



E
CMS Experiment at LHC CERN
Data recorded: Mon May 28 01:16:20 2012 CES
Run/Event: 195099 / 35438125
(lumi section: 65)
Orbit/Crossing: 16992111 / 2295

THE STANDARD MODEL AND EXPERIMENTAL ANOMALIES

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Raw $\Sigma E_T \sim 2$ TeV

14 jets with $E_T > 40$ GeV

Estimated PU ~ 50

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i \bar{\psi} D^\mu \psi \\ & + \bar{\chi}_i U_{ij} \chi_j \phi + h.c. \\ & + |D_\mu \phi|^2 - V(\phi) \\ & + \sqrt{-g} R\end{aligned}$$

Утренний кофе теоретика

Higgs sector

SUSY Higgs Bosons

SM

$$H = \begin{pmatrix} H^0 \\ H^- \end{pmatrix} = \begin{pmatrix} v + \frac{S+iP}{\sqrt{2}} \\ H^- \end{pmatrix} = \exp(i \frac{\vec{\xi} \vec{\sigma}}{2}) \begin{pmatrix} v + \frac{S}{\sqrt{2}} \\ 0 \end{pmatrix}$$

$4=2+2=3+1$

$$H \rightarrow H' = \exp(i \frac{\vec{\alpha} \vec{\sigma}}{2}) H \xrightarrow{(\vec{\alpha} = -\vec{\xi})} H' = \begin{pmatrix} v + \frac{S}{\sqrt{2}} \\ 0 \end{pmatrix}$$

MSSM

$8=4+4=3+5$

$$H_1 = \begin{pmatrix} H_1^0 \\ H_1^- \end{pmatrix} = \begin{pmatrix} v_1 + \frac{S_1+iP_1}{\sqrt{2}} \\ H_1^- \end{pmatrix}, \quad H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix} = \begin{pmatrix} H_2^+ \\ v_2 + \frac{S_2+iP_2}{\sqrt{2}} \end{pmatrix},$$

$$v_1^2 + v_2^2 = v^2, \quad v_2/v_1 \equiv \tan\beta$$

NMSSM

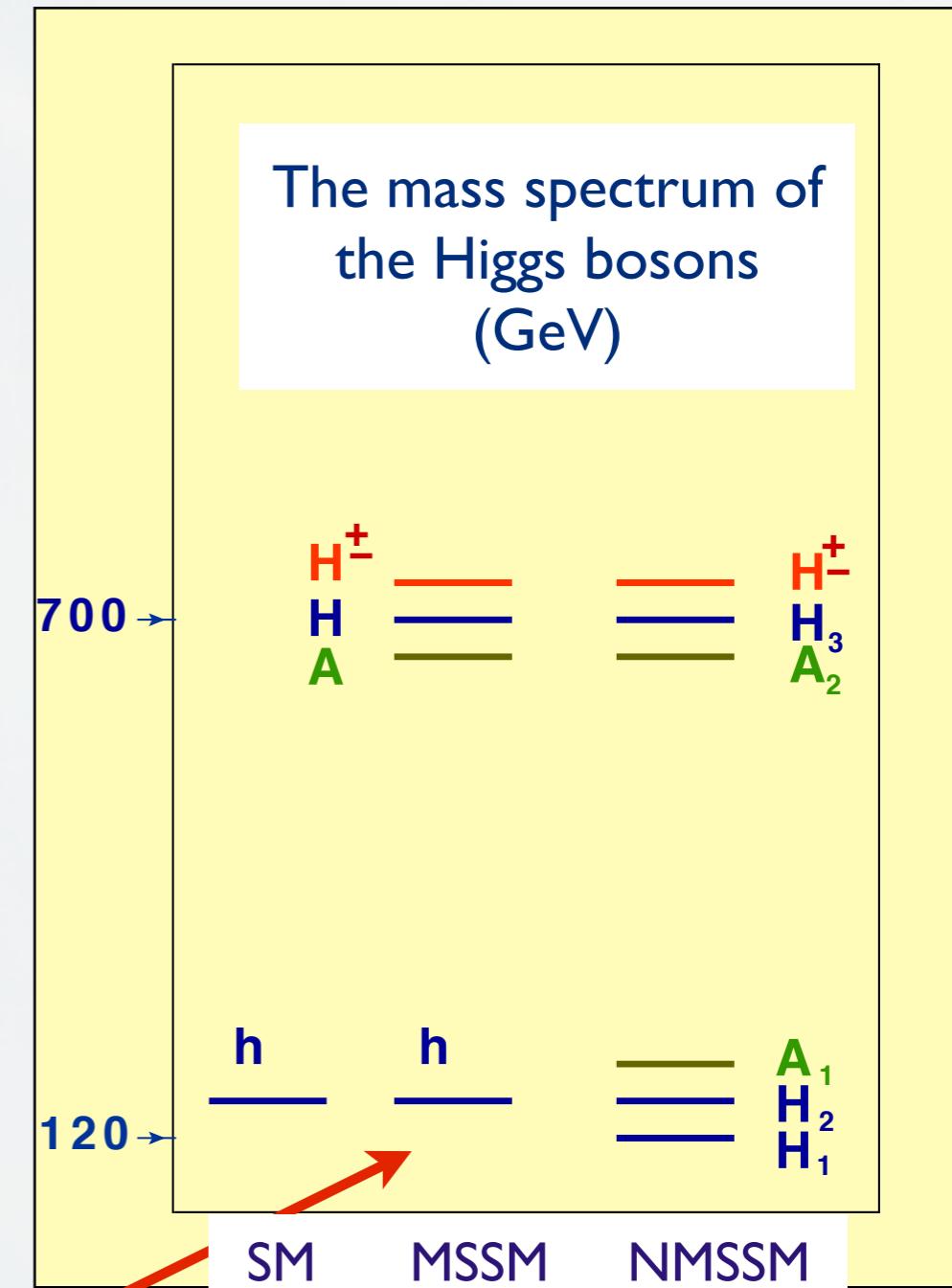
+2

$S + P$



The Higgs Sector: Alternatives

Model	Particle content
SM	h CP-even
2HDM/MSSM	h, H CP-even A CP-odd H^\pm
NMSSM	H_1, H_2, H_3 CP-even A_1, A_2 CP-odd H^\pm
Composite	h CP-even + excited states

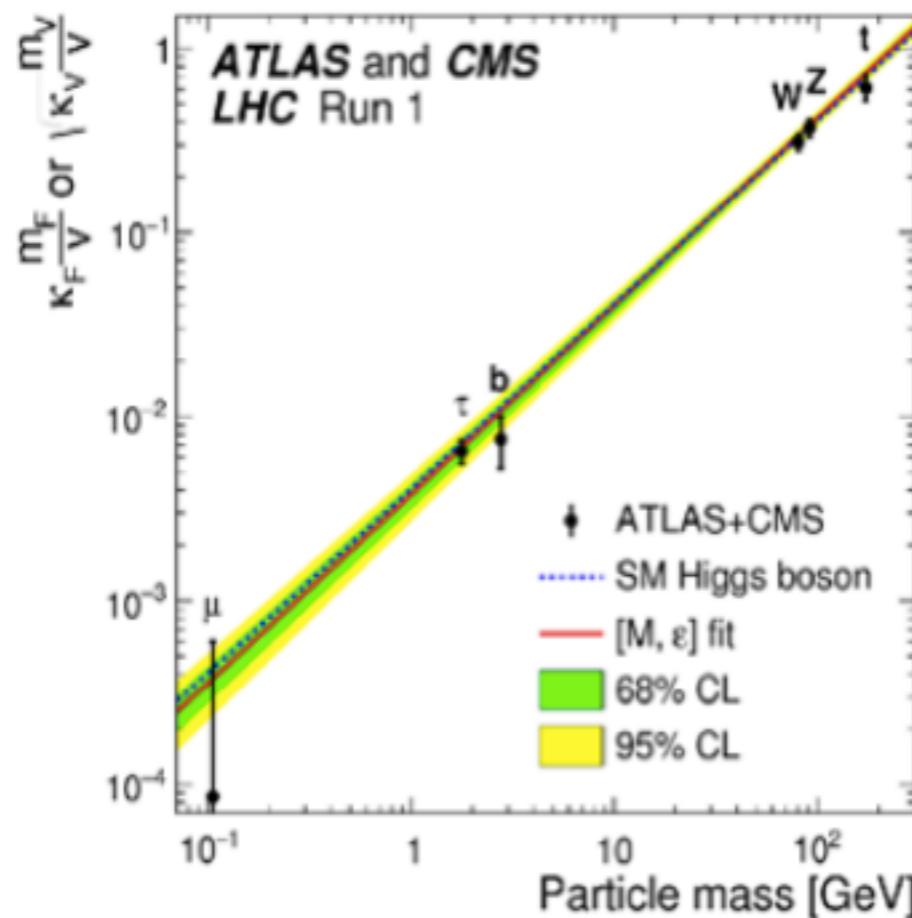


It may well be that we see one of these states

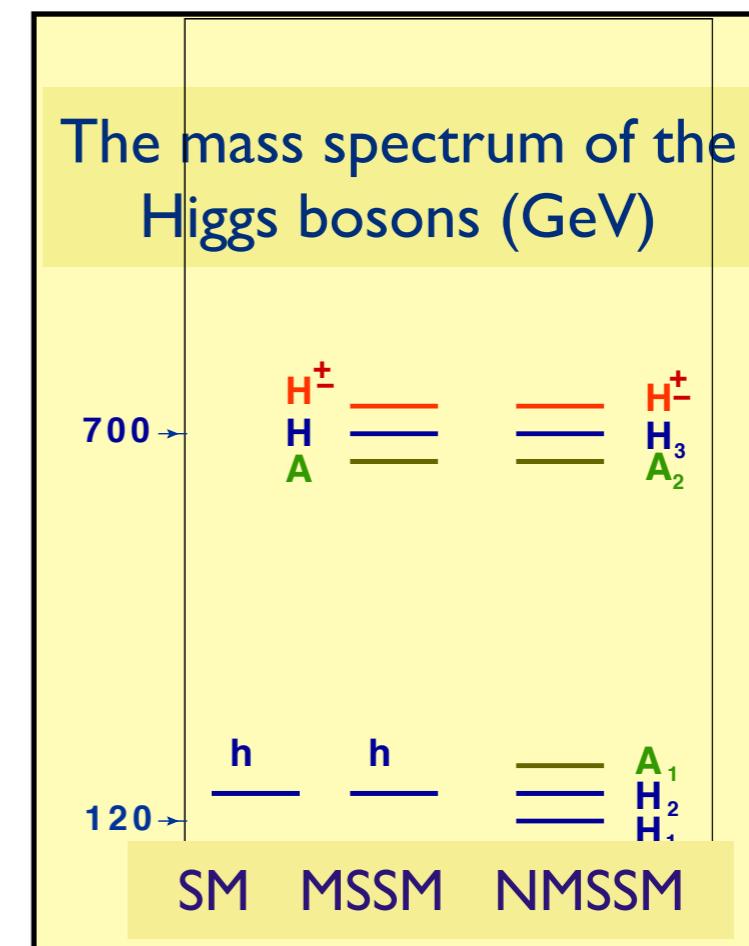
One has to check the presence or absence of heavy Higgs bosons

How to probe?

- Probe deviations from the SM Higgs couplings

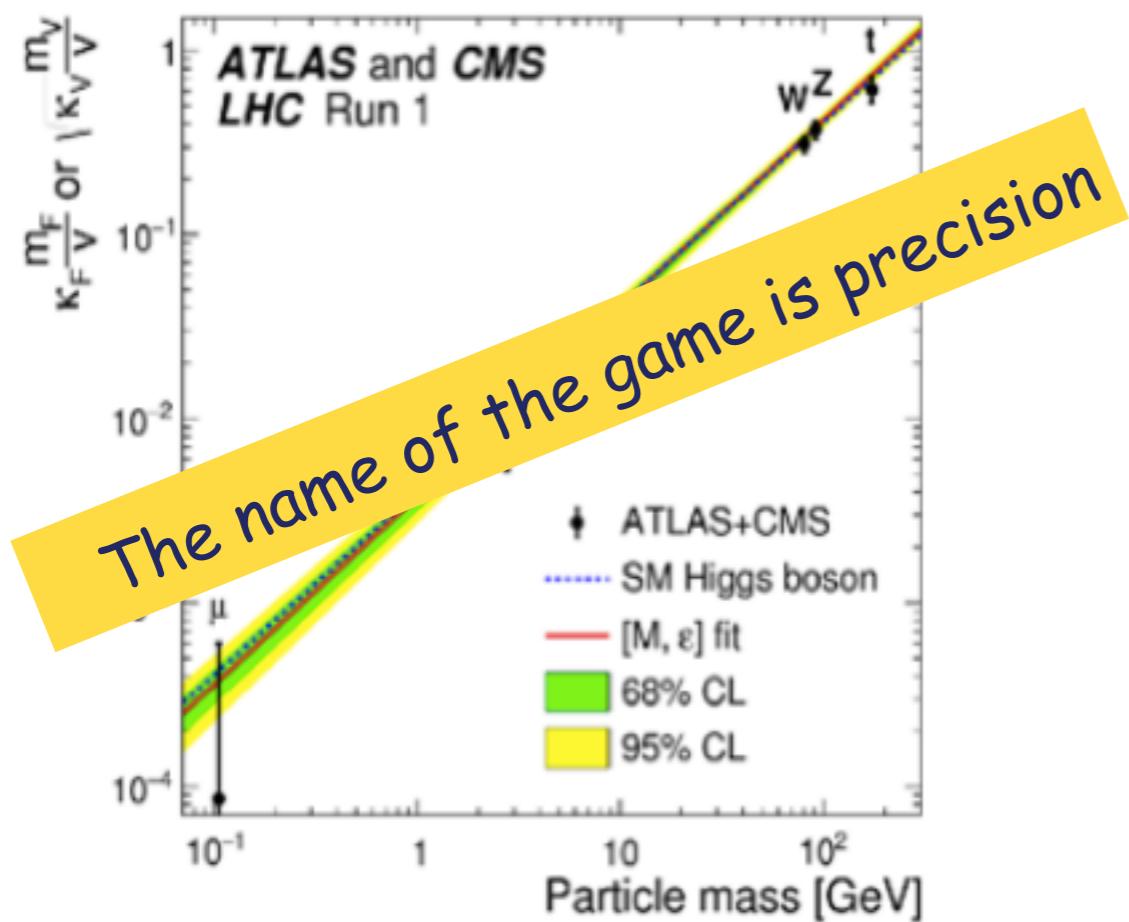


- Perform direct search for additional scalars

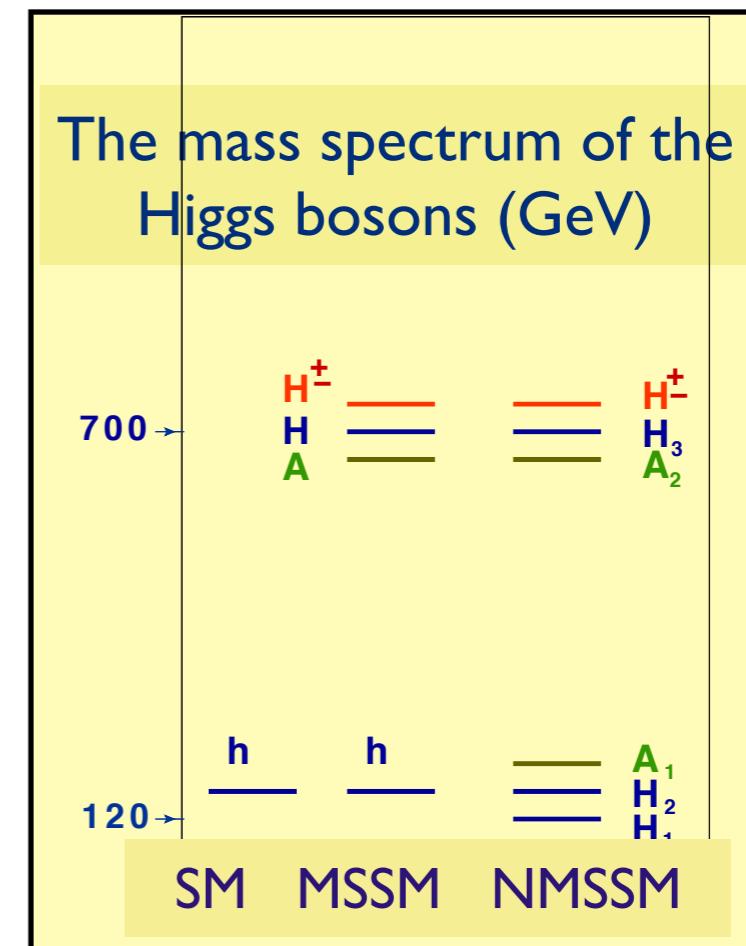


How to probe?

- Probe deviations from the SM Higgs couplings

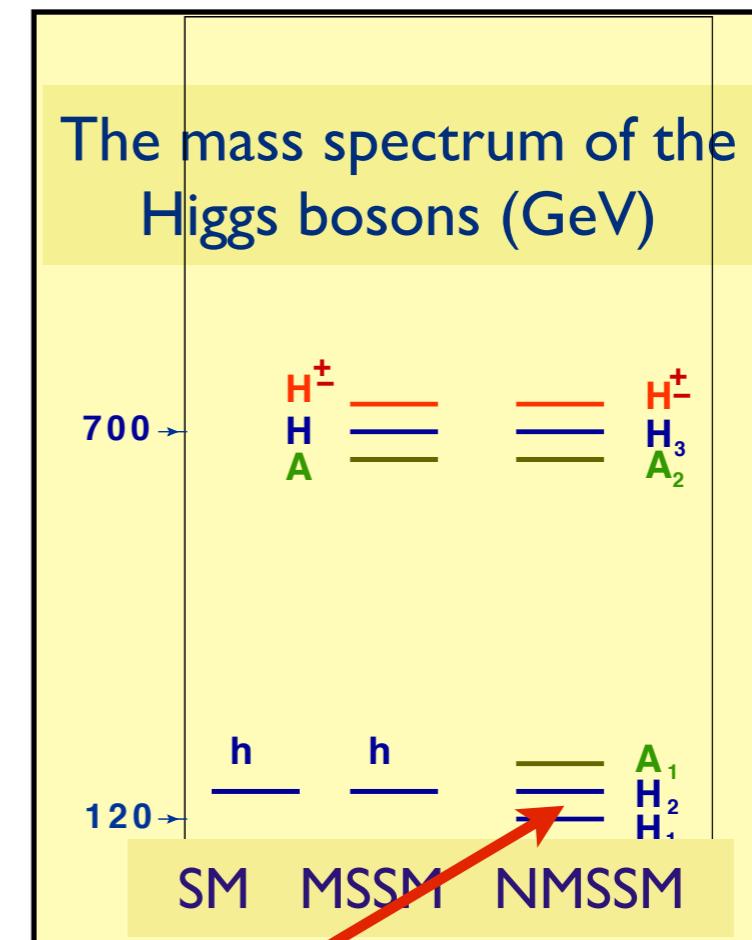
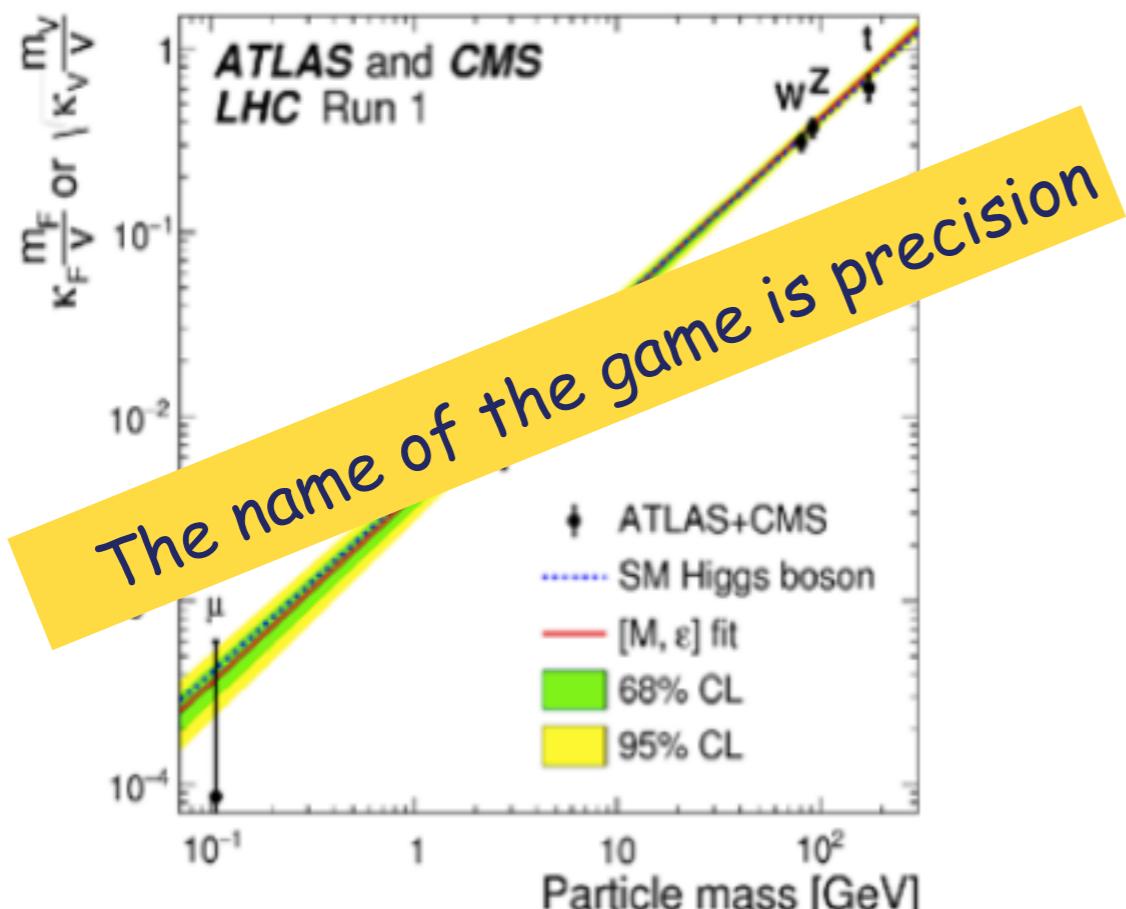


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- Probe deviations from the SM Higgs couplings
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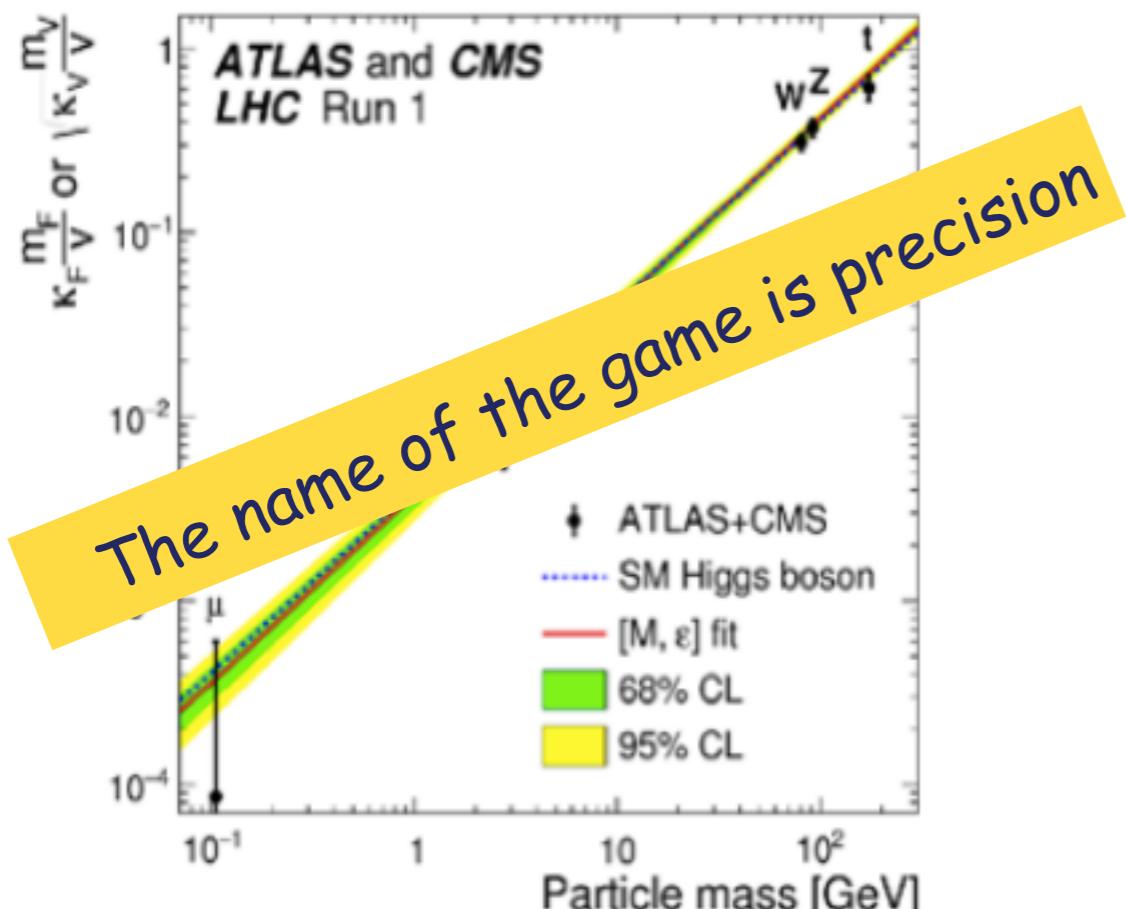


We may have found one of these states

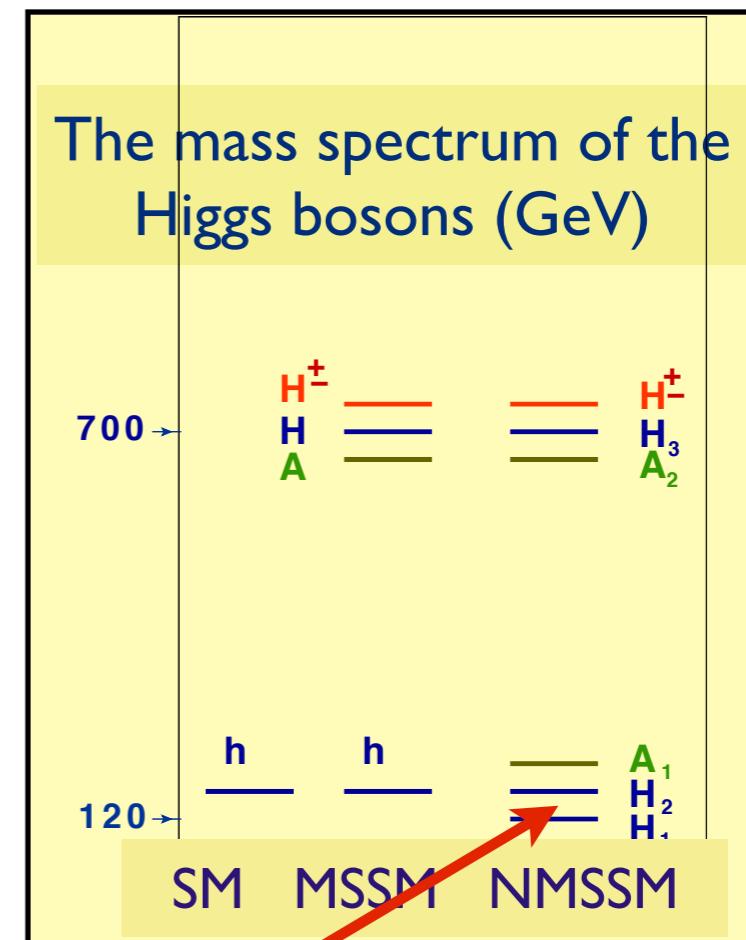


How to probe?

- Probe deviations from the SM Higgs couplings



- Perform direct search for additional scalars



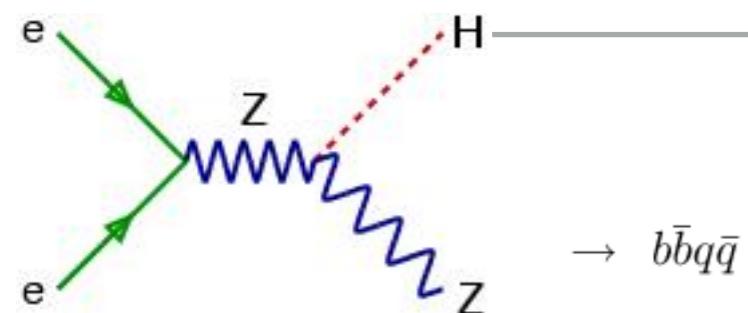
We may have found one of these states

One has to check the presence or absence of heavy Higgs bosons



PRECISION PHYSICS OF THE HIGGS BOSONS

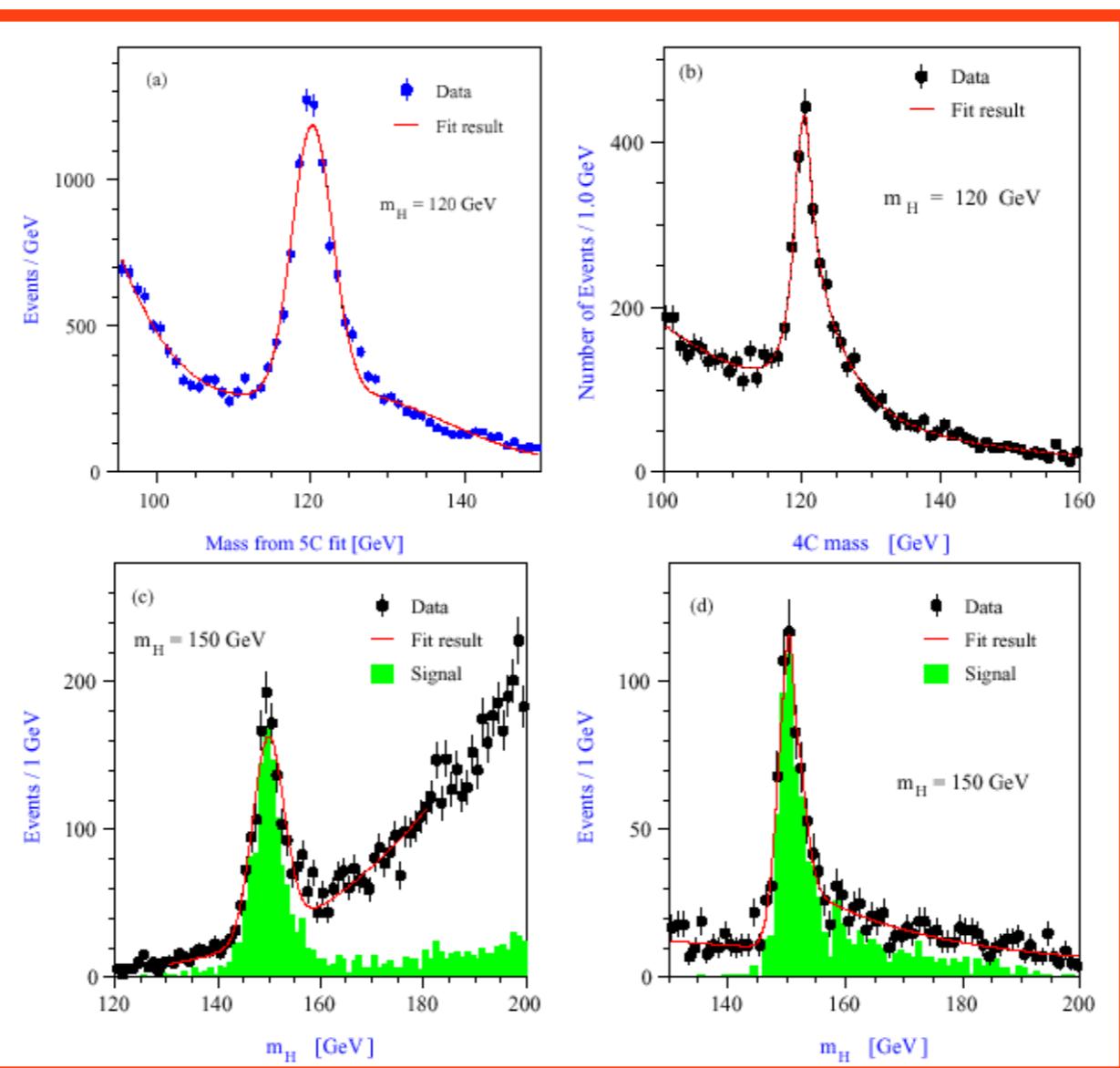
7



$ee \rightarrow HZ$ diff. decay channels

Int Linear Collider

$\rightarrow W^+W^-q\bar{q}$



$\rightarrow q\bar{q}\ell^+\ell^-$

$\Delta m_H = 40$ MeV

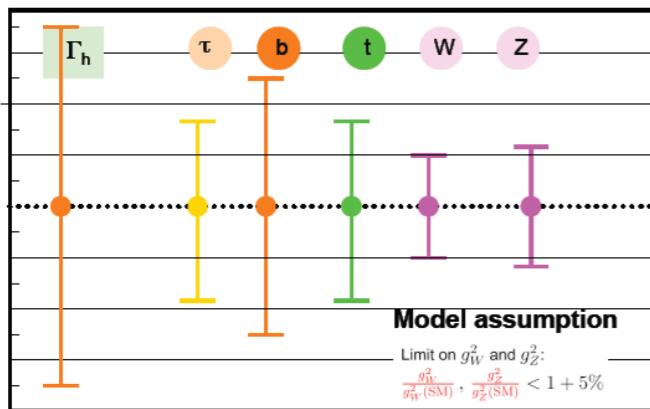
$\rightarrow W^+W^-\ell^+\ell^-$

$\Delta m_H = 70$ MeV

Coupling Precision

LHC 300 $\text{fb}^{-1} \times 2$

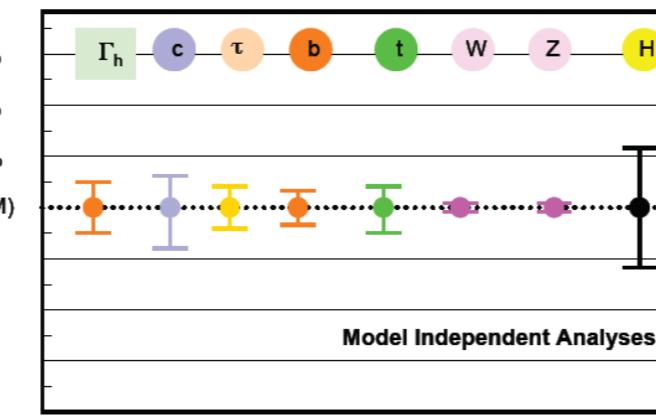
Deviation from SM value



Coupling Precision

ILC

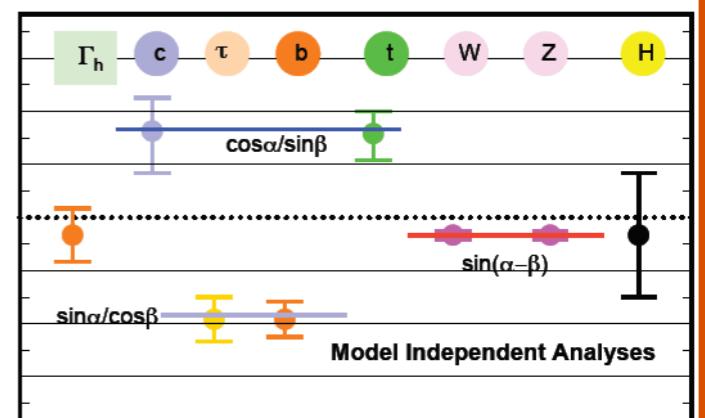
Deviation from SM value



SUSY or 2HDM

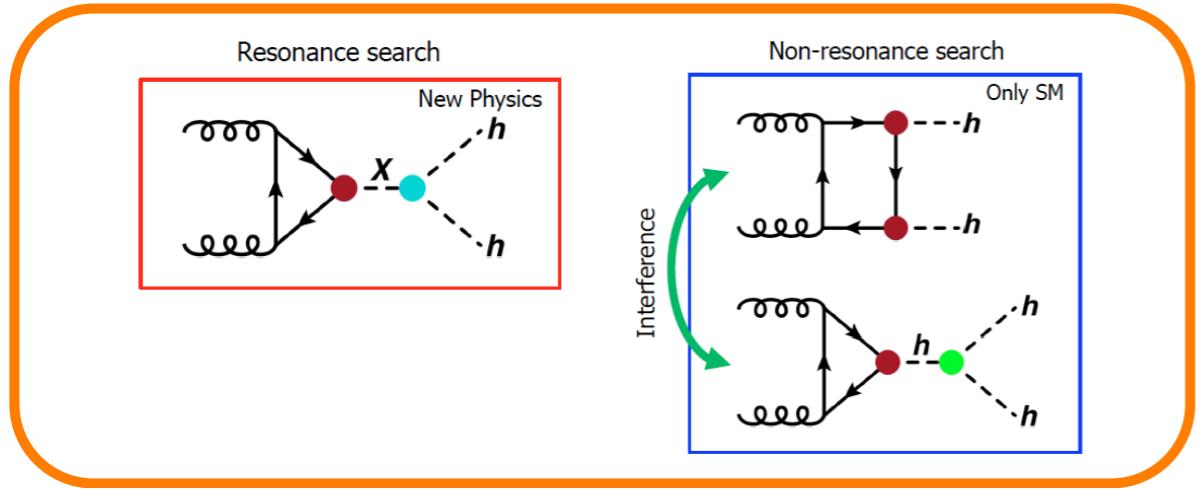
ILC

Deviation from SM value

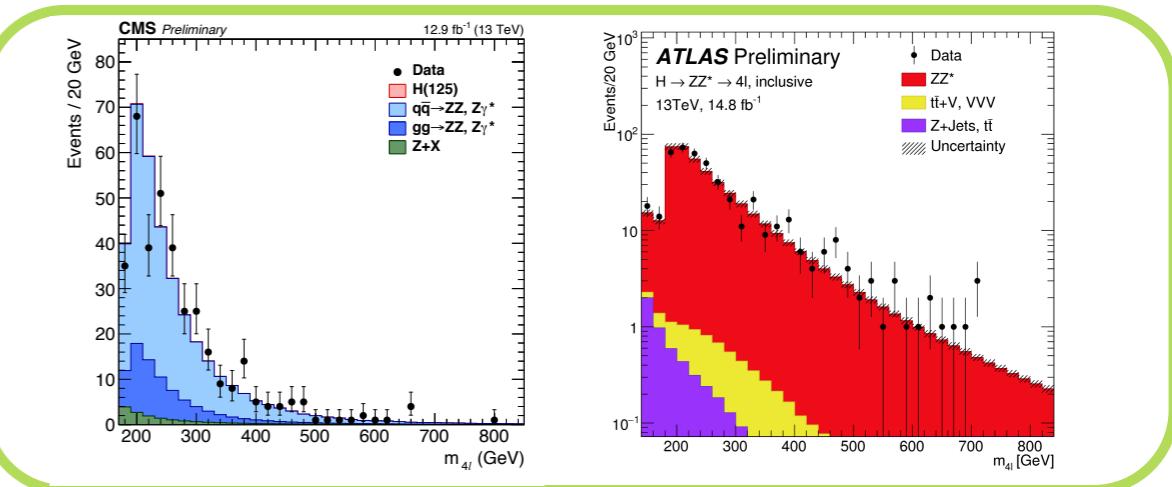


EXTRA HIGGS BOSONS

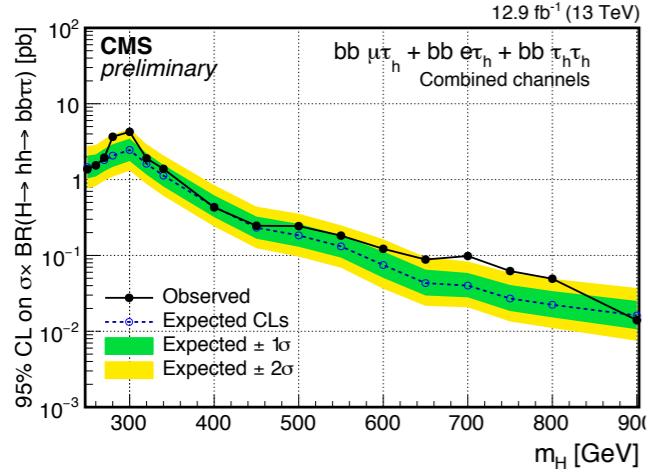
Higgs \rightarrow hh \rightarrow bb $\tau\tau$



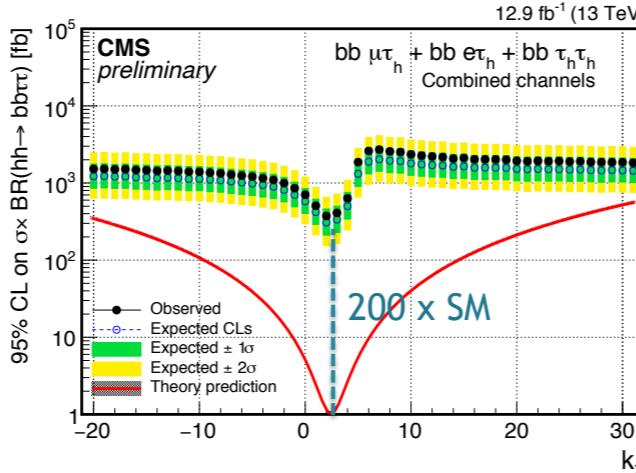
Heavy Higgs \rightarrow ZZ \rightarrow 4l



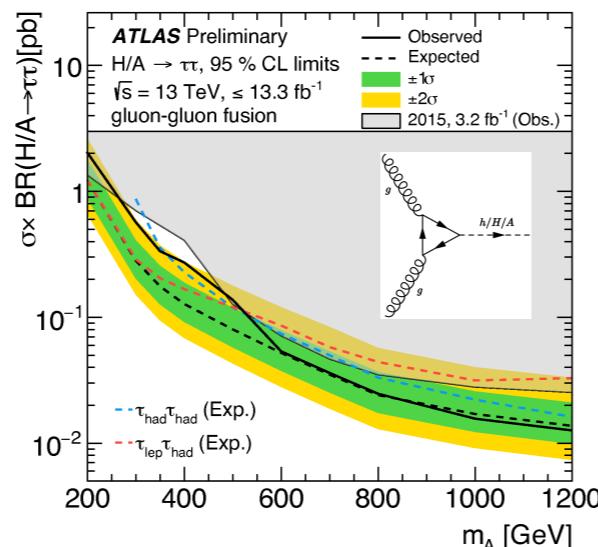
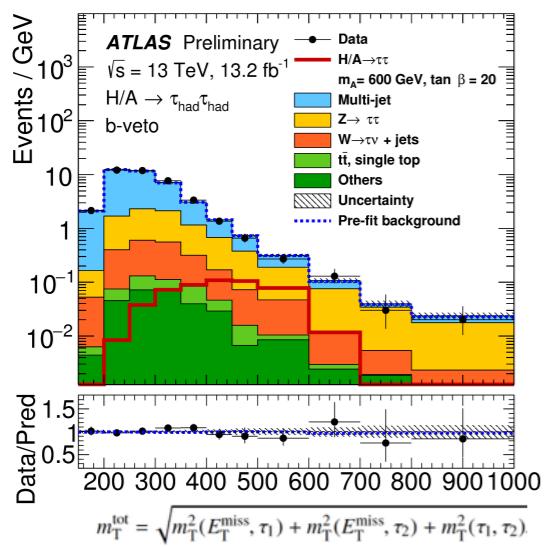
Resonant



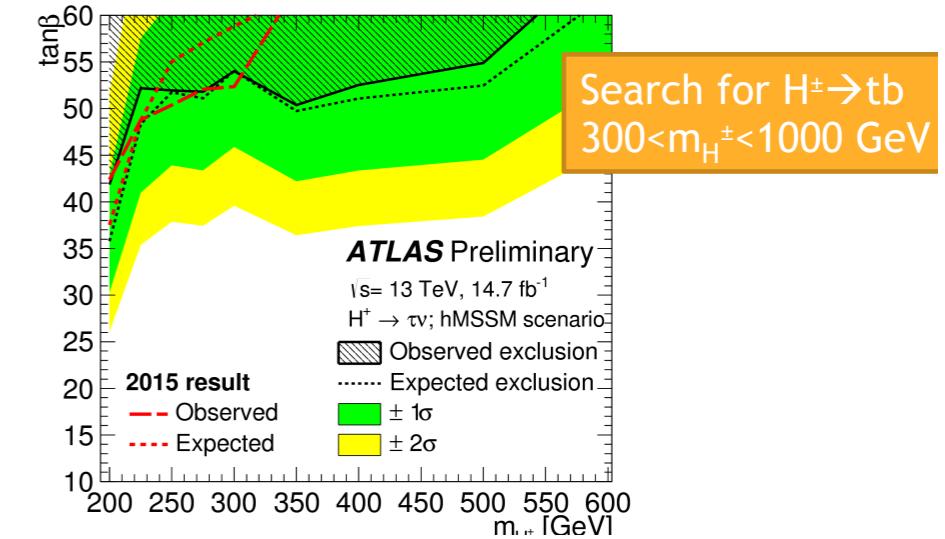
Non-Resonant



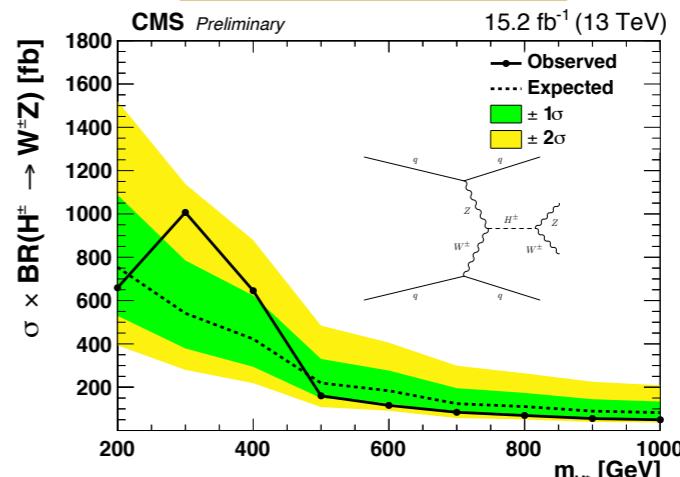
Heavy Higgs \rightarrow tau tau



Charged Higgs

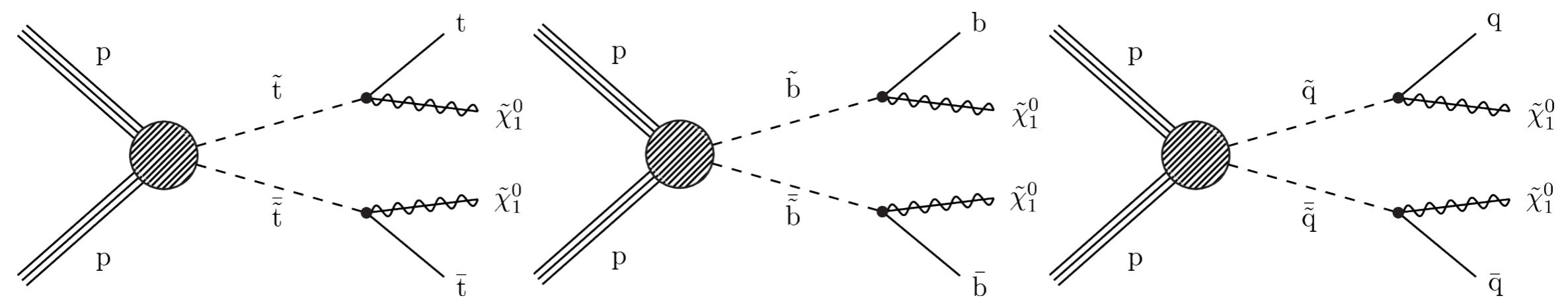
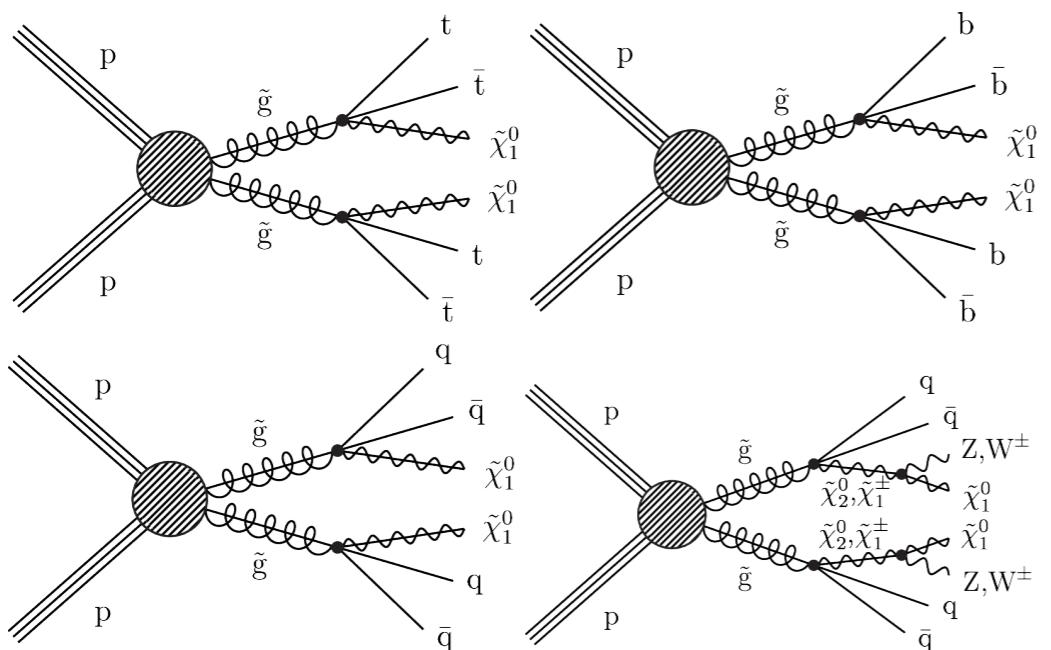


Search for H \pm WZ

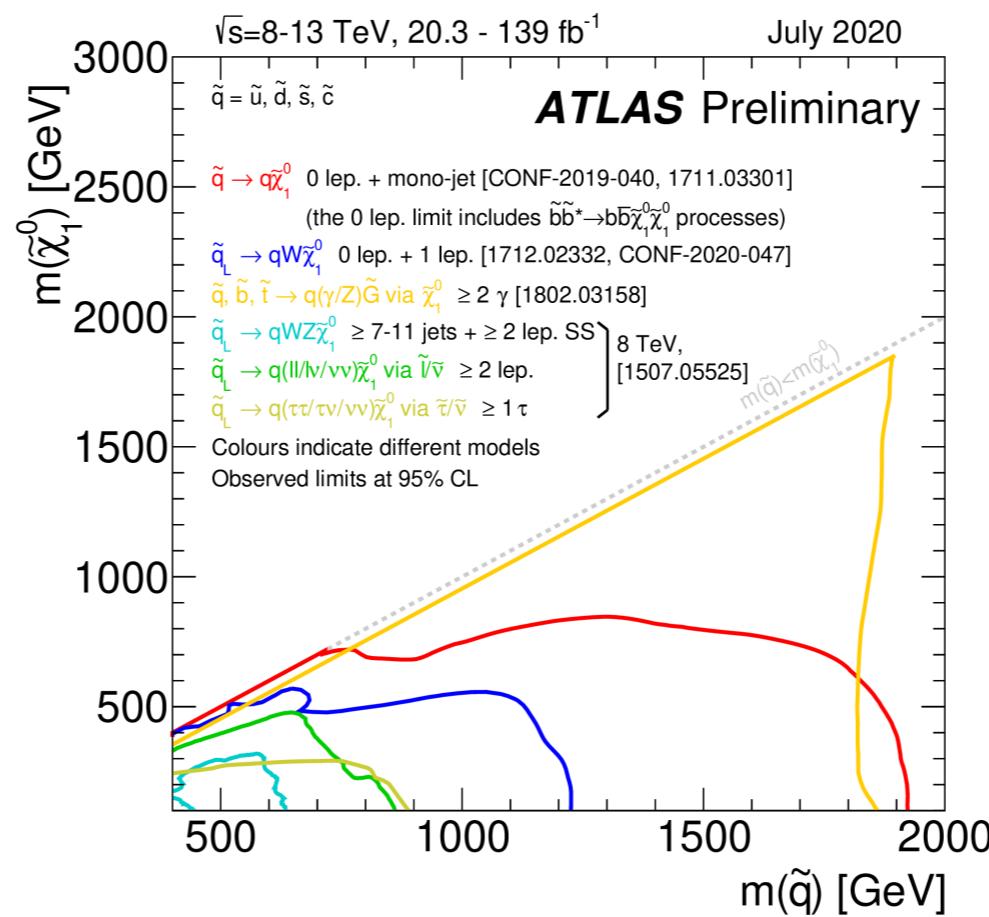
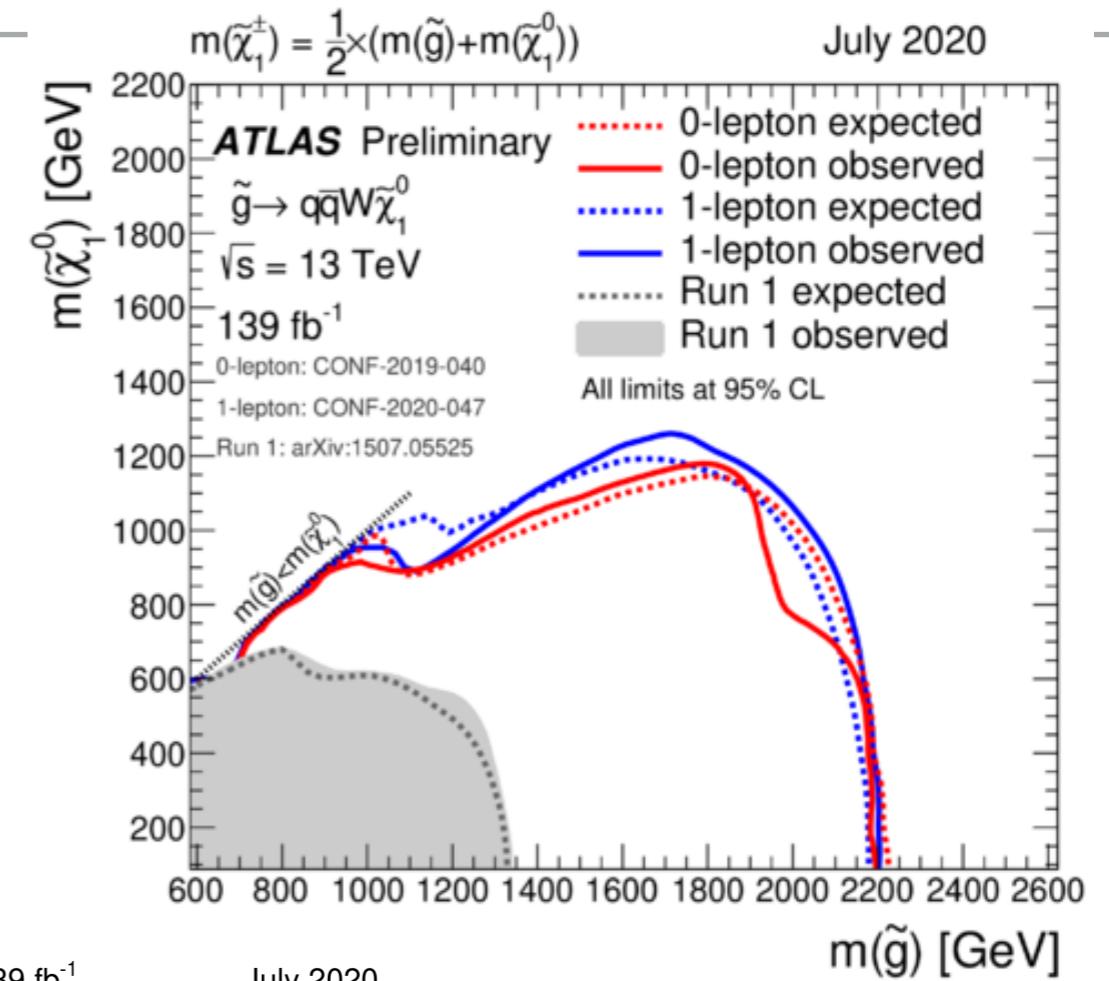
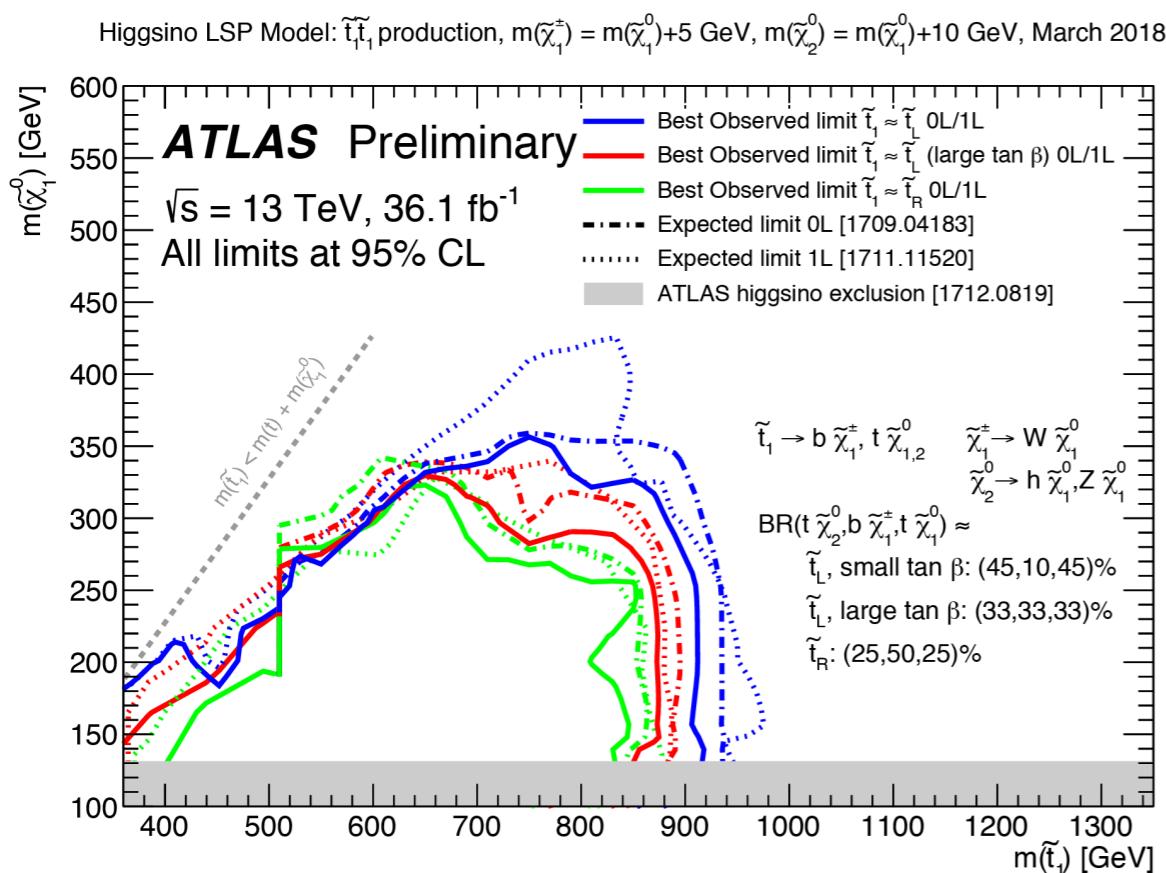


SUSY Searches

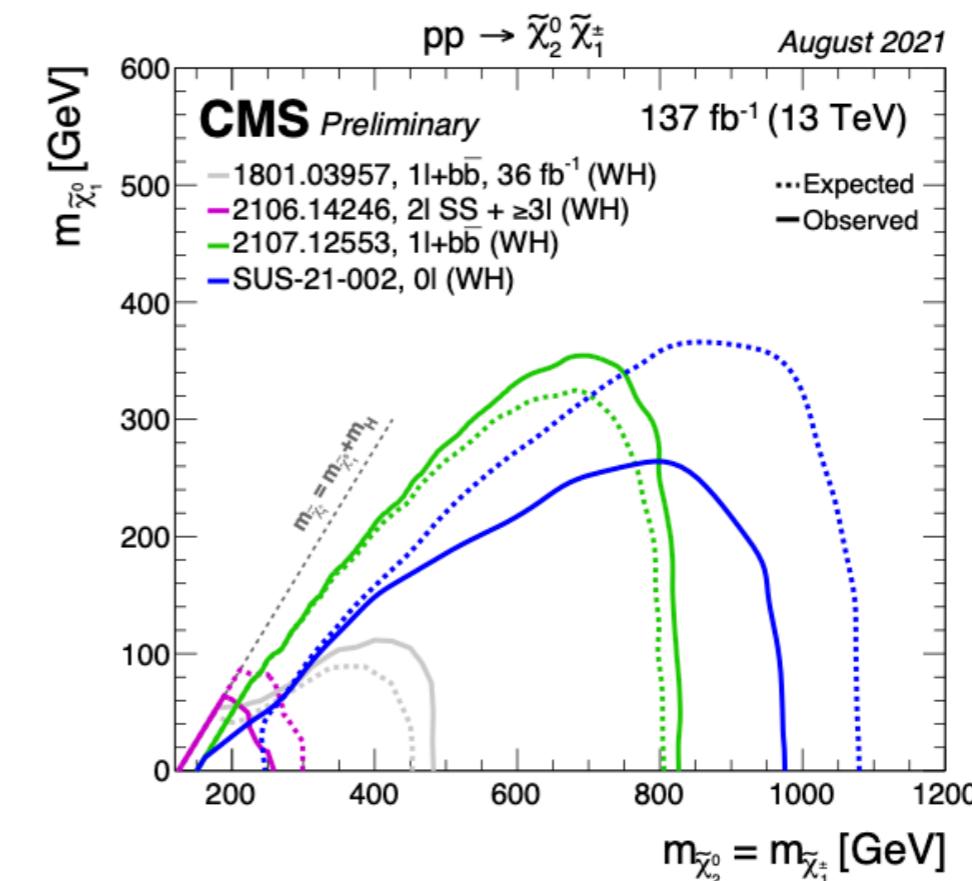
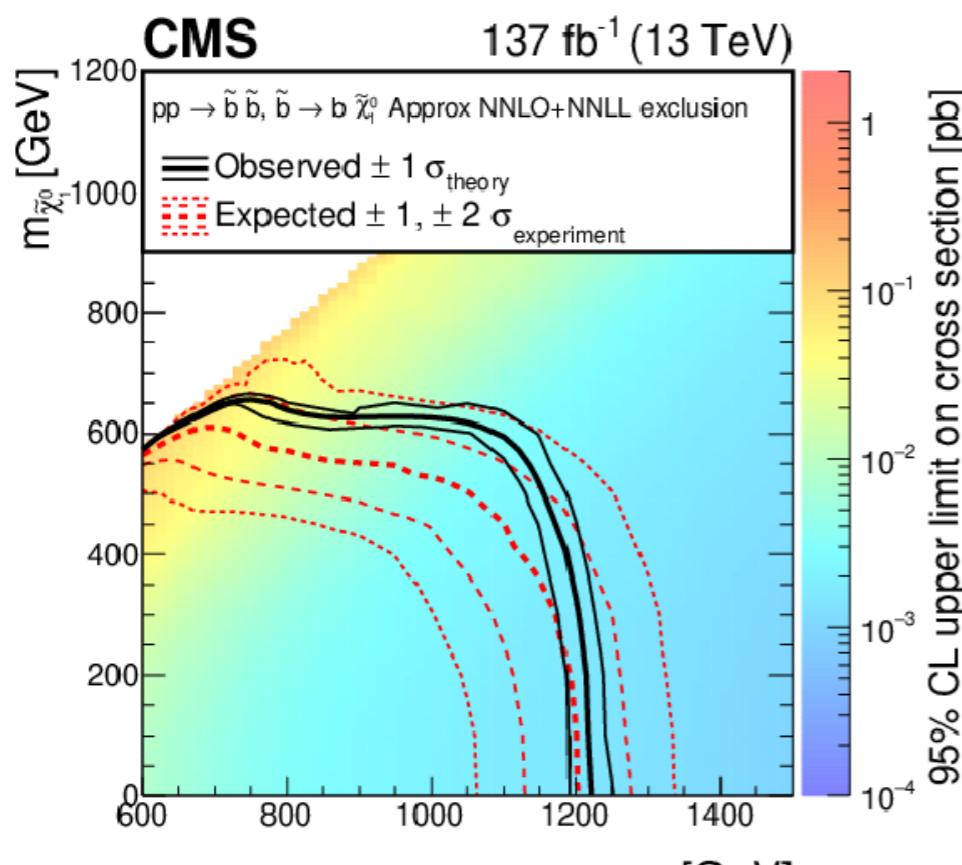
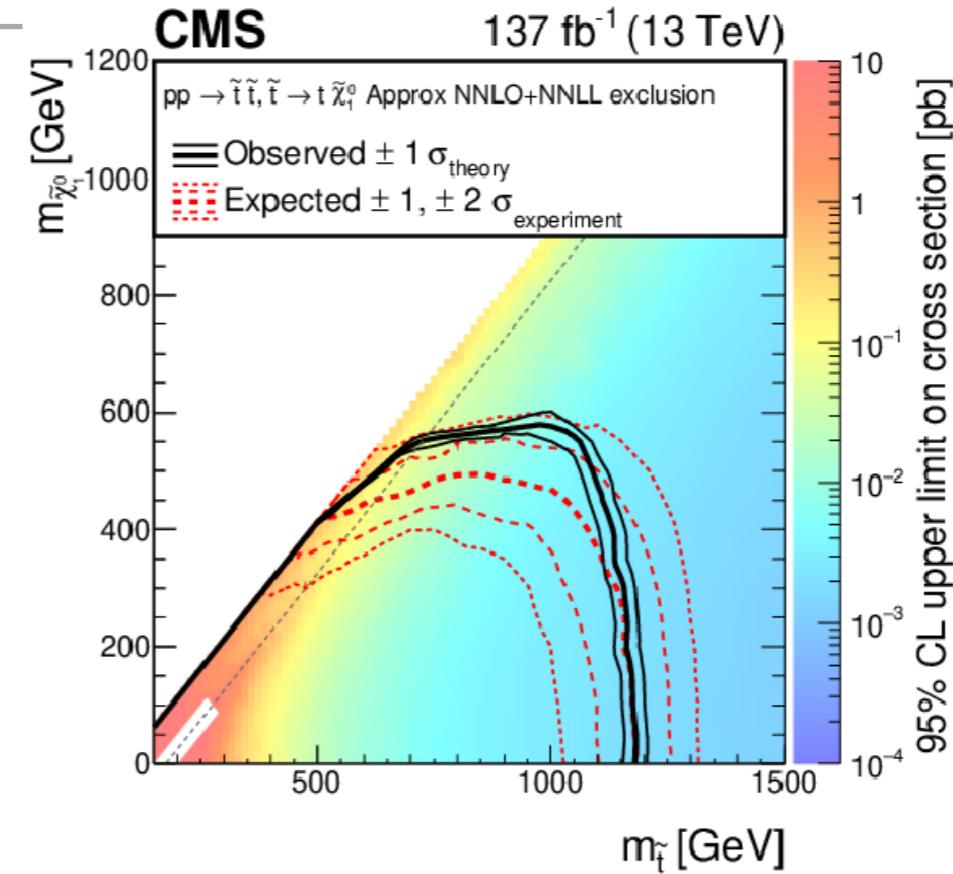
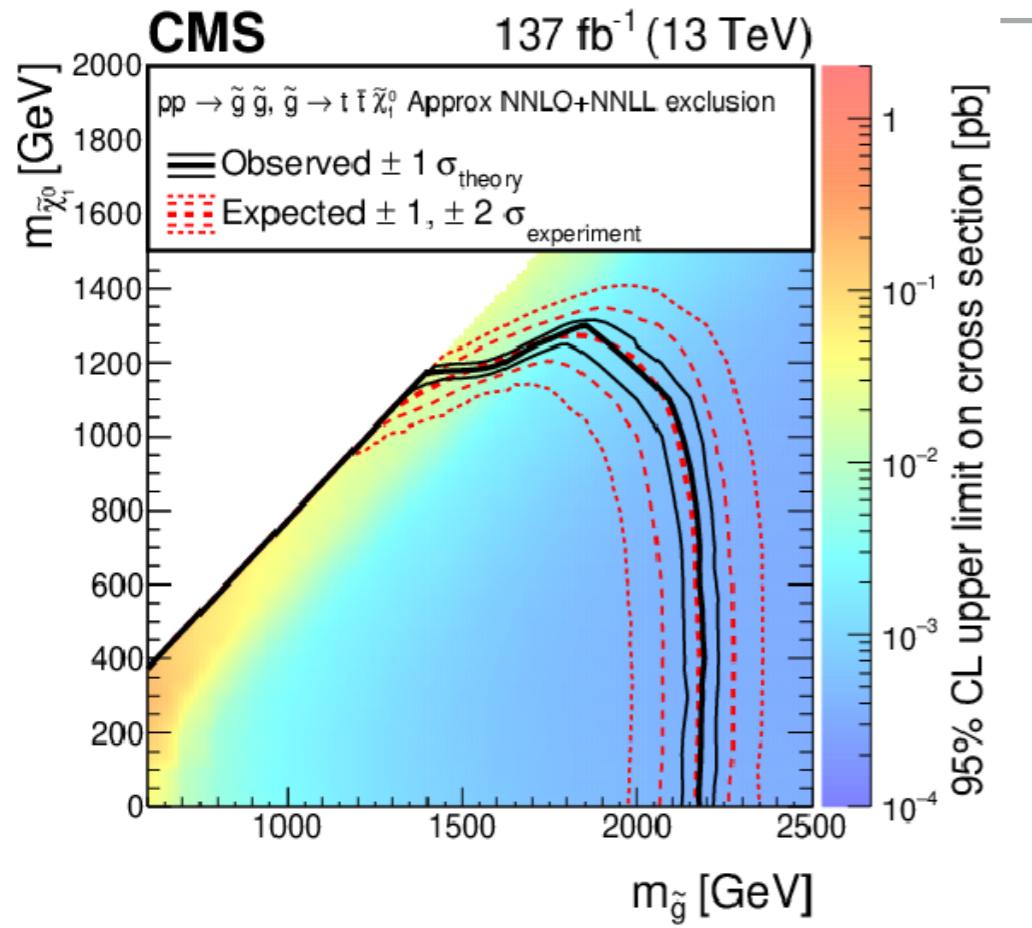
SUSY PRODUCTION



SUSY LIMITS

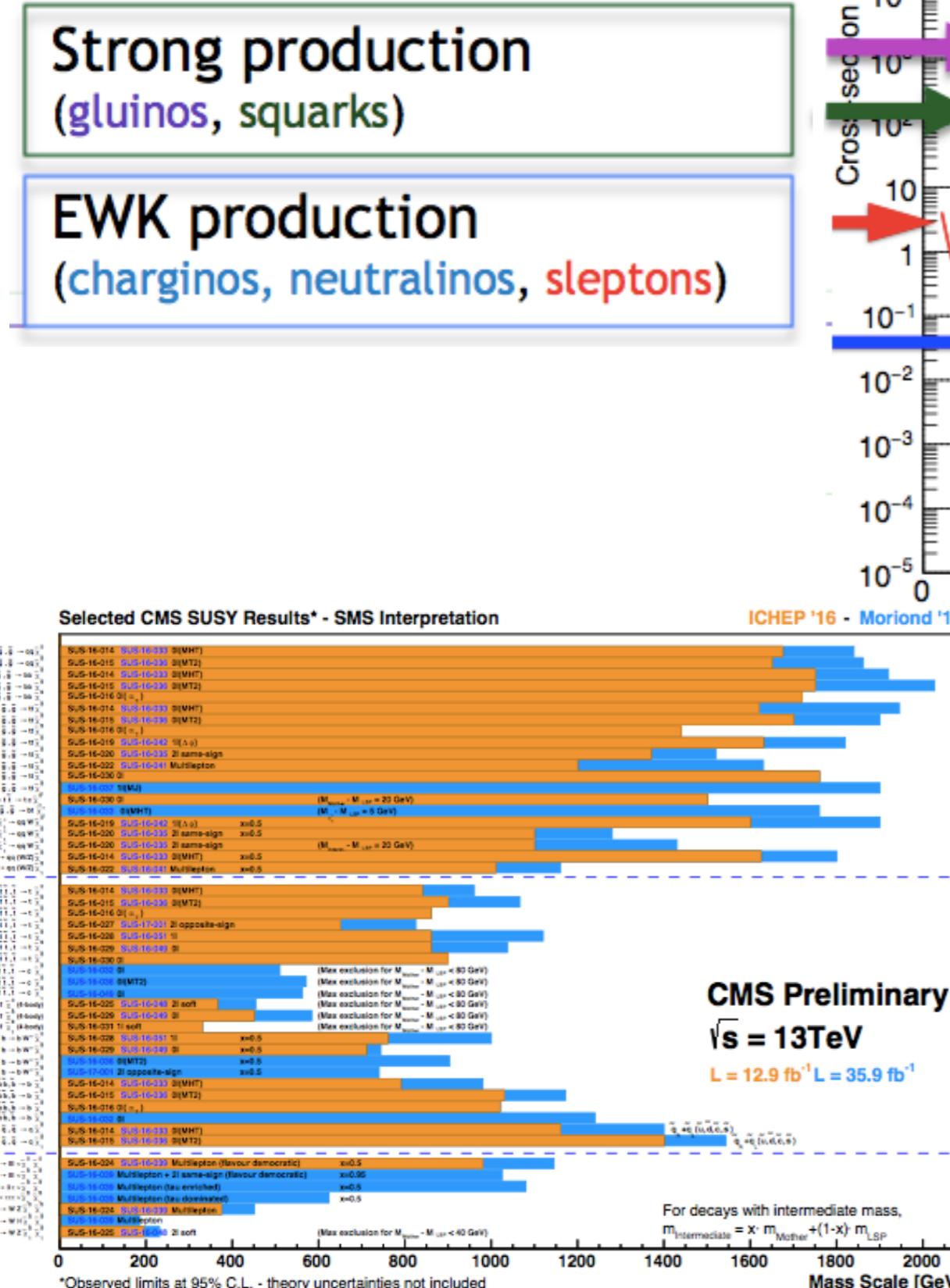


SUSY LIMITS



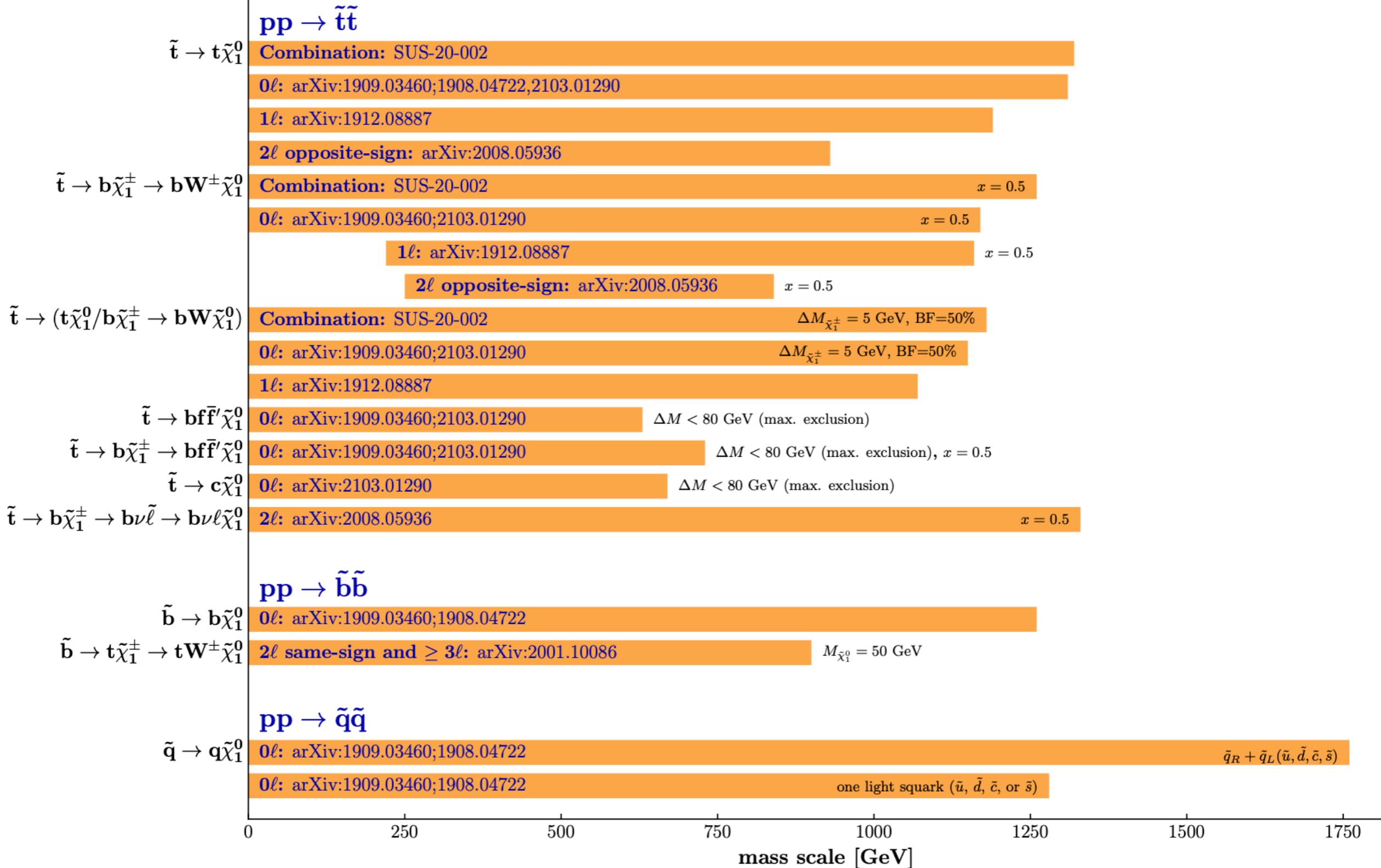
SUSY SEARCHES

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Overview of SUSY results: squark pair production

137 fb^{-1} (13 TeV)



Selection of observed limits at 95% C.L. (theory uncertainties are not included). Probe **up** to the quoted mass limit for light LSPs unless stated otherwise. The quantities ΔM and x represent the absolute mass difference between the primary sparticle and the LSP, and the difference between the intermediate sparticle and the LSP relative to ΔM , respectively, unless indicated otherwise.

Overview of SUSY results: electroweak production

137 fb⁻¹ (13 TeV)

$pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm$

$pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow \ell \tilde{\nu} \ell \tilde{\ell} \rightarrow \ell \nu \ell \ell \tilde{\chi}_1^0 \tilde{\chi}_1^0$

2 ℓ same-sign and 3 ℓ : SUS-19-012

flavour democratic, $x = 0.5$

2 ℓ same-sign and $\geq 3\ell$: SUS-19-012

flavour democratic, $x = 0.05$

2 ℓ same-sign and $\geq 3\ell$: SUS-19-012

flavour democratic, $x = 0.95$

$pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow \tilde{\tau} \nu \ell \tilde{\ell} \rightarrow \tau \nu \ell \ell \tilde{\chi}_1^0 \tilde{\chi}_1^0$

2 ℓ same-sign and 3 ℓ/τ_h : SUS-19-012

τ enriched, $x = 0.5$

3 ℓ/τ_h : SUS-19-012

τ enriched, $x = 0.05$

3 ℓ/τ_h : SUS-19-012

τ enriched, $x = 0.95$

$pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow \tilde{\tau} \nu \tau \tilde{\tau} \rightarrow \tau \nu \tau \tau \tilde{\chi}_1^0 \tilde{\chi}_1^0$

$\geq 3\ell/\tau_h$: SUS-19-012

τ dominated, $x = 0.5$

$pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow WH \tilde{\chi}_1^0 \tilde{\chi}_1^0$

2 ℓ same-sign and $\geq 3\ell/\tau_h$: SUS-19-012

1 ℓ +jets: SUS-20-003

$pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow WZ \tilde{\chi}_1^0 \tilde{\chi}_1^0$

2 ℓ opposite-sign: arXiv:2012.08600

2 ℓ same-sign and 3 ℓ : SUS-19-012

2 ℓ and 3 ℓ soft: SUS-18-004 $\Delta M = 5\text{--}10$ GeV

$pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm / \tilde{\chi}_1^0 \tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm / \tilde{\chi}_2^0 \rightarrow (W^*/Z^*) \tilde{\chi}_1^0$

2 ℓ and 3 ℓ soft: SUS-18-004 higgsino simplified model, $\Delta M = 5\text{--}10$ GeV

$pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_1^\pm$

2 ℓ opposite-sign: arXiv:1807.07799 $M_{\tilde{\chi}_1^0} = 1$ GeV

$pp \rightarrow \tilde{\chi}_1^\pm \tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow (\tilde{\ell} \nu / \ell \tilde{\nu}) \rightarrow \ell \nu \tilde{\chi}_1^0$

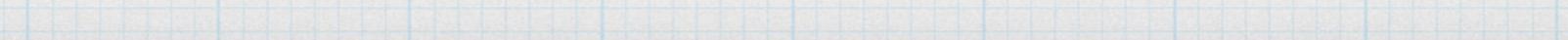
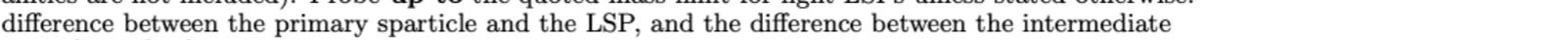
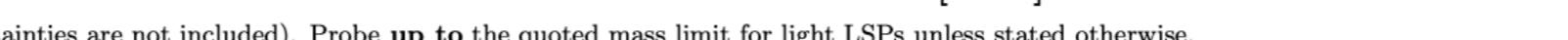
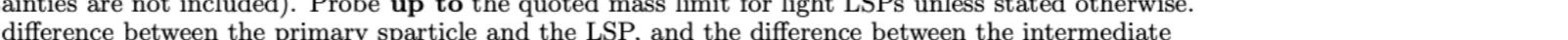
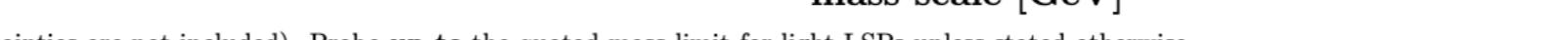
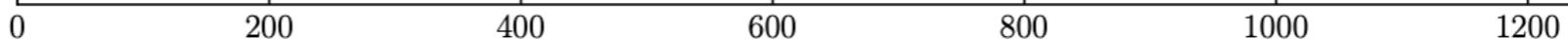
2 ℓ opposite-sign: arXiv:1807.07799

$BF(\tilde{\ell} \nu) = 50\%, x = 0.5$

$pp \rightarrow \tilde{\ell} \tilde{\ell}$

$pp \rightarrow \tilde{\ell}_{L/R} \tilde{\ell}_{L/R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_1^0$

e⁺e⁻, $\mu^+\mu^-$: arXiv:2012.08600

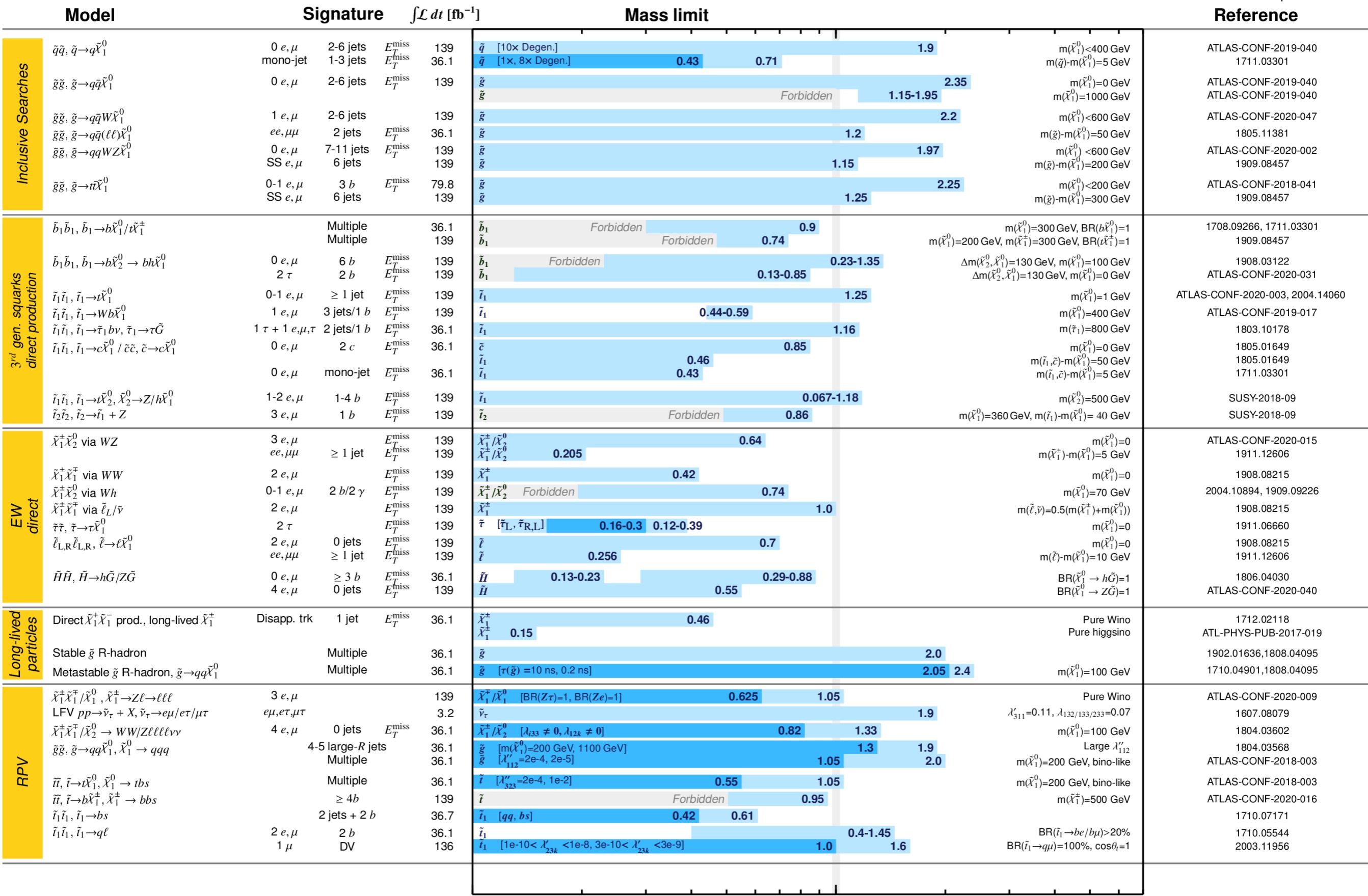


ATLAS SUSY Searches* - 95% CL Lower Limits

July 2020

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}$



*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹ 1 Mass scale [TeV]

Neutrinos

DIRAC OR MAJORANA?

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$$\nu_D = \begin{pmatrix} \nu_L \\ \nu_R \end{pmatrix} \quad \nu_{M_1} = \begin{pmatrix} \xi_1 \\ \xi_1^* \end{pmatrix}, \quad \nu_{M_2} = \begin{pmatrix} \xi_2 \\ \xi_2^* \end{pmatrix}$$

Mass matrix

$$\mathcal{M} = \begin{pmatrix} L & R \\ 0 & m_D \\ m_D^* & M \end{pmatrix} \begin{pmatrix} L \\ R \end{pmatrix}$$

Majorana term



Mass eigenvalues

$$m_1 \approx \frac{m_D^* m_D}{M} \quad m_2 \approx M$$

Light

Heavy

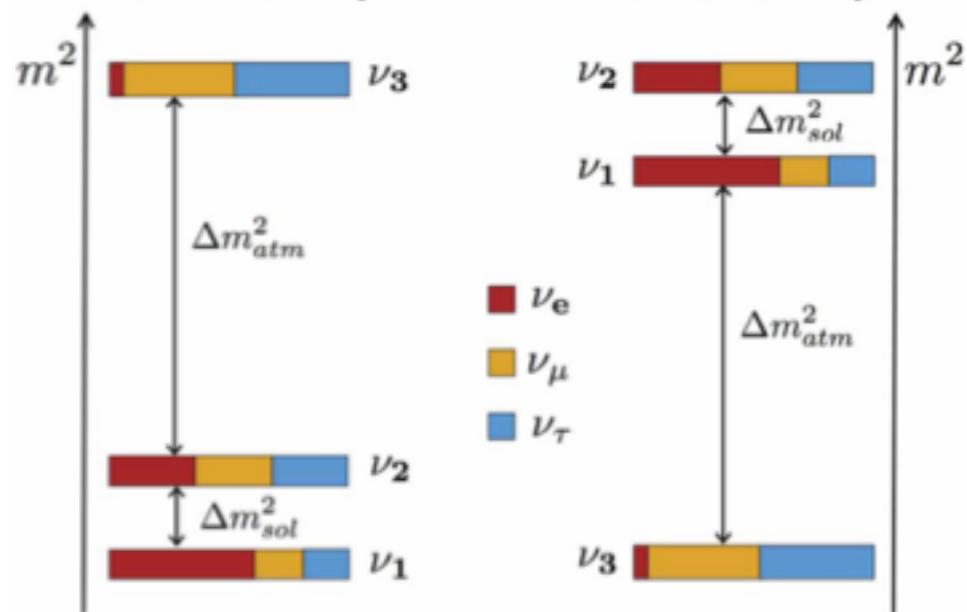


$$\nu_D \neq \nu_D^* \\ m_{\nu_L} = m_{\nu_R}$$



$$\nu_M = \nu_M^* \\ m_{\nu_{M_1}} \neq m_{\nu_{M_2}}$$

Neutrino Physics



- ➊ Absolute value of neutrino masses ?
- ➋ Mass hierarchy?
- ➌ Dirac or Majorana?
- ➍ Fourth sterile neutrino?
- ➎ Neutrino dark matter?

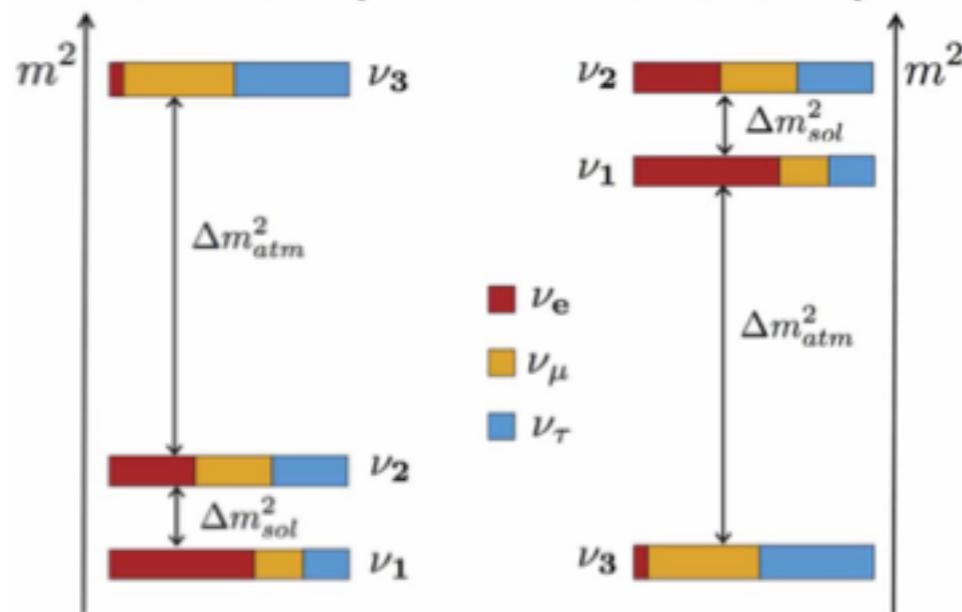
$$0.06 \text{ eV} < \sum m_\nu < 0.12 \text{ eV}$$

-OSC CMB

PMNS-matrix parameters are measured with high accuracy of few %

parameter	best fit $\pm 1\sigma$	3σ range
Δm_{21}^2 [10 $^{-5}$ eV 2]	$7.55^{+0.20}_{-0.16}$	7.05–8.14
$ \Delta m_{31}^2 $ [10 $^{-3}$ eV 2] (NO)	2.50 ± 0.03	2.41–2.60
$ \Delta m_{31}^2 $ [10 $^{-3}$ eV 2] (IO)	$2.42^{+0.03}_{-0.04}$	2.31–2.51
$\sin^2 \theta_{12}/10^{-1}$	$3.20^{+0.20}_{-0.16}$	2.73–3.79
$\sin^2 \theta_{23}/10^{-1}$ (NO)	$5.47^{+0.20}_{-0.30}$	4.45–5.99
$\sin^2 \theta_{23}/10^{-1}$ (IO)	$5.51^{+0.18}_{-0.30}$	4.53–5.98
$\sin^2 \theta_{13}/10^{-2}$ (NO)	$2.160^{+0.083}_{-0.069}$	1.96–2.41
$\sin^2 \theta_{13}/10^{-2}$ (IO)	$2.220^{+0.074}_{-0.076}$	1.99–2.44
δ/π (NO)	$1.32^{+0.21}_{-0.15}$	0.87–1.94
δ/π (IO)	$1.56^{+0.13}_{-0.15}$	1.12–1.94

Neutrino Physics



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$$0.06 \text{ eV} < \sum m_\nu < 0.12 \text{ eV}$$

$\nu\text{-OSC}$ CMB

PMNS-matrix parameters are measured with high accuracy of few %

- ➊ Normal hierarchy favoured at 3.1 σ
- ➋ Nonzero CP phase favoured
- ➌ Upper octant favoured

parameter	best fit $\pm 1\sigma$	3σ range
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BEYOND THE STANDARD MODEL: THE MASS SPECTRUM AND MIXINGS

- Mass spectrum?

$$m_{quark} = y_{quark} \cdot v$$

$$m_{lepton} = y_{lepton} \cdot v$$

$$m_W = g/\sqrt{2} \cdot v$$

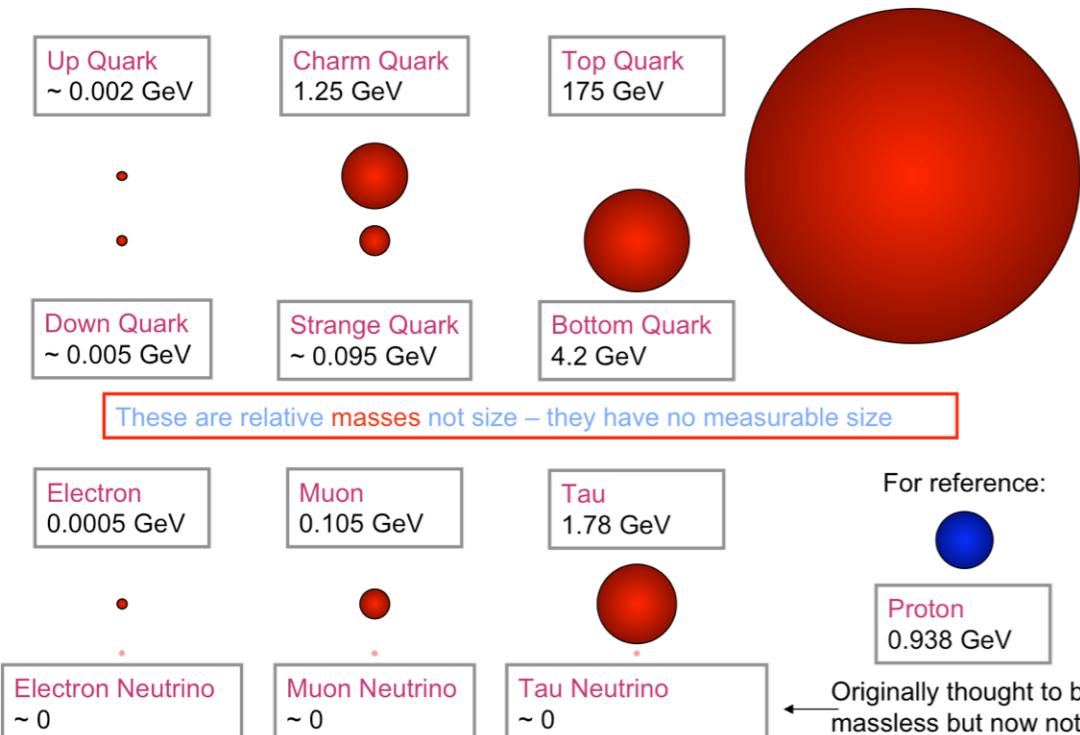
$$m_Z = \sqrt{g^2 + g'^2}/\sqrt{2} \cdot v$$

$$m_H = \sqrt{\lambda} \cdot v$$

SM

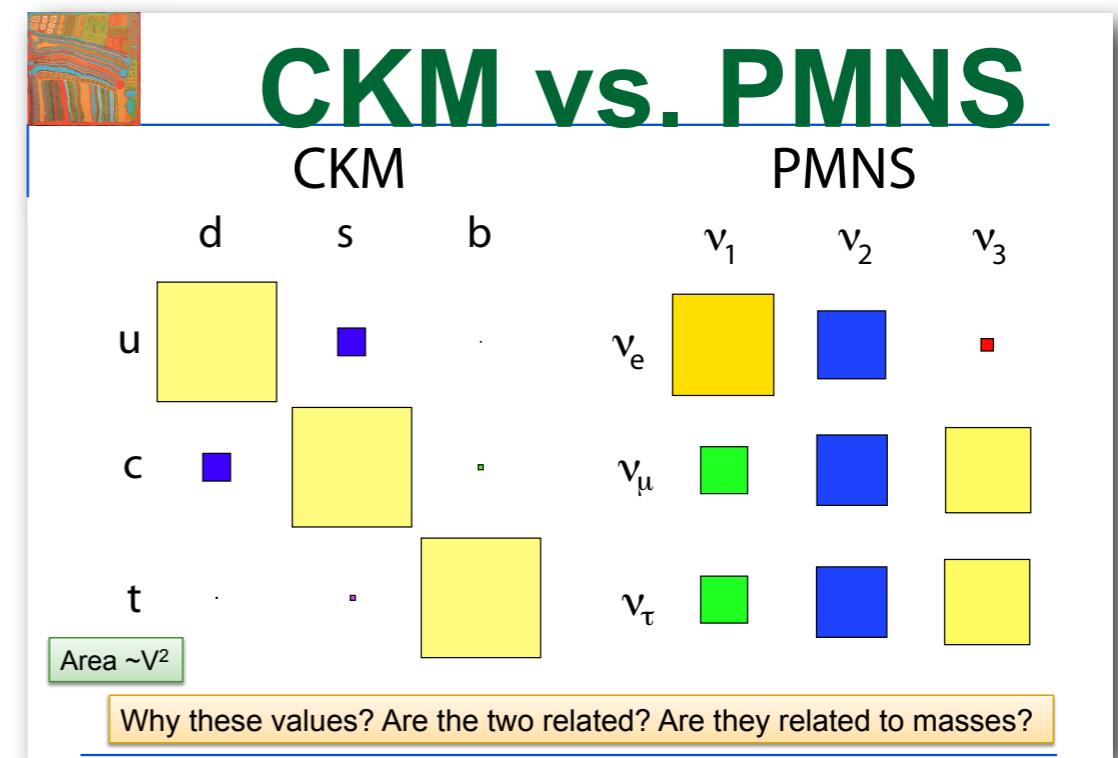
$$m_\gamma = 0$$

$$m_{gluon} = 0$$



- Mixing Matrices?

- Quark-Lepton Symmetry
- Strong difference in parameters



BEYOND THE STANDARD MODEL: THE MASS SPECTRUM AND MIXINGS

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$$m_{quark} = y_{quark} \cdot v$$

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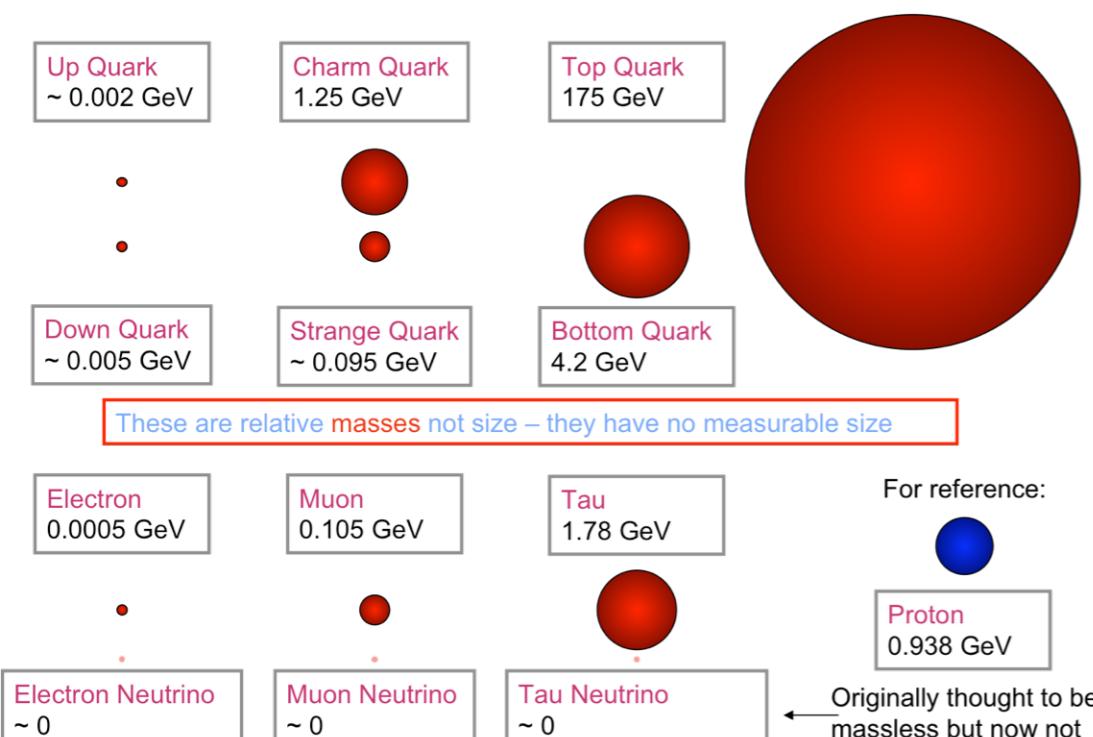
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SM

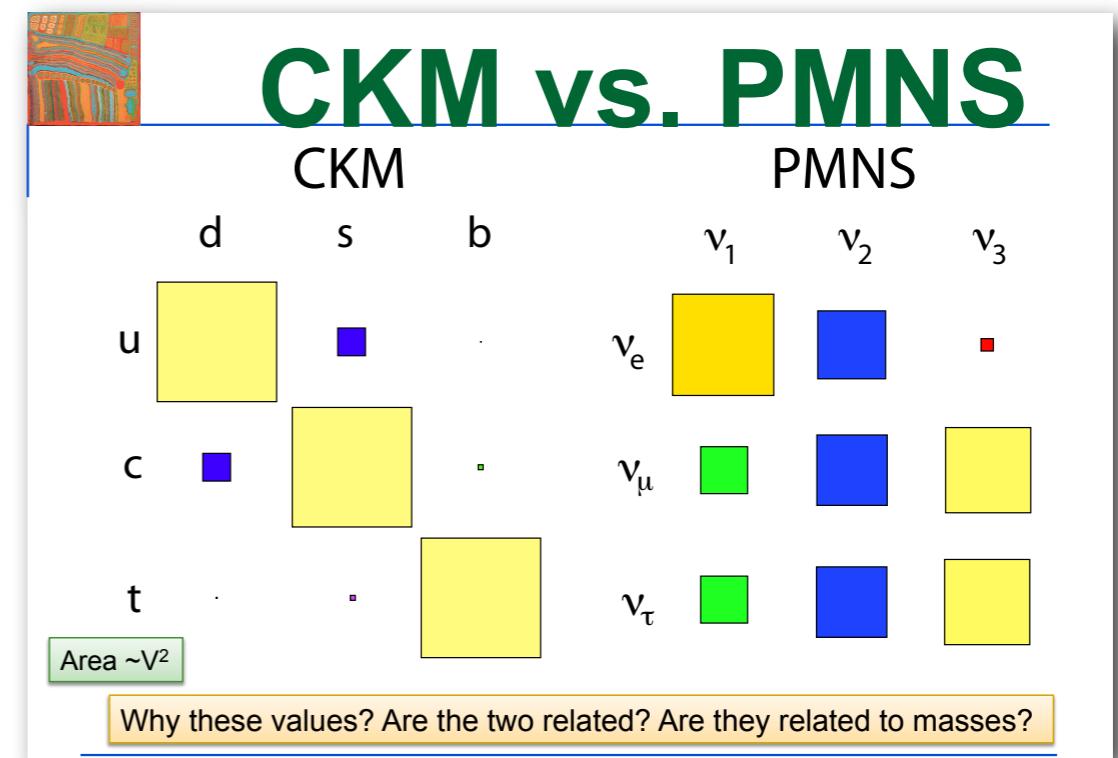
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- Mixing Matrices?

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- Strong difference in parameters



- What are the CKM and PMNS phases?
- Where lies the source of CP violation: in quark or lepton sector?

BEYOND THE STANDARD MODEL: THE MASS SPECTRUM AND MIXINGS

- Mass spectrum?

$$m_{quark} = y_{quark} \cdot v$$

$$m_{lepton} = y_{lepton} \cdot v$$

$$m_W = g/\sqrt{2} \cdot v$$

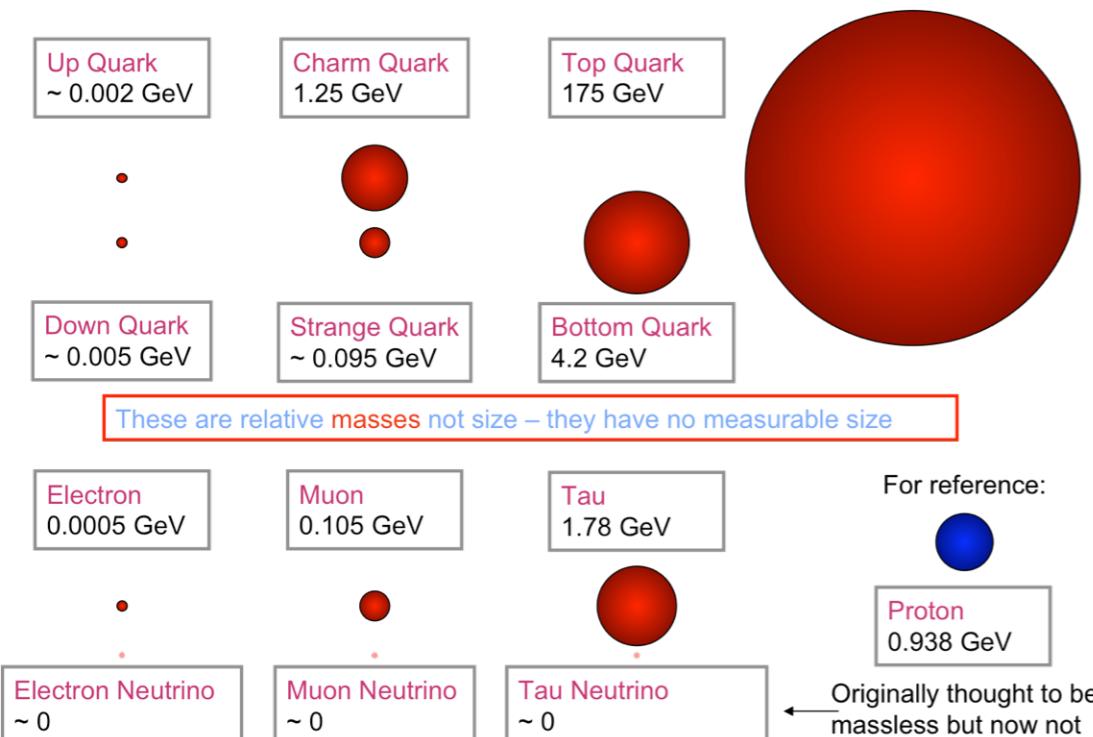
$$m_Z = \sqrt{g^2 + g'^2}/\sqrt{2} \cdot v$$

$$m_H = \sqrt{\lambda} \cdot v$$

SM

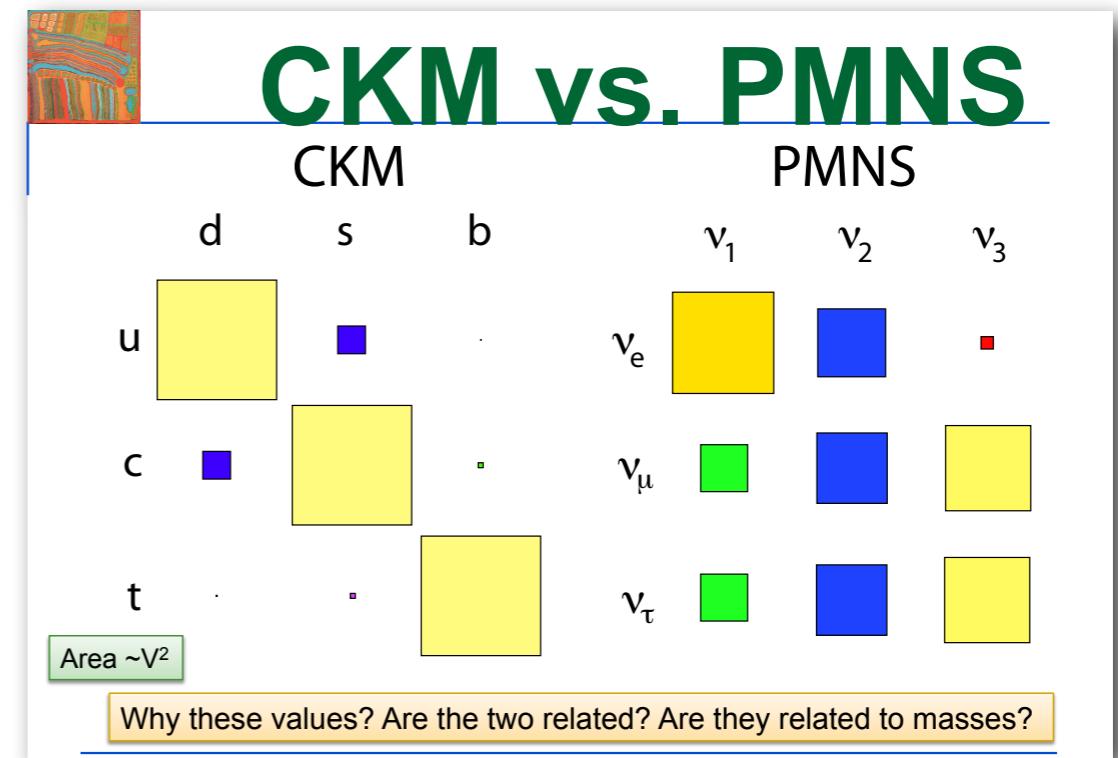
$$m_\gamma = 0$$

$$m_{gluon} = 0$$



- Mixing Matrices?

- Quark-Lepton Symmetry
- Strong difference in parameters

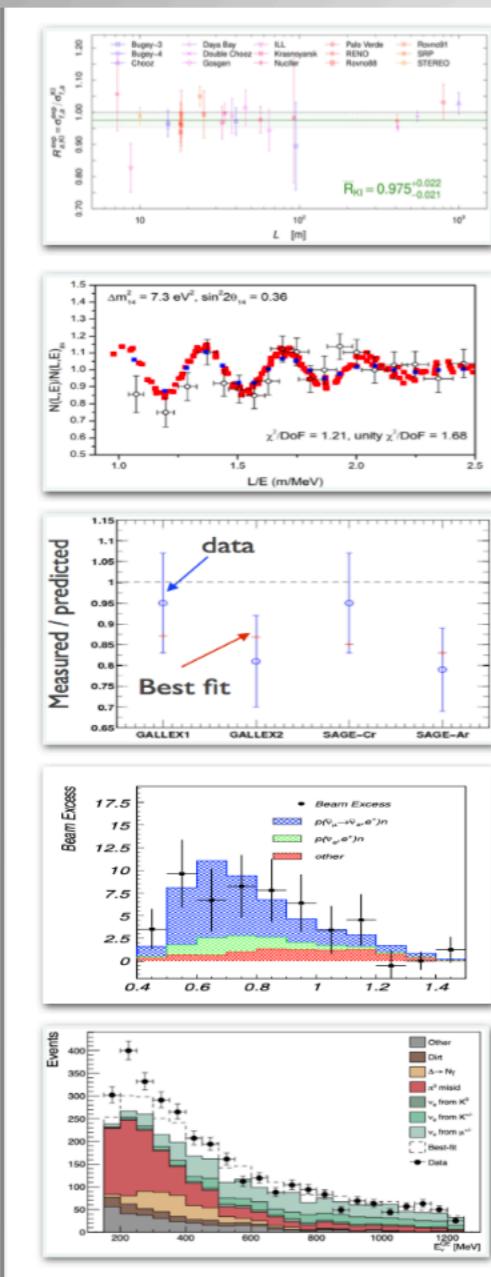


- What are the CKM and PMNS phases?
- Where lies the source of CP violation: in quark or lepton sector?

$$J_{CP} = \frac{1}{8} \sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \sin \delta$$

Neutrino Anomalies

Albert De Roeck



reactor flux anomaly
resolved with new input data
to flux calculation



reactor spectra
is there really an anomaly? -> DANSS



gallium anomaly
unresolved, recently reinforced BEST



LSND
unresolved



MiniBooNE
unresolved μ BooNe excluded some explanations
resolvable by next-gen. SBL experiments

More details
in the
backup

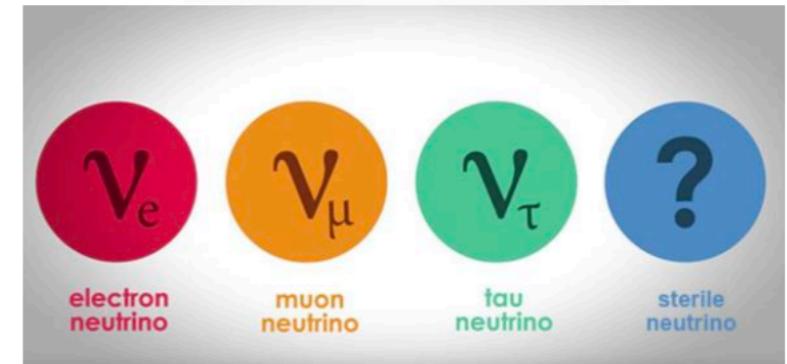
- Jury still out on many of these anomalies. No clear picture emerging yet.
- Simple sterile neutrino would not fit all the data. Tensions on all sides...
- Future: Reactor experiments continuing or new ones (eg JSNS²) or new experiments at the FNAL short neutrino baseline... (ICARUS, SBND)

Sterile Neutrinos

Several anomalies around in the community since some years...
Additional sterile neutrinos as a possible candidate explanation

Very generic extension of SM

- can be leftover of extended gauge multiplet



Useful phenomenological tool

- can explain ν masses (seesaw mechanism, $m \sim \text{TeV} \dots M_{\text{Pl}}$)
- can explain cosmic baryon asymmetry (leptogenesis, $m \gg 100 \text{ GeV}$)
- can explain dark matter ($m \sim \text{keV}$)
- can explain oscillation anomalies ($m \sim \text{eV}$)

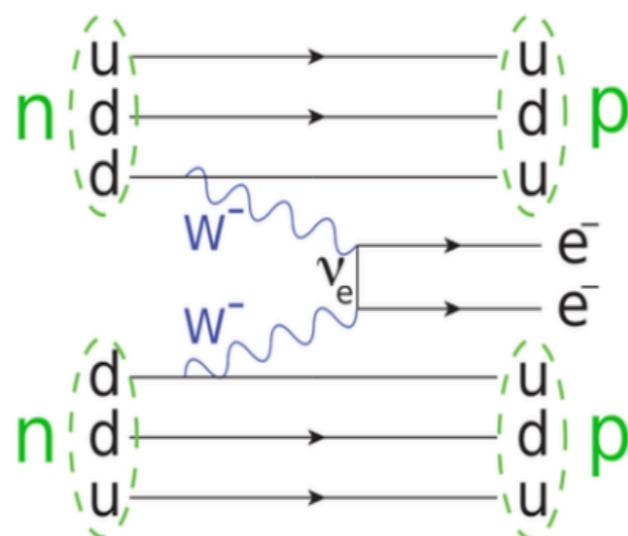
Promote mixing matrix to 4×4 , oscillation formula unchanged:

$$P_{\alpha \rightarrow \beta} = \sum_{j,k} U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* \exp [-i(E_j - E_k)T]$$

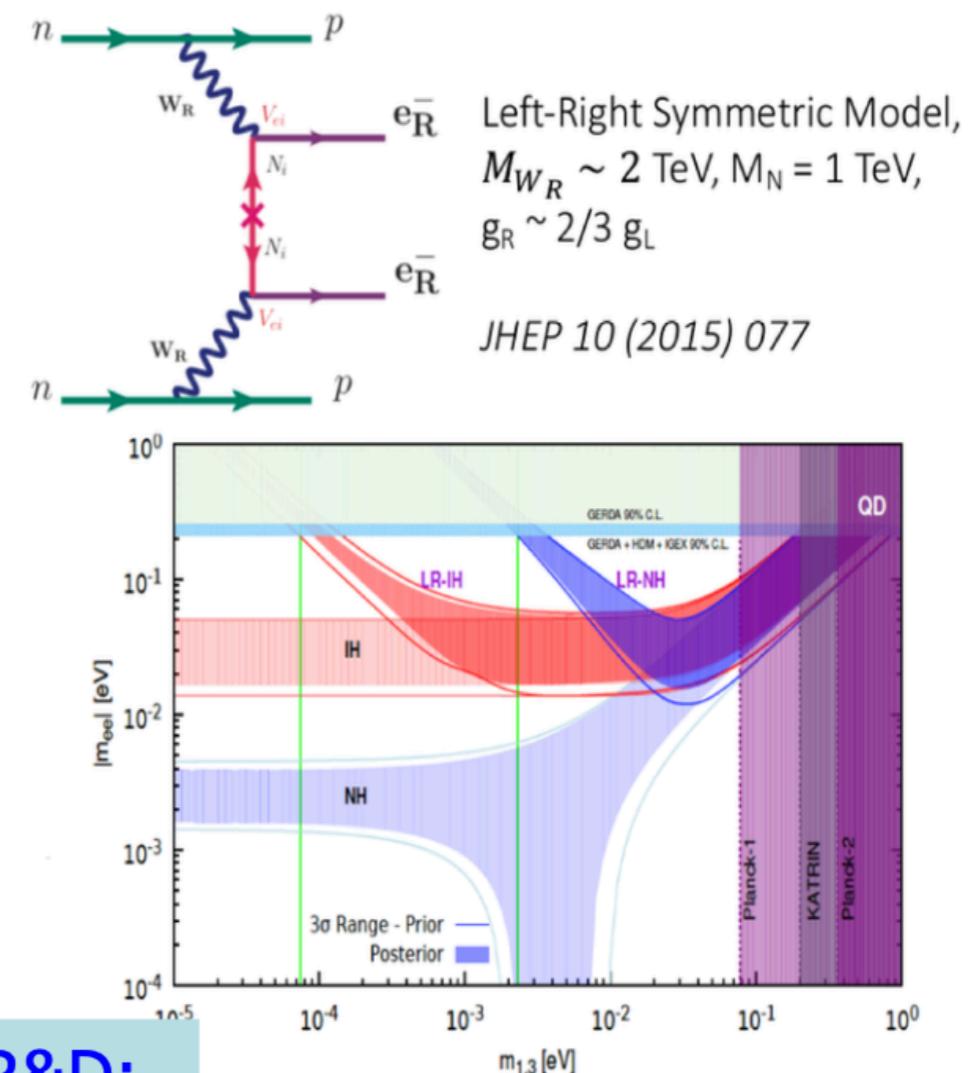
Neutrinoless Double Beta Decay

The question is still unanswered:

Are neutrinos their own antiparticles?



Ton scale OnuBB experiments will cover the inverted hierarchy by 2035



Many experiments operating, planned or in R&D:
LEGEND, SNO+, NEXT, CUPID, THEIA...

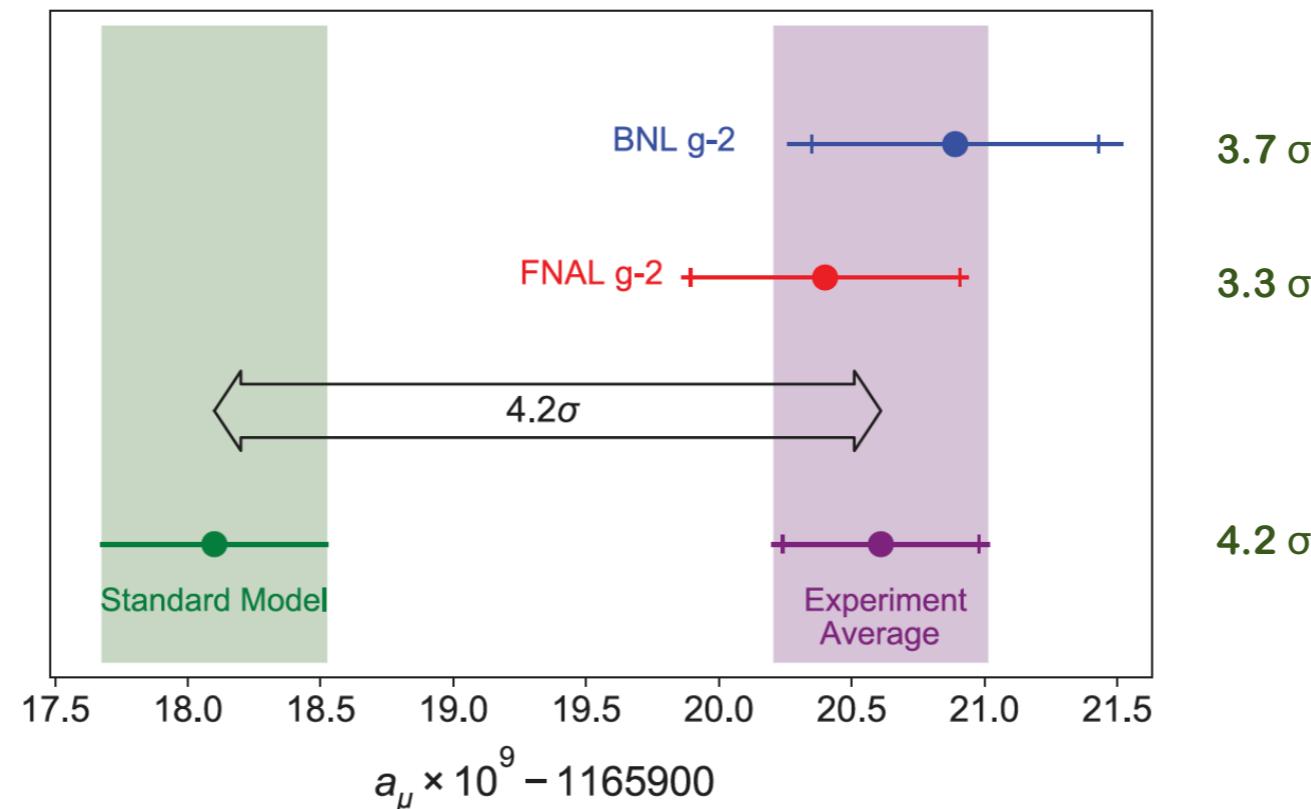
Summary

- In ~10 years from now, oscillation will mostly be understood: mass hierarchy and CP phase will be known
- Neutrino absolute masses may be measured in ~20 years, through cosmology, beta decays and double beta decays
- Majorana neutrino nature maybe determined in ~30 years
- Sterile neutrinos are unlikely the cause of reactor and Ga anomalies, but still possible for the LSND anomaly
- A new era of astrophysics with multi-messengers, including neutrinos
- Many new projects will start within 10 years
- A bright future

g-2

Experimental status

- April 7th 2021: Muon $g - 2$ experiment at FNAL confirms BNL!



$$a_\mu^{\text{EXP}} = (116592089 \pm 63) \times 10^{-11} [0.54\text{ppm}] \quad \text{BNL E821}$$

$$a_\mu^{\text{EXP}} = (116592040 \pm 54) \times 10^{-11} [0.46\text{ppm}] \quad \text{FNAL E989 Run 1}$$

$$a_\mu^{\text{EXP}} = (116592061 \pm 41) \times 10^{-11} [0.35\text{ppm}] \quad \text{WA}$$

- FNAL aims at 16×10^{-11} . First 4 runs completed, 5th in progress.
- Muon $g - 2$ proposal at J-PARC: Phase-1 with similar BNL precision.

“Old muon g-2 puzzle”

New Physics for the muon $g - 2$: at which scale?

- Δa_μ discrepancy at $\sim 4.2\sigma$ level:

$$\Delta a_\mu = a_\mu^{\text{EXP}} - a_\mu^{\text{SM}} \equiv a_\mu^{\text{NP}} = (2.51 \pm 0.59) \times 10^{-9}$$

$$\Delta a_\mu \equiv a_\mu^{\text{NP}} \approx (a_\mu^{\text{SM}})_{\text{weak}} \approx \frac{m_\mu^2}{16\pi^2 v^2} \approx 2 \times 10^{-9}$$

- ▶ NP is at the weak scale ($\Lambda \approx v$) and weakly coupled to SM particles.*
- ▶ NP is very light ($\Lambda \lesssim 1$ GeV) and feebly coupled to SM particles.
- ▶ NP is very heavy ($\Lambda \gg v$) and strongly coupled to SM particles.

*Favoured by the *hierarchy problem* and by a WIMP DM candidate but disfavoured by the LEP and LHC bounds (supersymmetry being the most prominent example).

[For a through compilation of models, see Athron, Balazs, Jacob, Kotlarski, Stockinger, Stockinger-Kim, '21.]

$\Lambda \approx v$: SUSY and the muon ($g - 2$)

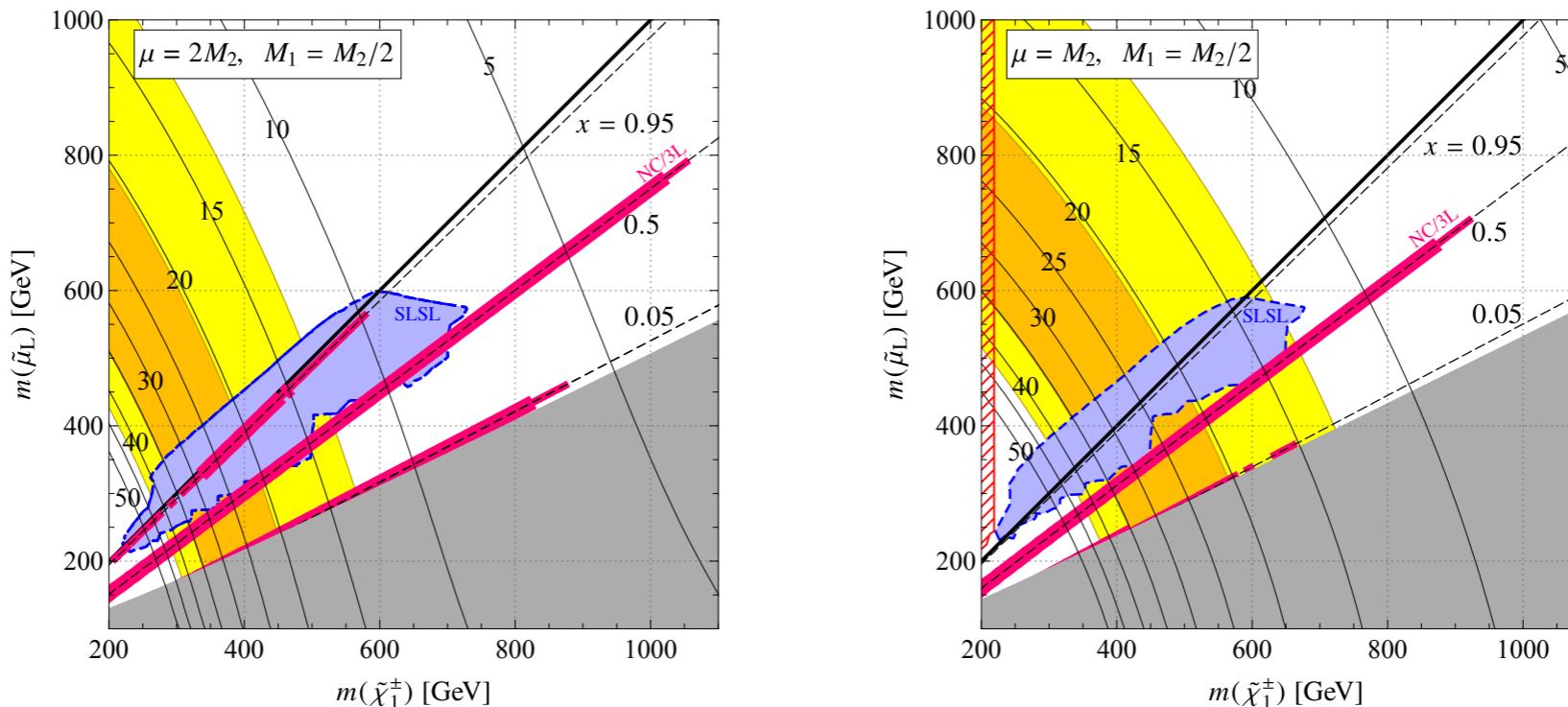
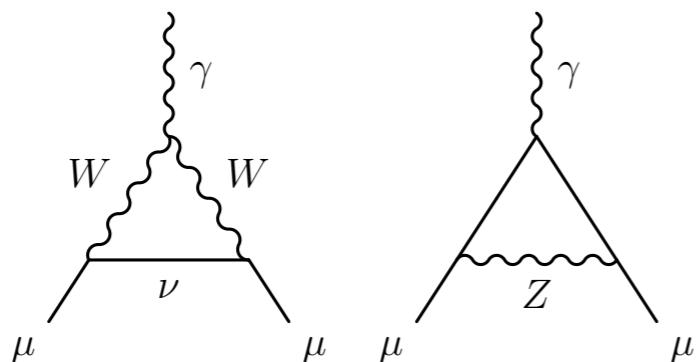
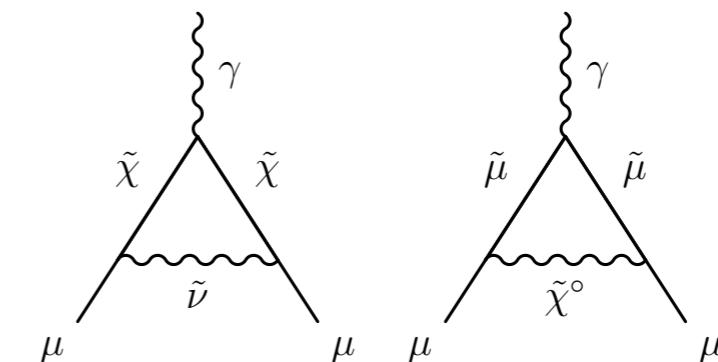


Figure: LHC Run 2 bounds on SUSY scenario for the muon $g - 2$ anomaly for $\tan \beta = 40$. Orange (yellow) regions satisfy the muon $g - 2$ anomaly at the 1σ (2σ) level [Endo et al., '20].



$$(a_\mu^{\text{SM}})_{\text{weak}} \approx \frac{g^2 m_\mu^2}{32\pi^2 M_W^2} \approx 2 \times 10^{-9}$$



$$a_\mu^{\text{SUSY}} \approx \frac{g^2 m_\mu^2 \tan \beta}{32\pi^2 \tilde{m}^2} \approx 2 \times 10^{-9}$$

$\tilde{m} = 500 \text{ GeV} \& \tan \beta = 40$

$\Lambda \lesssim 1 \text{ GeV}$: Axion-like Particles and the muon ($g - 2$)

Axion-like Particle effective Lagrangian

$$\mathcal{L} = e^2 C_{\gamma\gamma} \frac{a}{\Lambda} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{c_{\mu\mu}}{2} \frac{\partial^\nu a}{\Lambda} \bar{\mu} \gamma_\nu \gamma_5 \mu$$

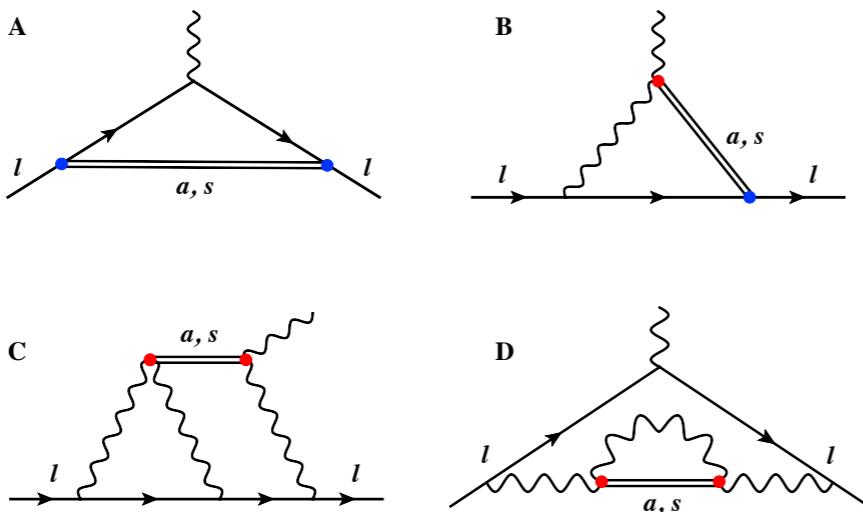


Figure: Contributions of a scalar ‘s’ and a pseudoscalar ‘a’ ALP to the $(g - 2)_\ell$.

[Marciano, Masiero, PP, Passera ’16]

[Cornella, P.P., Sumensari ’19]

$$\Delta a_\mu = \frac{m_\mu^2}{\Lambda^2} \left[\frac{12\alpha^3}{\pi} C_{\gamma\gamma}^2 \ln^2 \frac{\Lambda^2}{m_\mu^2} - \frac{(c_{\mu\mu})^2}{16\pi^2} h_1 \left(\frac{m_a^2}{m_\mu^2} \right) - \frac{2\alpha}{\pi} c_{\mu\mu} C_{\gamma\gamma} \ln \frac{\Lambda^2}{m_\mu^2} \right]$$

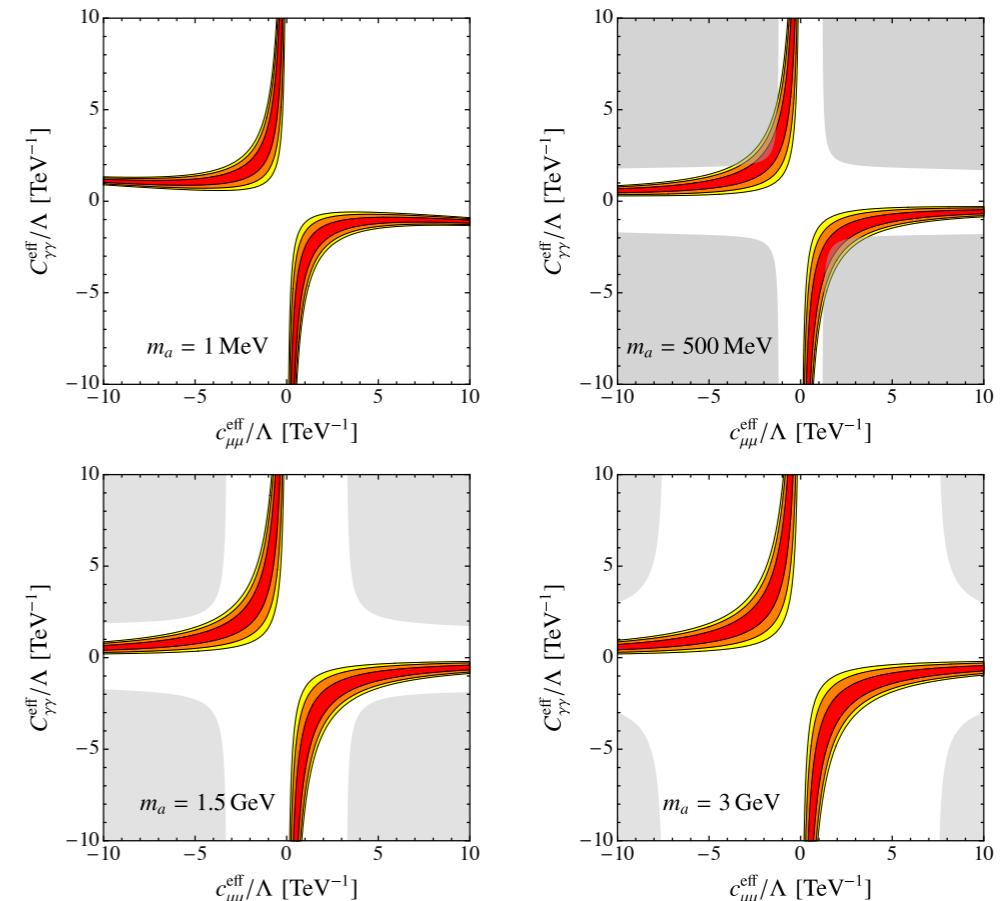


Figure: Δa_μ regions favoured at 68% (red), 95% (orange) and 99% (yellow) CL. Gray regions are excluded by the BaBar search $e^+e^- \rightarrow \mu^+\mu^- + \mu^+\mu^-$ [Bauer, Neubert, Thamm, ’17]

Breakdown of SM contributions

- a_μ from WP20 (w/o BMWc lattice result)

[Colangelo EPS-HEP2021 proceeding]

Contribution	Value $\times 10^{11}$	References
Experiment (E821)	116 592 089(63)	Ref. [3]
Experiment (FNAL)	116 592 040(54)	Ref. [1]
Experiment (World-Average)	116 592 061(41)	
HVP LO (e^+e^-)	6931(40)	Refs. [6–11]
HVP NLO (e^+e^-)	-98.3(7)	Ref. [11]
HVP NNLO (e^+e^-)	12.4(1)	Ref. [12]
HVP LO (lattice, $udsc$)	7116(184)	Refs. [13–21]
HLbL (phenomenology)	92(19)	Refs. [22–34]
HLbL NLO (phenomenology)	2(1)	Ref. [35]
HLbL (lattice, uds)	79(35)	Ref. [36]
HLbL (phenomenology + lattice)	90(17)	
QED	116 584 718.931(104)	Refs. [37, 38]
Electroweak	153.6(1.0)	Refs. [39, 40]
HVP (e^+e^- , LO + NLO + NNLO)	6845(40)	
HLbL (phenomenology + lattice + NLO)	92(18)	
Total SM Value	116 591 810(43)	
Difference: $\Delta a_\mu := a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$	251(59)	



HVP LO is the bottle-neck of the SM prediction

HLO contribution from $e^+e^- \rightarrow \text{hadrons}$

- dominated by $e^+e^- \rightarrow \pi^+\pi^-$ channel (70% of the full hadronic)

$$(a_\mu^{\text{HVP}})_{e^+e^-} = \frac{\alpha}{\pi^2} \int_{m_{\pi^0}^2}^{\infty} \frac{ds}{s} K(s) \text{Im } \Pi_{\text{had}}(s) = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \sigma_{\text{had}}(s)$$

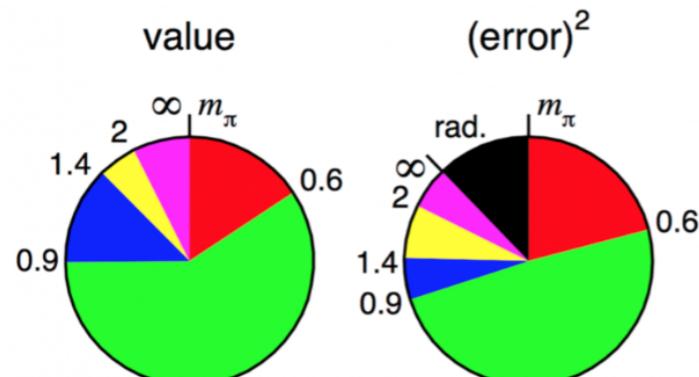
dispersion relations

optical theorem

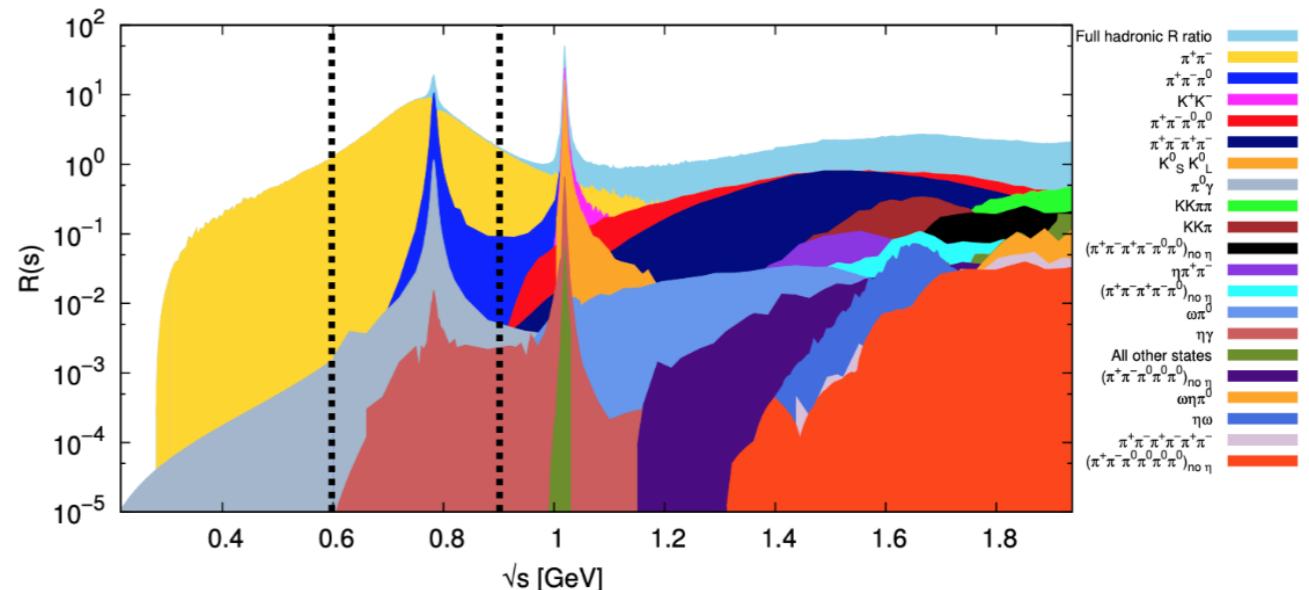
kernel function

$$K(s) \approx m_\mu^2/3s \quad \text{for} \quad \sqrt{s} \gg m_\mu$$

$$\text{Im } \langle \gamma \gamma \rangle \sim \left| \langle \gamma \gamma \rangle \right|^2 \sim \sigma(e^+e^- \rightarrow \gamma^* \rightarrow \text{hadrons})$$



Keshavarzi, Nomura, Teubner 2018



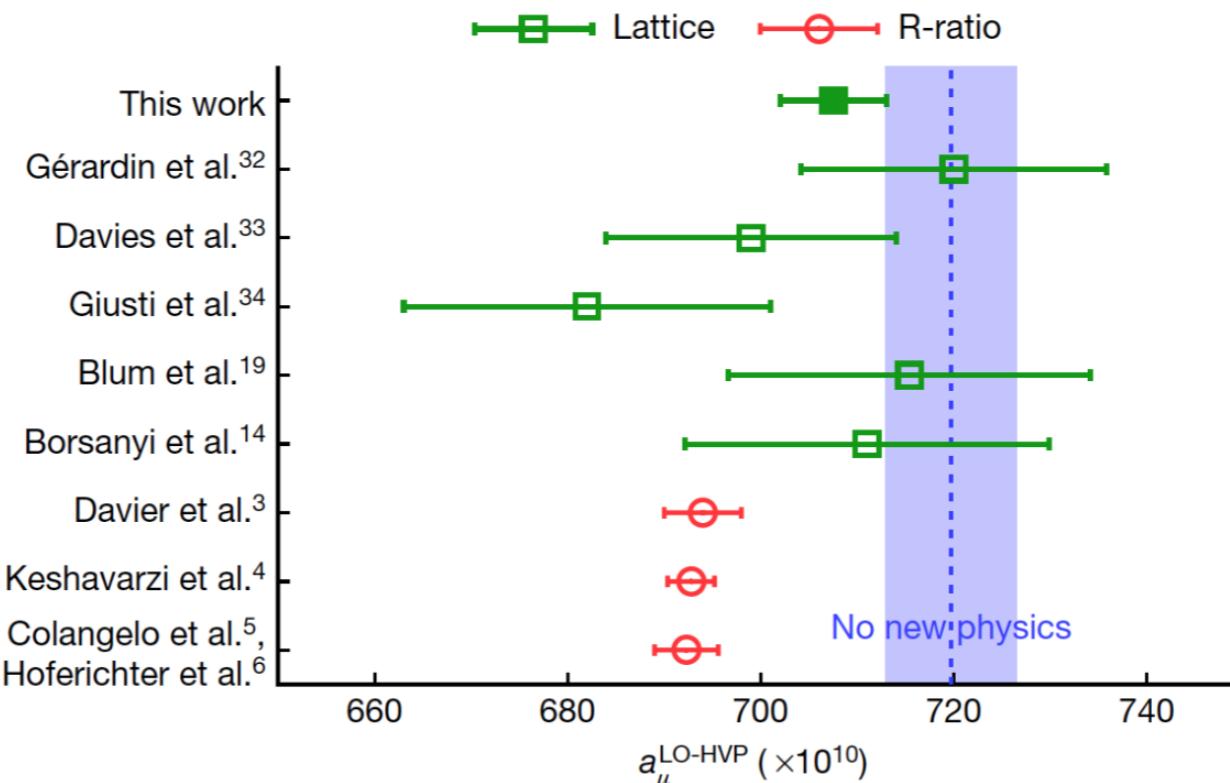
$$a_{\mu,e^+e^-}^{\text{HLO}} = 6931(40) \times 10^{-11} (0.6\%) \quad [\text{WP20}]$$

HLO contribution from lattice QCD

- Great progress also in lattice QCD, where spacetime is modeled as a discrete grid of points. The BMW collaboration reached a 0.8% precision!

$$a_\mu^{\text{HLO}} = 7075(23)_{\text{stat}}(50)_{\text{syst}} [55]_{\text{tot}} \times 10^{-11}$$

- 2–2.5 σ tension with the “data-driven” evaluations.



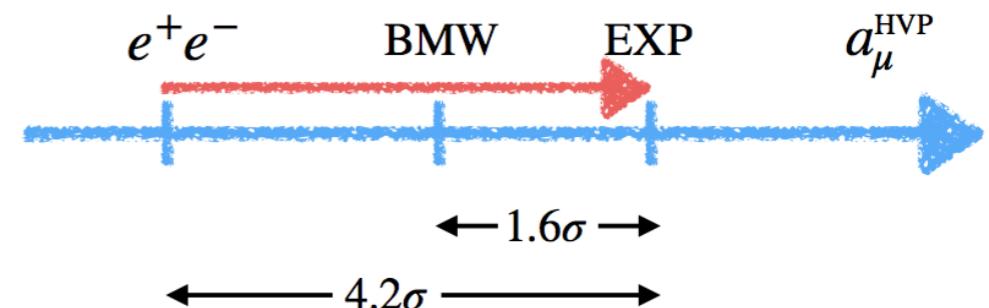
Borsanyi et al (BMWc), Nature 2021

“New muon g-2 puzzle”

$$(a_\mu^{\text{HVP}})_{\text{EXP}} = a_\mu^{\text{EXP}} - a_\mu^{\text{SM, rest}}$$

$$(a_\mu^{\text{HVP}})_{e^+e^-}^{\text{WP20}} = 6931(40) \times 10^{-11}$$

$$(a_\mu^{\text{HVP}})_{\text{BMW}} = 7075(55) \times 10^{-11}$$



“new puzzle”: if BMW is correct, the “old” g-2 discrepancy (4.2σ) would be basically gone

→ however, this brings in a new tension with e^+e^- data (2.2σ)

Here, NP in $\sigma_{\text{had}}(e^+e^- \rightarrow \text{hadrons})$ such that

[LDL, Masiero, Paradisi, Passera 2112.08312]

1. $(a_\mu^{\text{HVP}})_{e^+e^-}^{\text{WP20}} \approx (a_\mu^{\text{HVP}})_{\text{EXP}}$
2. the approximate agreement between BMW and EXP is not spoiled
3. w/o a direct contribution a_μ^{NP} (i.e. NP not in muons)

Consequences of the BMW result

- **Can Δa_μ be due to missing contributions in $\sigma(e^+e^- \rightarrow \text{had})$?**
 - ▶ An upward shift of $\sigma(s)$ also induces an increase of $\Delta\alpha_{\text{had}}^{(5)}(M_Z)$ defined by:

$$\alpha(M_Z) = \frac{\alpha}{1 - \Delta\alpha(M_Z) - \Delta\alpha_{\text{had}}^{(5)}(M_Z) - \Delta\alpha_{\text{top}}(M_Z)}$$

$$a_\mu^{\text{HLO}} \simeq \frac{m_\mu^2}{12\pi^3} \int_{4m_\pi^2}^\infty ds \frac{\sigma(s)}{s}, \quad \Delta\alpha_{\text{had}}^{(5)} = \frac{M_Z^2}{4\pi\alpha^2} \int_{4m_\pi^2}^\infty ds \frac{\sigma(s)}{M_Z^2 - s}$$

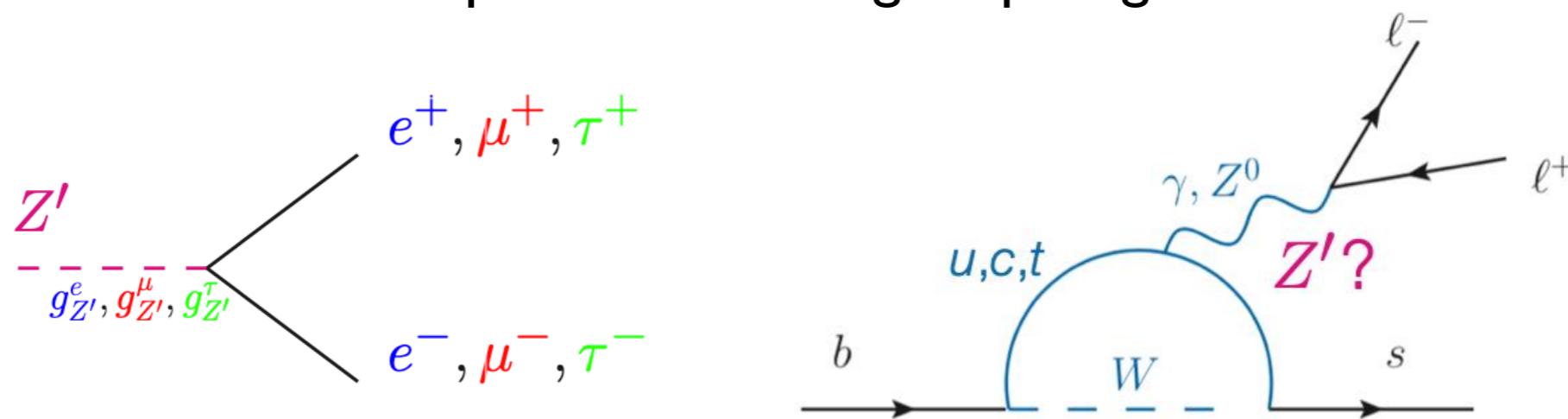
$$\text{Im } \text{---} \sim \left| \text{---} \right|^2 \sim \sigma(e^+e^- \rightarrow \gamma^* \rightarrow \text{hadrons})$$

- **A change in $\sigma(e^+e^- \rightarrow \text{had})$ is strongly disfavoured by:**
 - ▶ **EW-fit for $\sqrt{s} \gtrsim 1 \text{ GeV}$** [Marciano, Passera, Sirlin, '08, Keshavarzi, Marciano, Passera, Sirlin, '20, Crivellin, Hoferichter, Manzari, Montull, '20]. A shift of $\sigma(e^+e^- \rightarrow \text{had})$ to accomodate the Δa_μ anomaly would necessarily require new physics to show up in the EW-fit!
- **A check of the BMW results by other lattice QCD (LQCD) coll. is worth.**
- **LQCD coll. should provide $\Delta\alpha_{\text{had}}^{\text{LQCD}}$ to be compared with $\Delta\alpha_{\text{had}}^{e^+e^-}$.**

Flavor Anomaly

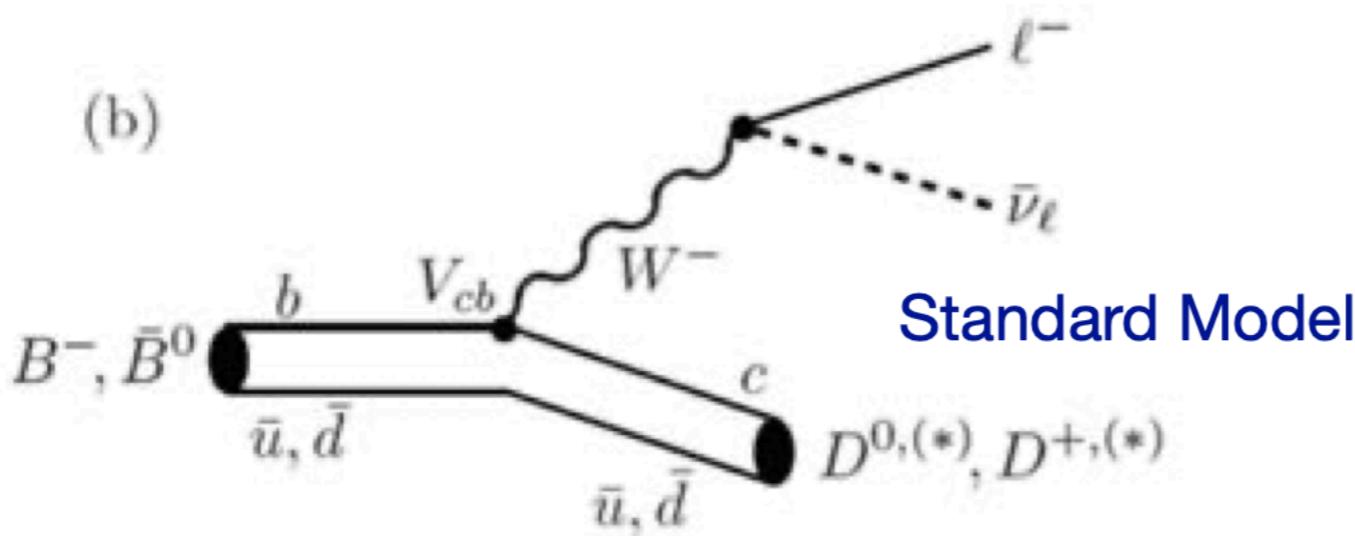
Lepton Flavour Universality (LFU)

- LFU is a cornerstone of the SM : charged leptons (e, μ, τ) couple in a universal way to the SM gauge bosons
- If NP couples in a non-universal way to the three lepton families, then we might see differences in rates of rare decays involving different lepton pairs (e.g. e/μ or μ/τ)
- Hence - LFU is tested in $b \rightarrow s \ell^+ \ell^-$ transitions. These are FCNC's with amplitudes involving loop diagrams



LEPTON (NON) UNIVERSALITY (?!)

Charged currents at tree level

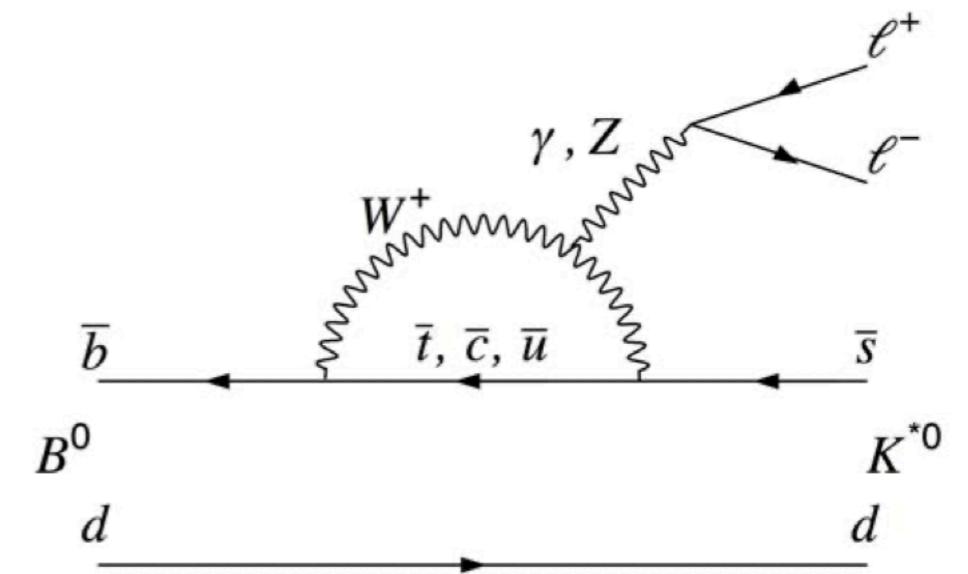


$$R_{D^{(*)}} = \frac{BR(B \rightarrow D^{(*)}\tau\nu_\tau)}{BR(B \rightarrow D^{(*)}\mu\nu_\mu)} \quad 3.8\sigma$$

$$\mathcal{R}(D)_{SM} = 0.299 \pm 0.003$$

$$\mathcal{R}(D^*)_{SM} = 0.258 \pm 0.005$$

Neutral currents at one-loop level



$$R_{K^{(*)}} = \left. \frac{BR(B \rightarrow K^{(*)}\mu\mu)}{BR(B \rightarrow K^{(*)}ee)} \right| \quad 2.5\sigma$$

$$\mathcal{R}(K^*)_{SM} = 1.0$$

Several R-ratio measurements

- Compare the rates of $B \rightarrow X_s e^+ e^-$ and $B \rightarrow X_s \mu^+ \mu^-$

[where B is B^+ , B^0 , B_s^0 , Λ_b^0 and X_s is K^+ , K^{*0} , ϕ , $pK \dots$]

- This allows precise testing of lepton flavour universality
- We can construct the ratio :

$$R_X = \frac{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\mathcal{B}(B_q \rightarrow X_s \mu^+ \mu^-)}{dq^2} dq^2}{\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\mathcal{B}(B_q \rightarrow X_s e^+ e^-)}{dq^2} dq^2} = 1 \pm \mathcal{O}(1\%)$$

- Small theoretical uncertainties because hadronic uncertainties cancel
- This ratio is unity in the SM, neglecting lepton masses, with QED corrections at the % level

- Five different ratios published so far by LHCb:
- $X_s = K^+, K_S^0, K^{*0}, K^{*+}$ and pK^-

Several R-ratio measurements

- Compare the rates of $B \rightarrow X_s e^+ e^-$ and $B \rightarrow X_s \mu^+ \mu^-$

[where B is B^+ , B^0 , B_s^0 , Λ_b^0 and X_s is K^+ , K^{*0} , ϕ , $pK \dots$]

- This allows precise testing of lepton flavour universality
- We can construct the ratio :

$$\int_{q_{\min}^2}^{q_{\max}^2} \frac{d\mathcal{B}(B_q \rightarrow X_s \mu^+ \mu^-)}{dq^2} dq^2$$

$$R_X =$$

- Actually measure double ratios which significantly reduce systematic uncertainties:

$$R_X = \frac{\mathcal{B}(B_q \rightarrow X_s \mu^+ \mu^-)}{\mathcal{B}(B_q \rightarrow X_s J/\psi(\mu^+ \mu^-))} \cdot \frac{\mathcal{B}(B_q \rightarrow X_s J/\psi(e^+ e^-))}{\mathcal{B}(B_q \rightarrow X_s e^+ e^-)}$$

- Small theoretical uncertainties because hadronic uncertainties cancel

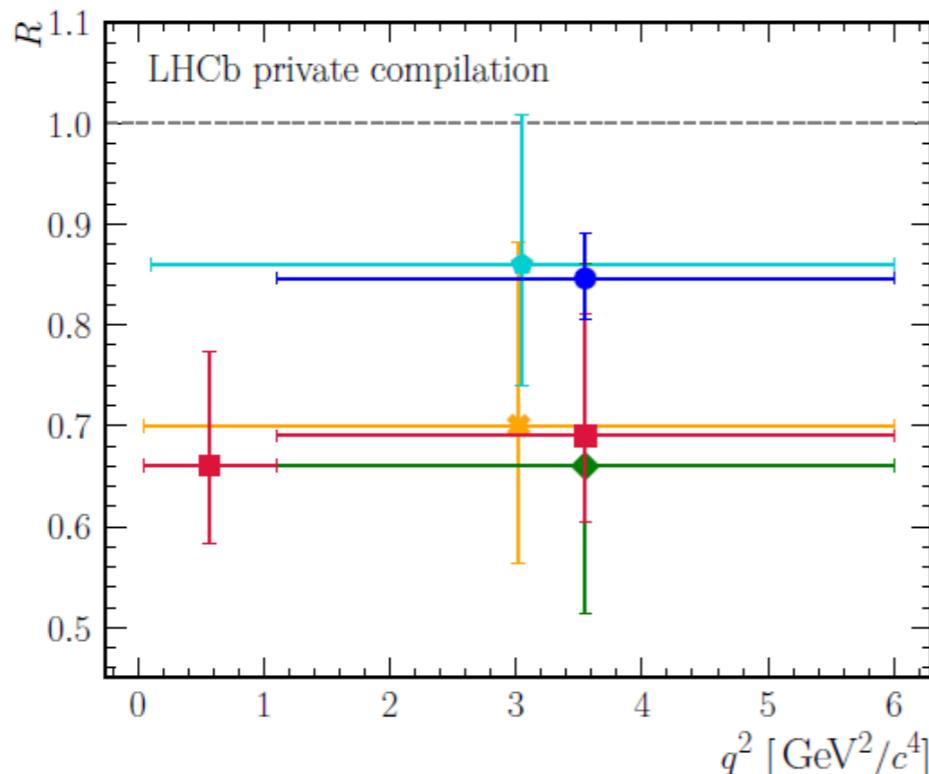
- Five different ratios published so far by LHCb:

$X_s = K^+, K_s^0, K^{*0}, K^{*+}$ and pK^-

LFU results : R_x

■ R_K [Nat. Phys. 18, 277–282 (2022)]
■ $R_{K_S^0}$ [PRL 128, No. 19]
■ $R_{K^{*+}}$ [PRL 128, No. 19]

■ R_{pK} [JHEP 05 (2020) 040]
■ $R_{K^{*0}}$ [JHEP 08 (2017) 055]



R_K (9/fb) 3.1 σ from SM
 $R_{K^{*+}}$ (9/fb) 1.4 σ NEW
 $R_{K_S^0}$ (9/fb) 1.5 σ NEW
 $R_{K^{*0}}$ low- q^2 (3/fb) 2.1 σ
 $R_{K^{*0}}$ central- q^2 (3/fb) 2.4 σ
 R_{pK} (5/fb) <1 σ

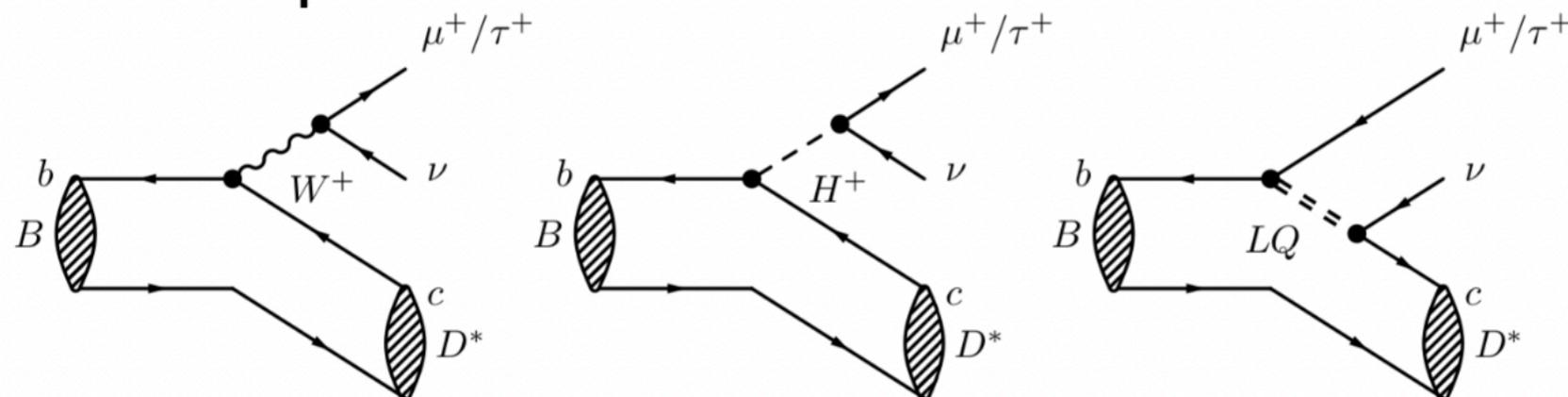
- All measurements have values less than unity
- The puzzle persists → we eagerly await Belle-II & CMS results
- LHCb is now focused on completing a combined analysis of R_K & R_{K^*} with the Run I+2 dataset. This work has led to a deeper understanding of systematics which will be reflected in the final result.

LFU studies in $B^0 \rightarrow D^{(*)-}\tau^+\nu_\tau$ decays

- Different class of decays (tree-level charged current with V_{cb} suppression)
- Not at all rare: $B(B^0 \rightarrow D^{*-}\tau^+\nu_\tau) \sim 1\%$, the problem is the background.
- Lepton-universality ratio $R(D^*)$:

$$R(D^*) = \frac{B(B^0 \rightarrow D^{*-}\tau^+\nu_\tau)}{B(B^0 \rightarrow D^{*-}\mu^+\nu_\mu)}$$

- may be sensitive to any NP model coupling preferentially to third generation leptons



- Ratios predicted theoretically at $\sim 1\%$:

$$R(D)_{SM} = 0.299 \pm 0.003$$

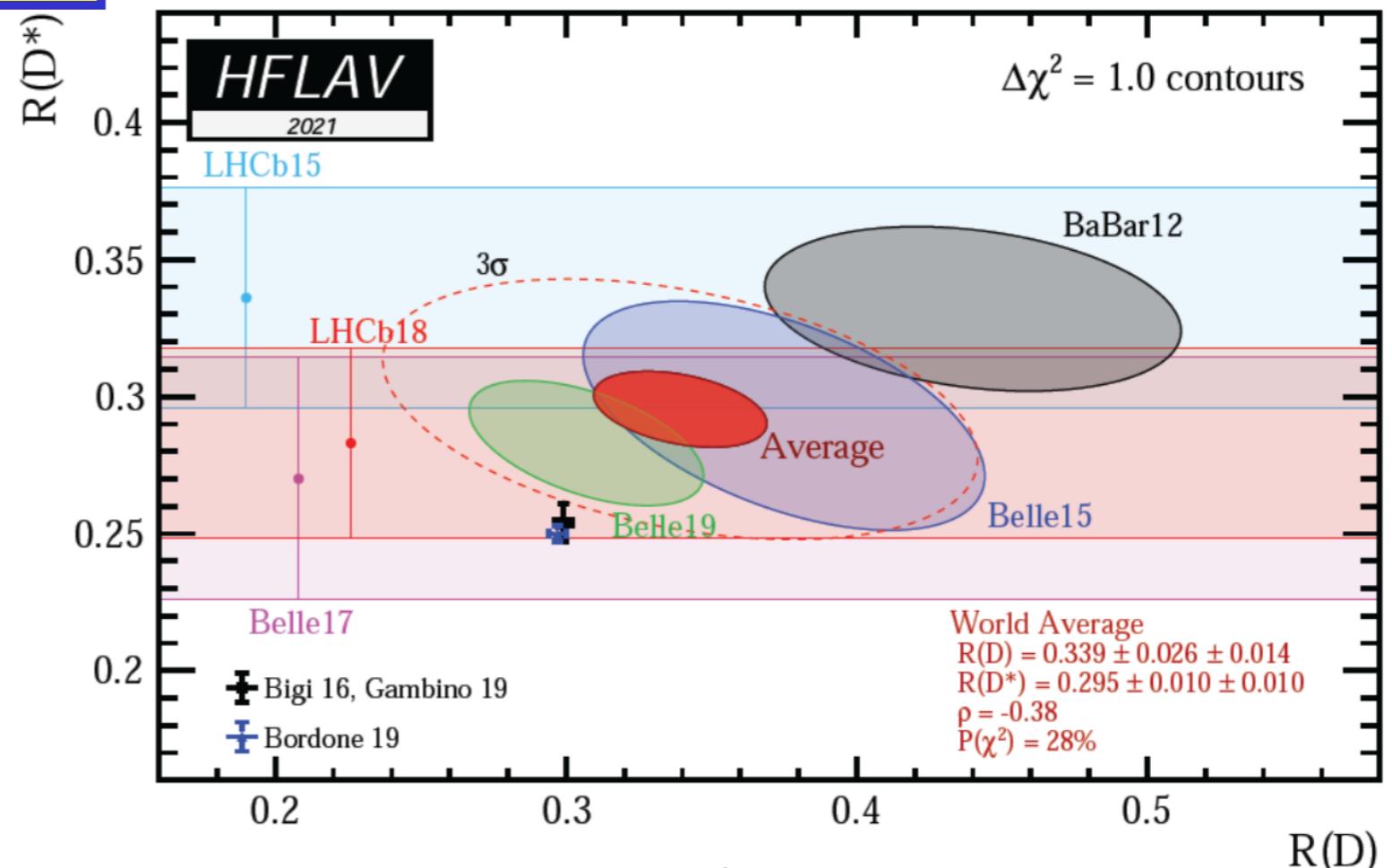
$$R(D^*)_{SM} = 0.258 \pm 0.005$$

- Anomalies first observed by Belle and BaBar

HFLAV 2019 average
of theoretical
predictions

$R(D)$ vs $R(D^*)$

- All experiments see an excess wrt SM predictions
- Combining $R(D)/R(D^*)$ average $\sim 3.4 \sigma$ tension with SM



- Intriguing as anomaly occurs in a tree-level SM process

- New LHCb result

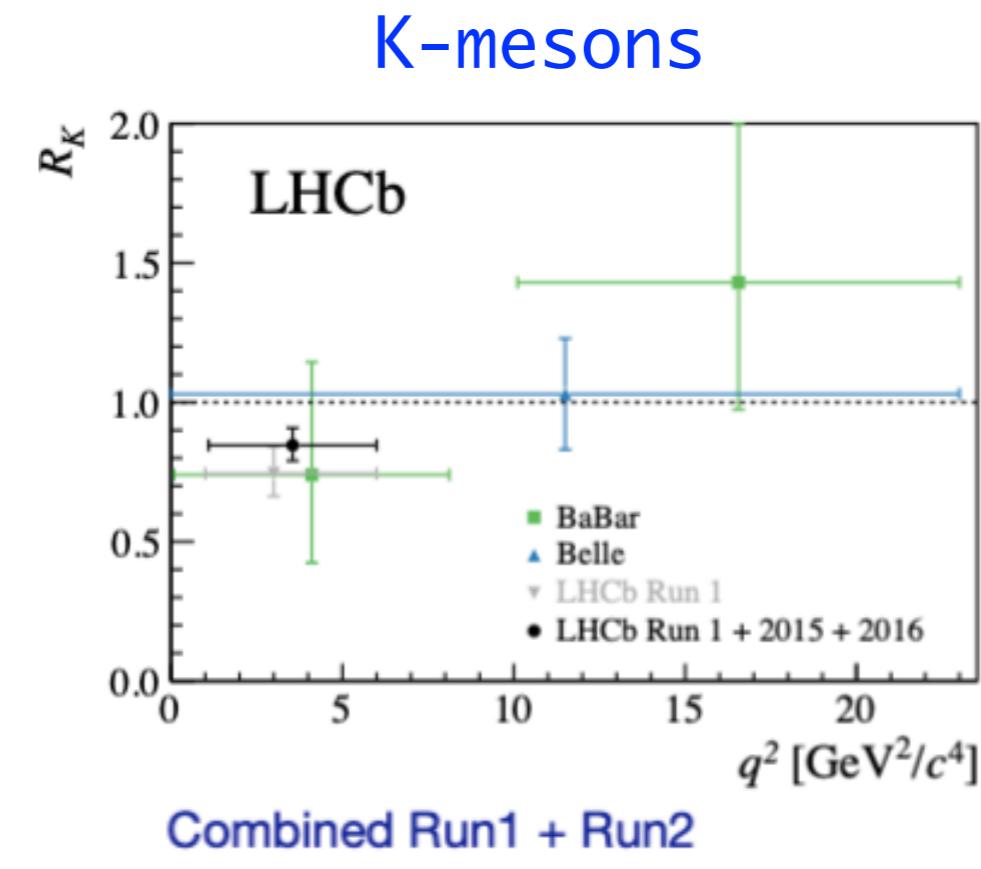
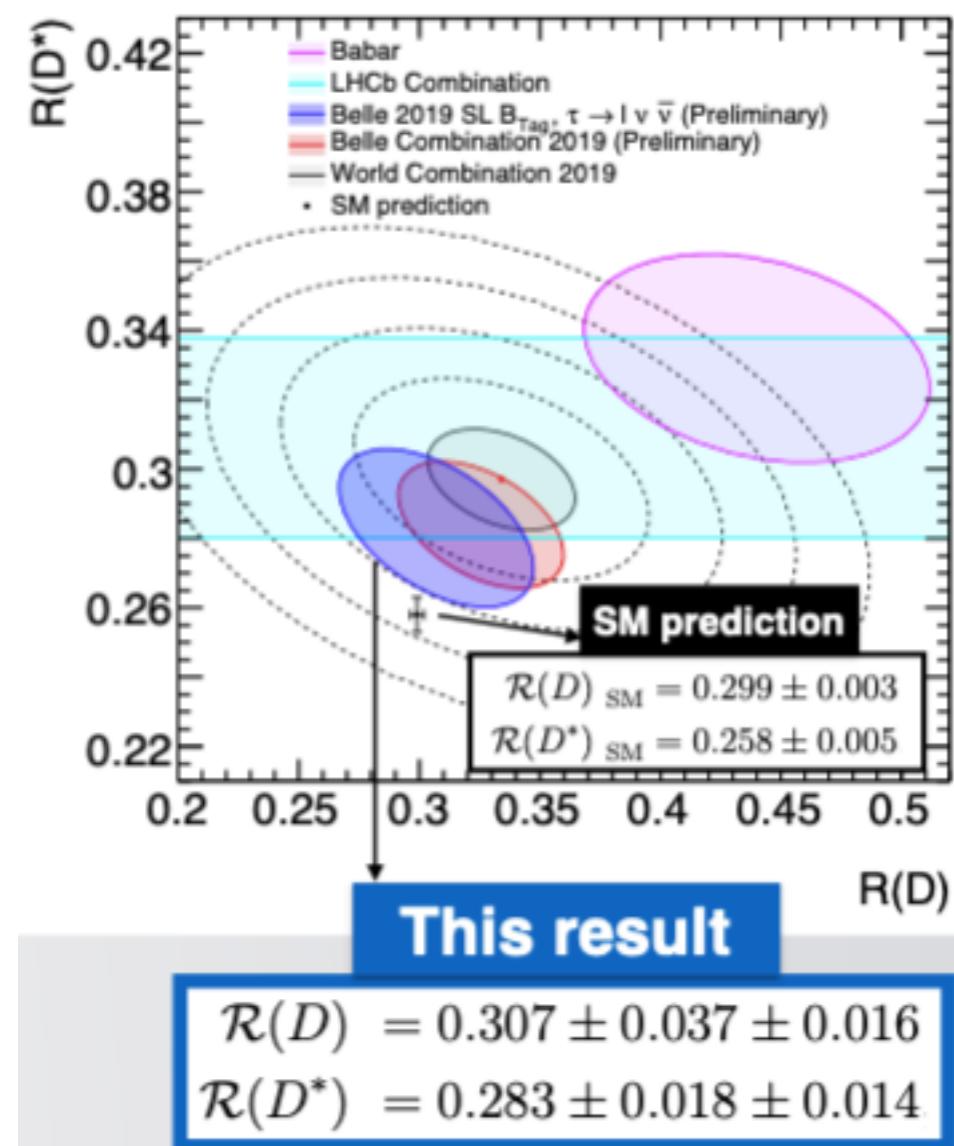
$$R(\Lambda_c) = 0.242 \pm 0.026 \pm 0.040 \pm 0.059(\text{ext})$$

arxiv:2201:03497

Measurement is consistent with SM ($\sim 1\sigma$ “low”) [SM=0.324±0.004].

THE STANDARD MODEL: THE STATUS AND OPEN QUESTIONS

Anomalies in B-meson decays: experiment \neq the SM predictions
 D-mesons

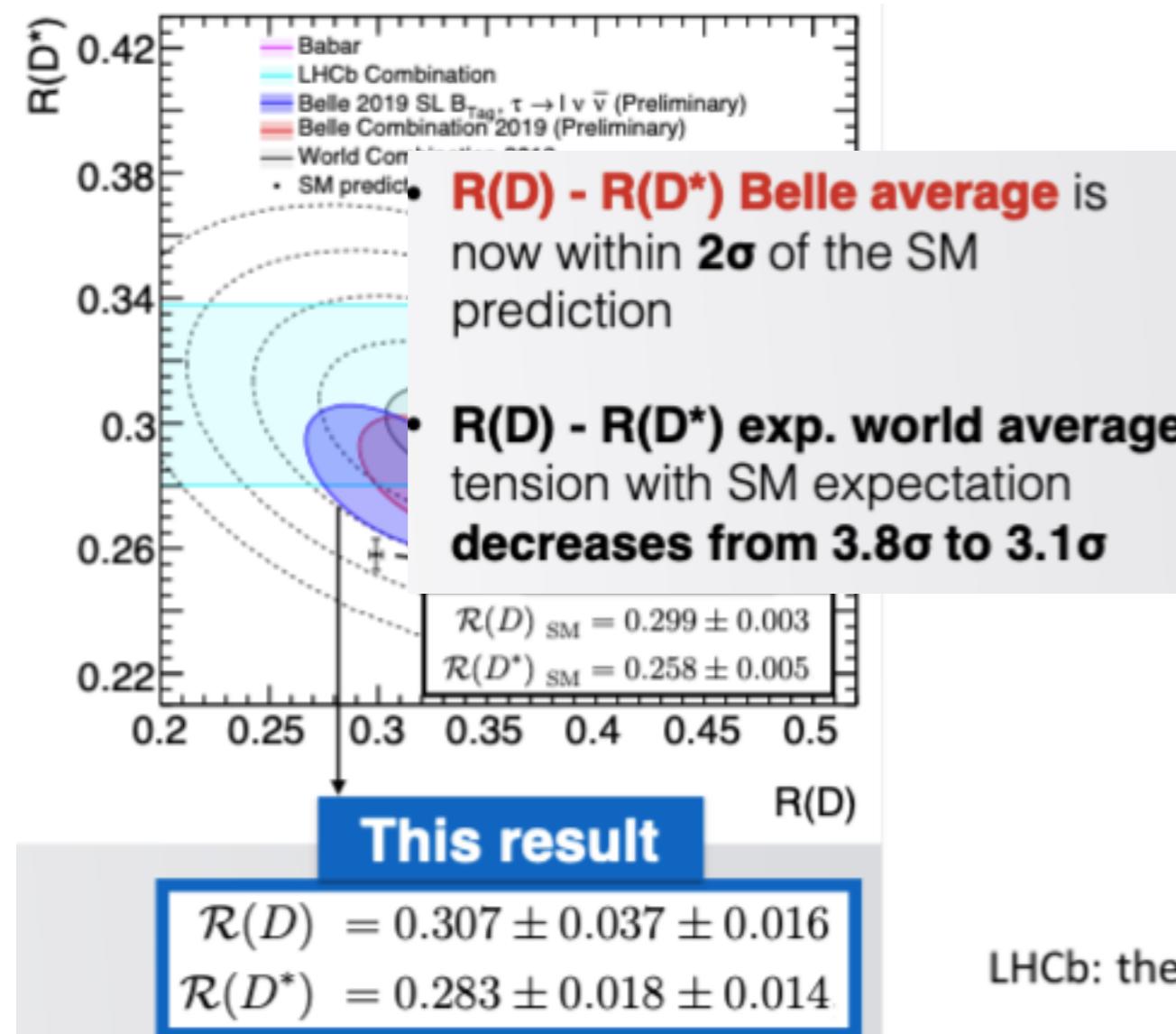


LHCb: the discrepancy present in $B_s \rightarrow \phi \mu \mu$ and $\Lambda_b \rightarrow \Lambda \mu \mu$

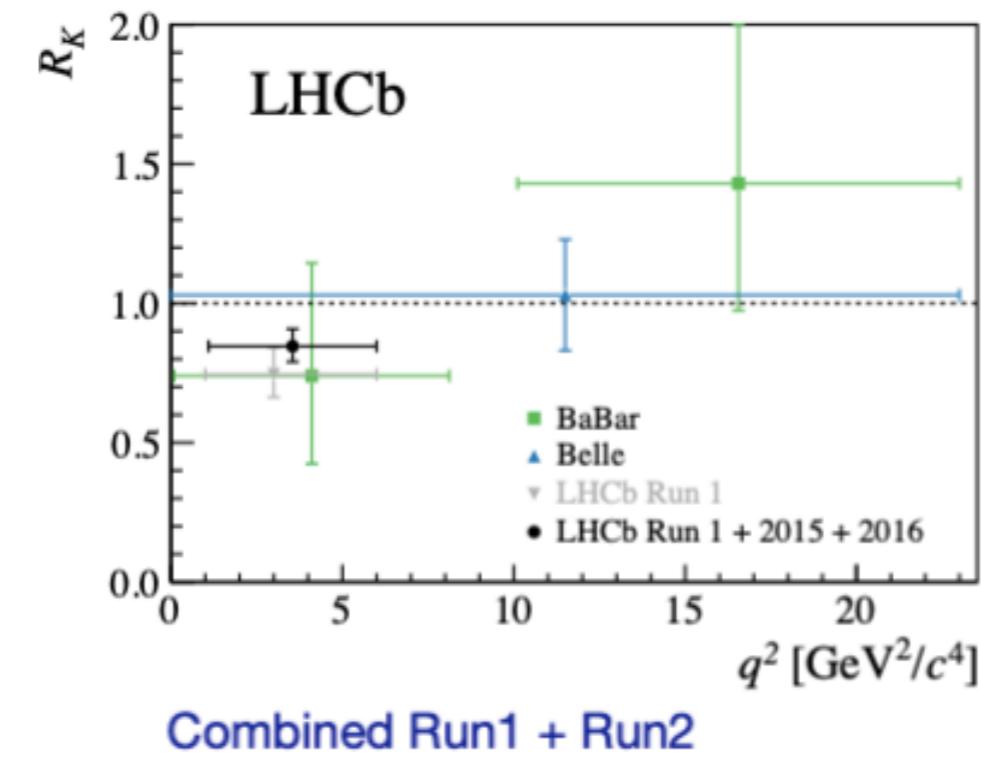
THE STANDARD MODEL: THE STATUS AND OPEN QUESTIONS

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D-mesons



K-mesons

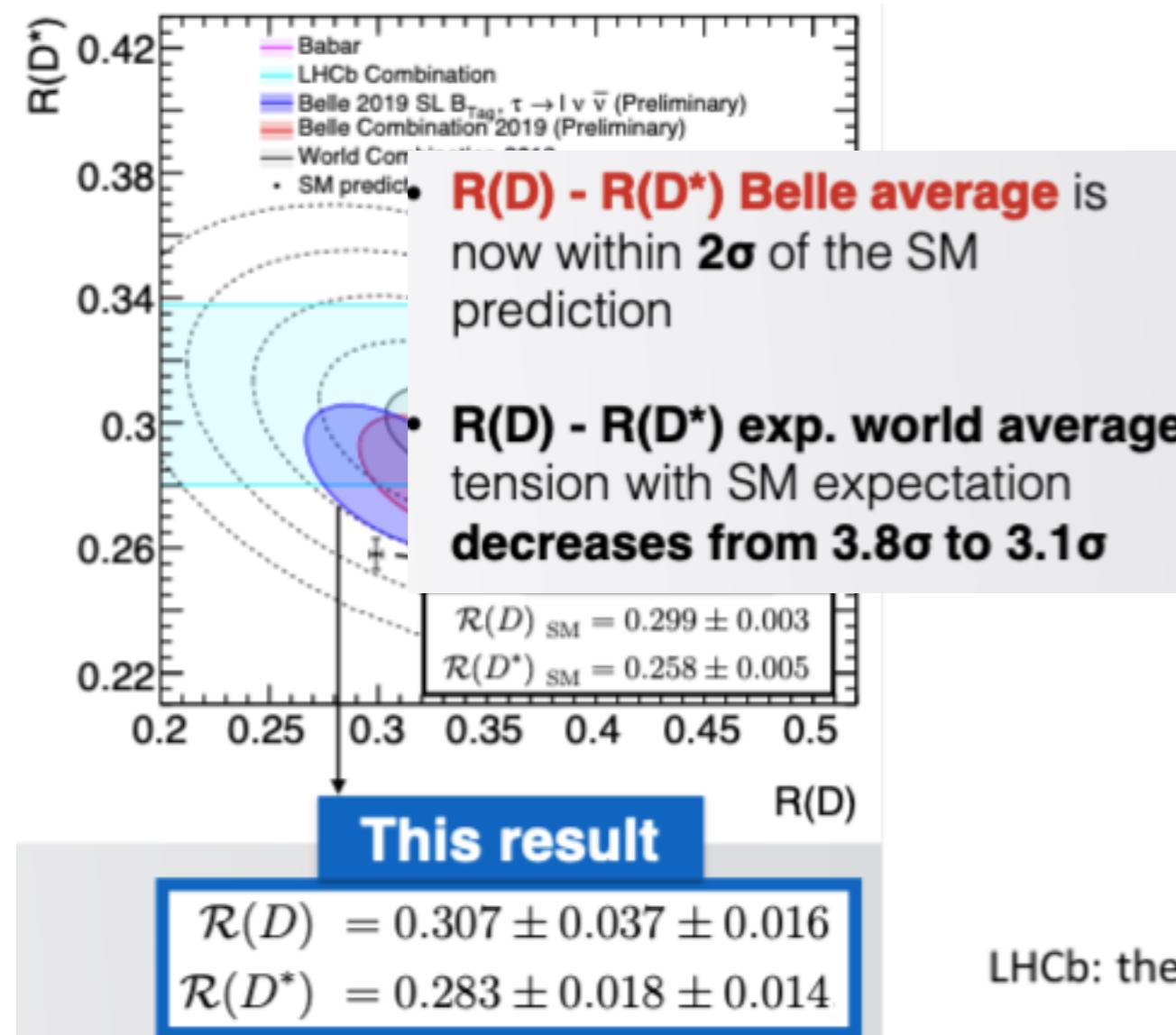


$$R_K = 0.846^{+0.060}_{-0.054} (\text{stat.})^{+0.016}_{-0.014} (\text{syst.})$$

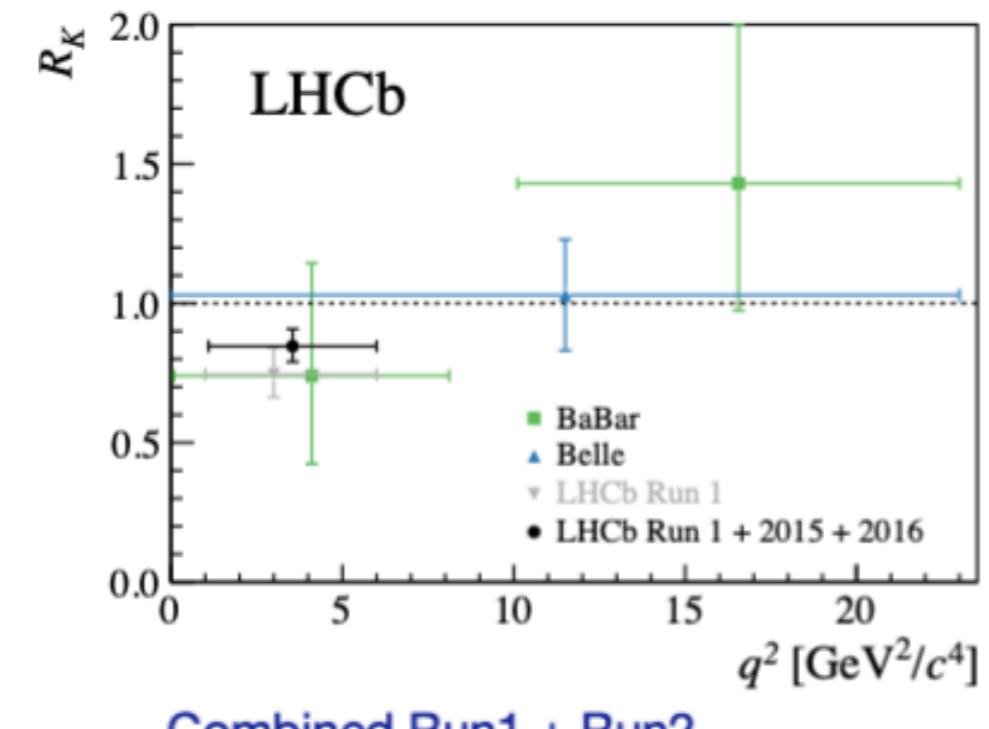
LHCb: the discrepancy present in $B_s \rightarrow \phi \mu \mu$ and $\Lambda_b \rightarrow \Lambda \mu \mu$

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 D-mesons



K-mesons

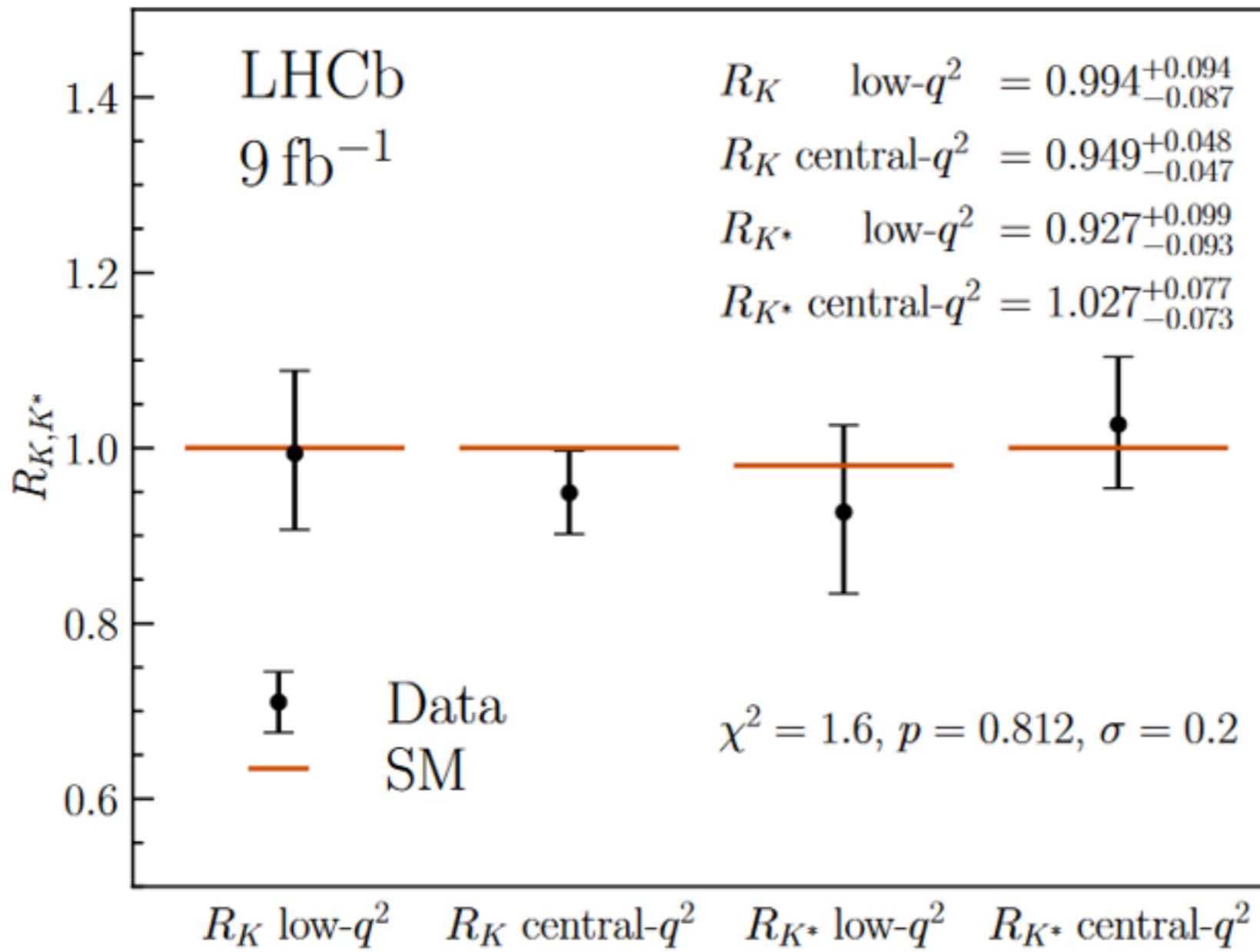


LHCb: the discrepancy present in $B_s \rightarrow \phi \mu \mu$ and $\Lambda_b \rightarrow \Lambda \mu \mu$

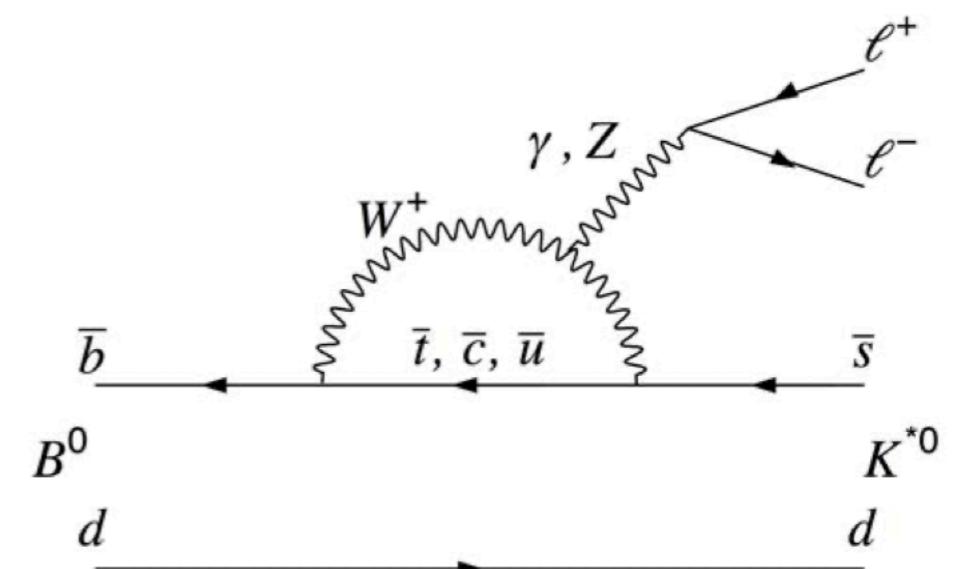
Discrepancy may increase but may decrease

Uncertainty of baryon contribution might be crucial!

THE LHCb RECENT RESULTS



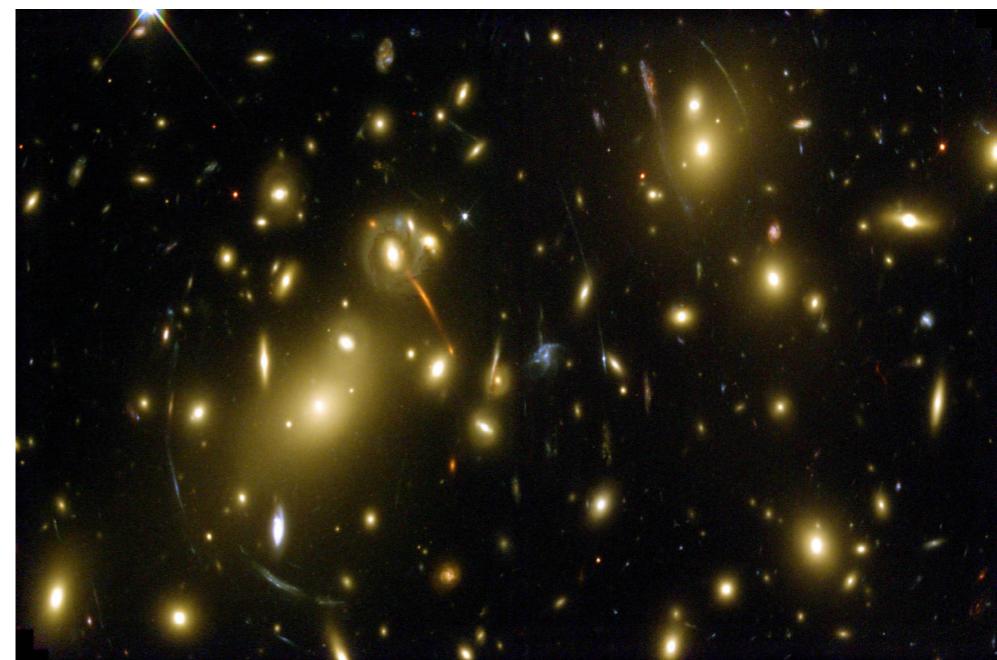
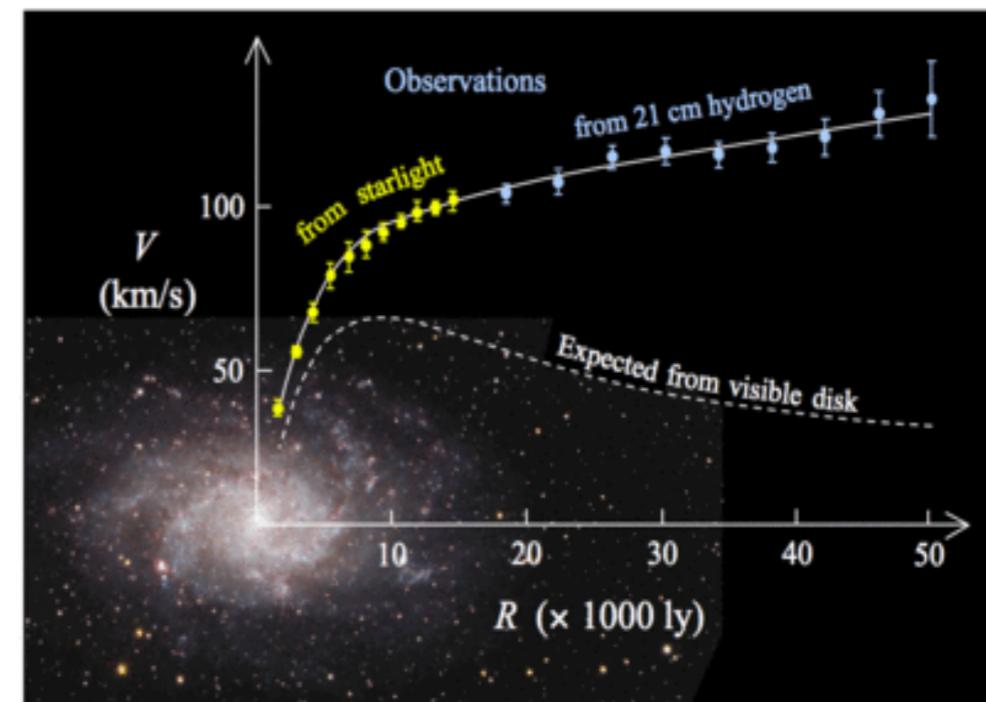
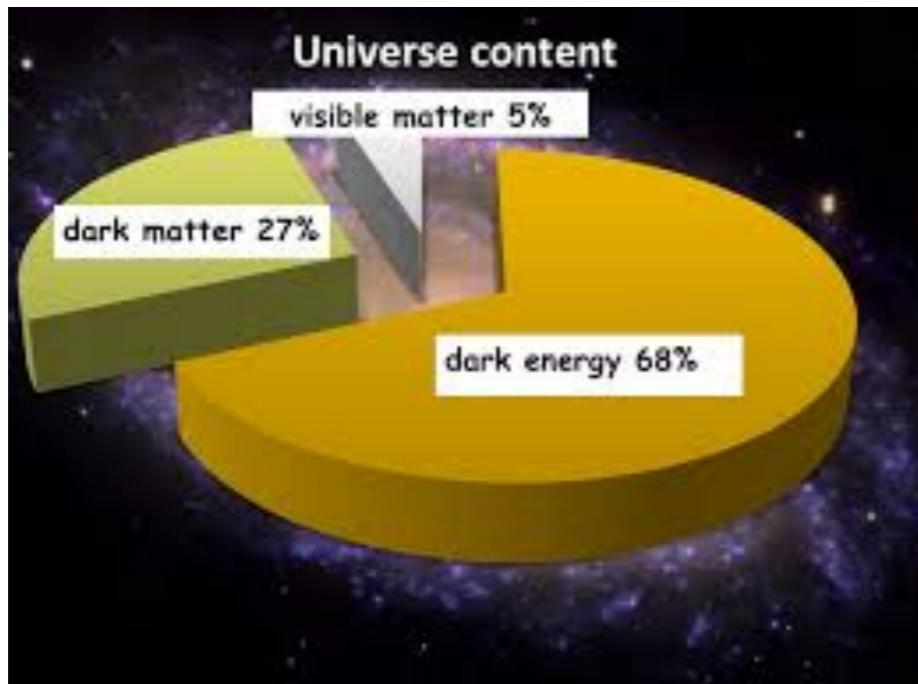
$$R_{K^{(*)}} = \left. \frac{BR(B \rightarrow K^{(*)}\mu\mu)}{BR(B \rightarrow K^{(*)}ee)} \right|$$



$$\mathcal{R}(K^*)_{SM} = 1.0$$

Dark matter

Тёмная материя



Главная проблема: 85% материи является тёмной и остаётся невидимой!

Совместимо ли это с Стандартной моделью?

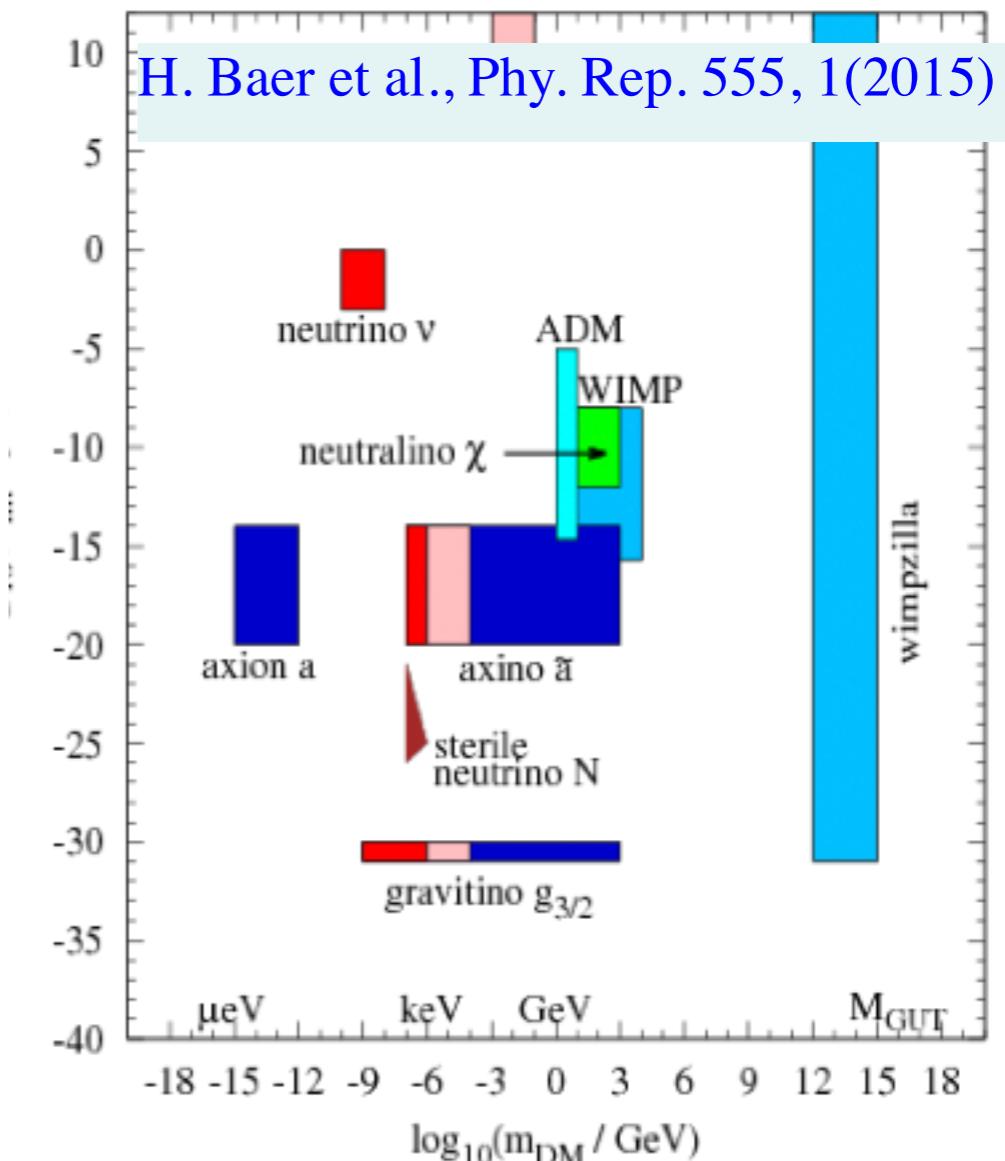
Требует ли это модификации СМ или добавление гравитации?

- **Много кандидатов в разбросах масс в несколько порядков**

- MOND (Problems: large scales, Bullet cluster)
- Primordial black holes (LIGO, but constraints)
- Fuzzy (very light bosons)
- Warm (KeV sterile)
- WIMP
- Axions/ALPs
- Dark sector
- Gravitinos
- Moduli
- Wimpzillas

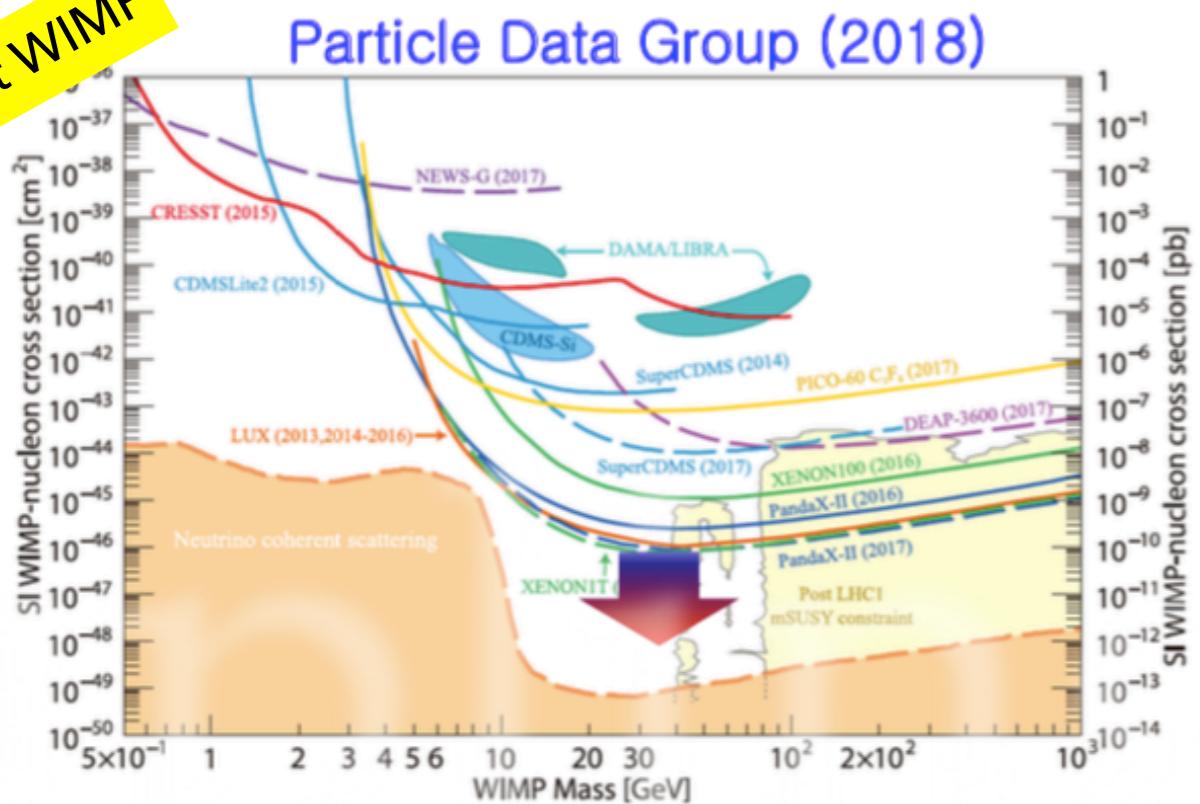


- Прямые, косвенные и коллайдерные поиски тёмной материи

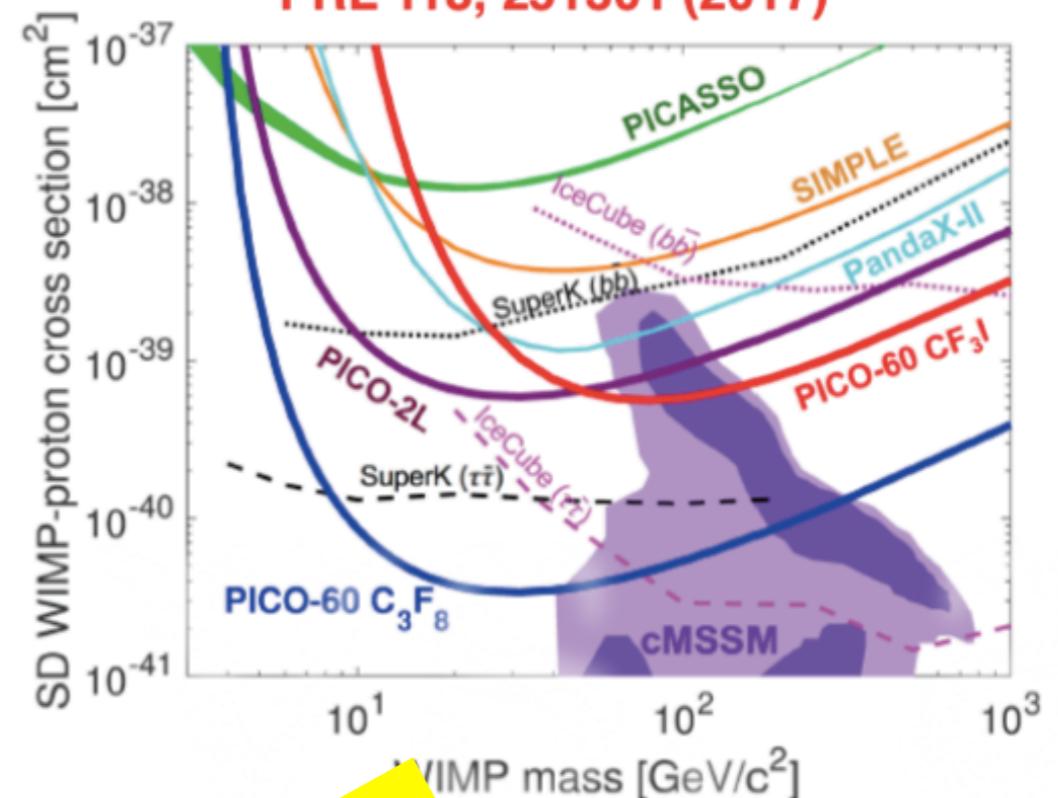


BEYOND THE STANDARD MODEL: DARK MATTER SEARCHES

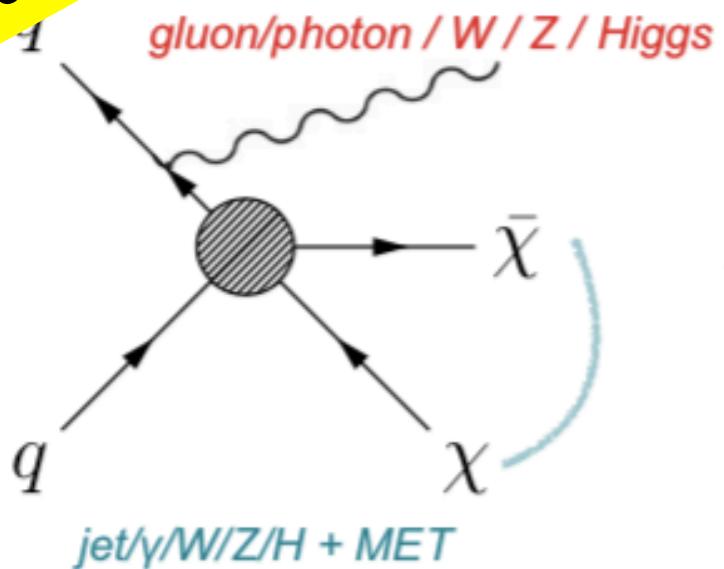
Direct WIMP



PRL 118, 251301 (2017)



Colliders WIMP

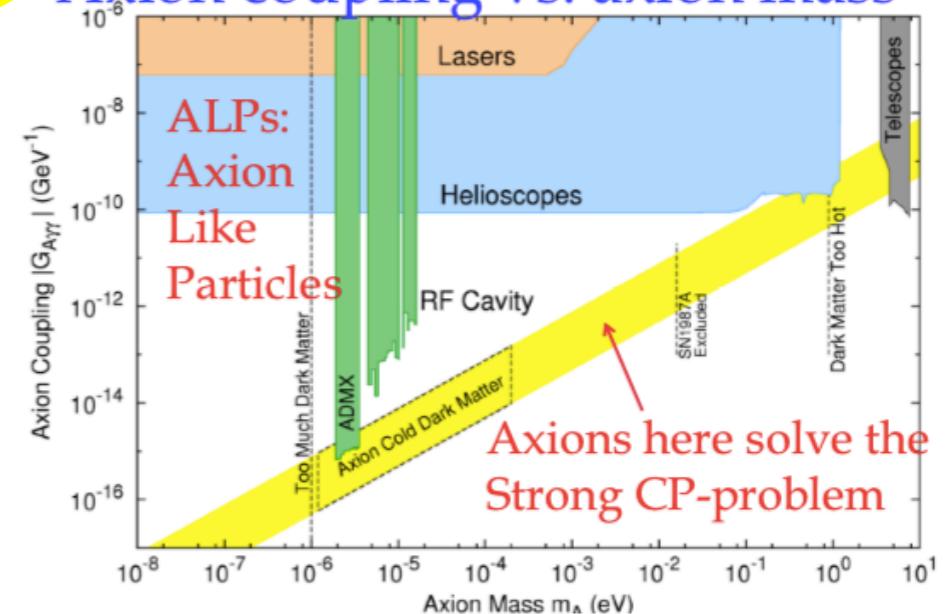


- mono-jet
 - most general signature, constraints on many models
- mono-photon
 - more challenging for background estimation
 - less powerful: EW vs. strong interaction
- mono-W/Z leptonic
 - clean signature and simple trigger
 - penalized by W/Z branching fraction
- mono-W/Z hadronic
 - larger statistics with larger background
- tt+MET/bb+MET and mono-top
 - more complicated experimentally
 - powerful in some scenarios
- mono-Higgs
 - powerful in some scenarios

D. del Re

Axion-likes

Axion coupling vs. axion mass



Y. Semertzidis

Baryon Asymmetry of the Universe

Барионная асимметрия Вселенной



- Барионное число сохраняется в СМ с экспоненциальной точностью
- Нарушение барионного числа имеет место в Теориях Великого Объединения и в моделях Пати-Салама (лептон = четвертый цвет)
Новые частицы = лептокварки, расширенный хиггсовский сектор

$$B = \frac{N_q - N_{\bar{q}}}{3}$$

~~B~~



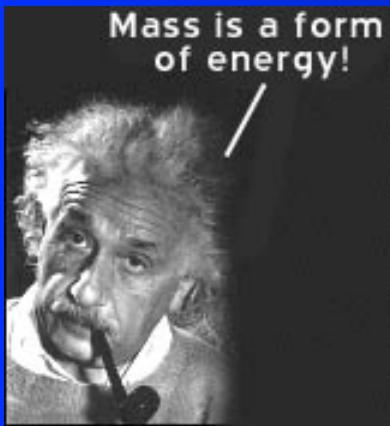
~~CP~~

- Нарушение CP инвариантности в СМ достигается за счёт фаз в матрицах смешивания CKM и PMNS
- ВАУ требует бо́льшего CP нарушения чем есть в СМ
- Возможен бариогенезис через лептогенезис
- В расширенных моделях (2HDM, SUSY, etc) существуют новые фазовые факторы

Gravity

Zagara № 3:

*Как проявляют
гравитацию?*



Общая теория Относительности

$$Action = \int d^4x \sqrt{-g} \left[\frac{c^4}{16\pi G} (R - 2\Lambda) + \mathcal{L}_M \right]$$

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = \frac{8\pi G}{c^2}T_{\mu\nu} \rightarrow R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + g_{\mu\nu}\Lambda = \frac{8\pi G}{c^2}T_{\mu\nu}$$

↑ ↑ ↑
тензор Риччи тензор энергии-импульса материи Космологическая постоянная
скалярная кривизна

Космологическая постоянная есть
вакуумная энергия = Λ^4

Приводит к антигравитации, что
порождает ускоренное расширение
Вселенной

Чтобы получить $\sim 70\%$ вклада в энергетический баланс
Вселенной Λ должна быть порядка 10^{-3} эв.

?

Квантование

$$g_{\mu\nu} = g_{\mu\nu}^{classic} + h_{\mu\nu}$$



метрика



квантовые флюктуации (гравитон)

Проблемы:

- Лишние степени свободы: духи
- Рост вероятностей с энергией: $\sim E^2/M_{Pl}^2$
- Наличие бесконечного числа бесконечностей: неперенормируемость

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Пути решения:

- Модификация сектора материи (суперсимметрия)
- Модификация гравитации (высшие члены по кривизне)
- Нелокальная теория (струна)
- Обуздание неперенормируемости

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?!
?

Решение пока отсутствует

Квантование

Модификация ОТО

$$R \rightarrow R + R^2 + R_{\mu\nu}R^{\mu\nu}$$

скалярная кривизна

тензор Риччи

Изменение космологических сценариев

Поляризация гравитационных волн

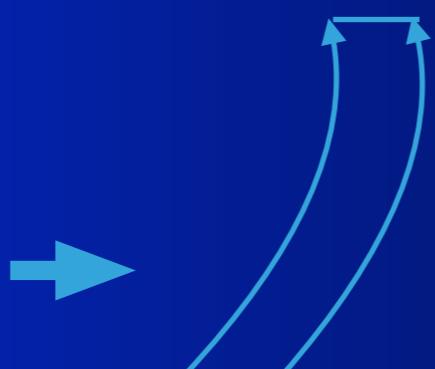
Новая парадигма: Теория струн

Мировая линия



$$X^\mu = X^\mu(\tau)$$

частица



открытая струна

Мировая поверхность

$$X^\mu = X^\mu(\tau, \sigma)$$



замкнутая струна

Частицы есть моды колебаний релятивистской струны

