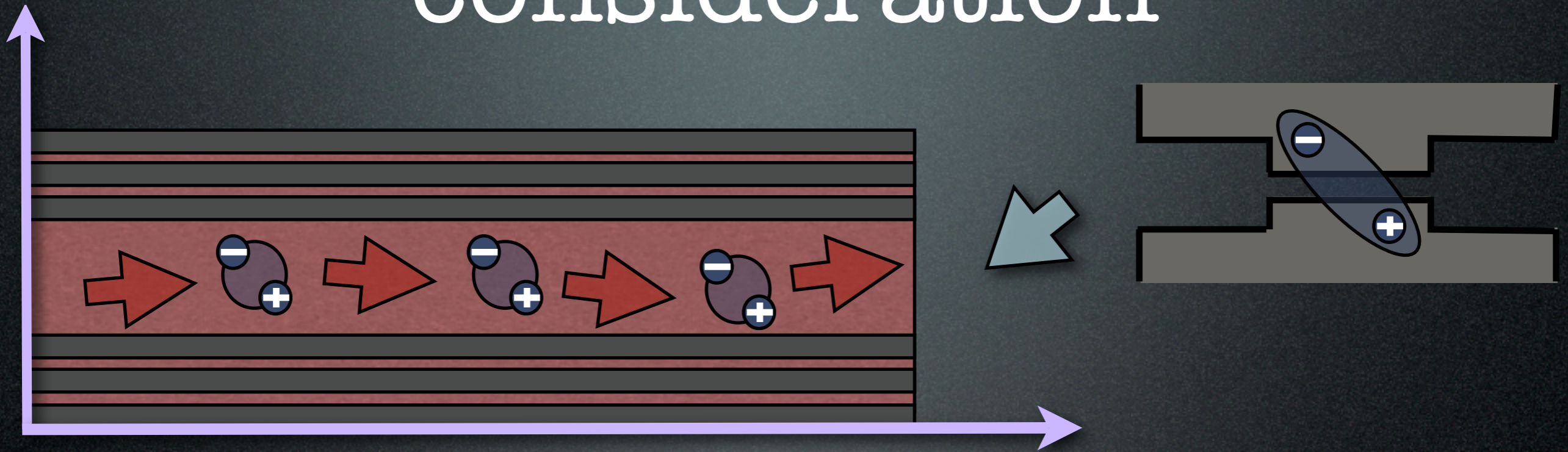


On the superfluid properties of polaritonic system

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System under consideration



$$\hat{H} = \frac{1}{V} \sum_p \left(\varepsilon_{ph} \hat{a}_p^\dagger \hat{a}_p + \varepsilon_{ex} \hat{b}_p^\dagger \hat{b}_p + \frac{\Omega}{2} (\hat{a}_p^\dagger \hat{b}_p + \hat{a}_p \hat{b}_p^\dagger) \right) + \hat{H}_{ex-ex}$$

$$\varepsilon_{ph}(p) = c \sqrt{\left(\frac{\pi n}{L}\right)^2 + p^2} \approx E_0 + \frac{p^2}{2m_{ph}} \quad \varepsilon_{ex}(p) = E_1 + \frac{p^2}{2m_{ex}}$$

System under consideration

Long range order and superfluidity

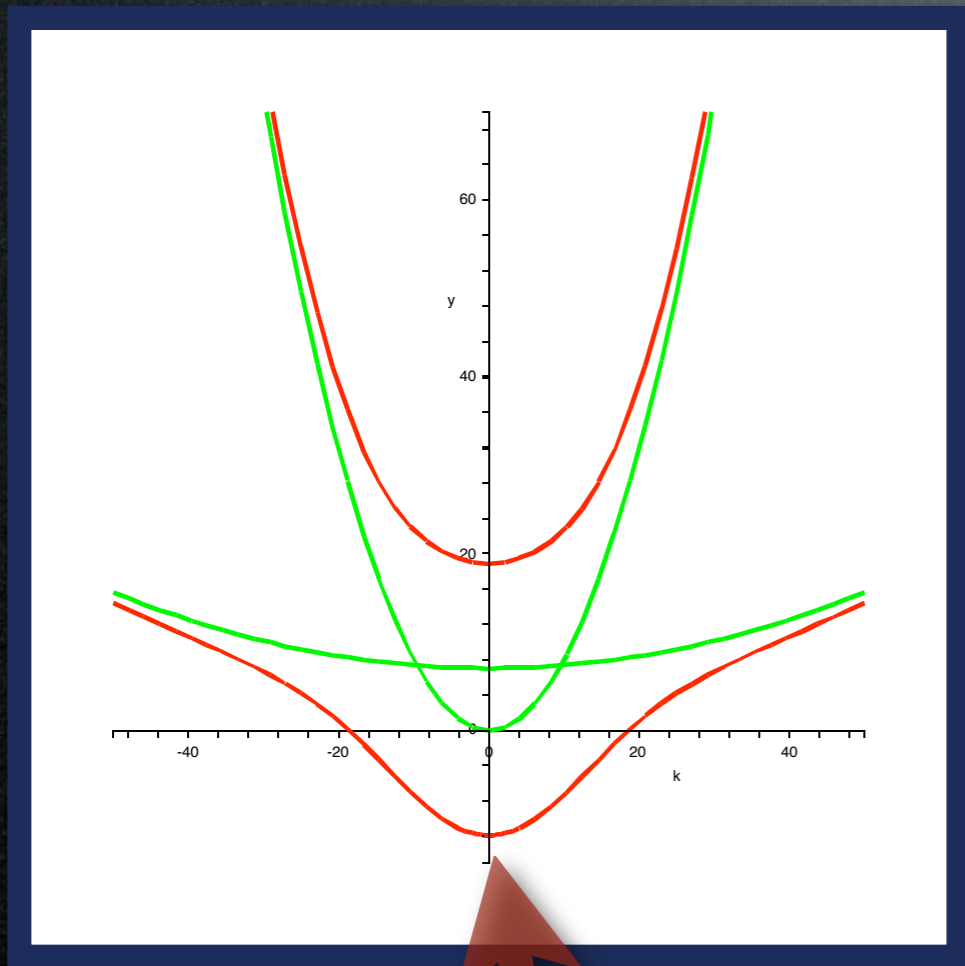
$$\rho(r - r') \sim \frac{1}{|r - r'|^\alpha}$$

? $\alpha = \frac{mT}{2\pi n_s}$?

System under consideration is two-dimensional, therefore there are no bose condensate at finite temperature and transition into superfluid state is of Berezinskii-Kosterlitz-Thouless type

$$\alpha_{crit} = \frac{1}{4}$$

? $T_c = \frac{2\pi n}{m \ln \Gamma}$



Small effective mass, therefore high transition temperature

Effective action

$$Z = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{-S}$$

$\psi = \psi_l + \psi_{sh}$

$\psi_l = \sqrt{\rho_0} e^{i\varphi} \quad \psi_{sh} = \tilde{\psi} e^{i\varphi}$

$$\partial_\mu (e^{i\varphi(x)} f(x)) = e^{i\varphi(x)} (\partial_\mu + i(\partial_\mu \varphi(x))) f(x)$$

$$S = S_0 + \int d\tau d^2x J_\mu(x, \tau) A_\mu(x, \tau) +$$
$$+ \frac{1}{2} \int d\tau d^2x d\tau' d^2x' \frac{\delta J_\mu(x, \tau)}{\delta A_\nu(x', \tau')} A_\mu(x, \tau) A_\nu(x', \tau') + \dots$$

$$A_\mu = \partial_\mu \varphi \quad J_\mu = \frac{\delta S}{\delta A_\mu} \quad K_{\mu\nu} = \frac{\delta J_\mu}{\delta A_\nu}$$

Effective action

$$S_{eff} = \frac{1}{2} \int d\tau d^2x d\tau' d^2x' \Pi_{\mu\nu}(x, \tau, x', \tau') A_\mu(x, \tau) A_\nu(x', \tau')$$
$$\Pi_{\mu\nu}(x, \tau, x', \tau') = \left\langle \frac{\delta J_\mu(x, \tau)}{\delta A_\nu(x', \tau')} \right\rangle - \langle J_\mu(x, \tau) J_\nu(x', \tau') \rangle$$

$$\Pi_{\mu\nu}(q, i\omega_n) = \langle K_{\mu\nu}(q, i\omega_n) \rangle - \langle J_\mu(q, i\omega_n) J_\nu(-q, -i\omega_n) \rangle$$

In the zero frequency and low momentum limit

$$\langle J_\mu(q, 0) J_\nu(-q, 0) \rangle = \chi_l \frac{q_\mu q_\nu}{q^2} + \chi_t \left(\delta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} \right), \quad q \rightarrow 0$$

One can show the following relation from the Sum Rule

$$K_{\mu\nu}(0, 0) = \chi_l \delta_{\mu\nu}$$

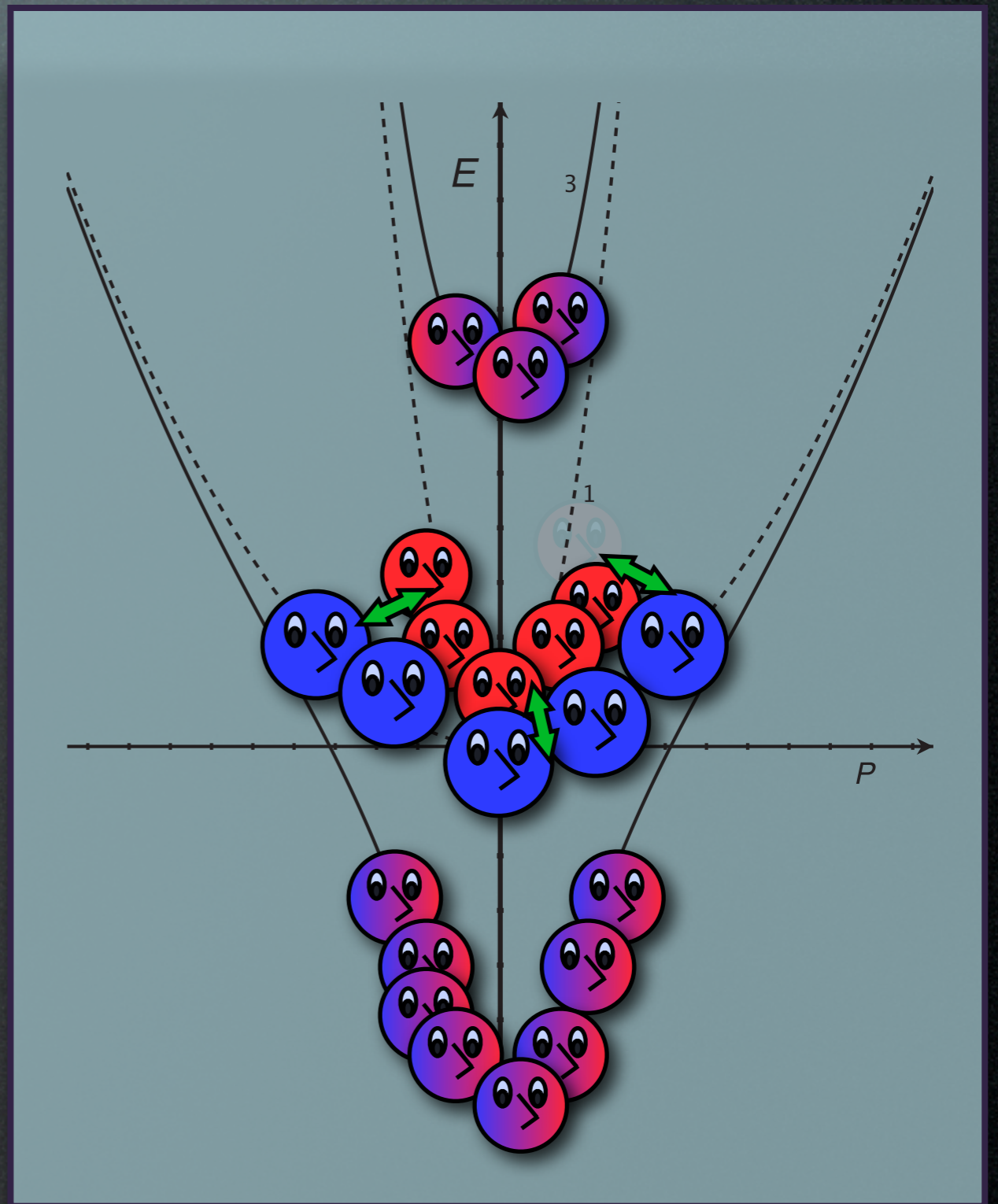
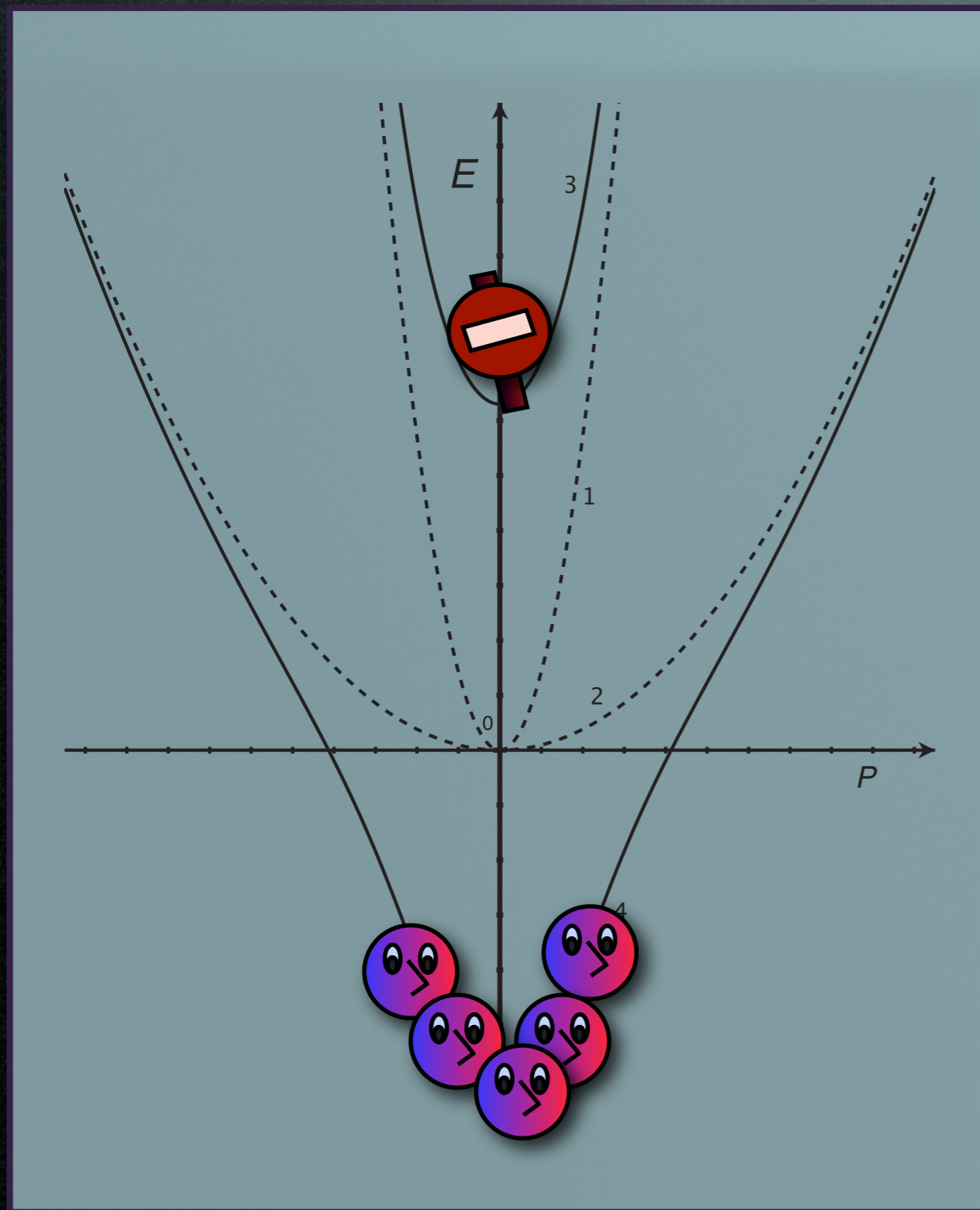
Effective action

$$\Pi_{\mu\nu}(q, 0) = (\chi_l - \chi_t) \left(\delta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} \right), \quad q \rightarrow 0$$

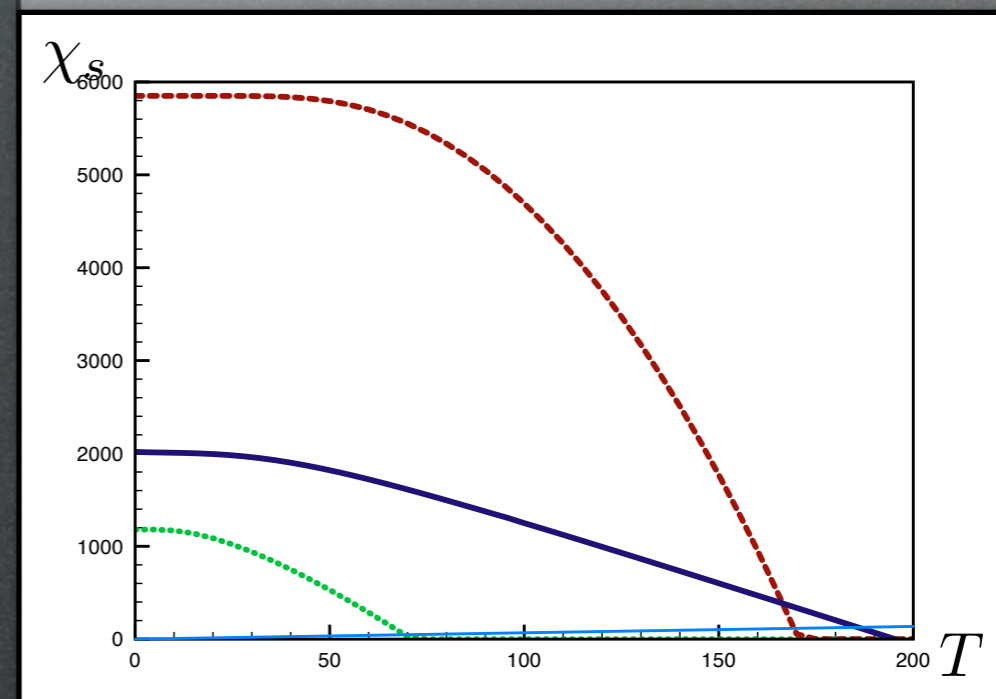
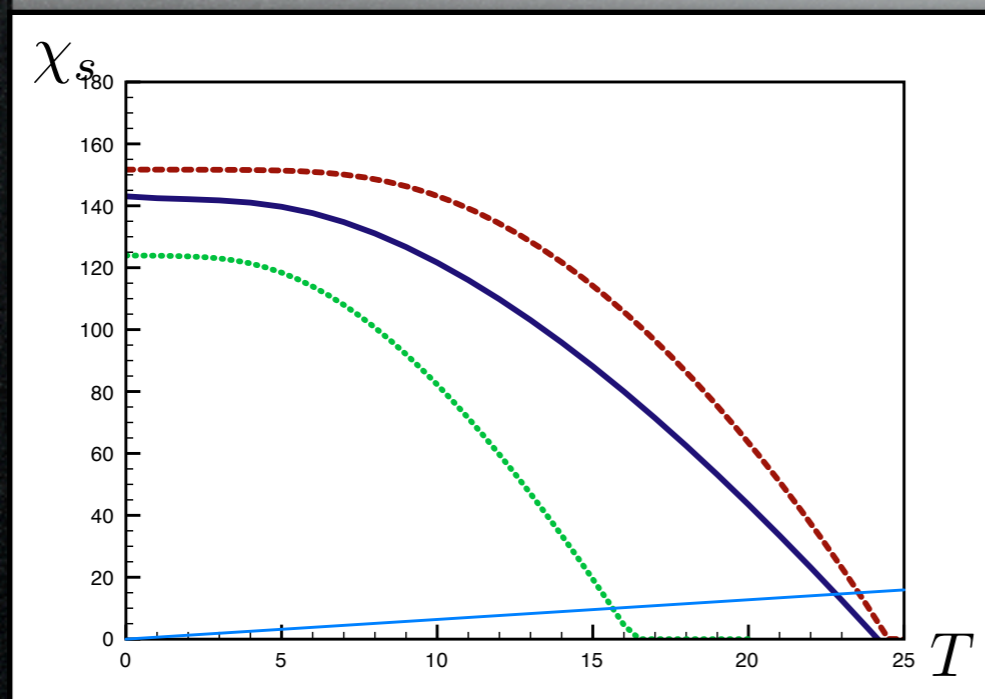
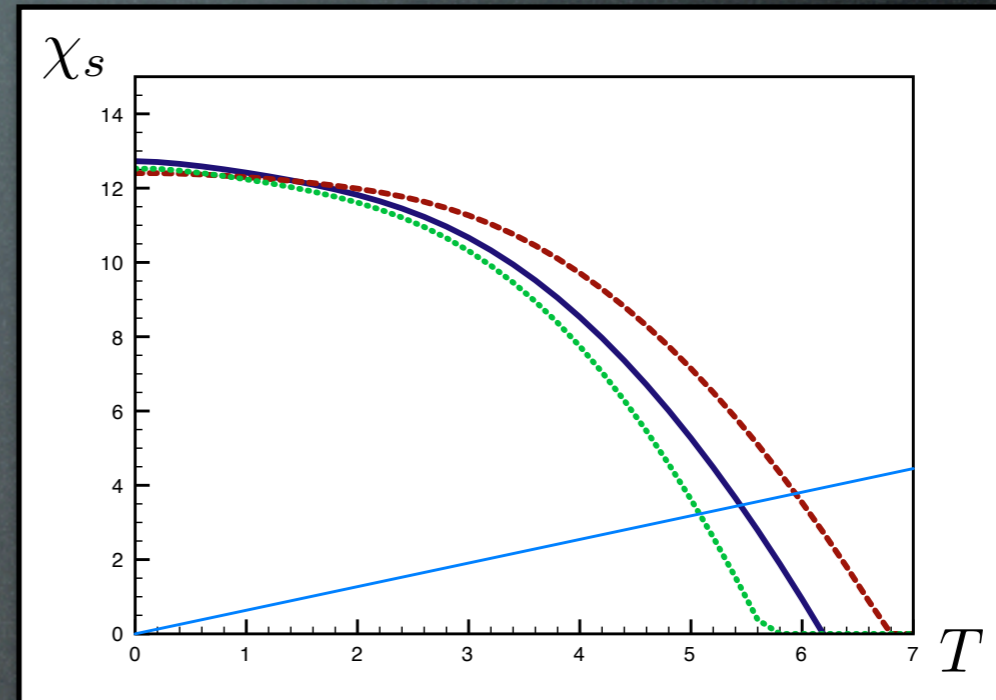
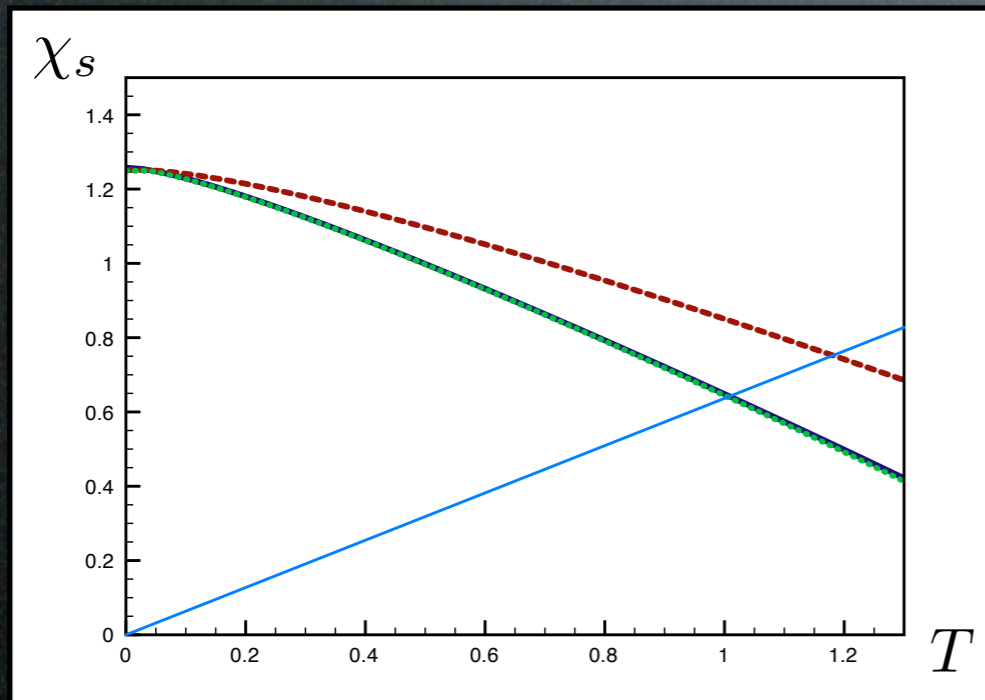
$$S_{eff} = \frac{\chi_s}{2T} \int (\partial_\mu \varphi_\nu(x))^2 d^2x$$

$$\chi_s = \chi_l - \chi_t$$

Considered approximations

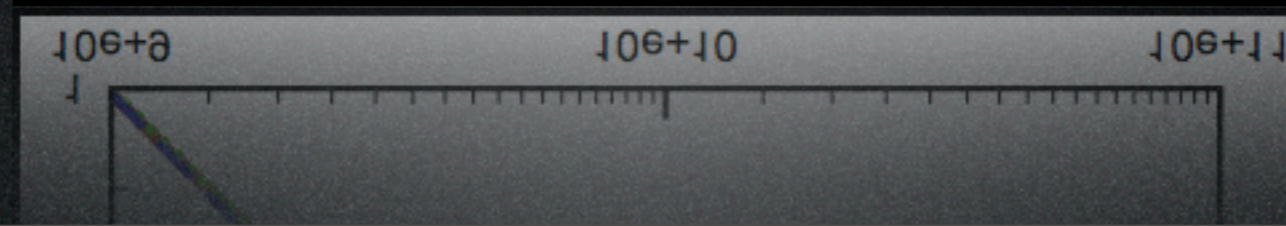
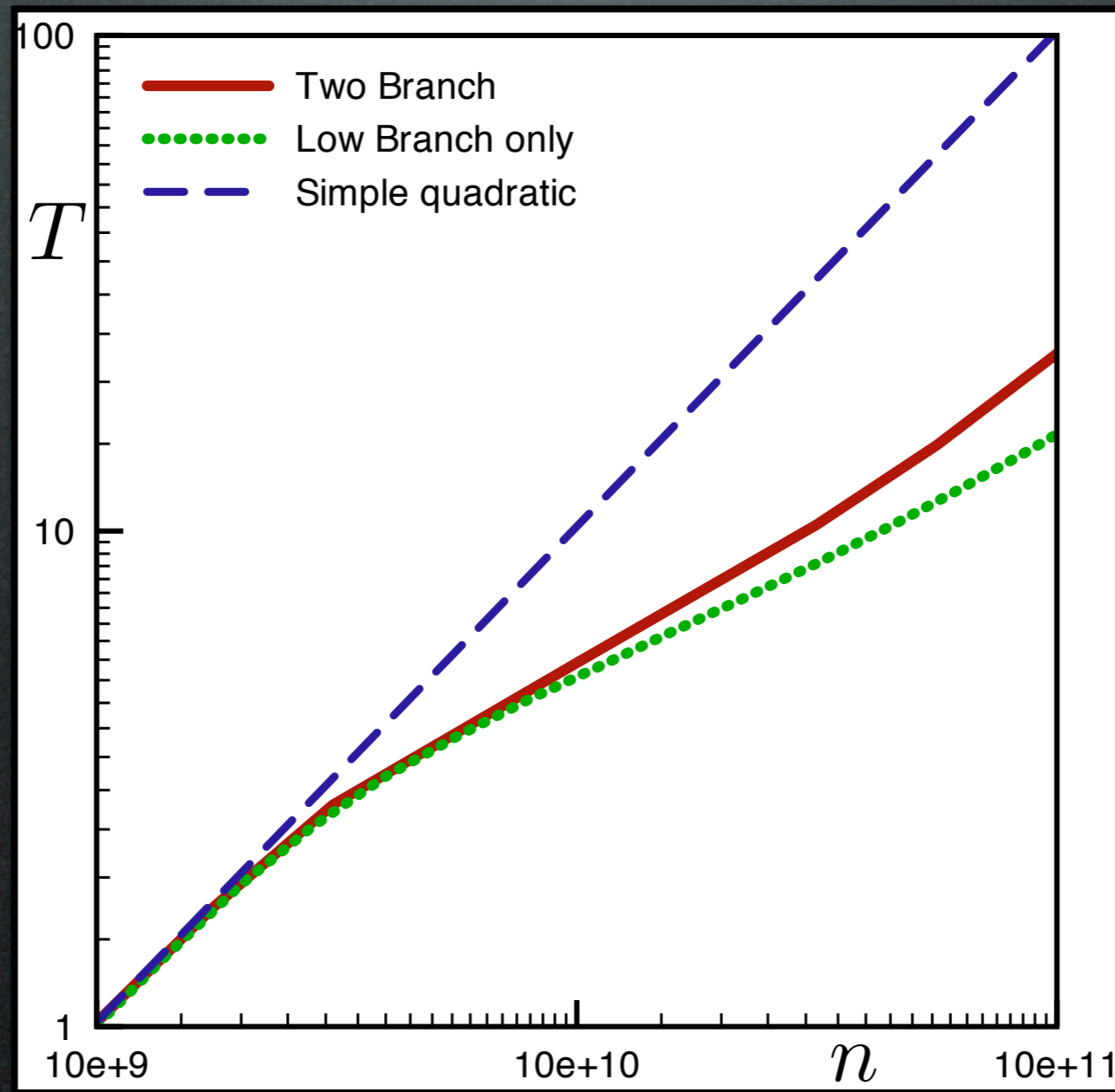


Result of calculations



χ_s

Transition temperature



Conclusions

- Theory of superfluidity in polaritonic system are developed;
- Superfluid density are calculated as function of temperature;
- Temperature of Berezinskii-Kosterlitz-Thouless transition are obtained for weakly interacting bose gas of polaritons.