

# Anisotropic flow fluctuations in heavy-ion collisions

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11.04.2017



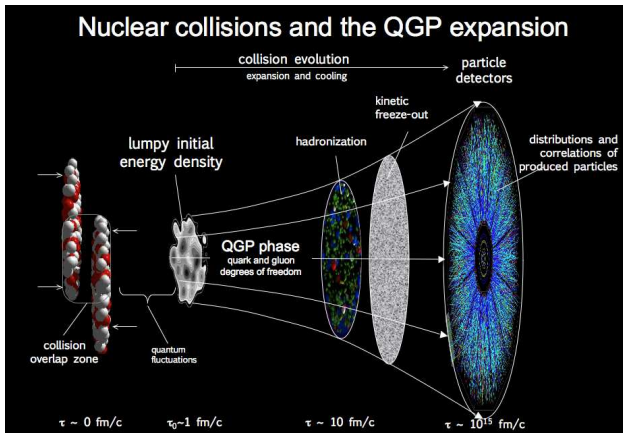
## Structure

- Introduction
- Motivation
- Analysis
- Results
- Conclusion

# Introduction



# Introduction: Evolution of HI collision



P. Sorensen, <http://arxiv.org/abs/arXiv:0905.0174>

Fig.1. Stages of evolution in non-central nuclear collision.

# Introduction: Flow

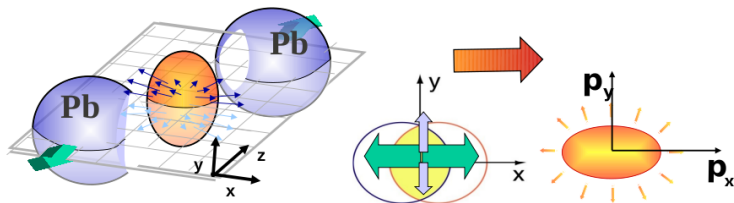


Fig.2. Formation of the overlap region:  
asymmetry in the initial geometry  $\rightarrow$  anisotropy in particle momenta distributions

Azimuthal anisotropy can be described by Fourier expansion of particles' angular distribution around the beam direction:

$$\frac{dN}{d\varphi} \propto 1 + 2 \sum_{n=1}^{\infty} v_n \cos(n(\varphi - \Psi_n))$$

where  $v_n$  and  $\varphi - \Psi_n$  represent magnitude and phase of the  $n^{\text{th}}$  - order anisotropy of a given event in the momentum space.

## Motivation



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- Azimuthal anisotropy is a powerful probe for collective properties of sub-nuclear matter created in heavy ion collision



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- Event-by-event (EbyE) analysis: studied **mean values**  $\implies$  can study **EbyE distributions** (like  $p(v_n)$ ). Why?
  - Observation of non-zero  $v_3$  at RHIC and LHC<sup>1</sup>  $\implies$  Participant eccentricity fluctuations ( $v_n \rightarrow p(v_n)$ )

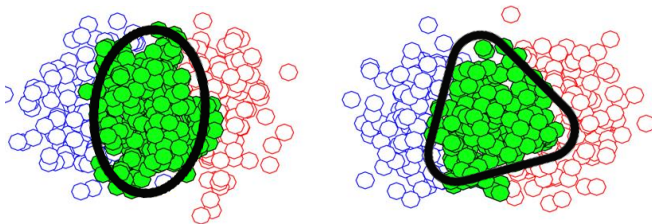


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- Multiparticle cumulants  $\implies$  moments of  $p(v_n)^2$ :
  - Observed  $v_2\{2\} > v_2\{4\} \approx v_2\{6\} \approx v_2\{8\}$ <sup>3</sup>:
  - Nature of fluctuations is Gaussian?

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<sup>4</sup>ATLAS Collaboration, JHEP 1311 (2013) 183



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  - Nature of fluctuations is Gaussian?
- Measurement of  $\rho(v_n)$  using EbyE unfolding<sup>4</sup>
  - Allows precise fluctuation studies
  - Cumulant extraction out of EbyE distributions
  - $\rho(\varepsilon_n)$  extraction

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- fluctuations of initial system distribution in the overlap area (dynamical fluctuations)

The method of EbyE unfolding analysis discussed here helps to get rid of the first two,  $\rightarrow$  leaves pure flow fluctuations connected with initial geometry of created system.



# Motivation: Skewness

- Splitting of higher-order cumulants was observed<sup>5</sup>

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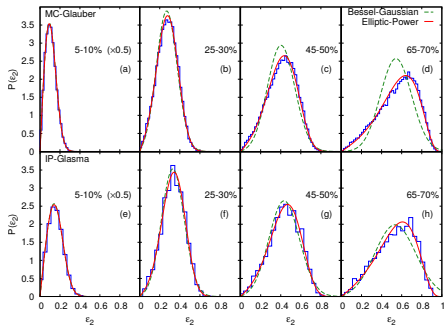
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# Motivation: Skewness

- Splitting of higher-order cumulants was observed<sup>5</sup>
- Can be explained by skewness of  $p(\varepsilon_2)$ <sup>6</sup>
- Suggests that initial state fluctuations are non-Gaussian



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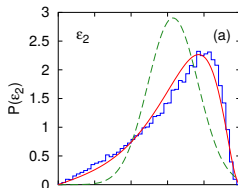
# Motivation: Extraction of $p(\varepsilon_n)$

Two common eccentricity parametrizations:

- Bessel-Gaussian<sup>7</sup>

$$p(\varepsilon_n|\varepsilon_0, \delta) = \frac{\varepsilon_n}{\delta^2} \text{Exp}\left[-\frac{\varepsilon_n^2 + \varepsilon_0^2}{2\delta}\right] I_0\left(\frac{\varepsilon_n \varepsilon_0}{\delta^2}\right)$$

$$\varepsilon_2 > \varepsilon_4 = \varepsilon_6 = \varepsilon_8$$



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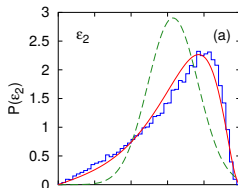
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- Elliptic Power Law<sup>8</sup>

$$p(\varepsilon_n|\varepsilon_0, \alpha) = \frac{2\alpha\varepsilon_n}{\pi} (1 - \varepsilon_0^2)^{\alpha+1/2} \int_0^\pi \frac{(1 - \varepsilon_n^2)^{\alpha-1} d\varphi}{(1 - \varepsilon_0 \varepsilon_n \cos \varphi)^{2\alpha+1}}$$

$$|\varepsilon_n| \leq 1, \varepsilon_2 > \varepsilon_4 > \varepsilon_6 > \varepsilon_8$$



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- data:

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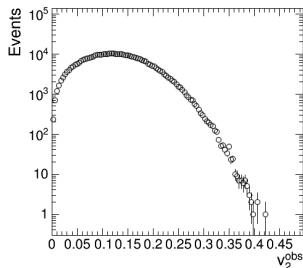
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- Construct response matrix

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# Analysis: Algorithm

- Construct response matrix
  - Data-driven approach (directly use the smearing function)
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- Construct response matrix
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- Construct unfolding matrix

$$M_{ij}^{iter} = \frac{A_{ji}c_i^{iter}}{\sum_{m,k} A_{mi}A_{jk}c_k^{iter}},$$
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- Perform D'Agostini Unfolding<sup>9</sup> (using RooUnfold package<sup>10</sup>)

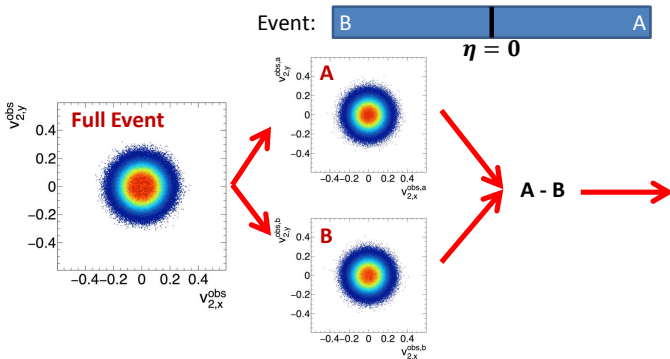
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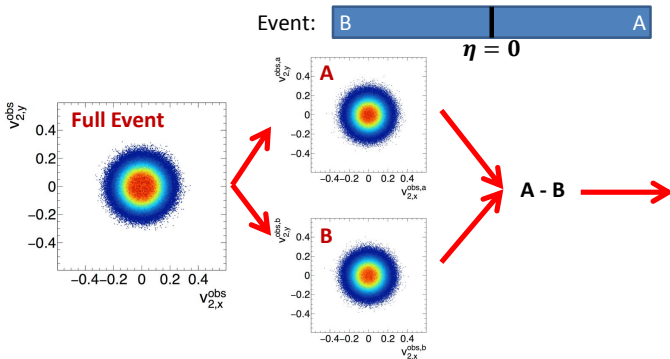
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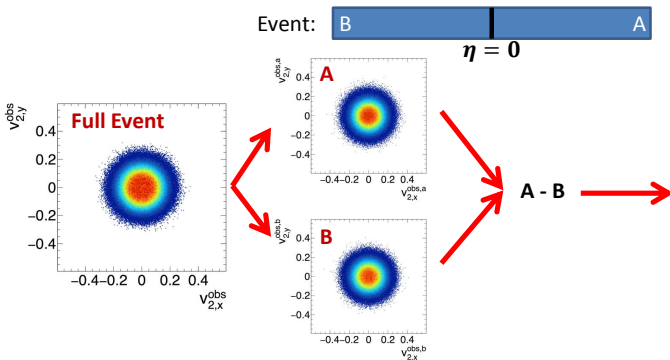


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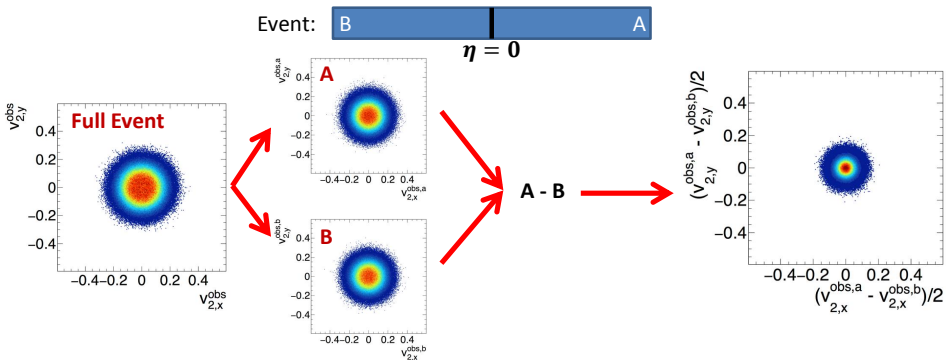
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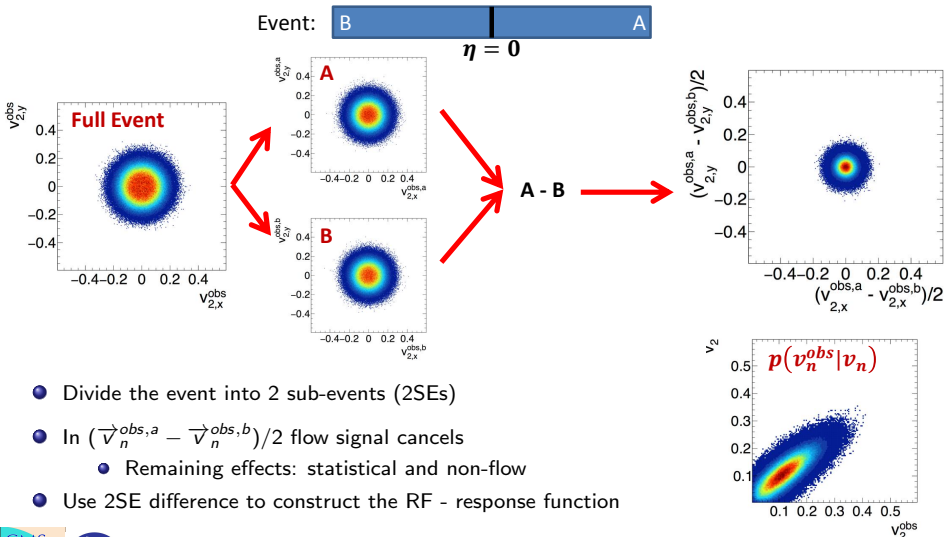


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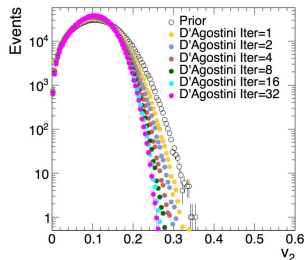


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Figure by James Castle, QM2017

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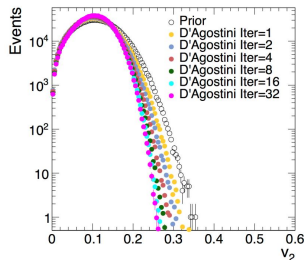
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- Need to choose "final distribution"



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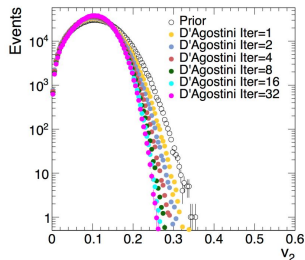
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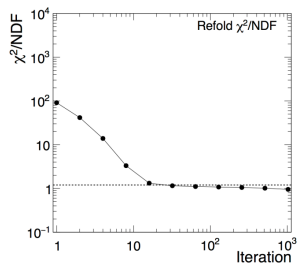
- Regularization:

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- Regularization:
  - Refolding
  - Cut-off criteria:  $\chi^2/NDF$



## Results



- Together with my colleagues from MSU and UiO performed the EbyE Unfolding analysis on HYDJET++<sup>11</sup>
- HYDJET++<sup>12</sup> is the event generator to simulate relativistic HI collisions as a superposition of two independent components:
  - The soft component is hydro-type state with preset freeze-out conditions
  - The hard state results from the in-medium multi-parton fragmentation with taking into account jet quenching effect

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<sup>11</sup>L.V. Bravina et al., Eur.Phys.J. C75 (2015) 588

<sup>12</sup>I.P. Lokhtin et al., Comp. Phys. Commun., 779 (2009)



## The direction and strength of the elliptic flow ( $v_2$ ):

- $\epsilon(b)$  - the spatial anisotropy - represents the elliptic modulation of the final freeze-out hyper-surface at a given impact parameter  $b$
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- $\epsilon(b)$  - the spatial anisotropy - represents the elliptic modulation of the final freeze-out hyper-surface at a given impact parameter  $b$
- $\delta(b)$  - the momentum anisotropy - deals with the modulation of flow velocity profile
- the relation between  $\delta(b)$  and  $\epsilon(b)$  has the form:

$$\delta = \frac{\sqrt{1 + 4B(\epsilon + B)} - 1}{2B}, \quad B = C(1 - \epsilon^2)\epsilon, \quad \epsilon = k\epsilon_0$$

# Modification of HYDJET++

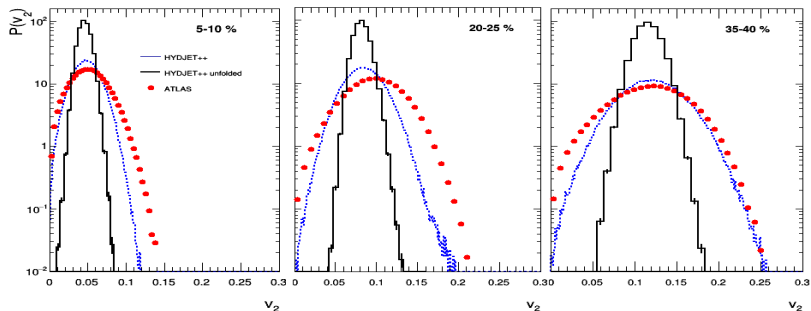


Fig. 3. EbyE  $p(v_2)$  distributions in original HYDJET++

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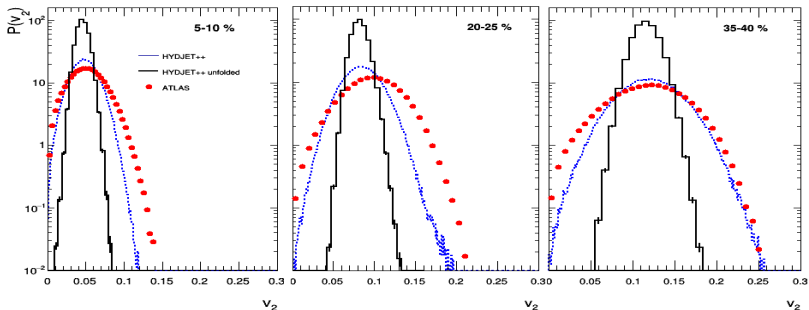


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Introducing additional “eccentricity” fluctuations in HYDJET++ model

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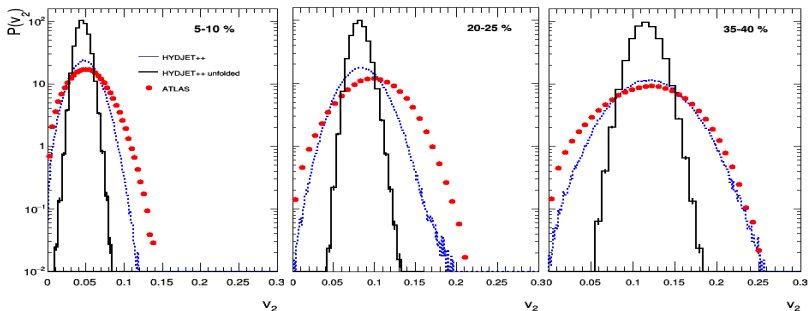


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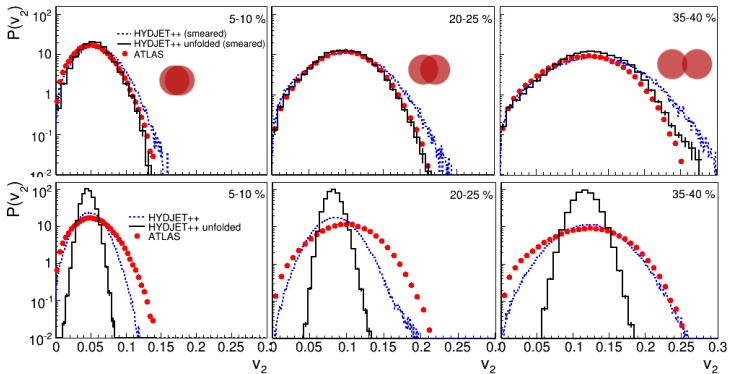
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The modification has been done by introducing EbyE Gaussian smearing of spatial anisotropy parameters  $\epsilon(b)$  and  $\epsilon_3(b)$ .

# Results: EbyE $p(v_2)$ in HYDJET++



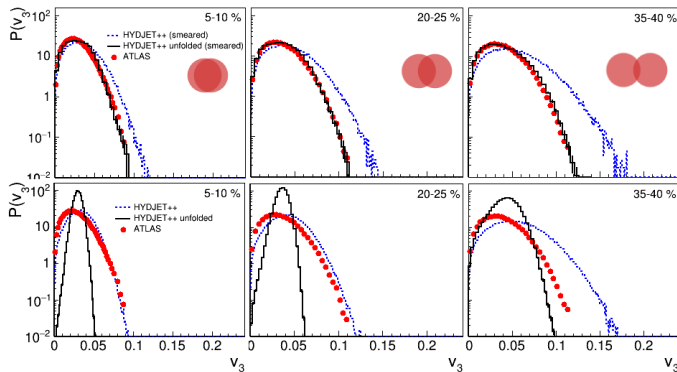
Eur.Phys.J. C75 (2015) 588

Fig.4. Comparison of  $p(v_2)$  in 5 – 10% (left), 20 – 25% (middle), 35 – 40% (right) centrality intervals in HYDJET++ with ATLAS data

Top/bottom row shows model results with/without additional smearing of spatial anisotropy parameters. ATLAS data are shown by full circles.



# Results: EbyE $P(v_3)$ in HYDJET++

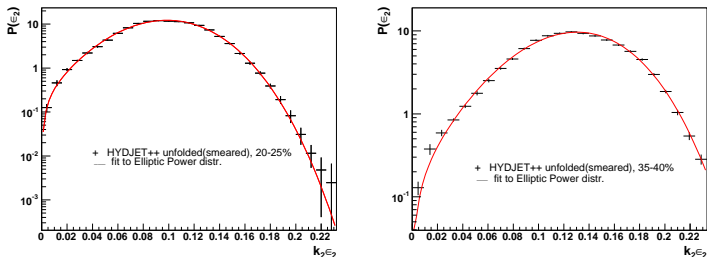


Eur.Phys.J. C75 (2015) 588

Fig.5. Comparison of  $p(v_3)$  in 5 – 10% (left), 20 – 25% (middle), 35 – 40% (right) centrality intervals in HYDJET++ with ATLAS data.

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# Results: EbyE $p(\varepsilon_2)$ in HYDJET++



Eur.Phys.J. C75 (2015) 588

Fig.6. The probability densities  $p(\varepsilon_n)$  obtained in HYDJET++ (with smearing and unfolding) of PbPb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV.

Left/right panel shows model results for centrality 20-25 % and 35-40 %.  
Full lines represent the fits to the Elliptic Power distribution

# Results: Cumulant ratios in CMS

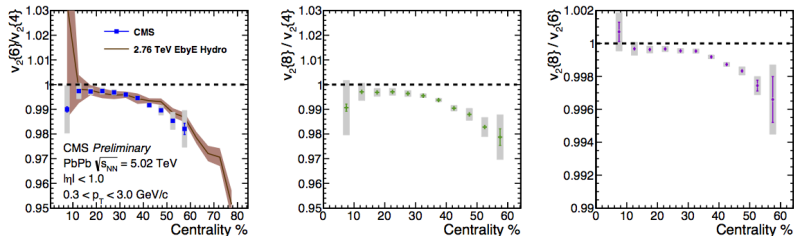


Fig. 7. Cumulant ratios

Figure by James Castle, QM2017



# Results: Cumulant ratios and skewness in CMS

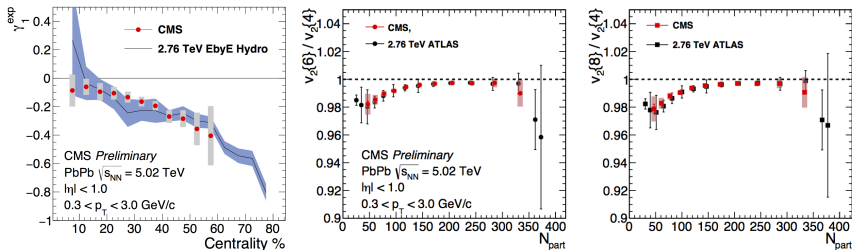


Fig. 8. Cumulant ratios and skewness

Figure by James Castle, QM2017

# Results: Fitting EbyE distributions in CMS

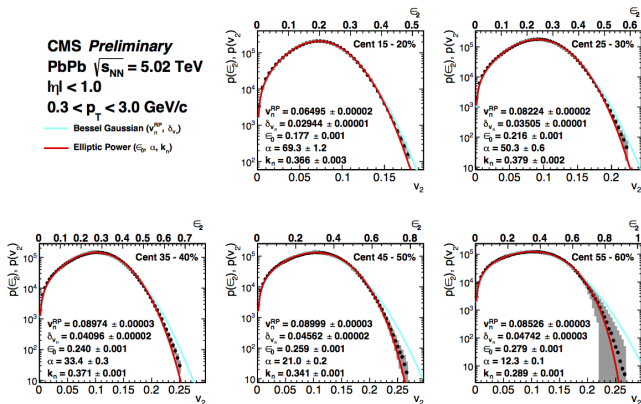


Fig. 9. Fitting with Elliptic Power and Bessel-Gaussian Parametrizations

Figure by James Castle, QM2017

## Conclusion



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- Using the EbyE unfolding with HI event generators  $\implies$  look into model's physics



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  - New set of parameters to explore:  $\langle v_n \rangle, \sigma_{v_n}, \langle v_n \rangle / \sigma_{v_n}$
  - Probe the initial state (extract  $p(\varepsilon_n)$ , skewness, ESE)
- Using the EbyE unfolding with HI event generators  $\implies$  look into model's physics
- Possibility of EbyE Unfolding for NICA energies?

Thank you for your attention!



# Back Up: Response Matrix

Constructing the Response Matrix:

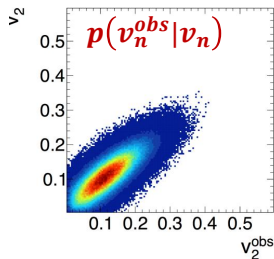
- Data-driven approach



# Back Up: Response Matrix

## Constructing the Response Matrix:

- Data-driven approach
- Prior  $p(\vec{v}_n)$  using  $p(\vec{v}_n^{obs})$  and  $p((\vec{v}_n^{obs,a} - \vec{v}_n^{obs,b})/2) = p(\vec{s}_n)$
- Sample a random  $\vec{v}_n$  from the prior
- Sample a random  $\vec{s}_n$  from the smearing function
- $\vec{v}_n^{obs} = \vec{v}_n + \vec{s}_n$
- Fill the response matrix with  $v_n$  and  $v_n^{obs}$



Constructing the Response Matrix:

- Analytical approach



## Constructing the Response Matrix:

- Analytical approach

- Using Bessel-Gaussian function:

$$p(v_n^{obs}|v_n^{true}) \approx \frac{v_n^{obs}}{\delta_{v_n}^2} \text{Exp}\left[-\frac{(v_n^{obs})^2 + (v_n^{true})^2}{2\delta_{v_n}}\right] I_0\left(\frac{v_n^{obs} v_n^{true}}{\delta_{v_n}^2}\right)$$

- Using Student's T function:  $p(v_n^{obs}|v_n^{true}, \delta_{v_n}, \nu) \approx$

$$v_n^{obs} \oint \left[1 + \frac{(v_n^{obs})^2 + (v_n^{true})^2 - 2v_n^{obs} v_n^{true} \cos \phi}{\nu \delta_{v_n}^2}\right]^{-\frac{\nu-1}{2}} d\phi$$



# Back-up: Possibilities of EbyE

- Extract cumulants from  $p(v_n)$  directly<sup>1</sup>:

$$v_n\{2\}^2 \equiv \langle v_n^2 \rangle$$

$$v_n\{4\}^4 \equiv -\langle v_n^4 \rangle + 2\langle v_n^2 \rangle^2$$

$$v_n\{6\}^6 \equiv (\langle v_n^6 \rangle - 9\langle v_n^4 \rangle\langle v_n^2 \rangle + 12\langle v_n^2 \rangle^3)/4$$

Where

$$\langle v_n^{2k} \rangle \equiv \int v_n^{2k} p(v_n) dv_n$$

- Measure skewness of  $p(v_n)$ :

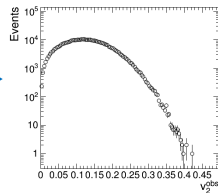
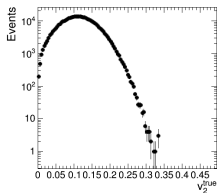
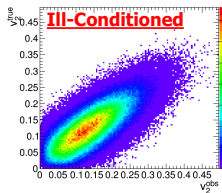
$$\gamma_1^{exp} = -6\sqrt{2}v_2\{4\}^4 \frac{v_2\{4\} - v_2\{6\}}{(v_2\{2\}^2 - v_2\{4\}^2)^{3/2}}$$

- Fit fluctuation models to data
- Bessel Gaussian
  - Elliptic Power
- Couple to ESE
- Extract detailed shape correlations

<sup>1</sup>Phys.Lett. B659 (2008) 537-541



# Back-up: Smearing



$$\begin{array}{ccc}
 \left[ \begin{array}{c} P(v_2^{obs} | v_2^{true}) \end{array} \right] & \left[ \begin{array}{c} v_2^{true} \end{array} \right] & = & \left[ \begin{array}{c} v_2^{obs} \end{array} \right] \\
 n_{true} \times n_{obs} & n_{true} \times 1 & & n_{obs} \times 1
 \end{array}$$



The altered radius of the freeze-out hyper-surface in azimuthal plane:

$$R(b, \phi) = R_{ell}(b, \phi)[1 + \epsilon_3(b) \cos[3(\phi - \Psi_3)]],$$

- $\epsilon_0 = b/2R_A$  - initial ellipticity,
- $R_A$  - the nuclear radius,
- $R_{ell}$  - former transverse radius of the fireball, which reproduces the elliptic deformation.