Kinematically Dependent Renormalization

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Based on arXiv:1804.08387 [hep-th]

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- **This is an attempt to shed some new light on non-renormalizable interactions with the aim to make sense of them at least in some cases**

- **As an example we consider maximally supersymmetric gauge theory in D=8 dimensions and focus on the on-shell scattering amplitudes**

- **The reason is that this case was studied in detail**

in collaboration with A. Borlakov, D.Tolkachev and D.Vlasenko

JHEP 1612 (2016) 154, arXiv:1610.05549v2 [hep-th]
Phys.Rev. D95 (2017) no.4, 045006 arXiv:1603.05501 [hep-th]
Phys.Rev. D97 (2018) no.12, 125008 arXiv:1712.04348 [hep-th],

and has important advantages

- **All analysis in performed within dimensional regularization**

Motivation

Maximal SYM

D=4 N=4
D=6 N=2
D=8 N=1
D=10 N=1

- Partial or total cancell... ...nces (all bubble an... ...cancel)
- First U... ...ams at D=4+6/L
- ...ual conformal symmetry
- ...ion structure of the integrands

All of them can be obtained from 10dim superstring by compactification on a torus

Bern, Dixon &Co 10
Drummond, Henn, Korchemsky, Sokatchev 10
Arkani-Hamed 12

D=4 N=8 Supergravity

- On-shell finite up to 8 loops
- Similar to higher dim SYM

Object: Helicity Amplitudes on mass shell with arbitrary number of legs and loops

The case: Planar limit $N_c \to \infty,\ g^2_{YM} \to 0$ and $g^2_{YM} N_c$ - fixed

The aim: to get all loop (exact) result

Study of higher dim SYM gives insight into quantum gravity

UV divergences in all Loops

Spinor-helicity formalism: S-matrix elements

D=4 N=4	**No UV div**	**IR div on shell**
D=6 N=2	**UV div from 3 loops**	**No IR div**
D=8 N=1	**UV div from 1 loop**	**No IR div**
D=10 N=1	**UV div from 1 loop**	**No IR div**

All these theories are non-renormalizable by power counting

The coupling g^2 **has dimension** $[g^2] = \frac{1}{M^{D-4}}$

The aim: to get all loop (exact) result for the leading (at least) divs

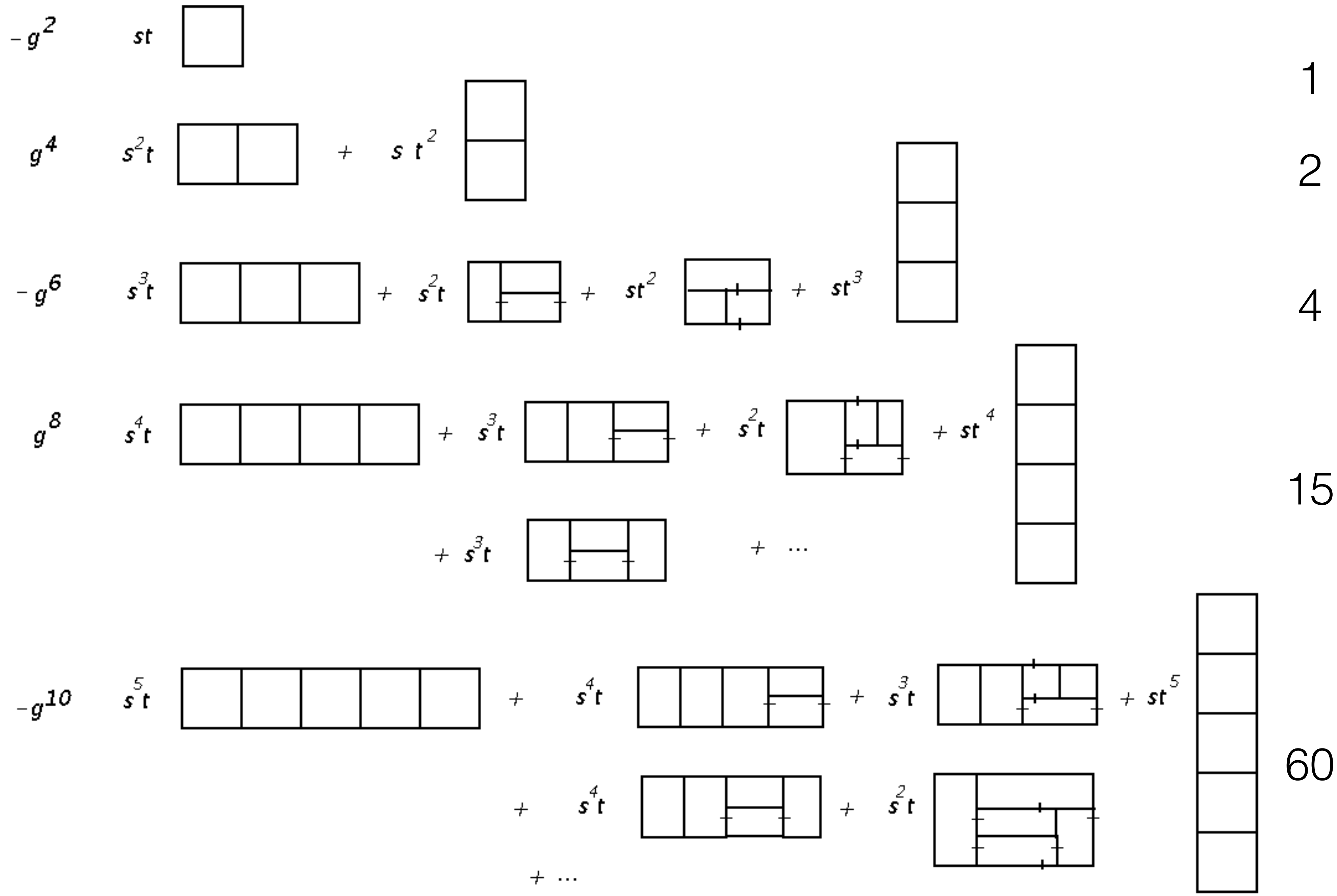
Perturbation Expansion for the 4-point Amplitudes for any D

A_4/A_4^{tree}

No bubbles
No Triangles

First UV div at L=[6/(D-4)] loops

IR finite



T. Dennen Yu-yin Huang 10 ,
S.Caron-Huot D.O'Connell 10

Universal expansion for any D in maximal SYM due to Dual conformal invariance

Leading Divergences from Generalized «Renormalization Group»

- **In renormalizable theories the leading divergences can be found from the 1-loop term due to the renormalization group, in particular, for a single coupling theory the coefficient of $1/\epsilon^n$ in n loops is**

$$\mathcal{R}'G = \sum_n \frac{a_n^{(n)}}{\epsilon^n} \qquad a_n^{(n)} = (a_1^{(1)})^n$$

- **In non-renormalizable theories the leading divergences can be also found from 1-loop due to locality and R-operation**

$$\mathcal{R}'G = 1 - \sum_\gamma K\mathcal{R}'_\gamma + \sum_{\gamma,\gamma'} K\mathcal{R}'_\gamma K\mathcal{R}'_{\gamma'} - ...,$$

$$\mathcal{R}'G_n = \frac{A_n^{(n)}(\mu^2)^{n\epsilon}}{\epsilon^n} + \frac{A_{n-1}^{(n)}(\mu^2)^{(n-1)\epsilon}}{\epsilon^n} + ... + \frac{A_1^{(n)}(\mu^2)^{\epsilon}}{\epsilon^n}$$

$$+ \frac{B_n^{(n)}(\mu^2)^{n\epsilon}}{\epsilon^{n-1}} + \frac{B_{n-1}^{(n)}(\mu^2)^{(n-1)\epsilon}}{\epsilon^{n-1}} + ... + \frac{B_1^{(n)}(\mu^2)^{\epsilon}}{\epsilon^{n-1}}$$

$+$lower order terms

Leading pole

SubLeading pole

$A_1^{(n)}, \ B_1^{(n)}$ **1-loop graph**

$B_2^{(n)}$ **2-loop graph**

SubLeading Divergences from Generalized «Renormalization Group»

- **In non-renormalizable theories the leading divergences can be also found from 1-loop due to locality and R-operation**

All terms like $(log\mu^2)^m/\epsilon^k$ **should cancel**

$$A_n^{(n)} = (-1)^{n+1}\frac{A_1^{(n)}}{n},$$

$$B_n^{(n)} = (-1)^n\left(\frac{2}{n}B_2^{(n)} + \frac{n-2}{n}B_1^{(n)}\right)$$

Leading pole from 1 loop diagrams

SubLeading pole from 2 loop diagrams

$$\mathcal{KR}'G_n = \sum_{k=1}^{n}\left(\frac{A_k^{(n)}}{\epsilon^n} + \frac{B_k^{(n)}}{\epsilon^{n-1}}\right) \equiv \frac{A_n^{(n)'}}{\epsilon^n} + \frac{B_n^{(n)'}}{\epsilon^{n-1}}.$$

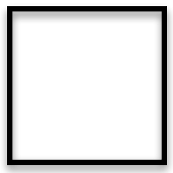
$$A_n^{(n)'} = (-1)^{n+1}A_n^{(n)} = \frac{A_1^{(n)}}{n},$$

$$B_n^{(n)'} = \left(\frac{2}{n(n-1)}B_2^{(n)} + \frac{2}{n}B_1^{(n)}\right)$$

Just like in renormalizable theories one can deduce the leading, subheading, etc divergences from 1, 2, etc diagrams

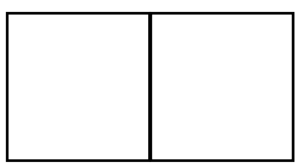
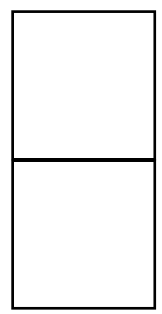
Kinematically dependent renormalization

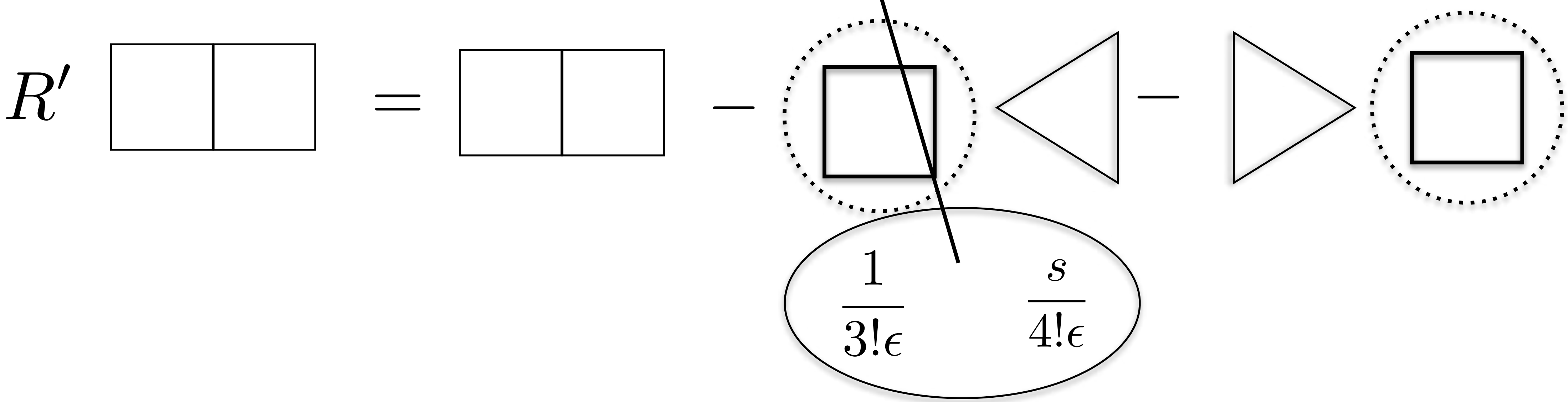
One-loop box

st  $= \frac{st}{3!}\frac{1}{\epsilon}$

Totally defined by 1 loop **Independent term**

Two-loop box

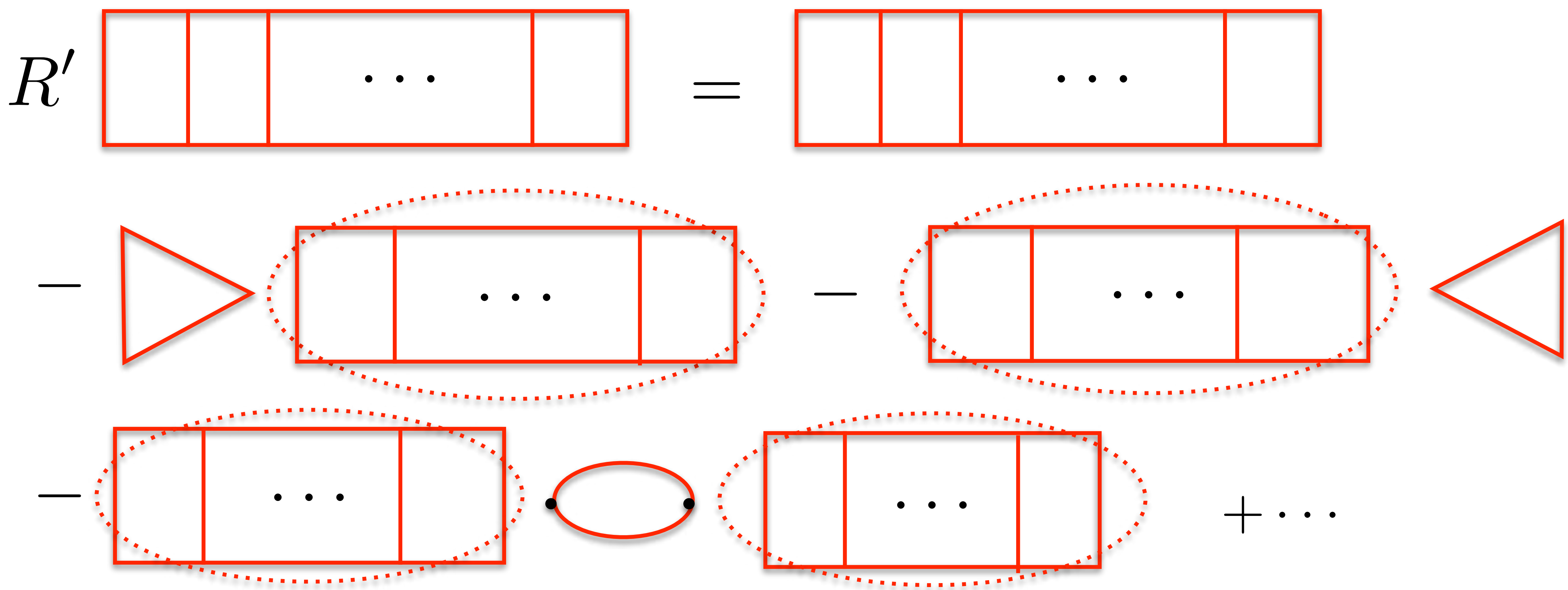
s^2t  + st^2  $= \frac{st}{3!4!}\left(\frac{s^2+t^2}{\epsilon^2} + \frac{27/4\ s^2 + 1/3\ st + 27/4\ t^2}{\epsilon}\right)$



This is true to all orders of PT like in renormalizable theories via the locality of the counnterterms due to the R-operation

Ladder diagrams (leading divs)

D=8 N=1 **Horizontal boxes**



Leading poles $A_n^{(n)} = s^{n-1} A_n$

1 loop box

$$A_1 = \frac{1}{3!}$$

$$A_2 = -\frac{1}{3!4!}$$

$$A_3 = \frac{2}{33!4!4!} + \frac{2}{35!3!3!}$$

$$nA_n = -\frac{2}{4!}A_{n-1} + \frac{2}{5!}\sum_{k=1}^{n-2} A_k A_{n-1-k}, \quad n \geq 3$$

Ladder diagrams (leading divs)

D=8 N=1 **Horizontal boxes** $A_n^{(n)} = s^{n-1} A_n$

$$nA_n = -\frac{2}{4!} A_{n-1} + \frac{2}{5!} \sum_{k=1}^{n-2} A_k A_{n-1-k}, \quad n \geq 3$$

$A_1 = 1/6$ **1 loop box**

Summation $\Sigma_m(z) = \sum_{n=m}^{\infty} A_n (-z)^n$

$$-\frac{d}{dz}\Sigma_3 = -\frac{2}{4!}\Sigma_2 + \frac{2}{5!}\Sigma_1\Sigma_1. \qquad \Sigma_3 = \Sigma_1 + A_1 z - A_2 z^2, \quad \Sigma_2 = \Sigma_1 + A_1 z, \quad A_1 = \frac{1}{3!}, \quad A_2 = -\frac{1}{3!4!}$$

$\Sigma_A \equiv \Sigma_1$ **Diff eqn**

$$\frac{d}{dz}\Sigma_A = -\frac{1}{3!} + \frac{2}{4!}\Sigma_A - \frac{2}{5!}\Sigma_A^2$$

$z = g^2 s^2/\epsilon$

$$\Sigma_A(z) = -\sqrt{5/3}\,\frac{4\tan(z/(8\sqrt{15}))}{1 - \tan(z/(8\sqrt{15}))\sqrt{5/3}} = \sqrt{10}\,\frac{\sin(z/(8\sqrt{15}))}{\sin(z/(8\sqrt{15}) - z_0)}$$

$$\Sigma(z) = -(z/6 + z^2/144 + z^3/2880 + 7z^4/414720 + \dots) \qquad z_0 = \arcsin(\sqrt{3/8})$$

All loop Exact Recurrence Relation

D=8 N=1

s-channel term $S_n(s,t)$ **t-channel term** $T_n(s,t)$ $T_n(s,t) = S_n(t,s)$

Exact relation for ALL diagrams

$$nS_n(s,t) = -2s^2 \int_0^1 dx \int_0^x dy\ y(1-x)\ \left(S_{n-1}(s,t') + T_{n-1}(s,t')\right)|_{t'=tx+yu}$$

$$+ \quad s^4 \int_0^1 dx\ x^2(1-x)^2 \sum_{k=1}^{n-2} \sum_{p=0}^{2k-2} \frac{1}{p!(p+2)!}\ \frac{d^p}{dt'^p}\left(S_k(s,t') + T_k(s,t')\right) \times$$

$$S_1 = \frac{1}{12},\ T_1 = \frac{1}{12} \qquad \times \frac{d^p}{dt'^p}\left(S_{n-1-k}(s,t') + T_{n-1-k}(s,t')\right)|_{t'=-sx}\ \left(tsx(1-x)\right)^p$$

summation $\Sigma_3(s,t,z) = \Sigma_1(s,t,z) - S_2(s,t)z^2 + S_1(s,t)z,\ \Sigma_2(s,t,z) = \Sigma_1(s,t,z) + S_1(s,t)z$

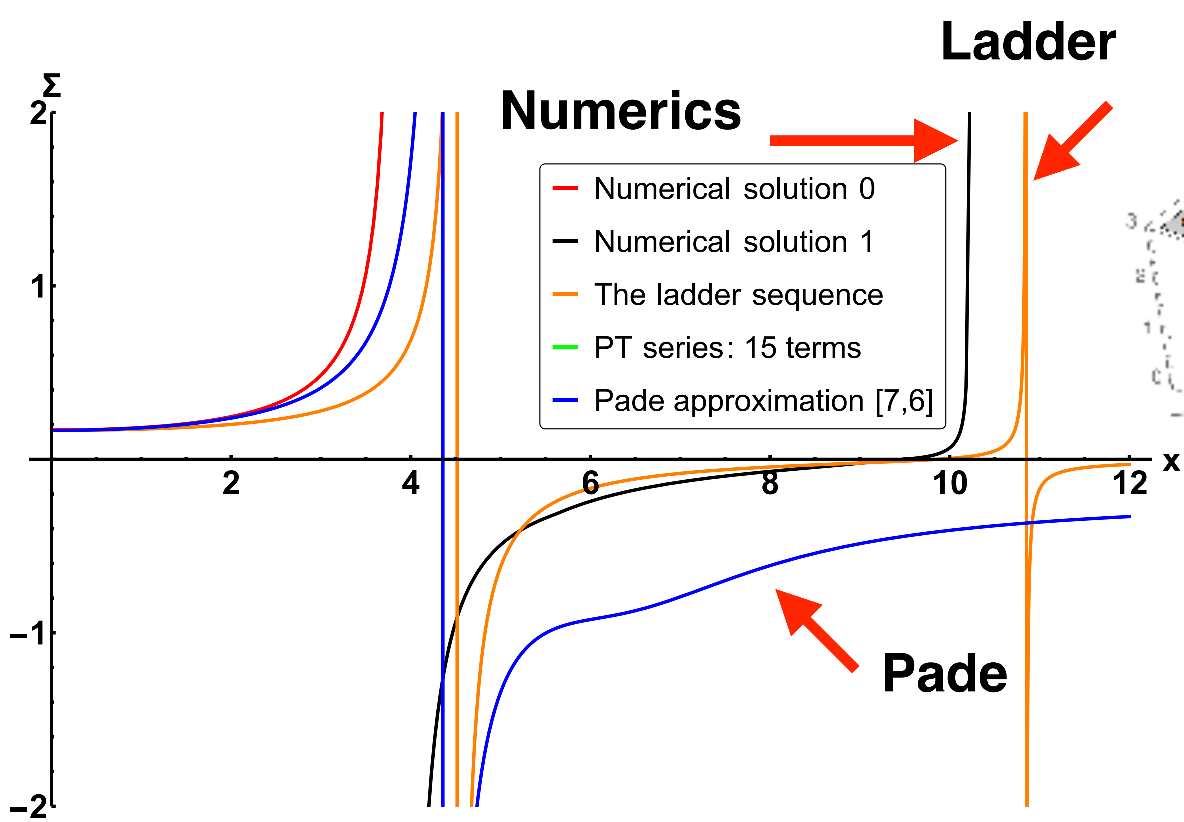
Diff eqn

$$\frac{d}{dz}\Sigma(s,t,z) = -\frac{1}{12} + 2s^2 \int_0^1 dx \int_0^x dy\ y(1-x)\ \left(\Sigma(s,t',z) + \Sigma(t',s,z)\right)|_{t'=tx+yu}$$

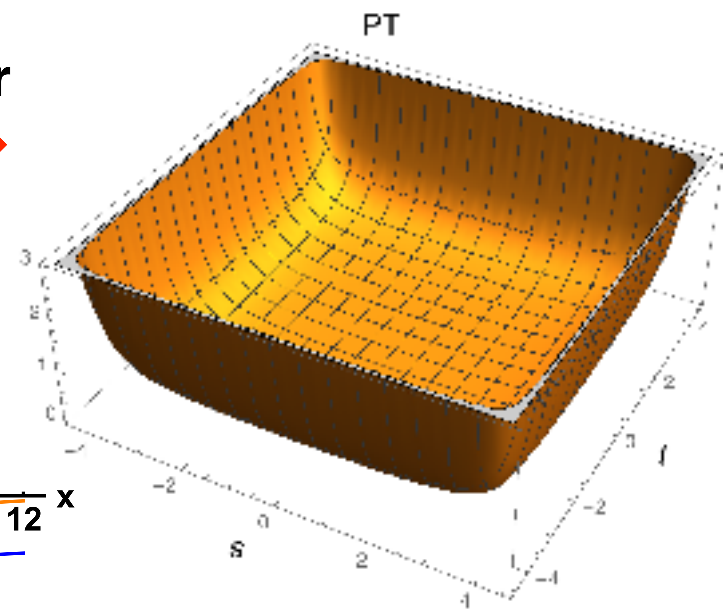
$$-s^4 \int_0^1 dx\ x^2(1-x)^2 \sum_{p=0}^{\infty} \frac{1}{p!(p+2)!}\left(\frac{d^p}{dt'^p}\left(\Sigma(s,t',z) + \Sigma(t',s,z)\right)|_{t'=-sx}\right)^2\ \left(tsx(1-x)\right)^p.$$

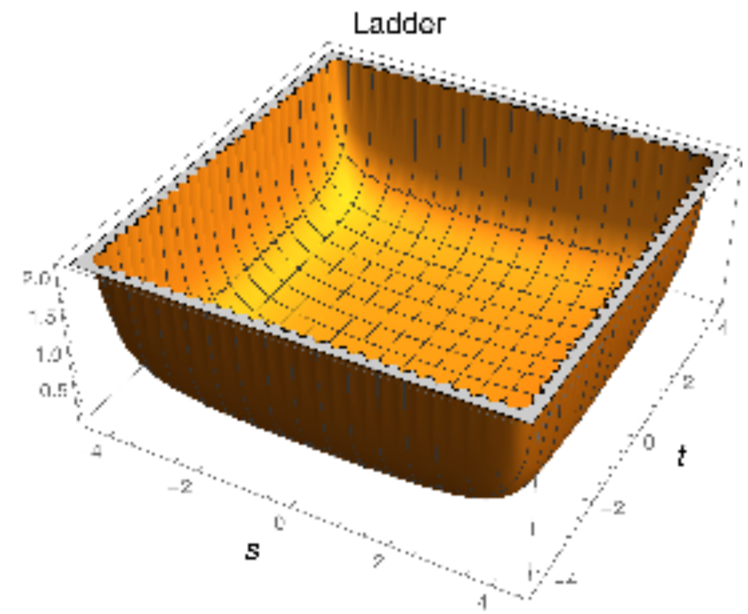
All loop Solution (leading divs)

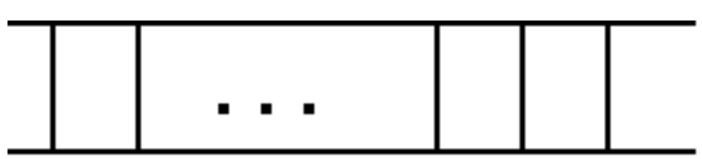
D=8 N=1



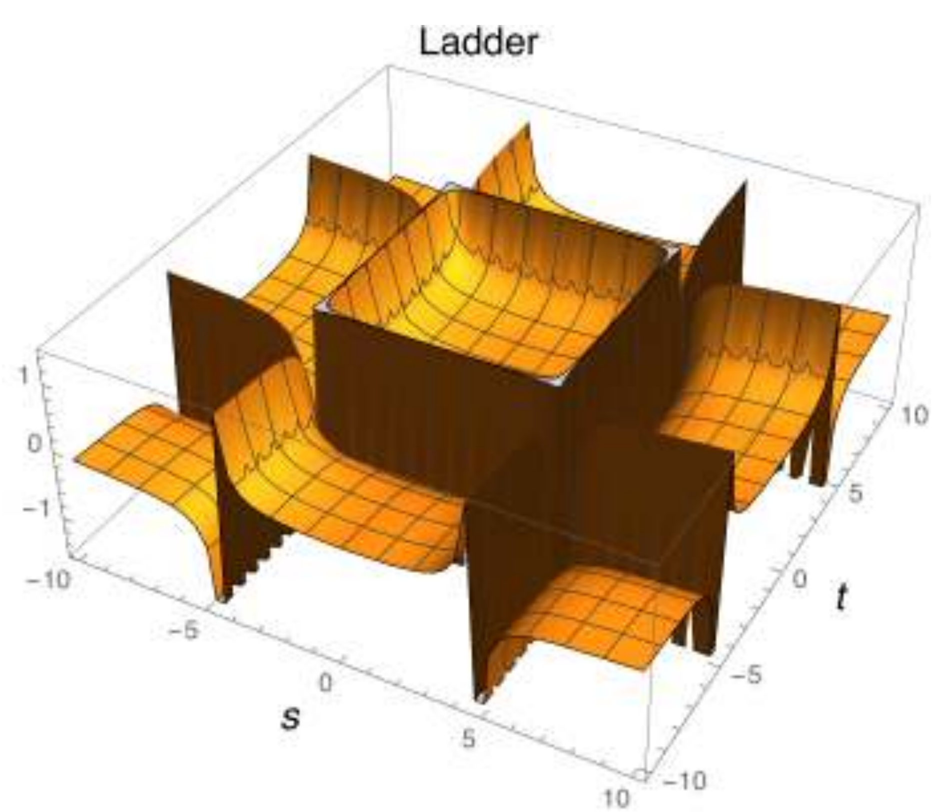
PT and Pade versus ladder for t=s







$$z = \frac{g^2}{\epsilon}$$



$$\Sigma_L(s,z) = -\sqrt{5/3}\frac{4\tan(zs^2/(8\sqrt{15}))}{1-\tan(zs^2/(8\sqrt{15}))\sqrt{5/3}}$$

Subleading divergences

$$\Sigma_L(z) + \epsilon\Sigma_{NL}(z) + \epsilon^2\Sigma_{NNL}(z) + \cdots$$

$$\Sigma(z) = \sum_n^\infty z^n F_n$$

$$D = 4 \quad N = 4 \quad z = g^2/\epsilon$$
$$D = 6 \quad N = 2 \quad z = g^2 s/\epsilon, z = g^2 t/\epsilon$$
$$D = 8 \quad N = 1 \quad z = g^2 s^2/\epsilon, z = g^2 st/\epsilon, ..$$
$$D = 10 \quad N = 1 \quad z = g^2 s^3/\epsilon, z = g^2 s^2 t/\epsilon, ..$$

D=8 N=1

sLadder case

$$\Sigma_{NL} = s\Sigma_{sB}(z) + t\Sigma_{tB}(z)$$

$$z = \frac{g^2 s^2}{\epsilon}$$

$$\Sigma'_{tB}(z) = \frac{5}{6}\left[e^{z/60}(2\cos(z/30) - \sin(z/30)) - 2\right]$$

$$\Sigma_{tB} = -\frac{1}{36}\left[60 + z + e^{z/60}(-(60+z)\cos(z/30) - 2(-15+z)\sin(z/30))\right]$$

Sum of Ladder diagrams (subleading divs)

$$\Sigma'_{sB} = \sum_{n=2}^{\infty} z^n B'_{sn}$$

$$\frac{d^2\Sigma'_{sB}(z)}{dz^2} + f_1(z)\frac{d\Sigma'_{sB}(z)}{dz} + f_2(z)\Sigma'_{sB}(z) = f_3(z)$$

Diff eqn

$$f_1(z) = -\frac{1}{6} + \frac{\Sigma_A}{15},$$

$$f_2(z) = \frac{1}{80} - \frac{\Sigma_A}{360} + \frac{\Sigma_A^2}{600} + \frac{1}{15}\frac{d\Sigma_A}{dz},$$

$$f_3(z) = \frac{2321}{5!5!2}\Sigma_A + \frac{11}{1800}\Sigma'_{tB} - \frac{47}{5!45}\Sigma_A^2 - \frac{1}{5!72}\Sigma_A\Sigma'_{tB} + \frac{23}{6750}\Sigma_A^3 + \frac{1}{1200}\Sigma_A^2\Sigma'_{tB}$$
$$-\frac{19}{36}\frac{d\Sigma_A}{dz} - \frac{1}{15}\frac{d\Sigma'_{tB}}{dz} + \frac{23}{225}\frac{d\Sigma_A^2}{dz} + \frac{1}{30}\frac{d(\Sigma_A\Sigma'_{tB})}{dz} - \frac{3}{32}$$

Solution to Diff eqn

smooth monotonic function

$$\Sigma'_{sB}(z) = \frac{d\Sigma_A}{dz}u(z) \qquad u(z) = \int_0^z dy \int_0^y dx \frac{f_3(x)}{d\Sigma_A(x)/dx}$$

Scheme dependence and arbitrariness of subtraction

subleading case

$$A_1' + B_{s1}' = \frac{1}{6\epsilon}(1 + c_1\epsilon) \qquad \Delta\Sigma_{sB}' = c_1 z \frac{d\Sigma_A'}{dz}. \quad \longrightarrow \quad z \to z(1 + c_1\epsilon).$$

sub-subleading case

$$A_2' + B_2' = \frac{s}{3!4!\epsilon^2}\left(1 - \frac{5}{12}\epsilon + 2c_1\epsilon + c_2\epsilon^2\right) \qquad \Delta\Sigma_{sC}' = c_2 z^2 \frac{d\Sigma_A'}{dz}.$$

$$\longrightarrow \quad z \to z(1 + c_1\epsilon) + z^2 c_2\epsilon^2.$$

$$\Delta\Sigma_{sC}' = -c_1^2\frac{z}{4!}\left(\frac{d\Sigma_A'}{dz} - 12\frac{d^2\Sigma_A'}{dz^2}\right) \quad \longrightarrow \quad z \to z(1 + c_1\epsilon) + z^2(c_2 + c_1^2/4!)\epsilon^2$$

Scheme dependence and arbitrariness of subtraction

sub-subleading case

linear term

$$A_2' + B_2' = \frac{s}{3!4!\epsilon^2}\left(1 - \frac{5}{12}\epsilon + 2c_1\epsilon + c_2\epsilon^2\right)$$

new contribution from subleading term

$$\Delta\Sigma_{sC}'(3-loop) = -\frac{719 c_1 s^2}{1036800\epsilon}$$

$$\Sigma_{sB}'(3-loop) = -\frac{71 s^2}{345600\epsilon^2}$$

$$\Sigma_{sB}'^{trunc}(3-loop) = -\frac{719 s^2}{3110400\epsilon^2}$$

the source of a problem

$$\Delta\Sigma_{sC}'(3-loop) = c_1 z \frac{d\Sigma_{sB}'^{trunc}}{dz}(3-loop)$$

$$z \to z(1 + c_1\epsilon) + z^2(c_2 - c_1^2/4!)\epsilon^2 + z^3 c_1^3/6!\epsilon^3 - z^4 c_1^4/4!6!\epsilon^4 + \ldots.$$

Kinematically dependent renormalization

- **R-operation is equivalent to**

renormalizable theories

nonrenormalizable theories

$$\bar{A}_4 = Z_4(g^2)\bar{A}_4^{bare}|_{g^2_{bare} -> g^2 Z_4}$$

$$g^2_{bare} = \mu^\epsilon Z_4(g^2) g^2.$$

$$Z = 1 - \sum_i KR'G_i$$

simple multiplication

operator multiplication

$$Z = 1 + \frac{g^2}{\epsilon} + g^4(\frac{1}{\epsilon^2} + \frac{1}{\epsilon}) + \ldots$$

$$Z = 1 + \frac{g^2}{\epsilon} st + g^4 st(\frac{s^2+t^2}{\epsilon^2} + \frac{s^2+st+t^2}{\epsilon}) + \cdots$$

$$\frac{g^2}{\epsilon}(D_\rho D_\sigma F_{\mu\nu})^2$$

Kinematically dependent renormalization

operator kinematically dependent renormalization

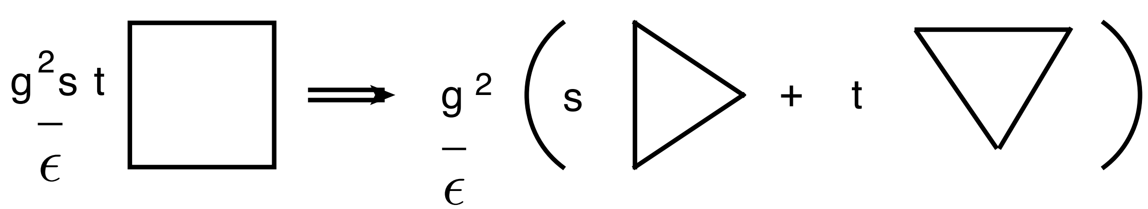
at 2 loops

$$\bar{A}_4 = 1 - \frac{g_B^2 st}{3!\epsilon} - \frac{g_B^4 st}{3!4!}\left(\frac{s^2+t^2}{\epsilon^2} + \frac{27/4s^2 + 1/3st + 27/4t^2}{\epsilon}\right) + \dots$$

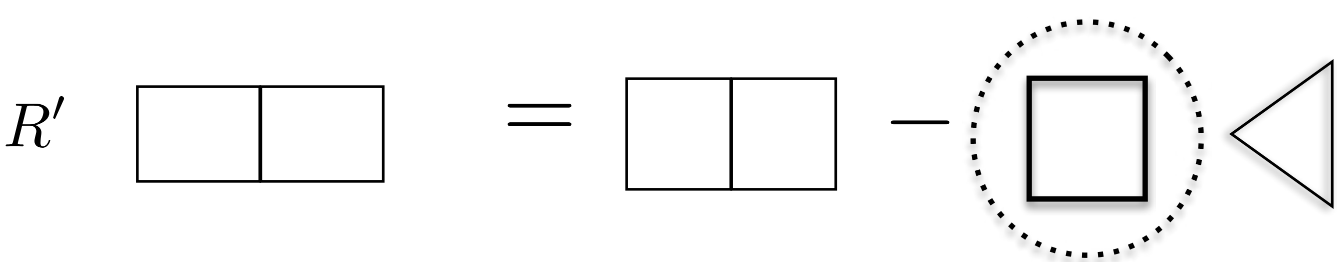
$$\bar{A}_4 = Z_4(g^2)\bar{A}_4^{bare}|_{g^2_{bare} -> g^2 Z_4}$$

$$Z_4 = 1 + \frac{g^2 st}{3!\epsilon} + \frac{g^4 st}{3!4!}\left(-\frac{s^2+t^2}{\epsilon^2} + \frac{5/12s^2 + 1/3st + 5/12t^2}{\epsilon}\right)$$

$$g_B^2 = g^2(1 + \frac{g^2}{3!\epsilon})$$



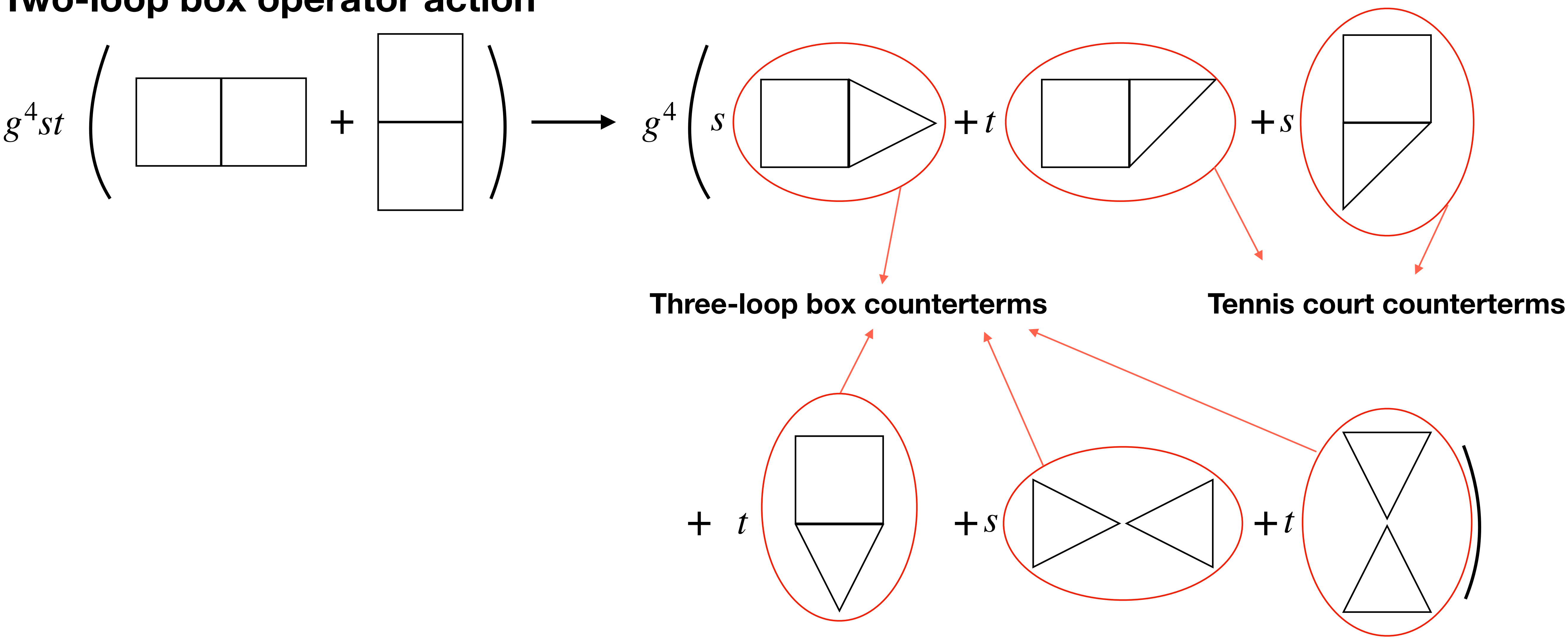
this is operator action!



compare with R-operation

Kinematically dependent renormalization

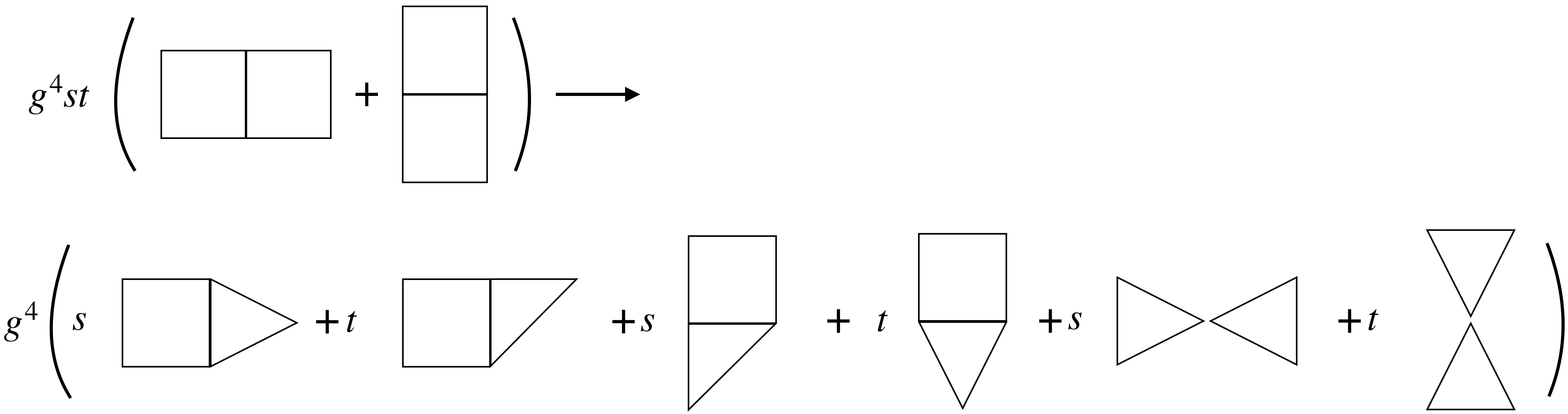
Two-loop box operator action



Z-operator reproduces R-operation like in renormalizable theories

Kinematically dependent renormalization

Two-loop box operator action



Three-loop box counterterms

Tennis court counterterms

Z-operator reproduces R-operation like in renormalizable theories

Kinematically dependent renormalization

renormalizable theories

scheme dependence

$$g^2 = zg'^2, \quad z = 1 + g'^2 c_1 + g'^4 c_2 + \ldots$$

infinite number of free parameters lead to a single multiplication constant -> redefinition of a single coupling

nonrenormalizable theories

scheme dependence

$$g^2 = zg'^2, \quad z = 1 + g'^2 st c_1 + g'^4 st(s^2 + t^2) c_2 + \ldots$$

infinite number of free parameters lead to a single multiplication constant acting as an operator -> redefinition of a series of couplings

Conclusions

- **The structure of UV divergences in non-renormalizable theories essentially copies that of renormalizable ones**

- **The main difference is that the renormalization constant Z depends on kinematics and acts like an operator rather than simple multiplication**

- **As a result, one can construct the higher derivative theory that gives the finite scattering amplitudes with a single arbitrary coupling g defined in PT within the given renormalization scheme.**

- **Transition to another scheme is performed by the action on the amplitude of a finite renormalization operator z that depends on kinematics.**

- **Assuming that one accepts these arguments, there is still a problem that at each order of PT the amplitude increases with energy, thus violating unitarity. However, apparently, this problem has to be addressed after summation of the whole PT series.**