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Lamb shift in muonic deuterium atom

R.N. Faustov¹, A.A. Krutov², A.P. Martynenko^{2,3},
R.N. Shamsutdinov³

¹Dorodnitsyn Computing Centre, Moscow, Russia,

²Samara State University, Samara, Russia,

³Samara State Aerospace University, Samara, Russia

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Outline

1. Effects of vacuum polarization in the one-photon interaction
2. Relativistic corrections with the vacuum polarization effects
3. Nuclear structure and vacuum polarization effects
4. Recoil corrections, muon self-energy and VP effects
5. Numerical results and discussion

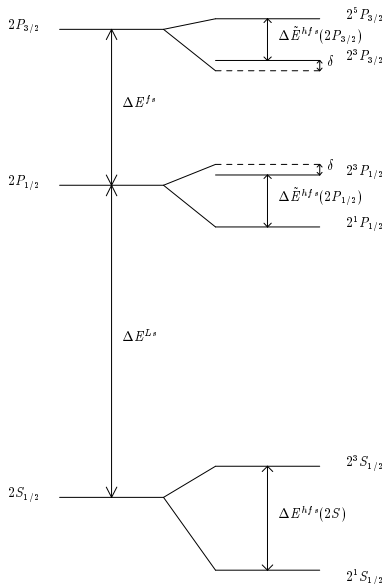


R. Pohl, A. Antognini, F. Nez et al., Nature **466**, 213 (2010).

The measurement of the (μp) Lamb shift transition energy ($2P_{3/2}^{F=2} - 2S_{1/2}^{F=1}$) at PSI (Paul Scherrer Institute) produced the result 49881.88 (76) GHz (206.2949 (32) meV).

- ▶ New value of the proton charge radius
 $r_p = 0.84184(36)(56)$ fm (experimental uncertainty of 0.76 GHz, the theoretical uncertainty 0.0049 meV).
- ▶ New value of proton radius differs by 5.0 standard deviations from the CODATA value $r_p = 0.8768(69)$ fm

The structure of S -wave and P -wave energy levels in (μp) .



The modification of the Coulomb potential due to the vacuum polarization in the coordinate representation

$$\frac{1}{k^2} \rightarrow \frac{\alpha}{3\pi} \int_1^\infty \rho(\xi) d\xi \frac{1}{k^2 + 4m_e^2 \xi^2}.$$

$$V_{VP}^C(r) = \frac{\alpha}{3\pi} \int_1^\infty d\xi \rho(\xi) \left(-\frac{Z\alpha}{r} e^{-2m_e \xi r} \right), \quad \rho(\xi) = \frac{\sqrt{\xi^2 - 1}(2\xi^2 + 1)}{\xi^4}$$

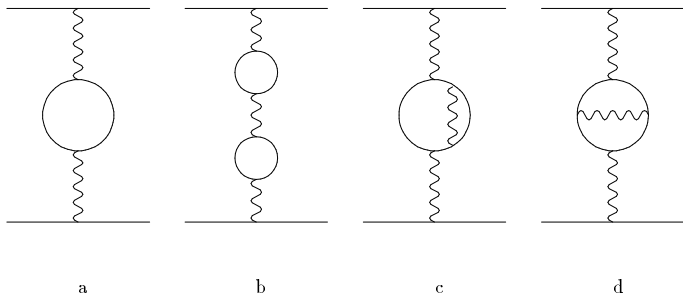


Figure: Effects of one-loop and two-loop vacuum polarization in the one-photon interaction.

$$\Delta E_{VP}(2S) = -\frac{\mu(Z\alpha)^2\alpha}{6\pi} \int_1^\infty \rho(\xi) d\xi \int_0^\infty x dx \left(1 - \frac{x}{2}\right)^2 e^{-x\left(1+\frac{2m_e\xi}{W}\right)} =$$

$$= -245.3194 \text{ meV}.$$

$$\Delta E_{VP}(2P) = -\frac{\mu(Z\alpha)^2\alpha}{72\pi} \int_1^\infty \rho(\xi) d\xi \int_0^\infty x^3 dx e^{-x\left(1+\frac{2m_e\xi}{W}\right)} =$$

$$= -17.6847 \text{ meV}.$$

$$\Delta E_{VP}(2P - 2S) = 227.6347 \text{ meV},$$

$$V_{VP-VP}^C(r) = \frac{\alpha^2}{9\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \left(-\frac{Z\alpha}{r}\right) \times$$

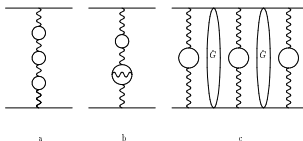
$$\times \frac{1}{(\xi^2 - \eta^2)} (\xi^2 e^{-2m_e\xi r} - \eta^2 e^{-2m_e\eta r})$$

$$\Delta V_{2-loop VP}^C = -\frac{2Z\alpha}{3r} \left(\frac{\alpha}{\pi}\right)^2 \int_0^1 \frac{f(v)dv}{(1-v^2)} e^{-\frac{2m_e r}{\sqrt{1-v^2}}},$$

$$\Delta E_{2-loop VP}(2P - 2S) = 1.3704 \text{ meV}.$$

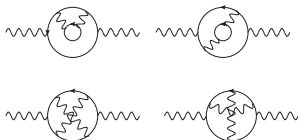
$$\Delta E_{VP-VP}(2P - 2S) = 0.2956 \text{ meV}.$$

Three loop VP contribution



$$\Delta E_{VP-VP-VP}(2P-2S) = 0.0005 \text{ meV},$$

$$\Delta E_{VP-2-loop VP}(2P-2S) = 0.0034 \text{ meV},$$



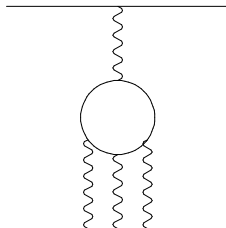
T. Kinoshita and M. Nio, PRL **82**, 3240 (1999); PRL **103**, 079901 (2009) (E)



S.G. Karshenboim et al., PRA **81**, 060501 (2010).

Our estimation of 3-loop VP contribution to the Lamb shift in (μd) is 0.0021 meV. Total 3-loop contribution in 1γ -interaction is 0.0060 meV

The Wichmann-Kroll correction



The particle interaction potential can be written in this case in the integral form:

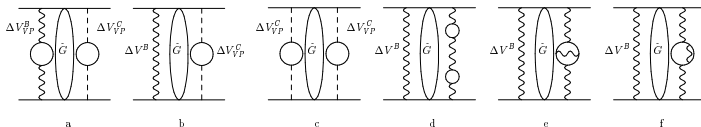
$$\Delta V^{WK}(r) = \frac{\alpha(Z\alpha)^3}{\pi r} \int_0^\infty \frac{d\zeta}{\zeta^4} e^{-2m_e\zeta r} \times$$

$$\times \left[-\frac{\pi^2}{12} \sqrt{\zeta^2 - 1} \theta(\zeta - 1) + \int_0^\zeta dx \sqrt{\zeta^2 - x^2} f^{WK}(x) \right].$$

The contribution to the Lamb shift:

$$\Delta E^{WK}(2P - 2S) = -0.0011 \text{ meV}.$$

The one-loop vacuum polarization corrections in the Breit interaction



$$\Delta V_{VP}^B(r) = \frac{\alpha}{3\pi} \int_1^\infty \rho(\xi) d\xi \sum_{i=1}^4 \Delta V_{i,VP}^B(r),$$

$$\Delta V_{1,VP}^B = \frac{Z\alpha}{8} \left(\frac{1}{m_1^2} + \frac{\delta_I}{m_2^2} \right) \left[4\pi\delta(\mathbf{r}) - \frac{4m_e^2\xi^2}{r} e^{-2m_e\xi r} \right],$$

$$\Delta V_{2,VP}^B = -\frac{Z\alpha m_e^2\xi^2}{m_1 m_2 r} e^{-2m_e\xi r} (1 - m_e\xi r),$$

$$\Delta V_{3,VP}^B = -\frac{Z\alpha}{2m_1 m_2} p_i \frac{e^{-2m_e\xi r}}{r} \left[\delta_{ij} + \frac{r_i r_j}{r^2} (1 + 2m_e\xi r) \right] p_j,$$

$$\Delta V_{4,VP}^B = \frac{Z\alpha}{r^3} \left(\frac{1}{4m_1^2} + \frac{1}{2m_1 m_2} \right) e^{-2m_e\xi r} (1 + 2m_e\xi r) (\mathbf{L}\boldsymbol{\sigma}_1)$$

In first order perturbation theory the potentials $\Delta V_{i,VP}^B(r)$ give the contributions of order $\alpha(Z\alpha)^4$ to the shift $(2P - 2S)$:

$$\Delta E_{1,VP}^B(2P - 2S) = -0.0353 \text{ meV},$$

$$\Delta E_{2,VP}^B(2P - 2S) = 0.0011 \text{ meV},$$

$$\Delta E_{3,VP}^B(2P - 2S) = 0.0012 \text{ meV},$$

$$\Delta E_{4,VP}^B(2P - 2S) = -0.0023 \text{ meV}.$$

Two-loop modification of the Breit Hamiltonian

$$\Delta V_{2-loop\ VP}^B(r) = \frac{\alpha^2(Z\alpha)}{12\pi^2} \left(\frac{1}{m_1^2} + \frac{\delta_I}{m_2^2} \right) \int_0^1 \frac{f(v)dv}{1-v^2} \times$$

$$\times \left[4\pi\delta(\mathbf{r}) - \frac{4m_e^2}{(1-v^2)r} e^{-\frac{2m_e r}{\sqrt{1-v^2}}} \right].$$

Corresponding $(2P - 2S)$ shift is the following:

$$\Delta E_{2-loop\ VP}^B(2P - 2S) = -0.0002 \text{ meV}.$$

VP corrections in second order perturbation theory of orders $\alpha^2(Z\alpha)^2$ and $\alpha(Z\alpha)^4$

$$\Delta E_{SOPT}^{VP} = \langle \psi | \Delta V_{VP}^C \tilde{G} \Delta V_{VP}^C | \psi \rangle + 2 \langle \psi | \Delta V^B \tilde{G} \Delta V_{VP}^C | \psi \rangle.$$

$$\begin{aligned} \Delta E_{SOPT}^{VP, VP}(2S) &= -\frac{\mu\alpha^2(Z\alpha)^2}{72\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \times \\ &\times \int_0^\infty \left(1 - \frac{x}{2}\right) e^{-x(1+\frac{2m_e\xi}{W})} dx \int_0^\infty \left(1 - \frac{x'}{2}\right) e^{-x'(1+\frac{2m_e\eta}{W})} dx' g_{2S}(x, x') = \\ &= -0.1750 \text{ meV}, \end{aligned}$$

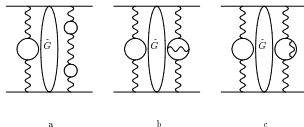
$$\begin{aligned} \Delta E_{SOPT}^{VP, VP}(2P) &= -\frac{\mu\alpha^2(Z\alpha)^2}{7776\pi^2} \int_1^\infty \rho(\xi) d\xi \int_1^\infty \rho(\eta) d\eta \times \\ &\times \int_0^\infty e^{-x(1+\frac{2m_e\xi}{W})} dx \int_0^\infty e^{-x'(1+\frac{2m_e\eta}{W})} dx' g_{2P}(x, x') = -0.0030 \text{ meV}, \end{aligned}$$

VP corrections in second order perturbation theory of orders $\alpha^2(Z\alpha)^2$ and $\alpha(Z\alpha)^4$

$$\Delta E_{SOPT}^{B,VP}(2P - 2S) = 0.0530 \text{ meV}.$$

$$\Delta E_{SOPT}^{VP,VP;\Delta V^B}(2P - 2S) = 0.0004 \text{ meV}.$$

$$\Delta E_{SOPT}^{VP,\Delta V_{VP}^B}(2P - 2S) = -0.00001 \text{ meV}.$$



$$\Delta E_{SOPT}^{VP-VP,VP}(2S) = -\frac{\mu\alpha^3(Z\alpha)^2}{108\pi^3} \int_1^\infty \rho(\xi)d\xi \int_1^\infty \rho(\eta)d\eta \int_1^\infty \rho(\zeta)d\zeta \int_0^\infty dx \left(1 - \frac{x}{2}\right) \times$$

$$\int_0^\infty dx' \left(1 - \frac{x'}{2}\right) e^{-x'(1+\frac{2m_e\zeta}{W})} \frac{1}{\xi^2 - \eta^2} \left[\xi^2 e^{-x(1+\frac{2m_e\xi}{W})} - \eta^2 e^{-x(1+\frac{2m_e\eta}{W})} \right] g_{2S}(x, x') = -0.0007 \text{ meV},$$

$$\Delta E_{SOPT}^{2-loop VP,VP}(2S) = -\frac{\mu\alpha^3(Z\alpha)^2}{18\pi^3} \int_0^1 \frac{f(v)dv}{1-v^2} \int_1^\infty \rho(\xi)d\xi \times$$

$$\times \int_0^\infty dx \left(1 - \frac{x}{2}\right) e^{-x(1+\frac{2m_e}{\sqrt{1-v^2}W})} \int_0^\infty dx' \left(1 - \frac{x'}{2}\right) e^{-x'(1+\frac{2m_e\xi}{W})} g_{2S}(x, x') = -0.0018 \text{ meV},$$

VP correction in third order perturbation theory

$$\delta E_{TOPT}^{(1)} = \langle \psi_n | \Delta V_{VP} \tilde{G}_n \delta V_{VP} \tilde{G}_n \delta V_{VP} | \psi_n \rangle$$

Another type of third order correction (V.G. Ivanov et al. PRD 80, 027702 (2009)):

$$\delta E_{TOPT}^{(2)} = \langle \psi_n | \Delta V_{VP} | \psi_n \rangle \langle \psi_n | \delta V_{VP} \tilde{G}_n \tilde{G}_n \delta V_{VP} | \psi_n \rangle$$

$$\Delta E_{TOPT}^{VP, VP, VP}(2P - 2S) = 0.0001 \text{ meV}$$

We present the charge radius corrections at two values of r_d :
 $r_d = 2.1424(21)$ fm (CODATA 2010) and $r_d = 2.130(9)$ fm (I. Sick):

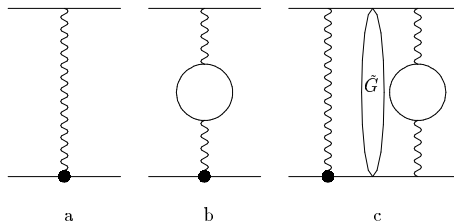


Figure: The leading order nuclear structure and vacuum polarization corrections.

$$\begin{aligned}\Delta E_{str}(2P - 2S) &= -\frac{\mu^3(Z\alpha)^4}{12}r_d^2 = -6.07313 \cdot r_d^2 = \\ &= -27.8749(-27.5532) \text{ meV}.\end{aligned}$$

VP and deuteron structure corrections

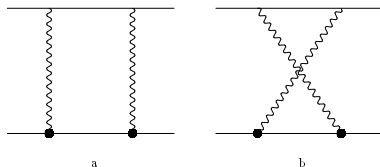
$$\Delta V_{str}^{VP}(r) = \frac{2Z\alpha^2}{9} r_d^2 \int_1^\infty \rho(\xi) d\xi \left[\delta(\mathbf{r}) - \frac{m_e^2 \xi^2}{\pi r} e^{-2m_e \xi r} \right].$$

$$\begin{aligned} \Delta E_{str}^{VP}(2P - 2S) &= -\frac{\mu^3 \alpha (Z\alpha)^4}{36\pi} r_d^2 \int_1^\infty \rho(\xi) d\xi \left[1 - \frac{16m_e^4 \xi^4}{(2m_e \xi + W)^2} \right] = \\ &= -0.01350 \cdot r_d^2 = -0.0620(-0.0612) \text{ meV}, \end{aligned}$$

$$\begin{aligned} \Delta E_{str,SOPT}^{VP}(2P - 2S) &= -\frac{\mu^3 \alpha (Z\alpha)^4}{36\pi} r_d^2 \int_1^\infty \rho(\xi) d\xi \times \\ &\times \frac{-12 + 23b_1 - 8b_1^2 - 4b_1^3 + 4b_1^4 + 4b_1(3 - 4b_1 + 2b_1^2) \ln b_1}{b_1^5} = \\ &= -0.020487 \cdot r_d^2 \text{ meV} = -0.0940(-0.0929) \text{ meV}, b_1 = 1 + \frac{2m_e}{W} \xi. \end{aligned}$$

$$\begin{aligned} \Delta E_{str}(2P - 2S) + \Delta E_{str}^{VP}(2P - 2S) + \Delta E_{str,SOPT}^{VP}(2P - 2S) &= \\ &= -6.10712 \cdot r_d^2 = -28.0309(-27.7074) \text{ meV}. \end{aligned}$$

Nuclear structure corrections of order $(Z\alpha)^5$



An investigation of elastic contribution to the Lamb shift and the deuteron polarizability contribution:



J.L. Friar, PRC **16**, 1540 (1977).



J.L. Friar, J. Martorell and D.W.L. Sprung, PRA **56**, 4579 (1997).



Y. Lu and R. Rosenfelder, PLB **319**, 7 (1993); W. Leidemann and R. Rosenfelder, PRC **51**, 427 (1995).



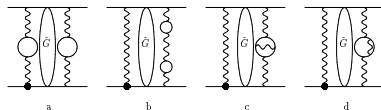
R.N. Faustov and A.P. Martynenko, PRA **67**, 052506 (2003); Phys. Atom. Nucl. **67**, 457 (2004).



K. Pachucki, PRL **106**, 193007 (2011).

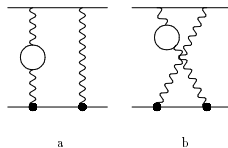
The elastic two-photon contribution is canceled by part of the deuteron polarizability contribution (K. Pachucki 2011): 1.680(16) meV

Two-loop vacuum polarization and nuclear structure corrections of order $\alpha^2(Z\alpha)^4$ in second order PT



$$\Delta E_{str,SOPT}^{VP,VP}(2P - 2S) = -9.5 \cdot 10^{-5} \cdot r_d^2 = -0.0004(0.0004) \text{ meV}.$$

Nuclear structure correction of order $\alpha(Z\alpha)^5$ from two-photon exchange diagrams with VP insertion



The muon matrix element P_{VP} for nonrelativistic two-photon exchange with an account of the vacuum polarization:

$$P_{VP} = \frac{2\alpha^3}{3\pi} \phi^2(0) \int_1^\infty \rho(\xi) d\xi \int \frac{d\mathbf{q}}{(2\pi)^3} \frac{(4\pi)^2}{q^2(q^2 + 4m_e^2\xi^2)} \frac{1}{E + \frac{q^2}{2m_1}} \left[e^{i\mathbf{q}(\mathbf{R}-\mathbf{R}')} - 1 + \frac{q^2}{6}(\mathbf{R}-\mathbf{R}')^2 \right]$$

Integrating over q and expanding the resulting expression over small parameter $\sqrt{2m_1E}|\mathbf{R} - \mathbf{R}'|$ we obtain:

$$P_{VP} = \frac{32\alpha^3}{3} m_1 \phi^2(0) |\mathbf{R} - \mathbf{R}'|^3 \int_1^\infty \rho(\xi) d\xi \left[\frac{a_\xi^3 - 3a_\xi^2 + 6a_\xi + 6e^{-a_\xi} - 6}{12a_\xi^4} - \right. \\ \left. - 2m_1 E |\mathbf{R} - \mathbf{R}'|^2 \frac{a_\xi^4 - 4a_\xi^3 + 12a_\xi^2 - 24a_\xi - 24e^{-a_\xi} + 24}{48a_\xi^6} \right], \quad a_\xi = 2m_e \xi |\mathbf{R} - \mathbf{R}'|.$$

Expanding second term at small a_ξ we obtain $(-\pi/240a_\xi)$. Integrating over ξ and summing over excited deuteron states we obtain the contribution to the Lamb shift:

$$\delta E_{pol}^{VP}(2P - 2S) = -\frac{m_1^2 \alpha^3 \phi^2(0)}{1024m_e} \left[\frac{1}{3} \langle \phi_D | R^2 H_D R^2 | \phi_D \rangle - \frac{4}{5} \langle \phi_D | R_i H_D R^2 R_i | \phi_D \rangle + \right. \\ \left. + \frac{2}{5} \langle \phi_D | (R_i R_j - \frac{1}{3} \delta_{ij} R^2) H_D (R_i R_j - \frac{1}{3} \delta_{ij} R^2) | \phi_D \rangle \right] = -0.0001 \text{ meV},$$

ϕ_D is the deuteron wave function: (I.B. Khriplovich and A.I. Milstein, JETP **98**, 181 (2004))

$$\phi_D(r) = \sqrt{\frac{\kappa}{2\pi}} \frac{1}{r} e^{-\kappa r},$$

$\kappa = 0.0457 \text{ GeV}$ is the inverse deuteron size.

Muon-line radiative correction to the nuclear size effect of order $\alpha(Z\alpha)^5$ (M.I. Eides et al., PRA **55**, 2447 (1997)):

$$\Delta E_{str}^{\alpha(Z\alpha)^5}(2P-2S) = 1.985 \frac{\alpha(Z\alpha)^5 \mu^3}{8} r_d^2 = 9.62 \cdot 10^{-4} \cdot r_d^2 = 0.0044(0.0044) \text{ meV}.$$

Nuclear structure corrections of order $(Z\alpha)^6$



J.L. Friar, Ann. Phys. **122**, 151 (1979).



J.L. Friar and G.L. Payne, PRA **56**, 5173 (1997).

$$\Delta E_{str}^{(Z\alpha)^6}(2P-2S) = \frac{(Z\alpha)^6}{12} \mu^3 \left\{ r_d^2 \left[\langle \ln \mu Z \alpha r \rangle + C - \frac{3}{2} \right] - \frac{1}{2} r_d^2 + \frac{1}{3} \langle r^3 \rangle \left\langle \frac{1}{r} \right\rangle - \right. \\ \left. - I_2^{rel} - I_3^{rel} - \mu^2 F_{NR} + \frac{1}{40} \mu^2 \langle r^4 \rangle \right\} = -21.28 \cdot 10^{-4} \cdot r_d^2 + 0.0029 = -0.0069(-0.0068) \text{ meV},$$

The recoil correction of order $(Z\alpha)^4 \mu^3/m_2^2$ (U.D. Jentschura, PRA **84**, 012505 (2011)):

$$\Delta E_{rec}(2P - 2S) = \frac{\mu^3(Z\alpha)^4}{12m_2^2} = 0.0672 \text{ meV}.$$

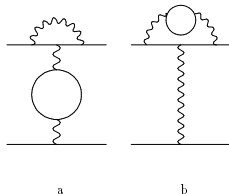


Figure: Radiative corrections with the vacuum polarization effects.

The recoil correction of fifth order in $(Z\alpha)$

$$\Delta E_{rec}^{(Z\alpha)^5} = \frac{\mu^3(Z\alpha)^5}{m_1 m_2 \pi n^3} \left[\frac{2}{3} \delta_{l0} \ln \frac{1}{Z\alpha} - \frac{8}{3} \ln k_0(n, l) - \frac{1}{9} \delta_{l0} - \frac{7}{3} a_n - \frac{2}{m_2^2 - m_1^2} \delta_{l0} (m_2^2 \ln \frac{m_1}{\mu} - m_1^2 \ln \frac{m_2}{\mu}) \right],$$

$$\ln k_0(2S) = 2.811769893120563, \ln k_0(2P) = -0.030016708630213,$$

$$a_n = -2 \left[\ln \frac{2}{n} + \left(1 + \frac{1}{2} + \dots + \frac{1}{n} \right) + 1 - \frac{1}{2n} \right] \delta_{l0} + \frac{(1 - \delta_{l0})}{l(l+1)(2l+1)}.$$

$$\Delta E_{rec}^{(Z\alpha)^5}(2P - 2S) = -0.0266 \text{ meV}.$$

The recoil correction of the sixth order in $(Z\alpha)$:

$$\Delta E_{rec}^{(Z\alpha)^6}(2P - 2S) = \frac{(Z\alpha)^6 m_1^2}{8m_2} \left(\frac{23}{6} - 4 \ln 2 \right) = 0.0001 \text{ meV}.$$

The energy contributions from radiative corrections to the lepton line, the Dirac and Pauli form factors and muon vacuum polarization:

$$\Delta E_{MVP,MSE}(2S) = \frac{\alpha(Z\alpha)^4}{8\pi} \frac{\mu^3}{m_1^2} \left[\frac{4}{3} \ln \frac{m_1}{\mu(Z\alpha)^2} - \frac{4}{3} \ln k_0(2S) + \frac{38}{45} + \frac{\alpha}{\pi} \left(-\frac{9}{4} \zeta(3) + \frac{3}{2} \pi^2 \ln 2 - \frac{10}{27} \pi^2 - \frac{2179}{648} \right) + 4\pi Z\alpha \left(\frac{427}{384} - \frac{\ln 2}{2} \right) \right] = 0.7647 \text{ meV},$$

$$\Delta E_{MVP,MSE}(2P) = \frac{\alpha(Z\alpha)^4}{8\pi} \frac{\mu^3}{m_1^2} \left[-\frac{4}{3} \ln k_0(2P) - \frac{m_1}{6\mu} - \frac{\alpha}{3\pi} \frac{m_1}{\mu} \left(\frac{3}{4} \zeta(3) - \frac{\pi^2}{2} \ln 2 + \frac{\pi^2}{12} + \frac{197}{144} \right) \right] = -0.0100 \text{ meV}.$$

Electron loop polarization insertion in the radiative photon can be expressed in terms of F'_1 and F_2

$$\Delta E_{rad+VP}(nS) = \frac{\mu^3}{m_1^2} \frac{(Z\alpha)^4}{n^3} \left[4m_1^2 F'_1(0) \delta_{l0} + F_2(0) \frac{C_{jl}}{2l+1} \right],$$

$$C_{jl} = \delta_{l0} + (1 - \delta_{l0}) \frac{[j(j+1) - l(l+1) - \frac{3}{4}] m_1}{l(l+1) \mu}.$$

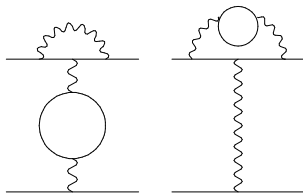
Two-loop contribution to the form factors $F'_1(0)$ and $F_2(0)$:

$$m_1^2 F'_1(0) = \left(\frac{\alpha}{\pi} \right)^2 \left[\frac{1}{9} \ln^2 \frac{m_1}{m_e} - \frac{29}{108} \ln \frac{m_1}{m_e} + \frac{1}{9} \zeta(2) + \frac{395}{1296} \right],$$

$$F_2(0) = \left(\frac{\alpha}{\pi} \right)^2 \left[\frac{1}{3} \ln \frac{m_1}{m_e} - \frac{25}{36} + \frac{\pi^2}{4} \frac{m_e}{m_1} - 4 \frac{m_e^2}{m_1^2} \ln \frac{m_1}{m_e} + 3 \frac{m_e^2}{m_1^2} \right].$$

Then the correction to the Lamb shift is equal to

$$\Delta E_{rad+VP}(2P - 2S) = -0.0018 \text{ meV}.$$



a

b

The estimation of the muon self-energy and electron vacuum polarization contribution in Fig.11(a)

$$\Delta E_{MSE}^{VP} = \frac{\alpha}{3\pi m_1^2} \ln \frac{m_1}{\mu(Z\alpha)^2} \left[\langle \psi_n | \Delta \cdot \Delta V_{VP}^C | \psi_n \rangle + 2 \langle \psi_n | \Delta V_{VP}^C \tilde{G} \Delta \left(-\frac{Z\alpha}{r} \right) | \psi_n \rangle \right].$$

The sum of all matrix elements leads to the following shift ($2P - 2S$):

$$\Delta E_{MSE}^{VP}(2P - 2S) = -0.0047 \text{ meV}.$$

TABLE I: Lamb shift ($2P_{1/2} - 2S_{1/2}$) in muonic deuterium atom.

Contribution to the splitting	$\Delta E(2P - 2S)$, meV	Equation, Reference
1	2	3
VP contribution of order $\alpha(Z\alpha)^2$ in one-photon interaction	227.6347	(6), [2]
Two-loop VP contribution of order $\alpha^2(Z\alpha)^2$ in one-photon interaction	1.6660	(9), (14), [2]
VP and MVP contribution in one-photon interaction	0.0001	(11), [2]
Three-loop VP contribution in one-photon interaction	0.0060	(17), (18), [29, 31]
The Wichmann-Kroll correction	-0.0011	(21), [2, 31]
Light-by-light contribution	0.0001	[31]
Relativistic and VP corrections of order $\alpha(Z\alpha)^4$ in first order PT	-0.0353	(27)-(30), [5]
Relativistic and two-loop VP corrections of order $\alpha^2(Z\alpha)^4$ in first order PT	-0.0002	(32)
Two-loop VP contribution of order $\alpha^2(Z\alpha)^2$ in second order PT	0.1720	(39)-(41), [31]
Relativistic and one-loop VP corrections of order $\alpha(Z\alpha)^4$ in second order PT	0.0530	(43), [5]
Relativistic and two-loop VP corrections of order $\alpha^2(Z\alpha)^4$ in second order PT	0.0004	(44)-(45)
Three-loop VP contribution in second order PT of order $\alpha^3(Z\alpha)^2$	0.0025	(46)-(47), [31]
Three-loop VP contribution in third order PT of order $\alpha^3(Z\alpha)^2$	0.0001	(48), [29, 31],
Nuclear structure contribution of order $(Z\alpha)^4$	-27.8749	(49), [2, 6]
Nuclear structure and polarizability contribution of order $(Z\alpha)^5$	1.6800	[51]
Nuclear structure and VP contribution in 1γ interaction of order $\alpha(Z\alpha)^4$	-0.0620	(51)
Nuclear structure and VP contribution in second order PT of order $\alpha(Z\alpha)^4$	-0.0940	(52)
Nuclear structure and two-loop VP contribution in 1γ interaction of order $\alpha^2(Z\alpha)^4$	-0.0005	(56)
Nuclear structure and two-loop VP contribution in second order PT of order $\alpha^2(Z\alpha)^4$	-0.0004	(57)

Table I (continued).

1	2	3
Nuclear structure and polarizability contribution of order $\alpha(Z\alpha)^5$ with VP correction	-0.0001	(60)
Nuclear structure contribution of order $\alpha(Z\alpha)^5$ with muon-line radiative correction	0.0044	(62), [61]
Nuclear structure contribution of order $(Z\alpha)^6$	-0.0069	(63), [22, 54]
Recoil correction of order $(Z\alpha)^4$	0.0672	(64), [5]
Recoil correction of order $(Z\alpha)^5$	-0.0266	(69), [2, 6, 55]
Recoil correction of order $(Z\alpha)^6$	0.0001	(70), [6]
Recoil correction to VP of order $\alpha(Z\alpha)^5$ (seagull term)	0.0002	[5]
Radiative-recoil corrections of orders $\alpha(Z\alpha)^5$, $(Z^2\alpha)(Z\alpha)^4$	-0.0026	(71), Table 9 [6]
Muon self-energy and MVP contribution	-0.7747	(72)-(73), [2, 6]
VP contribution to muon form factors $F_1^l(0)$, $F_2(0)$ of order $\alpha^2(Z\alpha)^4$	-0.0018	(78), [4, 6, 63]
VP correction to muon self-energy	-0.0047	(80), [4, 6]
HVP contribution	0.0129	[64, 65]
Total contribution	202.4139 ± 0.0573 ($r_d = 2.1424(21)$ fm) 202.7375 ± 0.2352 ($r_d = 2.130(9)$ fm)	

TABLE II: Fine structure of $2P$ -state in muonic deuterium atom.

Contribution to fine splitting ΔE^{fs}	Numerical value in meV	Equation, Reference
Contribution of order $(Z\alpha)^4$ $\frac{\mu^3(Z\alpha)^4}{32m_1^2} \left(1 + \frac{2m_1}{m_2}\right)$	8.83848	(82), [2, 6]
Muon AMM contribution $\frac{\mu^3(Z\alpha)^4}{16m_1^2} a_\mu \left(1 + \frac{m_1}{m_2}\right)$	0.01957	(82), [2, 6]
Contribution of order $(Z\alpha)^6$	0.00031	(82), [2, 6]
Contribution of order $(Z\alpha)^6 m_1/m_2$	-0.00001	(82), [2, 6]
Contribution of order $\alpha(Z\alpha)^4$ in first order PT $\langle \Delta V_{VP}^{fs} \rangle$	0.00346	(84)
Contribution of order $\alpha(Z\alpha)^4$ in second order PT $\langle \Delta V_{VP}^C \cdot \vec{G} \cdot \Delta V^{fs} \rangle$	0.00229	(85)
Contribution of order $\alpha(Z\alpha)^6$ $\frac{\alpha(Z\alpha)^6 \mu^3}{32\pi m_1^2} \left[\ln \frac{\mu(Z\alpha)^2}{m_1} + \frac{1}{5} \right]$	-0.00001	(82), [6]
VP Contribution from 1γ interaction of order $\alpha^2(Z\alpha)^4 \langle \Delta V_{P-V}^{fs} \rangle$	0.000003	(87)
VP Contribution from 1γ interaction of order $\alpha^2(Z\alpha)^4 \langle \Delta V_{2-loop,VP}^{fs} \rangle$	0.00002	(89)
VP Contribution in second order PT of order $\alpha^2(Z\alpha)^4$ $\langle \Delta V_{VP}^C \cdot \vec{G} \cdot \Delta V_{VP}^{fs} \rangle$	0.000002	Fig.4(a), $\Delta V^B \rightarrow \Delta V^{fs}$
VP Contribution in second order PT of order $\alpha^2(Z\alpha)^4$ $\langle \Delta V_{P-V}^C \cdot \vec{G} \cdot \Delta V^{fs} \rangle$	-0.000001	Fig.4(d), $\Delta V^B \rightarrow \Delta V^{fs}$
VP Contribution in second order PT of order $\alpha^2(Z\alpha)^4$ $\langle \Delta V_{2-loop,VP}^C \cdot \vec{G} \cdot \Delta V^{fs} \rangle$	0.000026	(90), Fig.4(e-f), $\Delta V^B \rightarrow \Delta V^{fs}$
Nuclear structure correction in 1γ interaction	-0.00028	(91), [22]
Summary contribution	8.86386	

