Neutrino-induced nucleosynthesis in supernovae

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Introduction

Two basic types of supernovae (SNe):

Thermonuclear SNe (Type Ia) Explosive carbon burning $({}^{12}C+{}^{12}C \otimes {}^{56}Ni)$ in a degenerate Chandrasekhar mass (~1.4M $_{\odot}$) white dwarf. The white dwarf turns out to be totally disrupted in the explosion — no stellar remnant is left!

The total energy of electromagnetic radiation is of $*6^{-}10^{49}$ erg and comes from the 56 Ni \otimes 56 Co \otimes 56 Fe decay. The explosion energy E_{exp} is of $*10^{51}$ erg and comes from the explosive carbon burning into 56 Ni , almost all the E_{exp} resides in the kinetic energy of expanding envelope.

Core-collapse SNe (all other Types but Ia)

The SN outburst is triggered by the gravitational collapse of the "iron" core of a mass $M_{Fe}=(1.2-2) M_{\odot}$ into a neutron star. About (10–15)% $M_{Fe}c^2$ is radiated in the form of neutrinos and antineutrinos of all the flavors (e, m, t):

$$E_{n\tilde{n}} = (3-5) \times 10^{53} erg$$

The explosion energy (kinetic energy of the envelope expansion):

$$E_{exp} = (0.5-2)^{-10^{51}}$$
erg

it comes from the shock wave created at the boundary between a new-born neutron star and the envelope to be expelled.

$$E_{exp} / E_{n\tilde{n}} \approx 3 \times 10^{-3}$$
 !!

Schematic Supernova «light curves»



Dynamics of gravitational collapse

Dynamical stability

Expanding force: $F_{\rm E} \sim 4 p R^2 P$

Contracting force: $F_{\rm G} \sim GM^2/R^2$ $\boldsymbol{r} = \boldsymbol{r}_0 \left(R_0 / R \right)^3$ $P = P_0 \left(\mathbf{r} / \mathbf{r}_0 \right)^{\mathbf{g}}$ $F_{\rm E} \sim 1/R^{3g-2}$ $F_{\rm G} \sim 1/R^2$ *Stability:* 3g - 2 > 2

Nadyozhin & Yudin (2005):



 56 Fe $\implies 26 p + 30 n$ Nuclear Statistical Equilibrium (NSE)

<u>Dynamics of gravitational collapse</u> *Initial "free fall". Accreting shock wave*



Dynamics of gravitational collapse



Collapse of a 2M_☉ iron stellar core (based on Nadyozhin 1977)

Dynamics of gravitational collapse Expulsion of supernova envelope (how?)



<u>Dynamics of gravitational collapse</u> Total energy liberated

The total energy liberated is equal to the neutron star gravitational binding energy E_g : $E_g = (M_{\text{Fe}} - M_g)c^2$

 $M_{\rm Fe}$ – initial mass of the iron core before collapse

 M_{g} – The gravitational mass of a neutron star measured by a distant observer that connected with M_{Fe} by (*Lattimer & Yahil 1989*)

$$M_{\text{Fe}} = M_g + 0.084 M_g^2 \longrightarrow E_g = 1.5 \times 10^{53} M_g^2 \text{ erg}$$

$$(M_{\text{Fe}}, M_g \text{ are in } M_{\odot})$$
V - total number of nucleons in the neutron star
$$\frac{M_{\text{Fe}} M_g E_g / 10^{53} \text{ erg } N / 10^{57} E_g / N \text{ MeV}}{1.10 \ 1.01 \ 1.54 \ 1.33 \ 73}$$

$$\frac{M_{\text{Fe}} M_g E_g / 10^{53} \text{ erg } N / 10^{57} E_g / N \text{ MeV}}{1.40 \ 1.27 \ 2.40 \ 1.67 \ 90}$$

$$\frac{1.70 \ 1.51 \ 3.41 \ 2.03 \ 105}{2.00 \ 1.74 \ 4.56 \ 2.38 \ 120}$$

$$\frac{2.30 \ 1.97 \ 5.84 \ 2.74 \ 133}$$

The properties of the Neutrino flux

Cumulative neutrino "light" curve (based on Nadyozhin 1978)



<u>The properties of the Neutrino flux</u> *Total energy radiated. Characteristic time.*

Total energy radiated : $E_{n\tilde{n}} \approx E_g = (2.4 - 4.6) \times 10^{53} \text{erg} \text{ (for } 1.4 \le M_{\text{Fe}} \le 2)$ $(10\% - 13\% \text{ of } M_{\text{Fe}}c^2)$

Nonthermal component: Mostly \mathbf{n}_{e} : $\mathbf{E}_{ne} \leq (1-2) \times 10^{52} \text{ erg} \approx 0.05 \text{ E}_{n\tilde{n}}$, $\mathbf{e}_{ne} \approx 15 - 20 \text{ MeV}$, $\Delta t \approx (0.01 - 0.02) \text{ sec}$

Thermal component: $E_{ne} \approx E_{nm} \approx E_{nt} \approx \frac{1}{3} E_{n\tilde{n}} \quad (ni \equiv n_i + \tilde{n_i}, i = e, m, t)$ $e_{ne,\tilde{n}e} \approx 10-12 \text{ MeV}, \quad e_{nm,\tilde{n}m} \approx e_{nt,\tilde{n}t} \approx 25 \text{ MeV},$ $\Delta t \approx (10-20) \text{ sec} \quad (e - \text{ individual energy of neutrinos}).$



Basics of the Neutrino Nucleosynthesis Light elements and p-nuclei



Cosmic abundance of elements

¹⁸⁰Ta exists in its excited (isomeric) state ($E_{ex} = 75.4 \text{keV}, J^p = 9^-$) abundance = $(0.00012)^{181}$ Ta $T_{1/2}(^{180}$ Ta) > 1.2×10^{15} y

> The ground state $(J^p = 1^+)$ **b**⁻decays: $T_{1/2} = 8.15$ h into ¹⁸⁰ W

Basics of the Neutrino Nucleosynthesis **Presupernova shell structure:**



A spatial scanning of Gamov's sub Coulomb-barrier penetration factor !! Basics of the Neutrino Nucleosynthesis Magnitude of the neutrino-induced transformations

$$\mathbf{n} + A \rightarrow B + c + \cdots$$



Basics of the Neutrino Nucleosynthesis Magnitude of the neutrino-induced transformations

$$\mathbf{n} + A \longrightarrow B + c + \cdots \quad (1)$$

$$\frac{\mathrm{d}n}{\mathrm{d}t} = -n \frac{L_{\mathbf{n}}(t)}{\langle \mathbf{e}_{\mathbf{n}} \rangle} \frac{\langle \mathbf{s}_{\mathbf{n}A} \rangle}{4\mathbf{p}r^2} \qquad n - \text{ density of species A}$$
$$\left|\frac{\mathrm{d}n}{n}\right| = \int_{0}^{\infty} \frac{L_{\mathbf{n}}(t)}{\langle \mathbf{e}_{\mathbf{n}} \rangle} \, \mathrm{d}t \frac{\langle \mathbf{s}_{\mathbf{n}A} \rangle}{4\mathbf{p}r^2} = N_{\mathbf{n}} \frac{\langle \mathbf{s}_{\mathbf{n}A} \rangle}{4\mathbf{p}r^2} = 10^{-2} - 10^{-5}$$
$$N_{\mathbf{n}} - \text{ total number of radiated neutrinos participating in reaction (1)}$$

Typical values:

$$N_{\mathbf{n}} \sim 10^{57}, \langle \mathbf{s}_{\mathbf{n}A} \rangle \approx 10^{-42} - 10^{-44} \text{ cm}^2, \ \mathcal{V} = 10^8 - 10^9 \text{ cm}$$

Basics of the Neutrino Nucleosynthesis The role of shock wave

$$T_{WW} = \left(\frac{3E}{4\pi a R_0^3}\right)^{1/4} = 2.37 \times 10^9 E_{51}^{0.25} (R_{09})^{-0.75} \text{ K}$$

$$t_u = 3.83 \times 10^{-3} \rho_0^{0.5} E_{51}^{-0.5} (R_{09})^{2.5} \text{ s}$$

$$T(t) = \frac{T_p}{1 + \xi_T t/t_u}, \quad T_p = \xi_p T_{WW},$$

$$\rho(t) = \rho_p \left(\frac{T}{T_p}\right)^3, \quad \rho_p = 7\rho_0,$$

$$R(t) = R_0 (1 + \xi_r t/t_u),$$

$$\Delta t_{SW} = \xi_{SW} E_{51}^{-0.38} (R_0/10^9 \text{ cm})^{1.4} \text{ s}$$

$$\mathbf{x} \sim 1 \text{ are adjustable parameters}$$

Nadyozhin & Deputovich (2002)

 $T_{\rm WW}$ – Weaver & Woosley (1980) temperature

Neutrino Nucleosynthesis of p-nuclei **p-process**

Creation of some 30 isotopes (from ⁷⁴Se up to ¹⁹⁶Hg) that can be produced by the neutron capture neither in the s-process nor in the r-process (bypassed or p-nuclei) The mechanisms of the production:

- The (p, γ) and (γ ,n) reactions at high temperatures (T₉=3-5)
- The neutrino resonance captures and scattering (GT & IA resonances)

The main neutrinoinduced process is

 $n_{e} + (A, Z) \rightarrow (A, Z+1)^{*} + e^{-}$ $(A, Z+1)^{*} \rightarrow (A-1, Z+1)_{bps} + n$



Gaponov Yu.V., Lyutostansky Yu. S., Sov. J. Nucl. Phys. (1974) Domogatskii G.V., Nadyozhin D.K., Sov. Astronomy (1978) Panov I.V., Astronomy Letters (1994)



The reaction $(A, Z) + \nu_e \Rightarrow (A, Z + 1)^* + e^- \Rightarrow (A - 1, Z + 1) + n + e^$ which goes through the GT & IA excited states of the nucleus (A, Z + 1) is of special importance for the neutrino-induced production of the p-nuclei.

Parents		p-nuclei		Final sum of yields
$^{75}\mathrm{As}$	(75, 33)	$^{74}\mathrm{Se}$	(74, 34)	$^{74}\mathrm{Sr} + ^{74}\mathrm{Rb} + ^{74}\mathrm{Kr} + ^{74}\mathrm{Br} + ^{74}\mathrm{Se} + 0.34^{74}\mathrm{As}$
$^{79}\mathrm{Br}$	(79, 35)	⁷⁸ Kr	(78, 36)	78 Y + 78 Sr+ 78 Rb+ 78 Y + 78 Kr
$^{85}\mathrm{Rb}$	(85, 37)	$^{84}\mathrm{Sr}$	(84, 38)	84 Nb + 84 Zr+ 84 Y + 84 Sr
$^{93}\mathrm{Nb}$	(93, 41)	$^{92}\mathrm{Mo}$	(92, 42)	92 Rh + 92 Ru+ 92 Tc + 92 Mo
$^{95}\mathrm{Mo}$	(95, 42)	$^{94}\mathrm{Tc}$	$(94, 43)^1$	$^{94}\mathrm{Tc}(eta^+)^{94}\mathbf{Mo}$
$^{99}\mathrm{Ru}$	(99, 44)	$^{98}\mathrm{Rh}$	$(98,45)^2$	$^{98}\mathrm{Rh}(\beta^+)^{98}\mathrm{Ru}$
$^{103}\mathrm{Rh}$	(103, 45)	$^{102}\mathrm{Pd}$	(102, 46)	102 In $+^{102}$ Cd $+^{102}$ Ag $+^{102}$ Pd $+0.2^{102}$ Rh
$^{107}\mathrm{Ag}$	(107, 47)	$^{106}\mathrm{Cd}$	(106, 48)	106 Sb $+^{106}$ Sn $+^{106}$ In $+^{106}$ Cd
$^{109}\mathrm{Ag}$	(109, 47)	$^{108}\mathrm{Cd}$	(108, 48)	108 Sb $+^{108}$ Sn $+^{108}$ In $+^{108}$ Cd $+0.97^{108}$ Ag
^{114}Cd	(114, 48)	113 In	(113, 49)	113 Te $+^{113}$ Sb $+^{113}$ Sn $+^{113}$ In
114 In	(114, 49)	$^{113}\mathrm{Sn}$	$(113,50)^3$	$^{113}{ m Sn}(eta^+)^{113}{ m In}$
$^{116}\mathrm{Sn}$	(116, 50)	$^{115}\mathrm{Sb}$	$(115,51)^4$	$^{115}{ m Sb}(eta^+)^{115}{ m Sn}$
113 In	(113, 49)	$^{112}\mathrm{Sn}$	(112, 50)	$^{112}\text{I} + ^{112}\text{Te} + ^{112}\text{Sb} + ^{112}\text{Sn} + 0.44^{112}\text{In}$
115 In	(115, 49)	$^{114}\mathrm{Sn}$	(114, 50)	$^{114}\text{I} + ^{114}\text{Te} + ^{114}\text{Sb} + ^{114}\text{Sn} + 0.995^{114}\text{In}$
$^{121}\mathrm{Sb}$	(121, 51)	$^{120}\mathrm{Te}$	(120, 52)	120 Cs $+^{120}$ Xe $+^{120}$ I $+^{120}$ Te
^{127}I	(127, 53)	$^{126}\mathrm{Xe}$	(126, 54)	126 La $+^{126}$ Ba $+^{126}$ Cs $+^{126}$ Xe $+0.437^{126}$ I
^{133}Cs	(133, 55)	^{132}Ba	(132, 56)	132 Pr + 132 Ce+ 132 La + 132 Ba +0.0187 132 Cs
¹³⁹ La	(139, 57)	$^{138}\mathrm{Ce}$	(138, 58)	138 Pm $+^{138}$ Nd $+^{138}$ Pr $+^{138}$ Ce
$^{153}\mathrm{Eu}$	$(153,\!63)$	^{152}Gd	(152, 64)	0.88^{152} Ho $+^{152}$ Dy $+^{152}$ Tb $+^{152}$ Gd $+0.279^{152}$ Eu
$^{159}\mathrm{Tb}$	$(159,\!65)$	$^{158}\mathrm{Dy}$	(153, 66)	158 Tm $+^{158}$ Er $+^{158}$ Ho $+^{158}$ Dy $+0.166^{158}$ Tb
$^{165}\mathrm{Ho}$	$(165,\!67)$	$^{164}\mathrm{Er}$	(164, 68)	164 Lu + 164 Yb+ 164 Tm + 164 Er +0.40 164 Ho
$^{169}\mathrm{Tm}$	(169, 69)	168 Yb	(168,70)	168 Ta $+^{168}$ Hf $+^{168}$ Lu $+^{168}$ Yb
^{175}Lu	(175,71)	$^{174}\mathrm{Hf}$	(174,72)	174 Re $+^{174}$ W $+^{174}$ Ta $+^{174}$ Hf
181 Ta	(181,73)	^{180}W	(180,74)	180 Ir $+^{180}$ Os $+^{180}$ Re $+^{180}$ W $+0.14^{180}$ Ta ^g
$^{185}\mathrm{Re}$	(185,75)	$^{184}\mathrm{Os}$	(184, 76)	184 Au $+^{184}$ Pt $+^{184}$ Ir $+^{184}$ Os
191 Ir	(191,77)	$^{190}\mathrm{Pt}$	(190, 78)	190 Tl $+^{190}$ Hg $+^{190}$ Au $+^{190}$ Pt
¹⁹⁷ Au	(197, 79)	¹⁹⁶ Hg	(196, 80)	$^{196}\text{Bi} + ^{196}\text{Pb} + ^{196}\text{Tl} + ^{196}\text{Hg} + 0.0695^{196}\text{Au}$

Stable parent nuclei

 $\label{eq:product} \begin{array}{ll} {}^{1}\beta^{+}\mbox{-decay to the p-nucleus } {}^{94}\mbox{Mo;} & {}^{2}\beta^{+}\mbox{-decay to the p-nucleus } {}^{98}\mbox{Ru} \\ {}^{3}\beta^{+}\mbox{-decay to the p-nucleus } {}^{113}\mbox{In;} & {}^{4}\beta^{+}\mbox{-decay to the p-nucleus } {}^{115}\mbox{Sn} \end{array}$

To find the final yield of a p-nucleus, marked in bold, one has to sum the yields of a number of ß-unstable nuclei that can exist in nuclear network

Unstable parent nuclei p-nuclei Final sum of yields Parents $^{93}\mathrm{Mo}$ $^{92}\mathrm{Tc}(\beta^+)^{92}\mathrm{Mo}$ $^{92}\mathrm{Tc}$ $(92,43)^1$ (93, 42) $^{95}\mathrm{Nb}$ 94 Rh + 94 Ru + 94 Tc + 94 Mo + 94 Nb ^{94}Mo (94, 42)(95,41)⁹⁶Ru $^{96}\mathrm{Ag} + ^{96}\mathrm{Pd} + ^{96}\mathrm{Rh} + ^{96}\mathrm{Ru}$ $^{97}\mathrm{Tc}$ (97, 43)(96, 44)⁹⁷Ru $^{96}\mathrm{Rh}$ $^{96}\mathrm{Rh}(\beta^+)^{96}\mathrm{Ru}$ (97, 44) $(96,45)^2$ $^{99}\mathrm{Tc}$ 98 Ru 98 Ag $+^{98}$ Pd $+^{98}$ Rh $+^{98}$ Ru $+^{98}$ Tc (99, 43)(98,44) $^{115}\mathrm{Sn}$ $^{115}\text{I} + ^{115}\text{Te} + ^{115}\text{Sb} + ^{115}\text{Sn}$ 116 In (116, 49)(115,50) 125 I $^{124}\mathrm{Xe}$ 124 Ba $+^{124}$ Cs $+^{124}$ **Xe** (125, 53)(124,54) $^{130}\mathrm{Ba}$ $^{131}\mathrm{Cs}$ 130 Pr + 130 Ce+ 130 La + 130 Ba +0.016 130 Cs (131,55)(130,56) 139 Ba 138 La 138 La (139, 56)(138,57) 137 La $^{136}\mathrm{Ce}$ (137, 57) 136 Nd $+^{136}$ Pr $+^{136}$ Ce (136, 58) $^{145}\mathrm{Pm}$ 144 Gd $+^{144}$ Eu $+^{144}$ Sm 144 Sm (145, 61)(144, 62) $^{157}\mathrm{Tb}$ $^{156}\mathrm{Dy}$ 156 Tm $+^{156}$ Er $+^{156}$ Ho $+^{156}$ Dy (157, 65)(156, 66) $^{162}\mathrm{Er}$ 163 Ho 162 Lu + 162 Yb+ 162 Tm + 162 Er (163, 67)(162, 68) 180 Ta ¹⁸⁰**Ta**^g and p-nucleus ¹⁸⁰**Ta**^{ex} ($E_{ex} = 75 \text{keV}$) (181,72)(180,73) $^{181}\mathrm{Hf}$

 ${}^{1}\beta^{+}$ -decay to the p-nucleus 92 Mo ${}^{2}\beta^{+}$ -decay to the p-nucleus 96 Ru

There is also a need in the cross-sections of a number of reactions involving muon and tau neutrinos and antineutrinos. Specifically, we need the total cross-section of the neutral-current excitation followed by the emission of a neutron for the following processes. The cross-section has to be averaged over the Fermi–Dirac spectra of zero chemical potential and with temperatures kT = 6, 8, 10 MeV.

Stable nuclei	Unstable nuclei
$ \begin{array}{c} ^{95}Mo(\nu,\nu')^{95}Mo^* \Longrightarrow {}^{94}Mo + n \\ $	$\begin{array}{c} {}^{97}\mathrm{Ru}(\nu,\nu'){}^{97}\mathrm{Ru}^* \Longrightarrow {}^{96}\mathrm{Ru} + \mathrm{n} \\ {}^{113}\mathrm{Sn}(\nu,\nu'){}^{113}\mathrm{Sn}^* \Longrightarrow {}^{112}\mathrm{Sn} + \mathrm{n} \\ {}^{114}\mathrm{In}(\nu,\nu'){}^{114}\mathrm{In}^* \Longrightarrow {}^{113}\mathrm{In} + \mathrm{n} \\ {}^{125}\mathrm{Xe}(\nu,\nu'){}^{125}\mathrm{Xe}^* \Longrightarrow {}^{124}\mathrm{Xe} + \mathrm{n} \\ {}^{127}\mathrm{Xe}(\nu,\nu'){}^{127}\mathrm{Xe}^* \Longrightarrow {}^{126}\mathrm{Xe} + \mathrm{n} \\ {}^{131}\mathrm{Ba}(\nu,\nu'){}^{131}\mathrm{Ba}^* \Longrightarrow {}^{130}\mathrm{Ba} + \mathrm{n} \\ {}^{137}\mathrm{Ce}(\nu,\nu'){}^{137}\mathrm{Ce}^* \Longrightarrow {}^{136}\mathrm{Ce} + \mathrm{n} \\ {}^{145}\mathrm{Sm}(\nu,\nu'){}^{145}\mathrm{Sm}^* \Longrightarrow {}^{144}\mathrm{Sm} + \mathrm{n} \\ {}^{157}\mathrm{Dy}(\nu,\nu'){}^{157}\mathrm{Dy}^* \Longrightarrow {}^{156}\mathrm{Dy} + \mathrm{n} \\ {}^{163}\mathrm{Er}(\nu,\nu'){}^{163}\mathrm{Er}^* \Longrightarrow {}^{162}\mathrm{Er} + \mathrm{n} \end{array}$

The p-nuclei are marked in bold, ν stands for ν_{μ} , $\tilde{\nu}_{\mu}$, ν_{τ} , and $\tilde{\nu}_{\tau}$.

He shell



Composition of a 15 M_{\odot} solar metallicity (Z_{\odot}) presupernova He-shell (*Woosley, Heger et al.*)

The key reactions:

Thermonuclear network (code I) $p + n \rightarrow D$; ⁴He(T,g)⁷Li; ⁴He(³He,g)⁷Be; (T = ³H) ⁷Li(a,g)¹¹B; ⁷Be(a,g)¹¹C; (a = ⁴He) ... ¹³C(n,g)¹⁴C, ¹³C(a,n)¹⁶O....

About 160 thermonuclear reactions and beta-decays connecting 40 nuclides from **n** and **p** through ²⁴Mg The rates are from the Fowler et al., Thielemann et al., Angulo et al. data bases.

Thermonuclear network (code II)

deals with ~1300 heavy nuclides up to the atomic number Z=60 (Nd)

The rates of the **n**, **p**, and α-captures are from Cowan et al. (1991). The β-decay rates are from Kratz et al. (1993).



Iterative adjustment

Nadyozhin, Panov, Blinnikov, AA 335, 207 (1998)

Neutrino-nuclear interactions included ⁴He(n,n'n)³He, ⁴He(n,n'p)³H $\frac{12}{C(n,n'n)} \frac{11}{C}$, $\frac{12}{C(n,n'p)} \frac{11}{B}$ ${}^{12}C(n,n'{}^{3}\text{He}){}^{9}\text{Be}, {}^{14}N(n,n'a){}^{10}\text{B}$ ${}^{16}O(n,n'n){}^{15}O$, ${}^{16}O(n,n'p){}^{15}N$ 20 Ne(*n*,*n*'**n**)¹⁹Ne, 20 Ne(*n*,*n*'**p**)¹⁹F $p(\tilde{\boldsymbol{n}}_{e}, e^{+})n$

 $n \equiv n_{m,t}, \tilde{n}_{m,t}$

The neutrino cross-sections from:

Woosley, Hartmann, Hoffman, Haxton (1990) Heger, Kolbe, Haxton, Langanke et al. (2003) Domogatsky, Svetlana Imshennik (1982) Domogatsky, Nadyozhin, Eramzhyan (1978) The cross-sections are averaged over the μ and τ neutrino-antineutrino spectra given by the Fermi-Dirac distributions with T=8 MeV and zero chemical potential ($\langle E_v \rangle = 25$ MeV).

The post-shock temporal behavior of temperature, density, and radius from: Nadyozhin, Deputovich (2002)



No thermonuclear processing ! (only v-induced interactions!)



(Y versus time t)

 $E_{exp} = 10^{51} erg$ $T_{peak} = 6.5 \ 10^8 K$ $r_{peak} = 5.3 \ 10^3$

Shock arrives at $\Delta t_{SW} \approx 8 \text{ s}$

Post shock characteristic time $\Delta t = 5 \text{ s}$

$$\mathbf{Y}_{j} \equiv \frac{m_{\mathrm{u}}}{\boldsymbol{r}} n_{j}$$

number of nuclides *j* per barion



Post-shock burning and recreation II (Y versus time t)



Bottom of He-shell: Final yields versus E_{exp} (Here Y9 = actual Y9 × 100)

Y3 = ${}^{3}\text{He} + {}^{3}\text{H}$ (T_{1/2} = 12.33 y) Y7 = ${}^{7}Li + {}^{7}Be (T_{1/2} = 53.29 d)$ $Y9 = {}^{9}Be$ $Y11 = {}^{11}B + {}^{11}C$ (T₁₂ = 20.4 m) $Y15 = {}^{15}N + {}^{15}O (T_{1/2} = 122 s)$ Y19 = ${}^{19}F + {}^{19}Ne(T_{1/2} = 17.3 \text{ s})$

Neutrino-driven weak r-process component in the helium shell of poor-metal stars



Composition of a 15 M_O poor-metal (Z=0.0001 Z_O) presupernova He-shell (Woosley, Heger et al.)



Composition of a 15 M_{\odot} solar metallicity (Z_{\odot}) presupernova He-shell (*Woosley, Heger et al.*)



On average 50 neutrons were captured per Fe-seed *Nadyozhin, Panov Astron. Lett.* **33**, 385 (2007)

<u>Main issues</u>

- 1. Creation of some 30 isotopes (from ⁷⁴Se up to ¹⁹⁶Hg) that can be produced neither in the s-process nor in the r-process (bypassed or p-nuclei)
- 2. Creation of ⁷Li, ¹¹B, ⁹Be (in He-shell and in CO-shell) There is no need in hypothetical low-energy (E<100 MeV) CR

⁶Li comes from CR;

⁷Li comes from BBN + CR + NN;

 9 Be comes from CR + NN;

 ^{10}B comes from CR;

¹¹B comes from CR + NN;

No problem with large cosmic ratios ⁷Li/⁶Li and ¹¹B/¹⁰B

- 3. Creation of rare isotopes such as ¹⁹F, ¹⁵N and some others
- 4. Possibility to reproduce at least a weak component of the r-process in He-shell
- 5. Neutrino induced acceleration of heavy nuclide fission

6. Evaporation of free neutrons due to the neutrino inelastic scattering off heavy nuclides (helps to drive the r-process)

For detailed information on the core-collapse supernovae see a remarkable review:

Woosley S.E., Heger A. and Weaver T.A. The evolution and explosion of massive stars *Rev. Modern Phys.* 74, 1015 (2002)

For general review of the origin of chemical elements in the framework of hystorical aspects see:

> Wallerstein G., Iben I., Parker P. et al. *Rev. Modern Phys.* 69, 195 (1997)

> > Presupernova models

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