

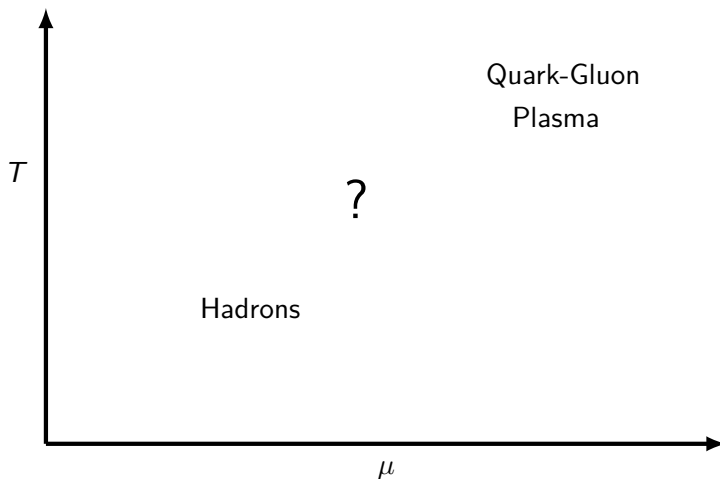
# Towards the QCD phase diagram using Complex Langevin

Benjamin Jäger

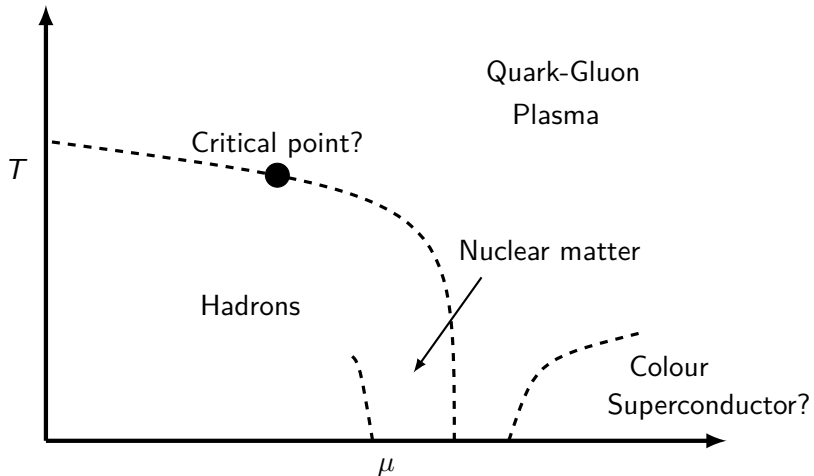
**ETH** zürich

In collaboration with G. Aarts, F. Attanasio, D. Sexty

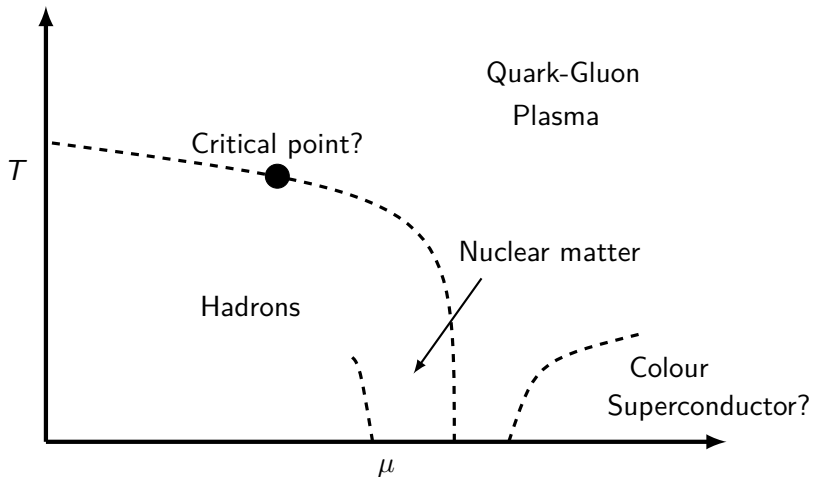
# Phase diagram for QCD



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- Sign problem  $\leftrightarrow \det D(\mu \neq 0) \in \mathbb{C}$

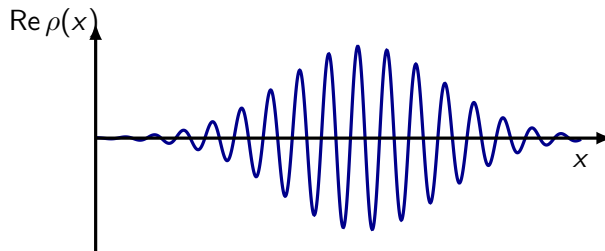
# The sign problem

- Finite chemical potential  $\rightarrow$  Sign problem

$$(\det D(\mu))^* = \det D(-\mu^*) \rightarrow \det D(\mu \neq 0) \in \mathbb{C}.$$

- Importance Sampling based Monte Carlo methods fail

$$\langle A \rangle = \frac{1}{Z} \int DU A(U) |\det D| e^{i\Theta} e^{-S_G(U)}$$



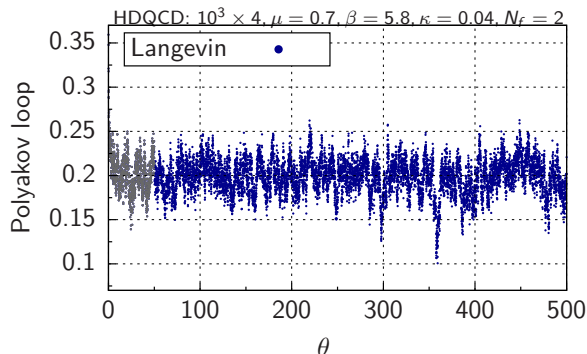
## Complex Langevin simulations

- Complexify degrees of freedom  $SU(3) \rightarrow SL(3, \mathbb{C})$

$$U_{x,\mu} = \exp \left[ i a \lambda^c \left( A_{x,\mu}^c + i B_{x,\mu}^c \right) \right]$$

- Evolve links according (1st order) Langevin equation

$$U_{x,\mu}(\theta + \varepsilon) = \exp \left[ i \lambda^a \left( -\varepsilon D_{x,\mu}^a S + \sqrt{\varepsilon} \eta_{x,\mu}^a \right) \right] U_{x,\mu}(\theta)$$



# Complex Langevin simulations

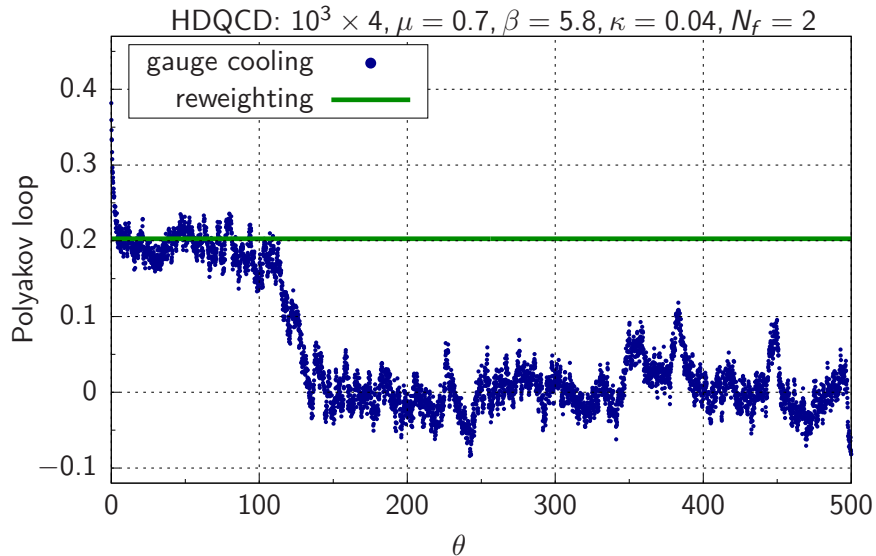
- However,  $SL(3, \mathbb{C})$  is **not** a compact group. . .
- Convergence  $\Leftrightarrow$ 
  - Action  $S$  is holomorphic
  - "Imaginary" direction of  $SL(3, \mathbb{C})$  falls off quickly enough
- Measure distance to  $SU(3)$  manifold

$$\text{unitnorm} = \text{Tr} \left( U_{x,\mu} U_{x,\mu}^\dagger - 1 \right)^2$$

- Gauge cooling is essential, but **sometimes** not sufficient. . .

$$U_{x,\mu} \rightarrow \Omega_x U_{x,\mu} \Omega_{x+\mu}^{-1}$$

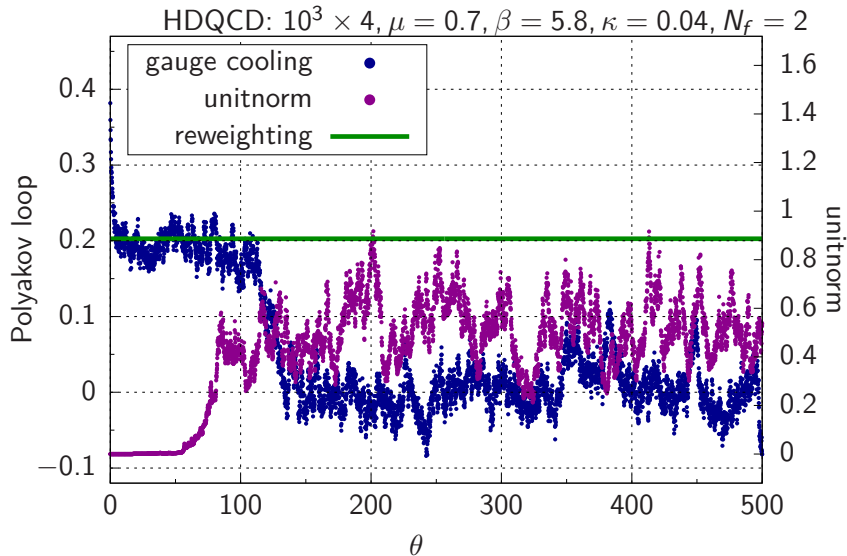
# Gauge cooling



- Tunneling to wrong results.



# Gauge cooling



- Tunneling to wrong results, when unitnorm grows too large.

## Dynamic stabilization

- Adding a trivial force to the Langevin dynamics

$$U_{x,\nu}(\theta + \varepsilon) = \exp \left[ i\lambda^a \left( \varepsilon K_{x,\nu}^a + i\varepsilon \alpha_{DS} M_x^a + \sqrt{\varepsilon} \eta_{x,\nu}^a \right) \right] U_{x,\nu}(\theta)$$

where

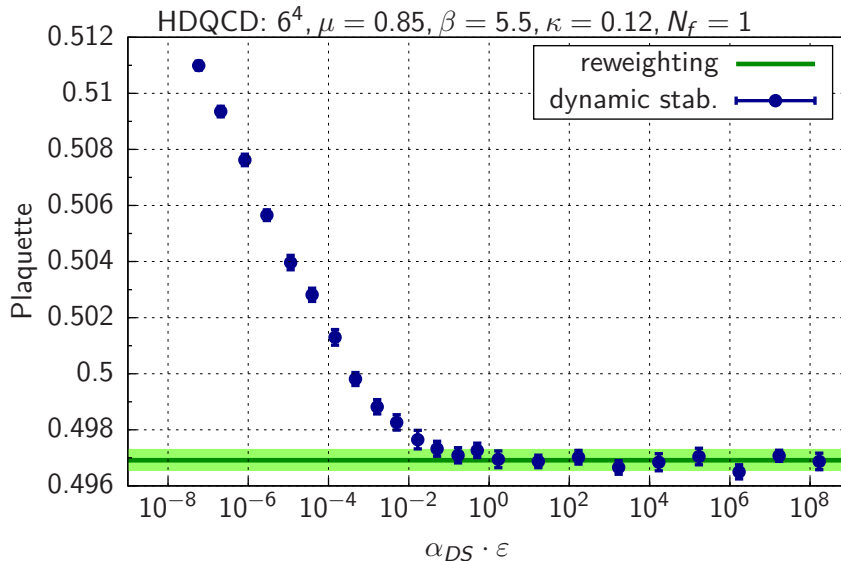
$$M_x^a = i b_x^a \left( \sum_c b_x^c b_x^c \right)^3 \text{ and } b_x^a = \text{Tr} \left[ \lambda^a \sum_\nu U_{x,\nu} U_{x,\nu}^\dagger \right].$$

- Expanding the force in terms of gauge fields  $A$  and  $B$

$$M_x^a \sim a^7 \left( \bar{B}_y^c \bar{B}_y^c \right)^3 \bar{B}_x^a + \mathcal{O}(a^8).$$

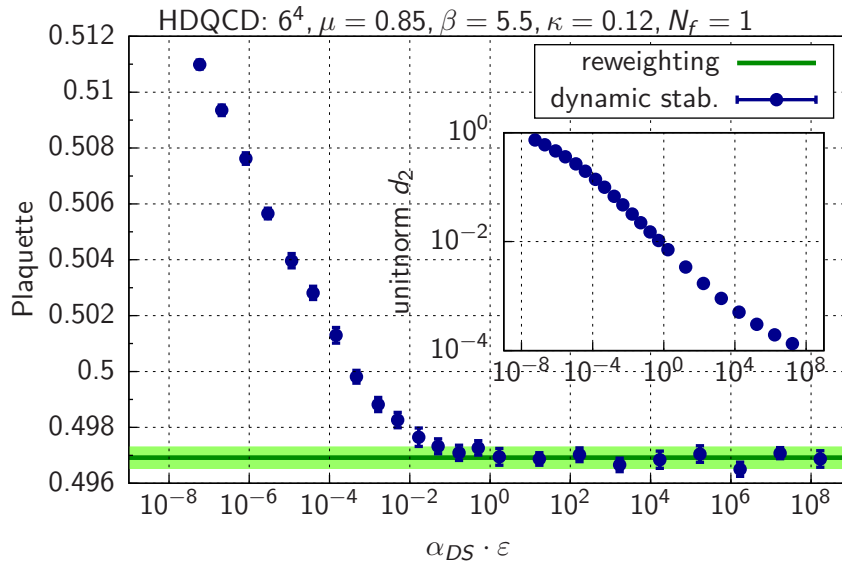
- Dynamic stabilization is numerically cheap and can be combined with gauge cooling (Here: 1 step)

# Dynamic stabilization



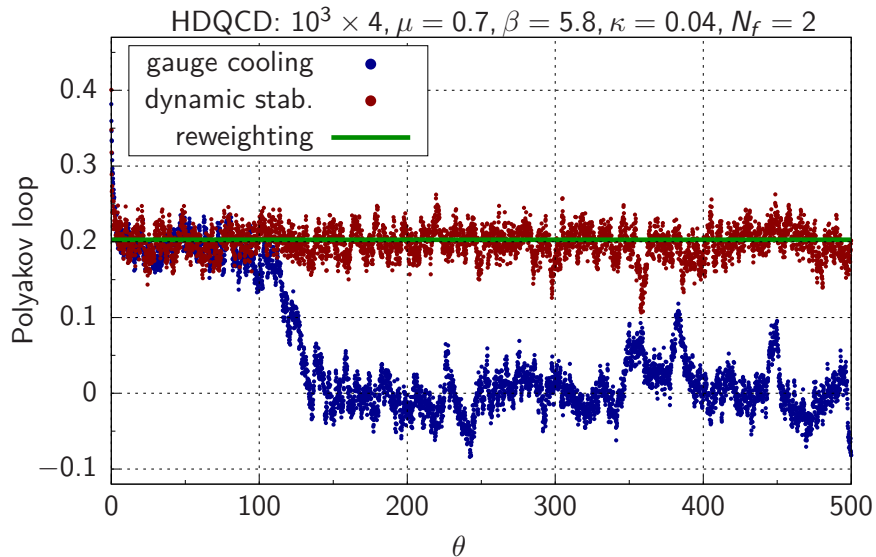
- For sufficient large  $\alpha_{DS}$  we find agreement with reweighting.

# Dynamic stabilization



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# Dynamic stabilization

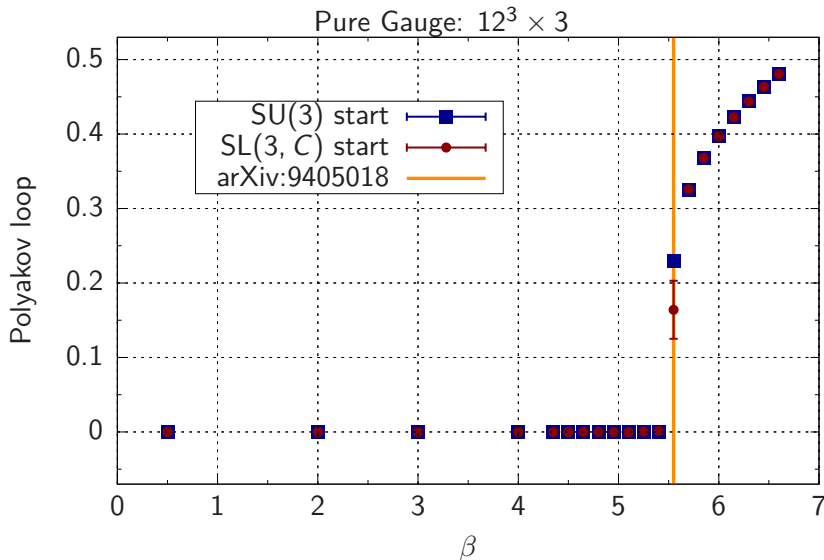


- Improved stability using dynamic stabilization

# Results

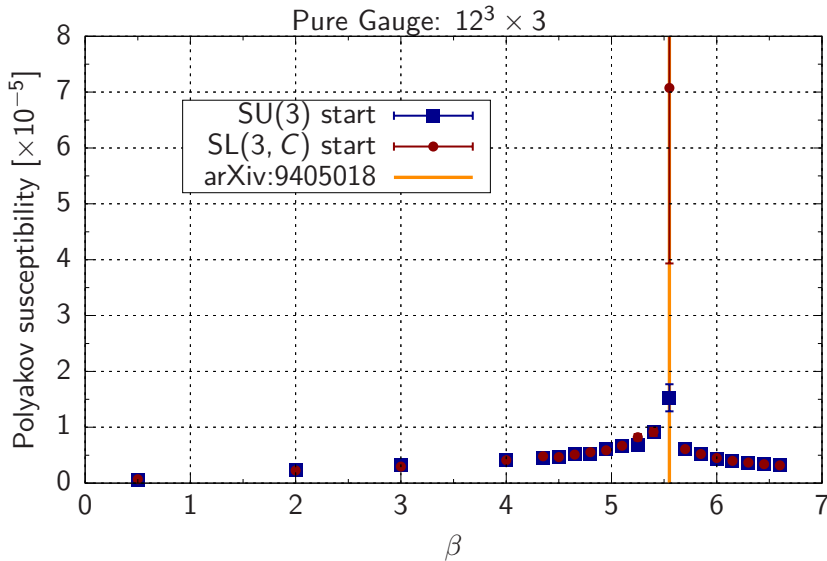
- **Pure Yang-Mills**  $\Leftrightarrow$  No sign problem
  - Checking Langevin simulations
  - Compare to standard Hybrid Monte Carlo methods
- **Heavy dense QCD (HDQCD)**
  - Heavy and dense approximation of QCD
  - Has a sign problem and phase transitions
  - Numerical cheap :)
- **Full (Staggered) QCD**
  - Including light dynamical fermions
  - Very preliminary results

# Pure Yang-Mills



- The correct transition is obtained, even for  $SL(3, \mathbb{C})$  start.

# Pure Yang-Mills



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# Heavy Dense QCD

- Here: QCD in the limit of **heavy quarks** (HDQCD).
- Fermion determinant simplifies with the Polyakov loops  $P_{\vec{x}}$  and  $P_{\vec{x}}^{-1}$  as

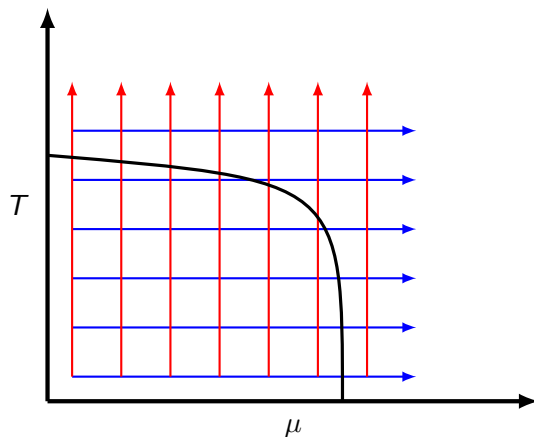
$$\det D(\mu) = \prod_{\vec{x}} \det(1 + C P_{\vec{x}})^2 \det(1 + C' P_{\vec{x}}^{-1})^2,$$

where the Polyakov loop  $P_{\vec{x}}$  is defined as

$$P_{\vec{x}} = \frac{1}{V} \sum_t U_0(\vec{x}, t)$$

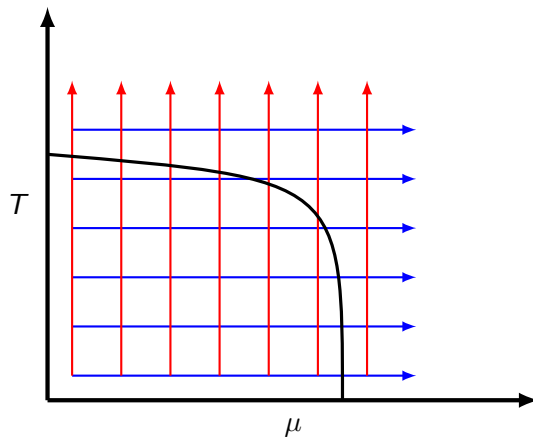
- For the gluonic part, we use the full Wilson gauge action.
- Map out the **phase diagram** for HDQCD.
- Expected transition:  $\mu_c^0 = -\ln(2\kappa)$

# Heavy Dense QCD



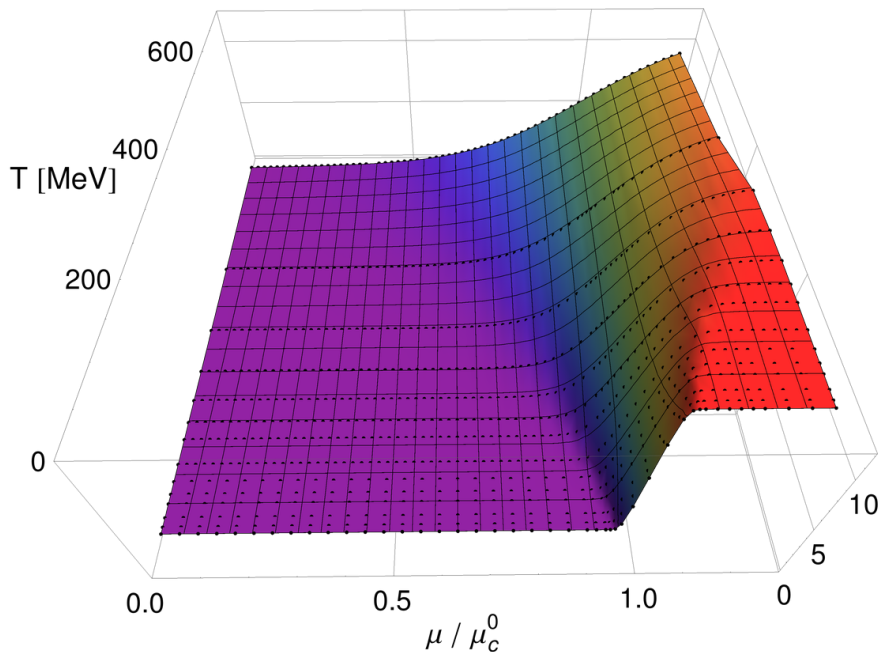
- Strategy:
  - Determine  $\mu$ -transition in Fermion density
  - Determine  $T$ -transition in Polyakov loop

# Heavy Dense QCD



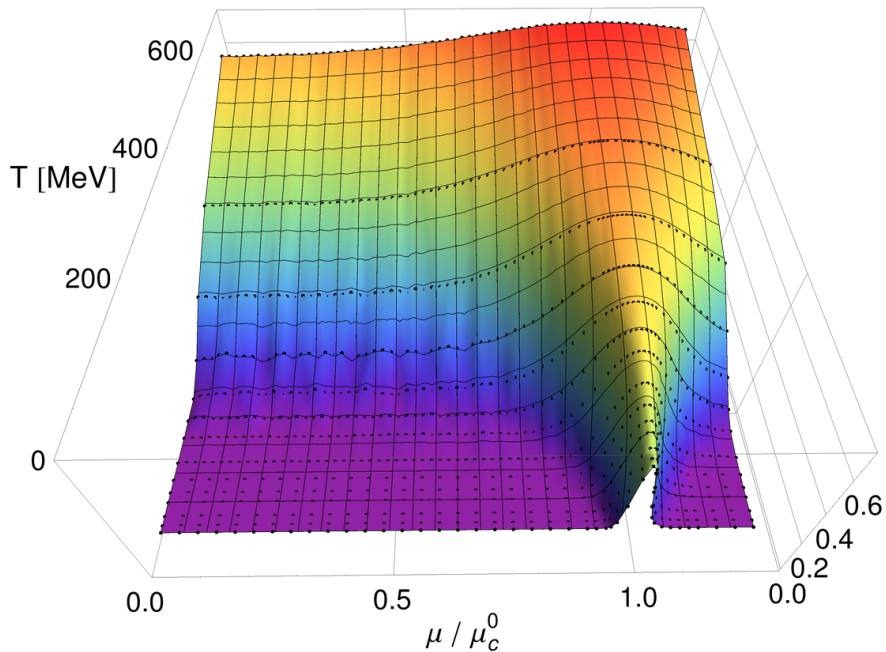
$\beta = 5.8$	$V = 6^3, 8^3, 10^3$	$a \sim 0.15 \text{ fm}$
$\kappa = 0.04$	$N_f = 2$	$\mu_c^0 = 2.53$
$N_\tau$	28 - 2	
$T \text{ [MeV]}$	48 - 671	

# Heavy Dense QCD



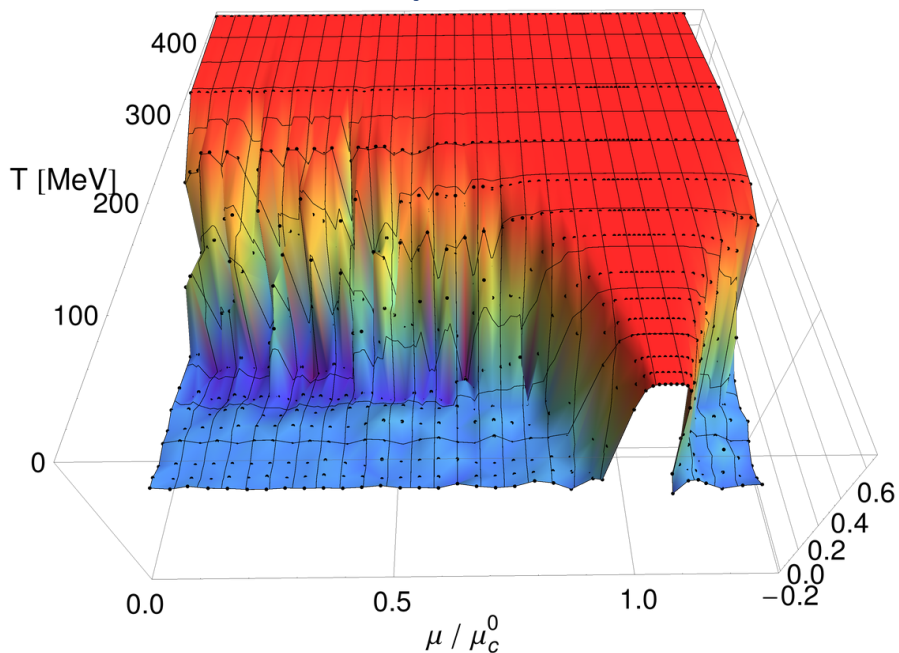
- Fermion density:  $n = \frac{1}{N_t} \frac{\partial \ln Z}{\partial \mu}$

# Heavy Dense QCD



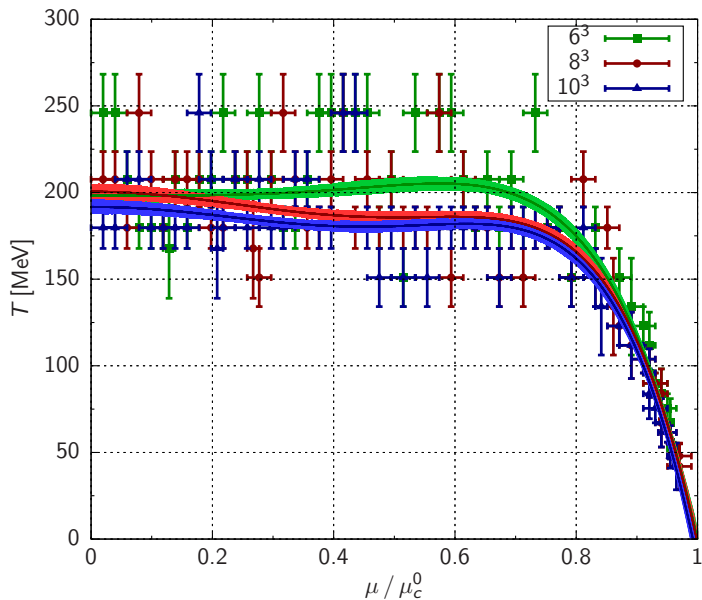
- Polyakov loop

# Heavy Dense QCD



- Binder cumulant of the Polyakov loop  $B = 1 - \frac{\langle P^4 \rangle}{3 \langle P^2 \rangle^2}$

# Heavy Dense QCD



- Fit the phase boundary using  $T_c(\mu) = \sum_k b_k (1 - \mu/\mu_c)^k$

## Full (Staggered) QCD

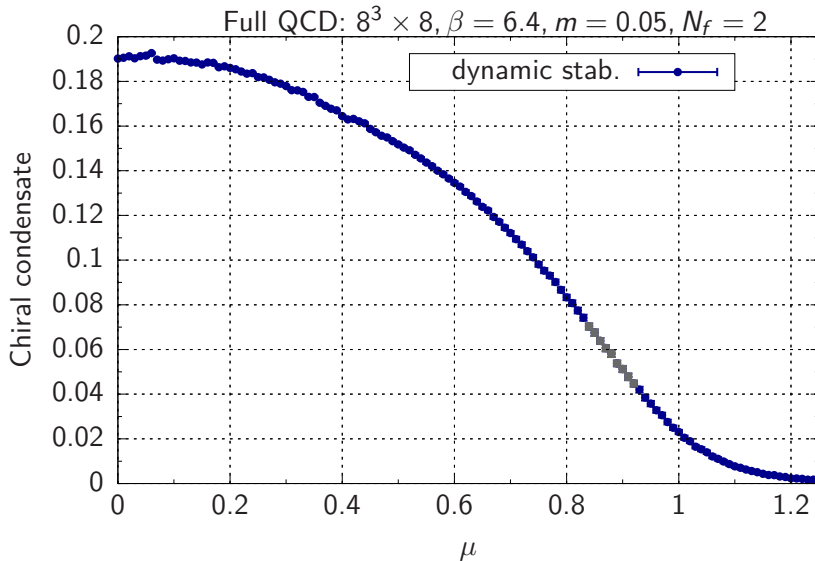
- Here: QCD using unimproved Staggered quarks
- Large gauge coupling  $\beta = 6.4$
- Small volume  $V = 8^3$
- (Small) pion masses  $\sim \mathcal{O}(700)$  MeV
- Rough inversion of Dirac operator  $M \sim 10^{-3}$

$$D_{x,\mu}^a S = D_{x,\mu}^a S_{YM} - \frac{N_f}{4} \text{Tr} \left[ M^{-1} D_{x,\mu}^a M \right]$$

- Every Langevin update needs inversion of the Dirac operator...

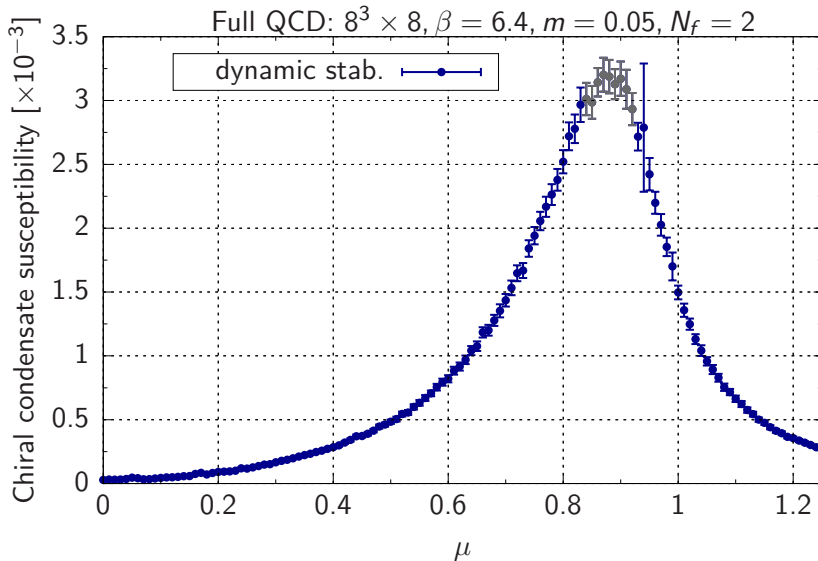


# Full (Staggered) QCD



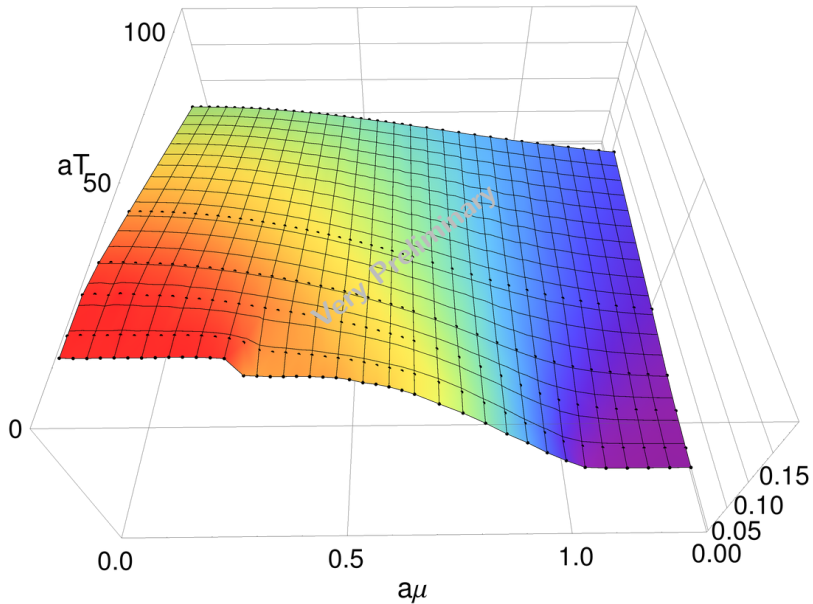
- Unimproved Staggered quarks

# Full (Staggered) QCD



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# Full (Staggered) QCD

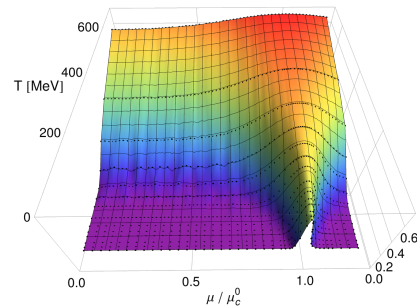


- Unimproved Staggered quarks

## Future work

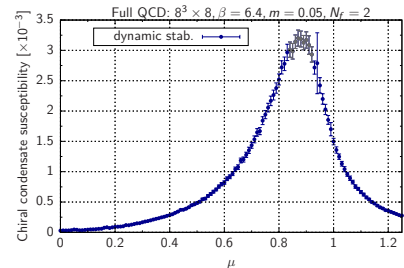
### Conclusion

- Complex Langevin simulation can be used to study the QCD phase diagram
- Dynamical stabilisation improves convergence
- Work on the convergence, especially around  $\mu_c$ .



### Future work

- Start proper Full QCD simulations to identify phase structure of QCD.
- A lot of work to be done!



Thank you for your attention!