD++



TPC: $-1.0 < \eta < 1.0$

FTPC: $2.8 < |\eta| < 3.8$

BBC : $2.2 < |\eta| < 5.0$

EEMC: $1 < \eta < 2$

BEMC: $-1 < \eta < 1$

FPD++/FPD:

$$\eta \sim 3.3/3.7$$

12/30/2005

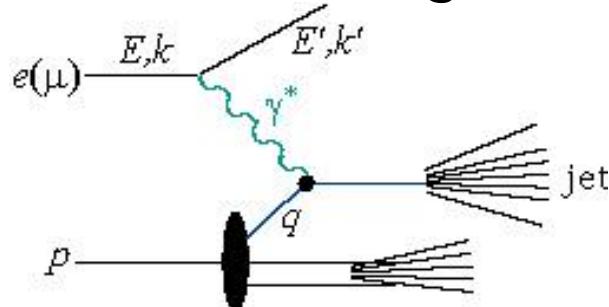


Mirko Planinić



Why Consider Forward Physics at a Collider?

Deep inelastic scattering



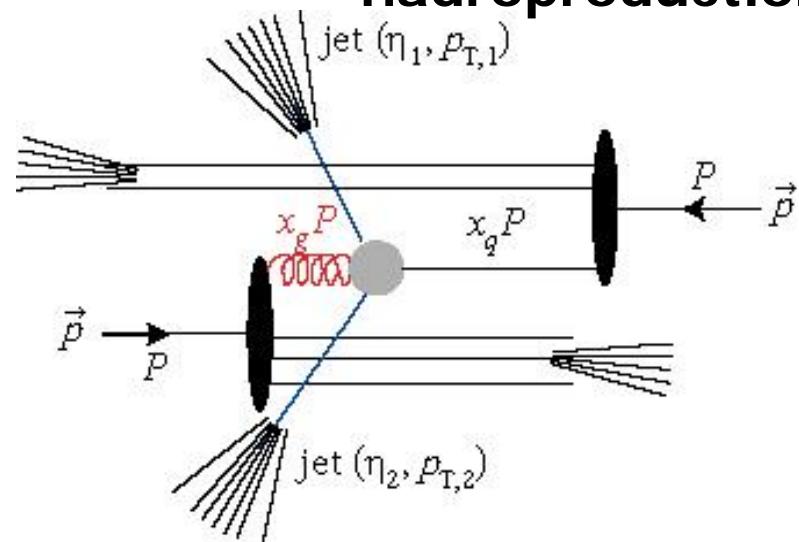
$$Q^2 = 2(EE' - \vec{k} \cdot \vec{k}')$$

$$\nu = E - E'$$

$$x = Q^2 / 2M\nu$$

Kinematics

Hard scattering hadroproduction



How can Bjorken x values be selected in hard scattering?

Assume:

1. Initial partons are collinear
2. Partonic int. is elastic $\Rightarrow p_{T,1} \approx p_{T,2}$

$$\Rightarrow x_q \approx p_T / \sqrt{s} (e^{+\eta_1} + e^{+\eta_2})$$

$$x_g \approx p_T / \sqrt{s} (e^{-\eta_1} + e^{-\eta_2})$$

Studying pseudorapidity, $\eta = -\ln(\tan\theta/2)$, dependence of particle production probes parton distributions at different Bjorken x values and involves different admixtures of gg, qg and qq' subprocesses.



Simple Kinematic Limits

Mid-rapidity particle detection:

$$\eta_1 \approx 0 \text{ and } \langle \eta_2 \rangle \approx 0$$

$$\Rightarrow x_q \approx x_g \approx x_T = 2 p_T / \sqrt{s}$$

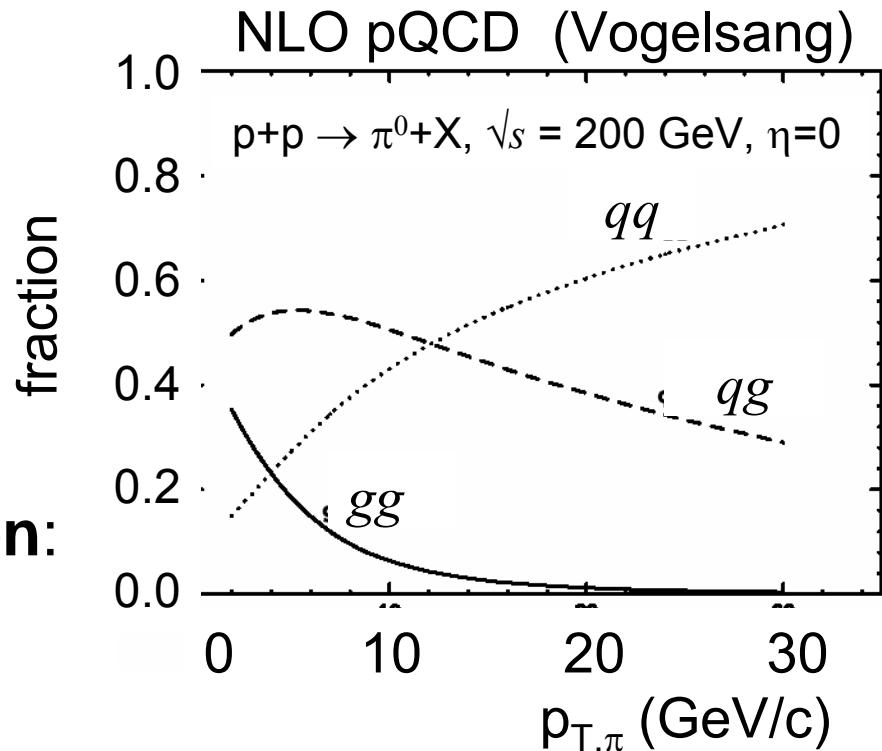
Large-rapidity particle detection:

$$\eta_1 \gg \eta_2$$

$$\Rightarrow x_q \approx x_T e^{\eta_1} \approx x_F \text{ (Feynman } x\text{), and}$$

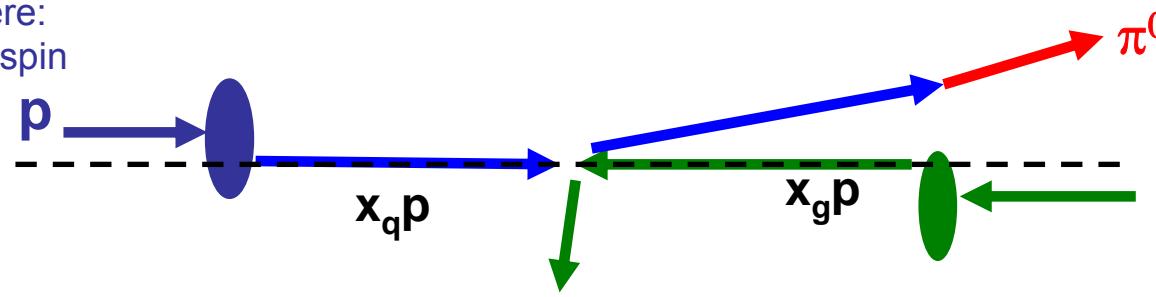
$$x_g \approx x_F e^{-(\eta_1 + \eta_2)}$$

⇒ Large rapidity particle production and correlations involving large rapidity particle probes low- x parton distributions using valence quarks



Recent asymmetry measurements at STAR

Polarization here:
valence quark spin
effects



Polarization here: low
x-gluons and other
partons

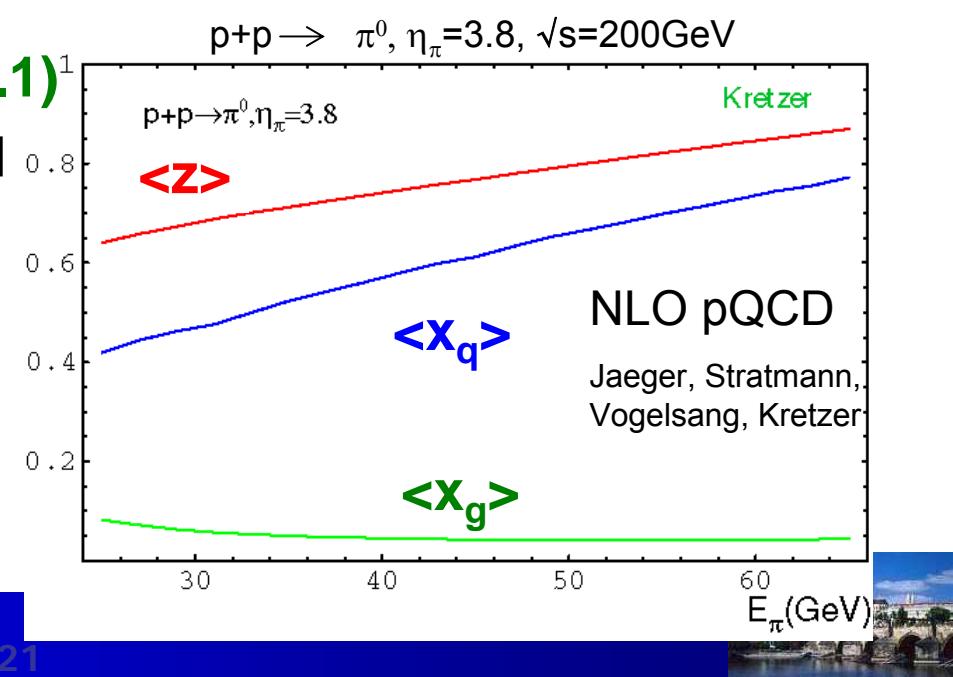
High rapidity π 's ($\eta_\pi \sim 4$) from asymmetric partonic collisions

Mostly high-x valence quark on low-x gluon

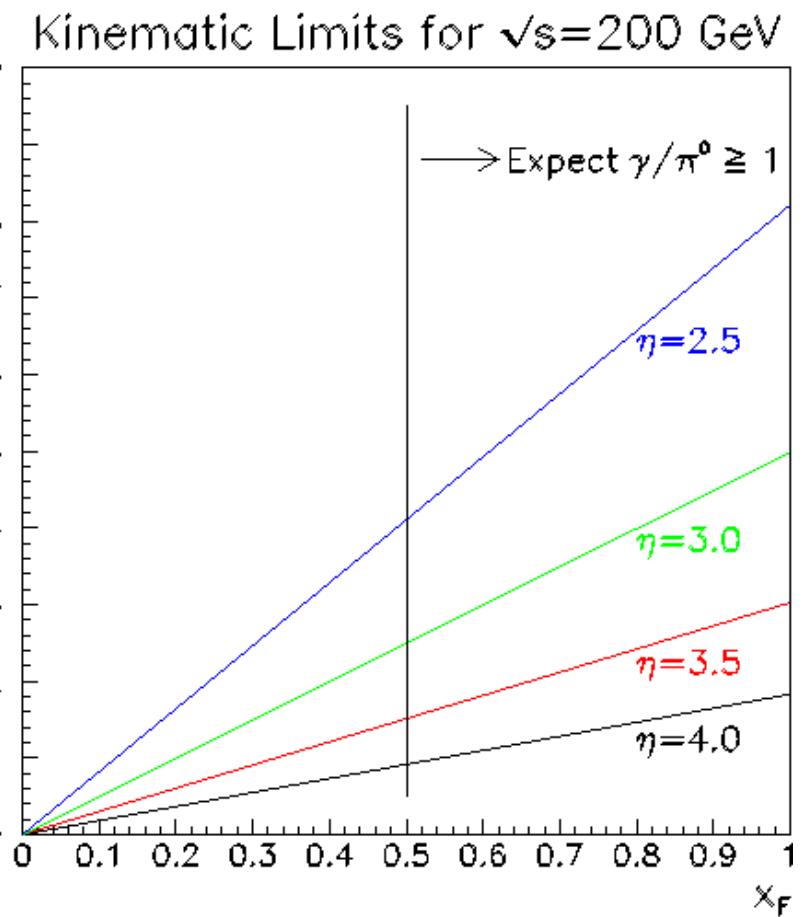
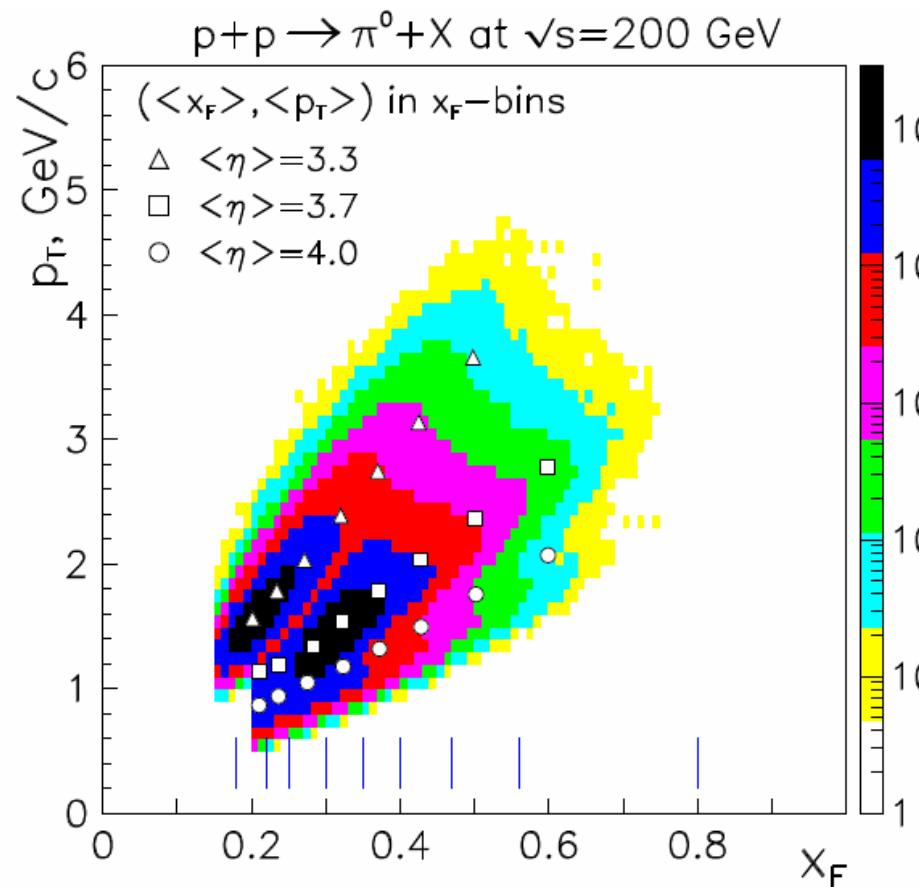
$$(0.3 < x_q < 0.7, \quad 0.001 < x_g < 0.1)$$

Fragmentation z nearly constant and
high $0.7 \sim 0.8$

- Large-x quark polarization is known to be large from DIS
- Directly couple to gluons = A probe of low x gluons

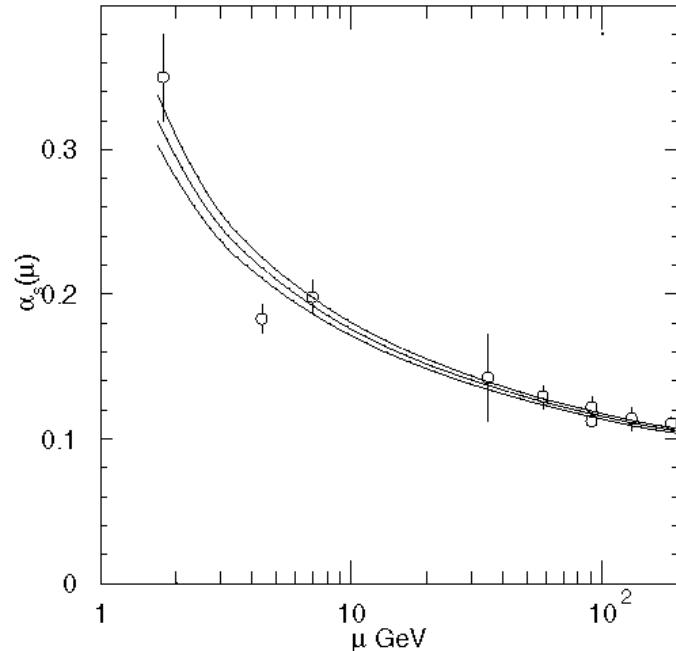


x_F and p_T range of FPD data



Possible Problems at Forward Angles

- Is it possible to access large enough p_T where NLO pQCD is applicable?



Although α_s does not vary much over accessible scales at RHIC, large η will primarily probe small $p_T \Rightarrow$ need to understand scale dependence of fixed order pQCD calculations.

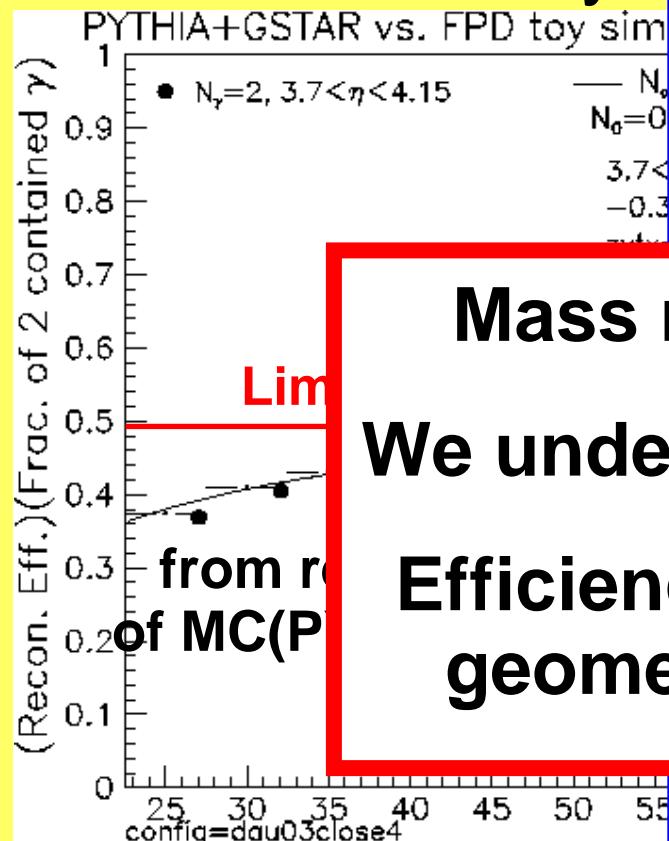
- Large x_F means high energy particles. Detection is best accomplished using electromagnetic + hadronic calorimetry + charge-sign determination from tracking through a magnetic field.
- For increasing p_T at large x_F , faced with increasingly steep falloff of $dN/d\eta$ distributions.



Di-photon Mass Reconstruction and calibration

Ph-glass reconstruction (no SMD)

π^0 reconstruction
efficiency



Categorization

MC & Data comparison



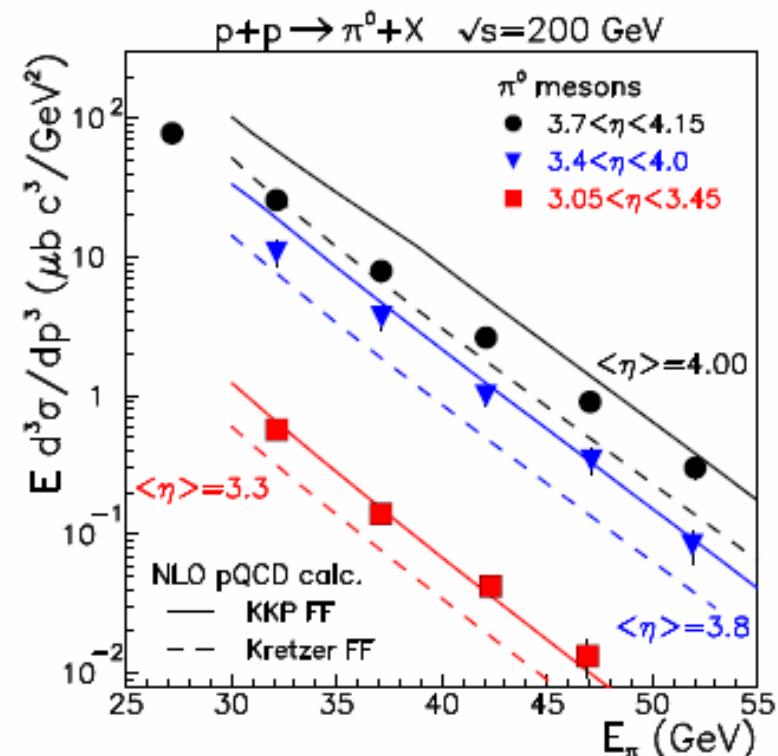
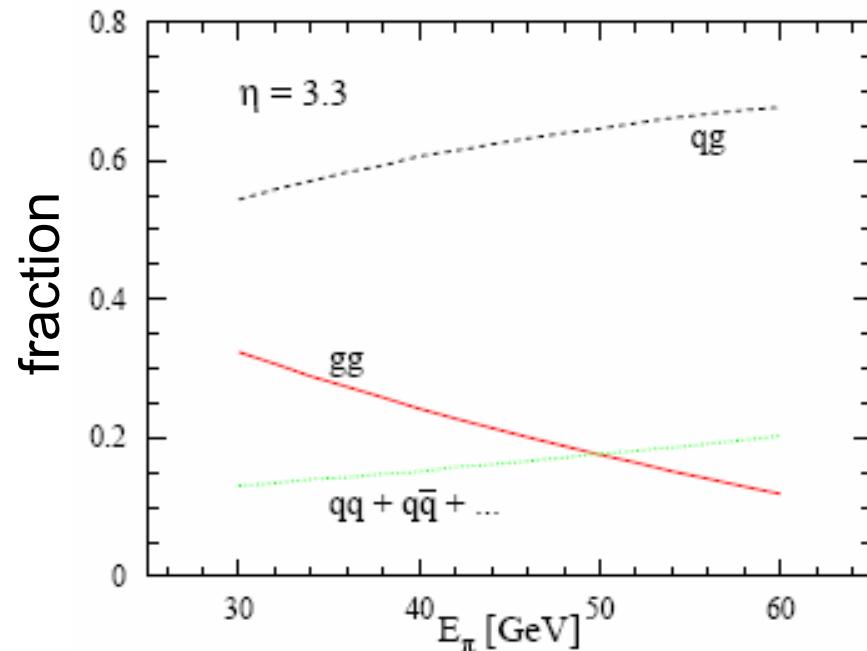
Mass resolution ~ 20MeV

We understand gain ~2% level

Efficiencies is almost purely
geometrically determined

Cross section measurements: Forward pion production

Phys. Lett. B603 (2004) 173



Forward production is dominated by asymmetric qg collisions

Data compares favorably to NLO pQCD at $\sqrt{s} = 200$ GeV in contrast to fixed-target or ISR energies



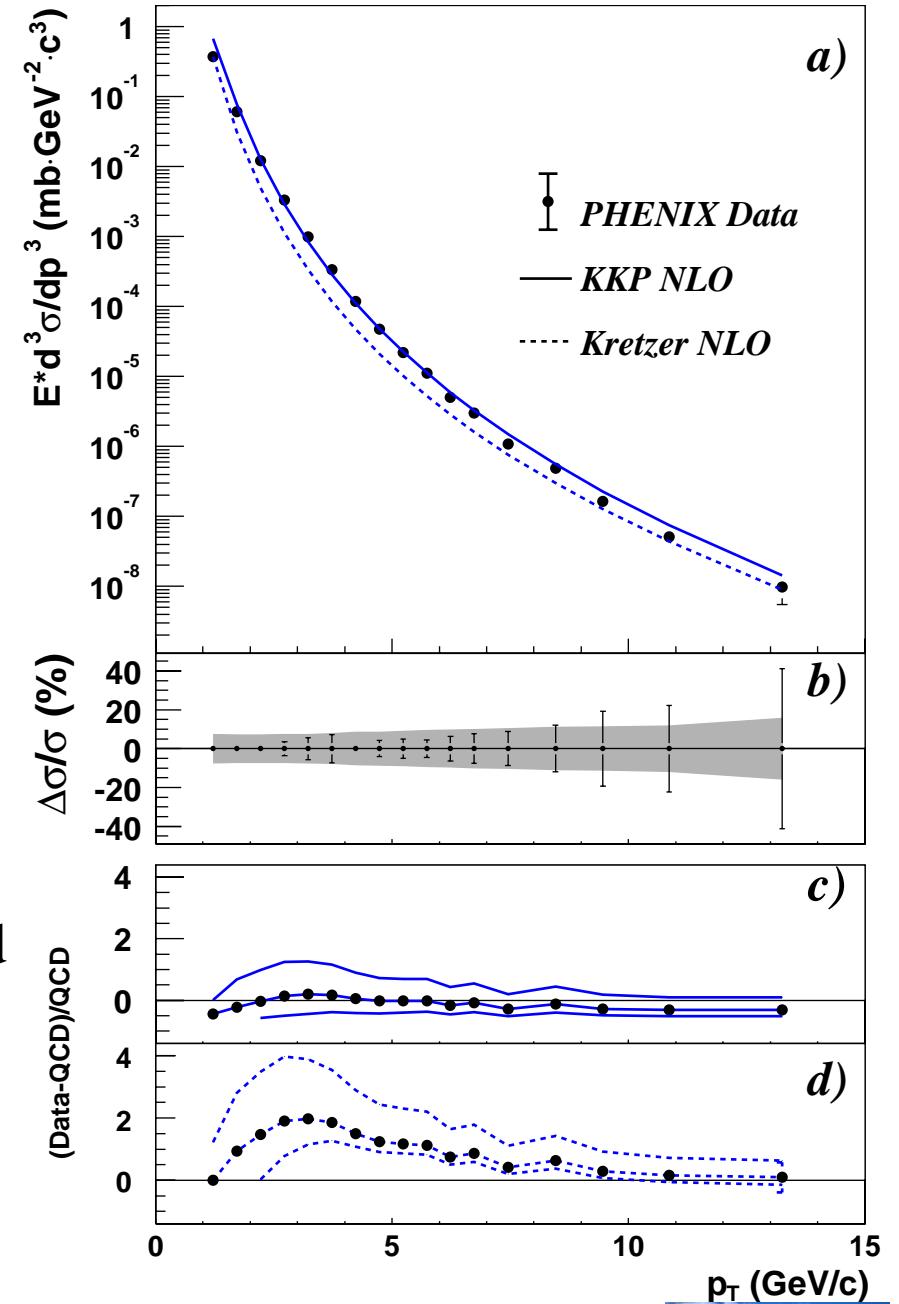
π^0 production at midrapidity

$$p + p \rightarrow \pi^0 + X, \sqrt{s} = 200 \text{ GeV}$$

S.S. Adler *et al.* (PHENIX), PRL 91, 241803 (2003).

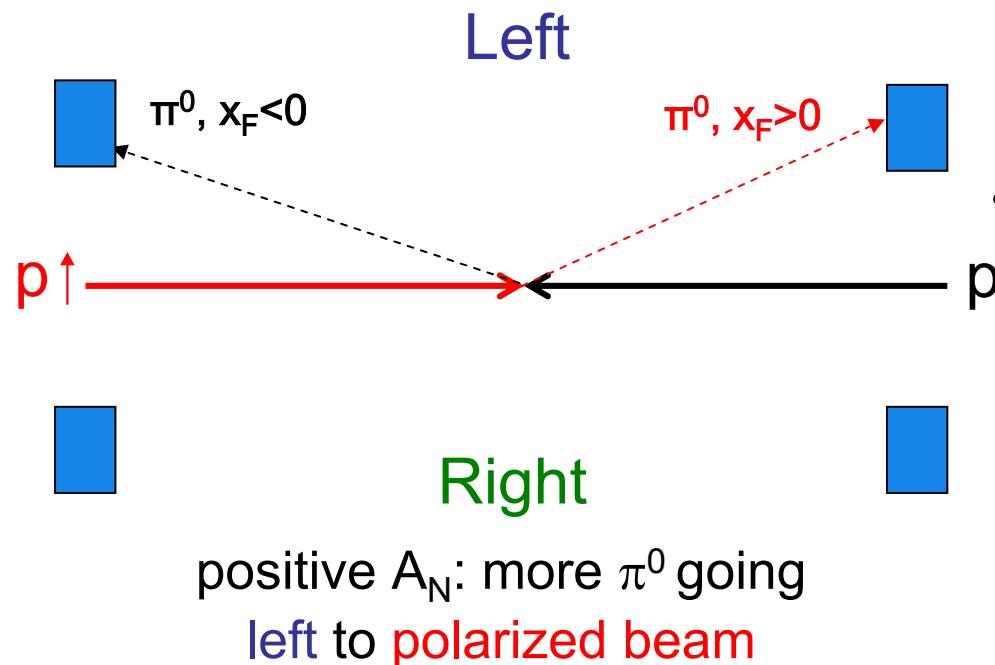
NLO pQCD calculation, using CTEQ5M PDF and KKP fragmentation functions is found to be consistent with data down to surprisingly low p_T .

Universality tests at collider energies yield comparable results.



Single Spin Asymmetry

- Definition: $A_N = \frac{d\sigma^{\uparrow} - d\sigma^{\downarrow}}{d\sigma^{\uparrow} + d\sigma^{\downarrow}}$
- $d\sigma^{\uparrow(\downarrow)}$ – differential cross section of π^0 when incoming proton has spin up(down)



Two methods of measurements:

- Single arm calorimeter:

$$A_N = \frac{1}{P_{beam}} \cdot \left(\frac{N^{\uparrow} - RN^{\downarrow}}{N^{\uparrow} + RN^{\downarrow}} \right) \quad R = \frac{L^{\uparrow}}{L^{\downarrow}}$$

R – relative luminosity (by BBC)

P_{beam} – beam polarization

- Two arm (left-right) calorimeter:

$$A_N = \frac{1}{P_{Beam}} \cdot \left(\frac{\sqrt{N_L^{\uparrow} \cdot N_R^{\downarrow}} - \sqrt{N_R^{\uparrow} \cdot N_L^{\downarrow}}}{\sqrt{N_L^{\uparrow} \cdot N_R^{\downarrow}} + \sqrt{N_R^{\uparrow} \cdot N_L^{\downarrow}}} \right)$$

No relative luminosity needed



Transverse spin runs at STAR with forward calorimetry: 2001→2006

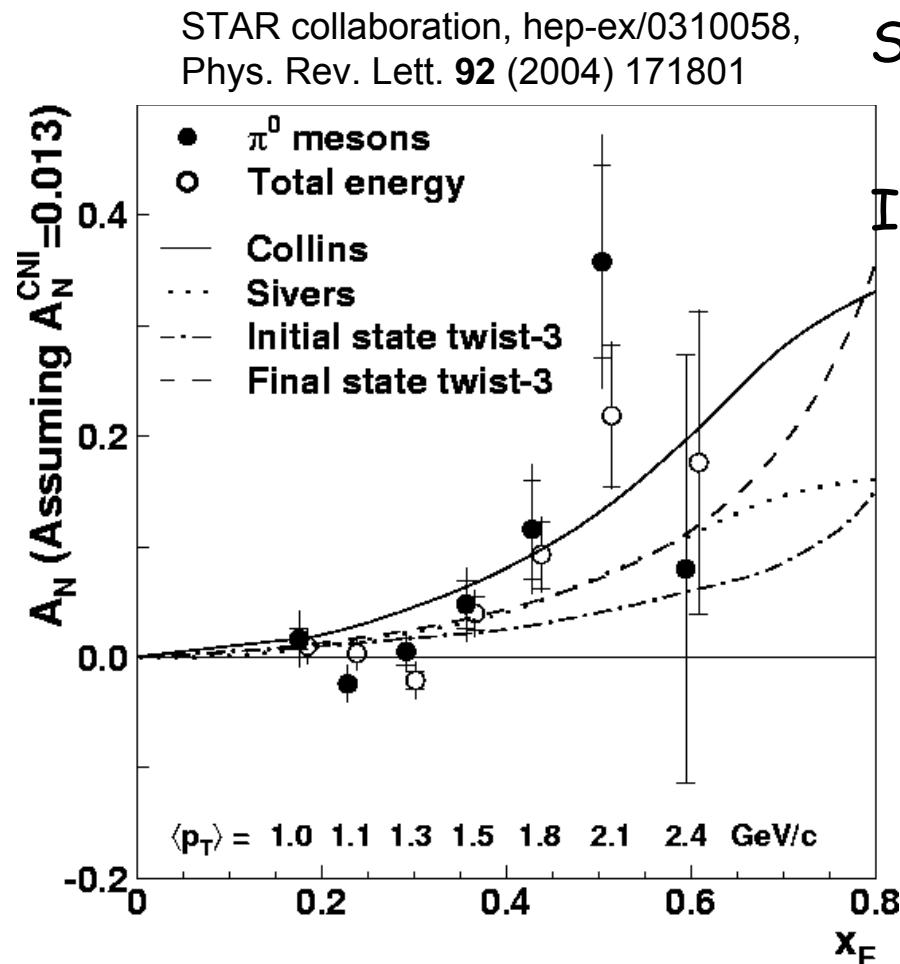
	Run2	Run3	Run5	Run6
detector	EEMC and FPD prototypes	6 matrices of FPD	full FPD (8 matrices)	East FPD West FPD++
P_{BEAM} , %	~15	~30	~45	~60
$\int L dt$, pb ⁻¹	0.15	0.25	0.1	6.8
$< \eta >$	3.8	$\pm 3.3/\pm 4.0$	$\pm 3.7/\pm 4.0$	-3.7/3.3

FOM ($P^2 L$) in Run 6 is ~50 times larger than from all the previous STAR runs



Large Analyzing Powers at RHIC

First measurement of A_N for forward π^0 production at $\sqrt{s}=200\text{GeV}$



Similar to FNAL E704 result at $\sqrt{s} = 20$ GeV

In agreement with several models including different dynamics:

- Sivers: spin and k_\perp correlation in initial state (related to orbital angular momentum?)
- Collins: Transversity distribution function & spin-dependent fragmentation function
- Qiu and Sterman (initial-state) / Koike (final-state) twist-3 pQCD calculations



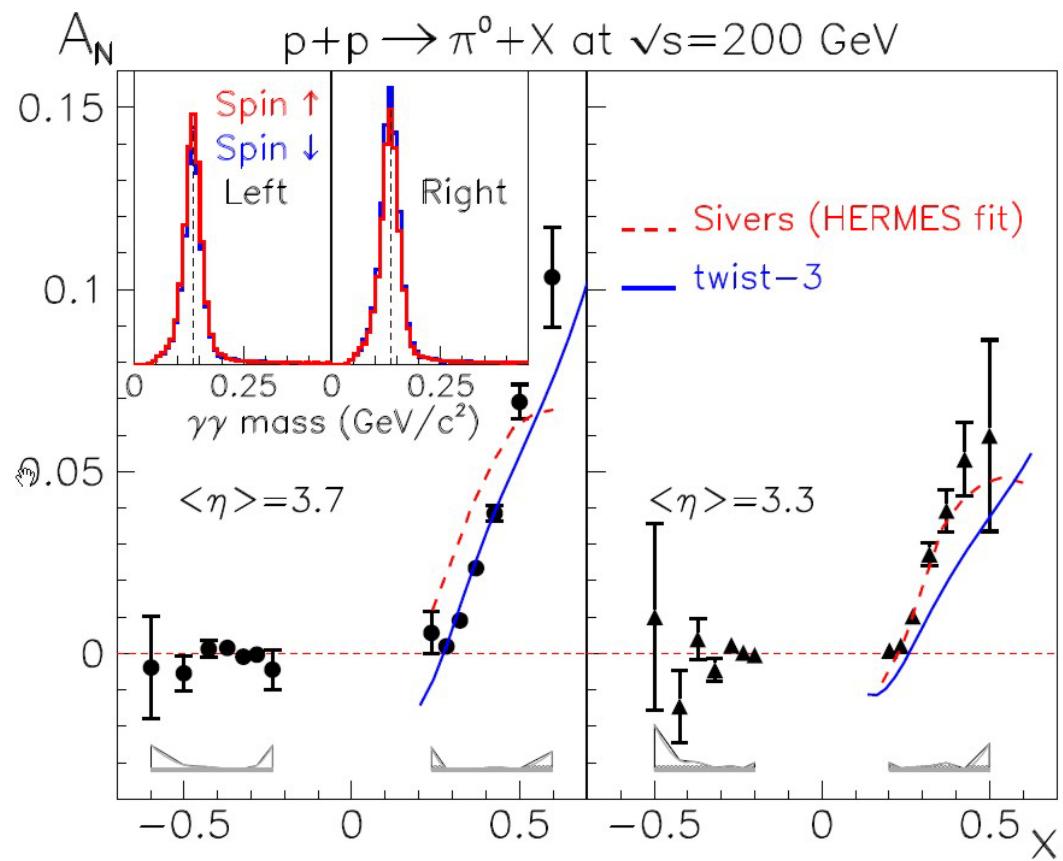
High Precision Analyzing Powers

(red line) M. Bogiljone, U. D'Alesio, F. Murgia,

PRD 77 (2008) 051502

(blue line) C. Kourvaris, J. Qiu, W. Vogelsang, F. Yuan,

PRD 74 (2006)



(2003 –
2006)

B.I. Abelev, *et al*,
hep-ex/0801.2990

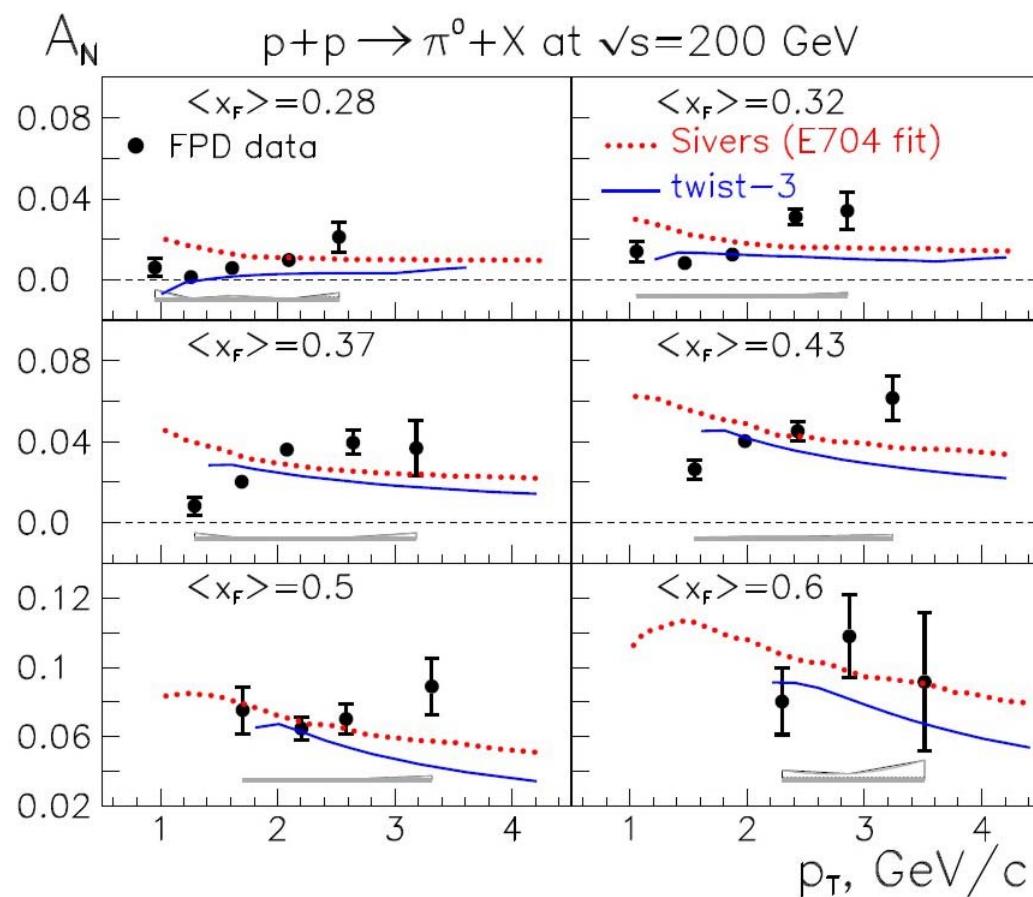
Null at $x_F < 0$ is natural since gluon Sivers function is probed where unp. gluon distribution is large.

→ Precision measurements at $\sqrt{s} = 200$ GeV provide stringent constraints on the models...



High Precision Analyzing Powers

Data
broken out
in X_F bins



(2003 –
2006)

B.I. Abelev, et al,
hep-ex/0801.2990

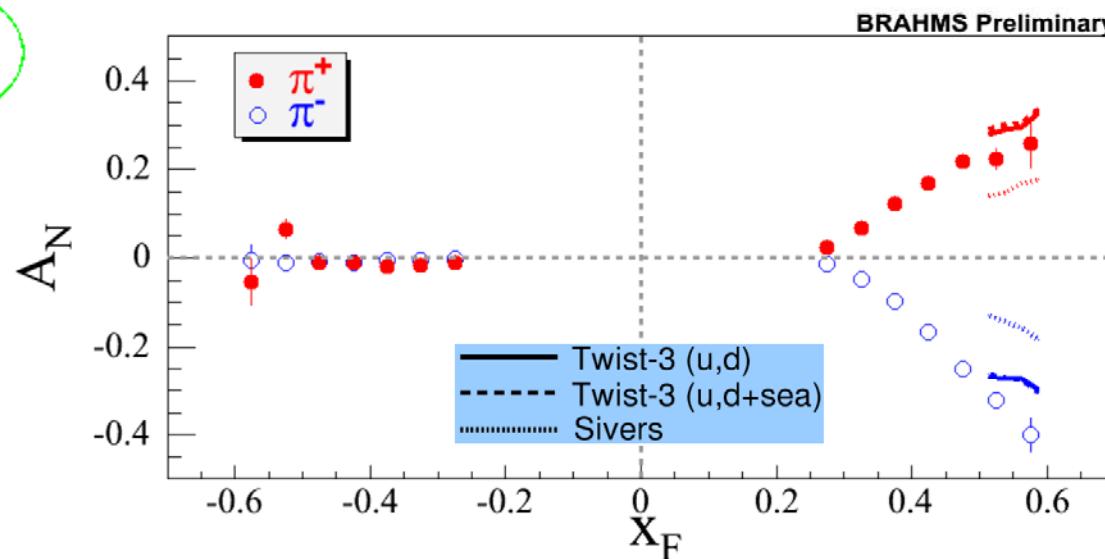
...but rising P_T dependence is not predicted by the same fits

→ No model fully describes the precision data

More experimental (and theoretical) work needed...



SSA of π^\pm at $\sqrt{s} = 62.4$ GeV

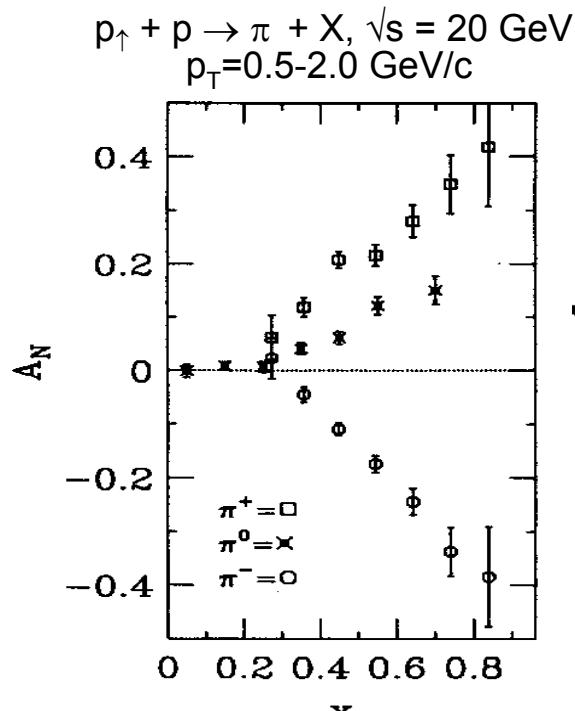


- Large $A_N(\pi^-)$: 40% at $x_F \sim 0.6$, $p_T \sim 1.3 \text{ GeV}/c$, $A_N(-x_F) \sim 0$.
 - Strong x_F -dependence
 - $|A_N(\pi^+)/A_N(\pi^-)|$ decreases with x_F
 - Sivers and Twist-3 calculations are compared with the data: Twist-3 calculations are in a better agreement with data.

Large asymmetries persist at high \sqrt{s}

Examples:

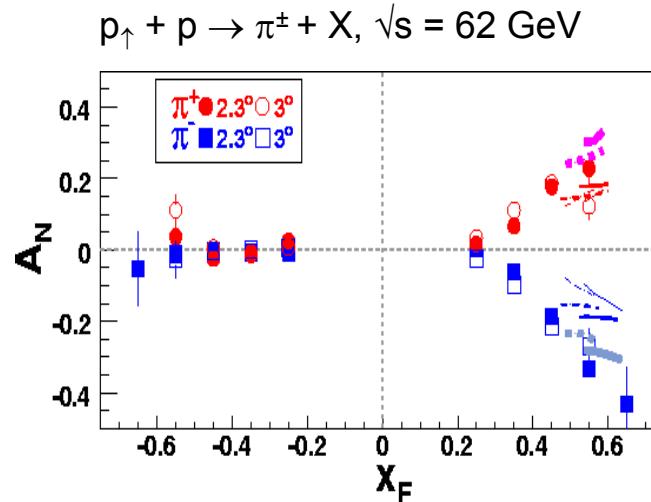
$\sqrt{s} = 20 \text{ GeV}$



π^0 : E704, Phys.Lett. B261 (1991) 201.
 π^{+-} : E704, Phys.Lett. B264 (1991) 462.

Fermilab, Fixed target, E704, 1991

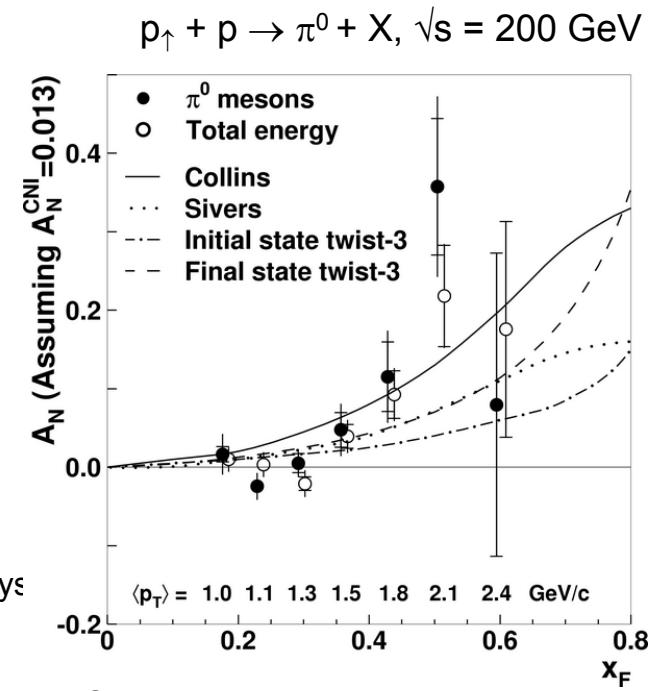
$\sqrt{s} = 62 \text{ GeV}$



Arsene et al. (BRAHMS), submitted to Phys Rev. Lett. [arXiv:nucl-ex/0801.1078]

RHIC, Brahms, 2007

$\sqrt{s} = 200 \text{ GeV}$



(STAR) Phys. Rev. Lett. 92 (2004)
171801

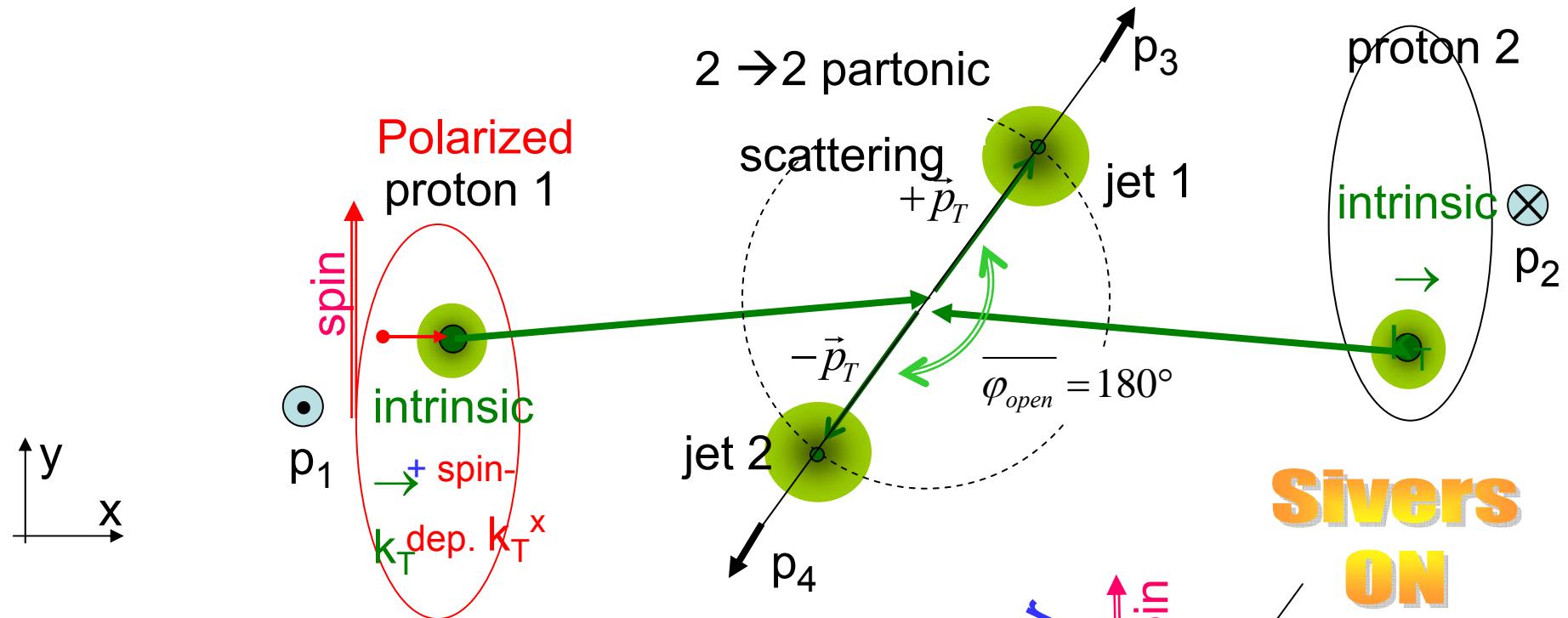
RHIC, STAR, 2004

Non-Perturbative cross section

Perturbative cross section

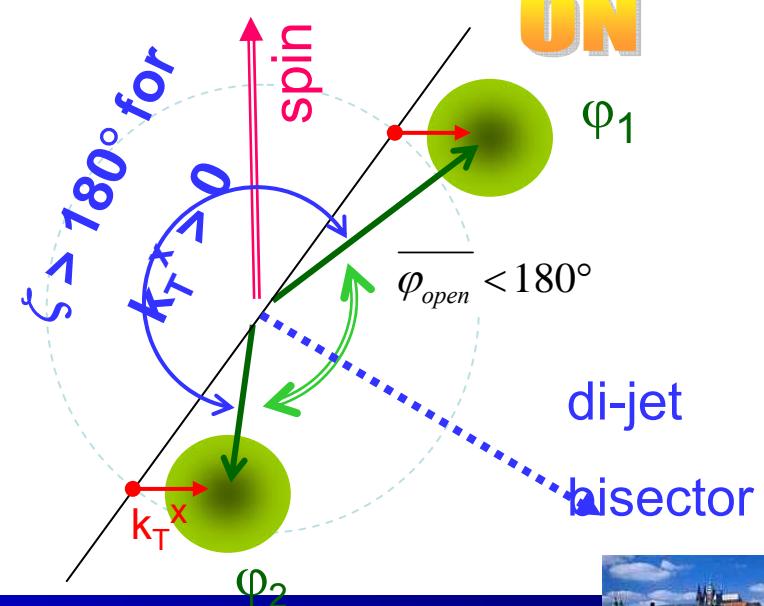


Sivers Mechanism of SSA



Spin dependent k_T^x offset

- deflects both jets in the same direction
- reduces average di-jet opening angle



STAR Results vs. Di-Jet Pseudorapidity Sum

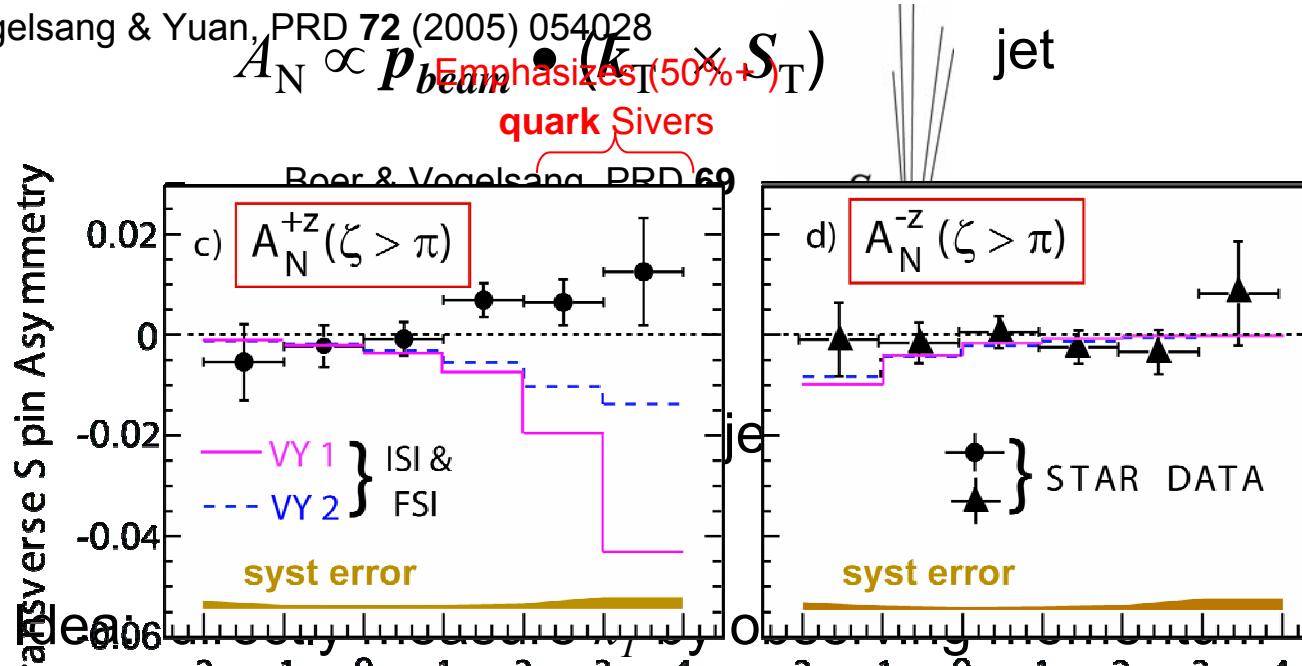
Run-6 Result

VY 1, VY 2 are calculations by

Vogelsang & Yuan, PRD 72 (2005) 054028

$$A_N \propto p_{beam} \cdot (k_T \times S_T)$$

Emphasizes (50%+)
quark Sivers



imbalance of a pair of jets produced in $p + p$ collision
 $n + n \rightarrow \ln(x_B^{+z}/x_B^{-z})$
 $n + n \rightarrow \ln(x_B^{+z}/x_B^{-z})$
and attempt to measure if k_T is correlated with
incoming proton spin

A_N consistent with zero

⇒~order of magnitude smaller in $pp \rightarrow$ di-jets than in semi-inclusive DIS
quark Sivers asymmetry!

arXiv:0705.4629



Mirko Planinić



Summary I

Transverse Single Spin Asymmetry (SSA) Measurements

- Feynman-x dependence of large-rapidity pion production shows large transverse SSA at RHIC energies, where cross sections are described by NLO pQCD
- Feynman-x dependence of large-rapidity transverse SSA are consistent with theoretical models (Sivers effect \Rightarrow orbital motion / twist-3 calculations)
- The p_T dependence of large-rapidity π^0 transverse SSA does not follow theoretical expectations



Summary II

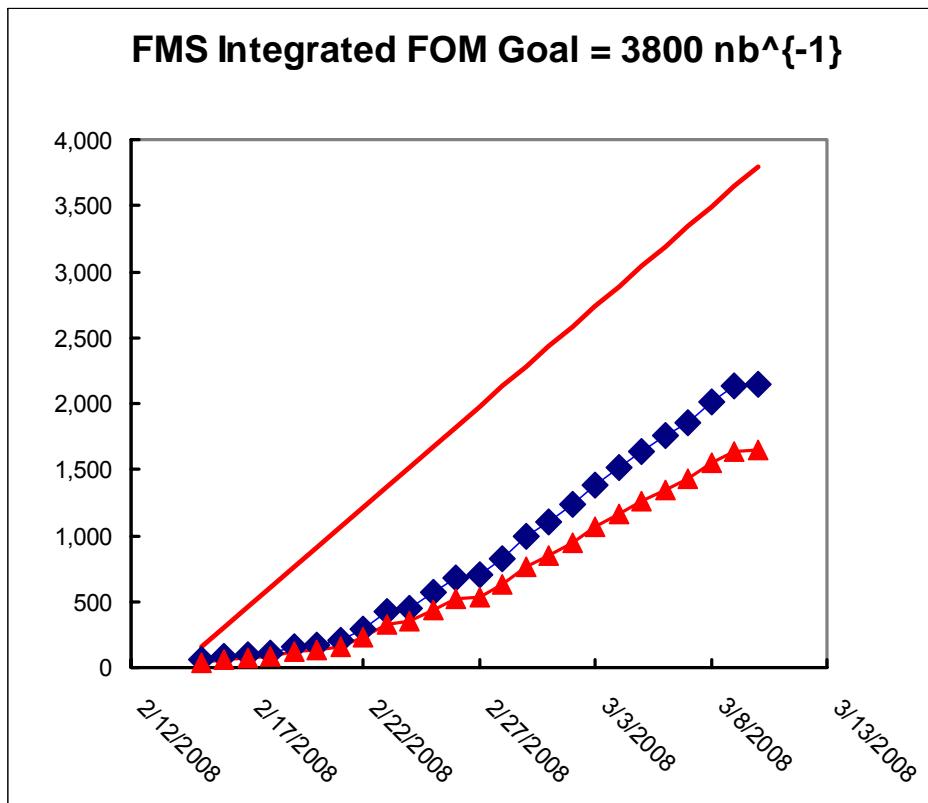
Transverse Single Spin Asymmetry (SSA) Measurements

- Direct measurement of spin-correlated k_T (Sivers effect) via midrapidity di-jet spin asymmetries completed in RHIC run 6 and found consistent with zero.
- Cancellations found in theory calculations subsequent to measurements also expect small di-jet spin asymmetries at midrapidity.



Run 8 data

Because of lower polarization than expected, figure of Merit (P^2L) fell short by roughly a factor of two.



Run 8 Integrated FMS transverse figure-of-merit.

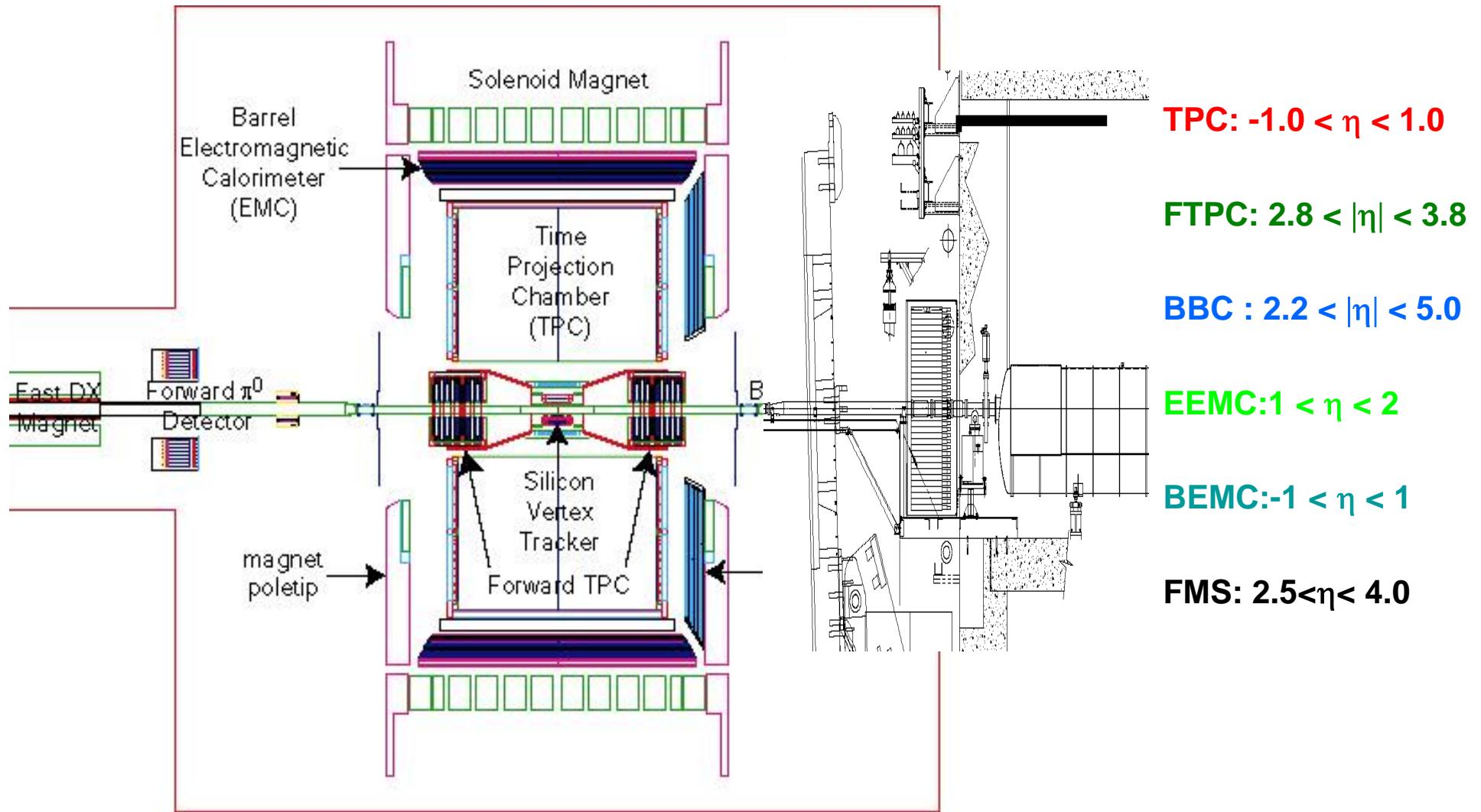
Only 43% of goal, after calibration from jet

However, the FMS provides roughly 20 times the coverage of previous runs in the forward region.





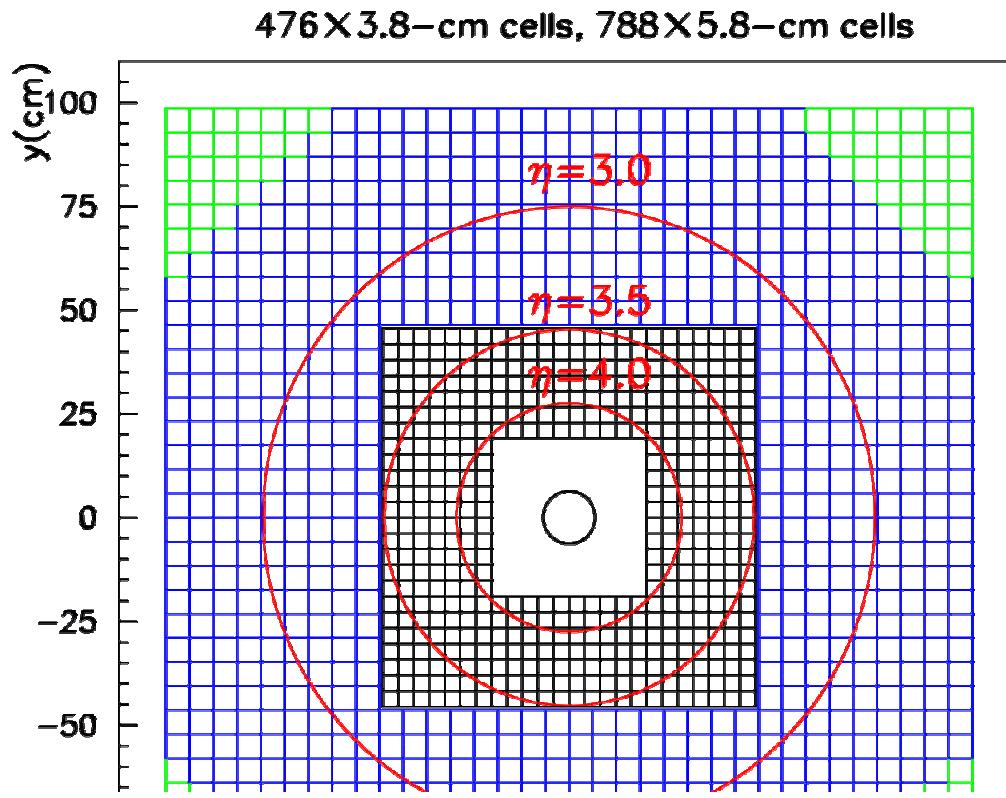
STAR detector layout with FMS



New forward detector for Run 8: FMS

40

FMS provides nearly 20x the coverage of previous forward detectors



North-half, view from
the hall



Nearly contiguous coverage for
 $2.5 < \eta < 4.0$



Run 5 FPD

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