

# Горение атомных ядер:

Гидростатическое горение

FBS2014

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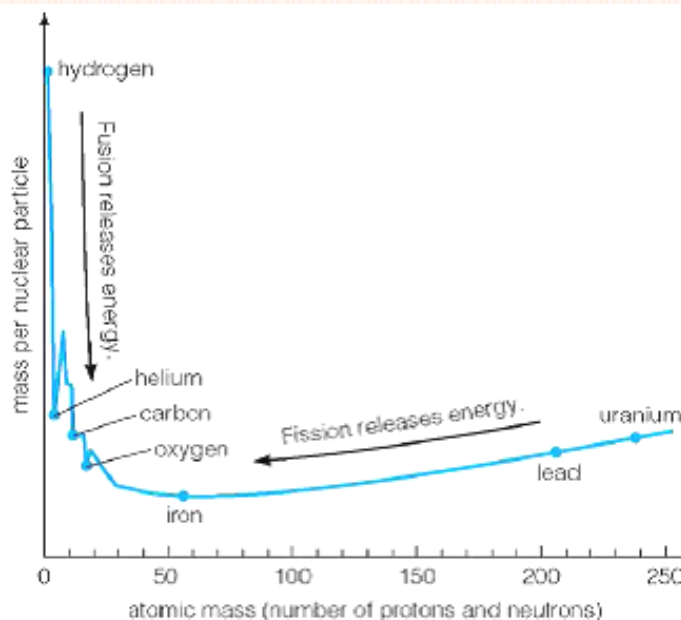
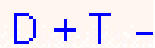
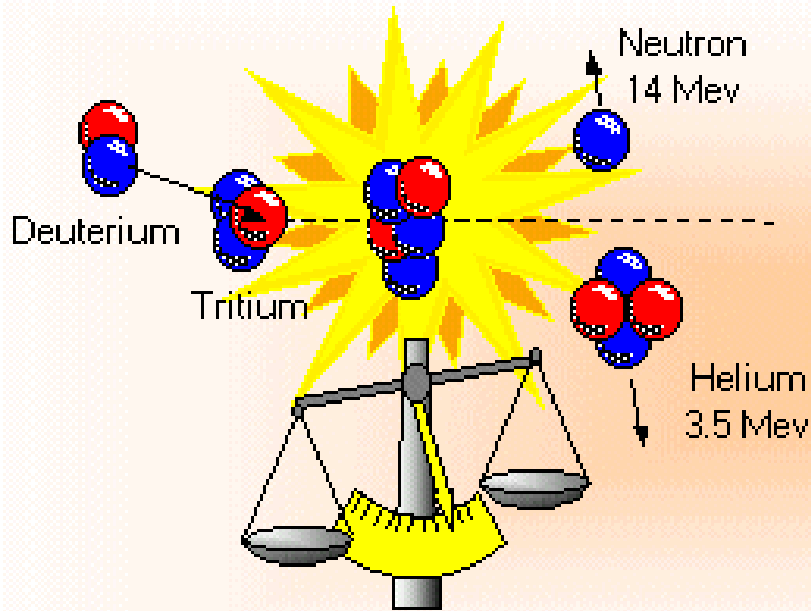
# *Energy vs Work*

## in Comparison

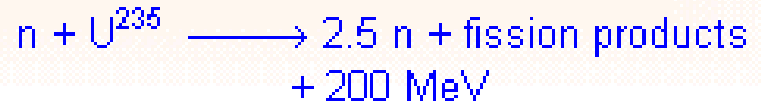
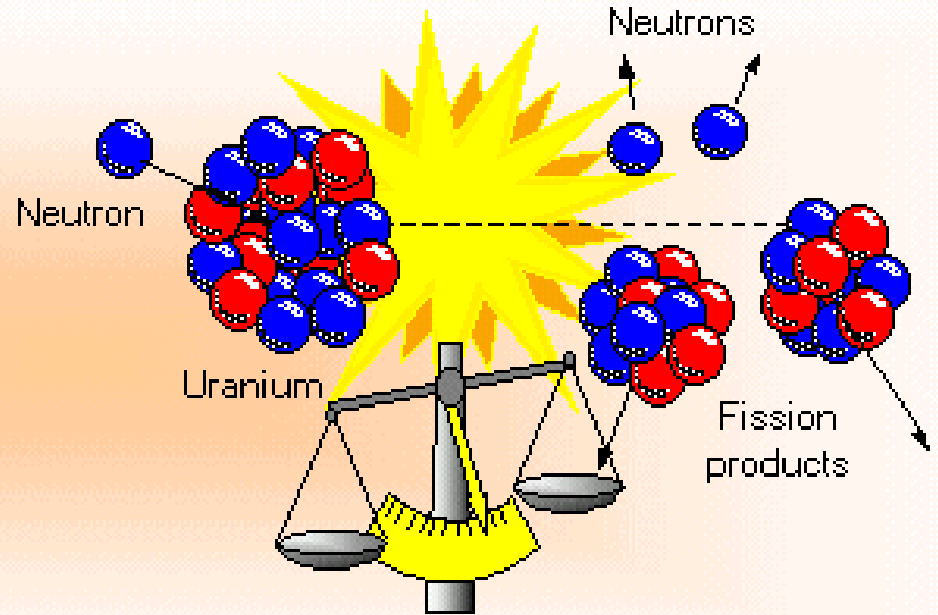
	– Energy (joules)
• Average daytime solar energy on Earth/m <sup>2</sup> sec	1.3 10 <sup>3</sup>
• Energy from metabolism of a candy bar	1 10 <sup>6</sup>
• Energy spend at 1 hour walking (adult)	1 10 <sup>6</sup>
• Kinetic energy of a car at 100 km/hr	1 10 <sup>6</sup>
• Daily energy needs of an adult	1 10 <sup>7</sup>
• Energy released at burning 1 liter oil	1.2 10 <sup>7</sup>
• Energy released at fission of 1kg <sup>235</sup> U	5.6 10 <sup>13</sup>
• Energy from fusion of hydrogen in 1l water	7 10 <sup>13</sup>
• Energy released by 1 megaton H-bomb	5 10 <sup>15</sup>
• Energy at major earthquake (magnitude 8)	2.5 10 <sup>16</sup>
• Annual energy generation by Sun	10 <sup>34</sup>
• Energy released at supernova (star explosion)	10 <sup>44</sup> --10 <sup>46</sup>

# Nuclear Reactors

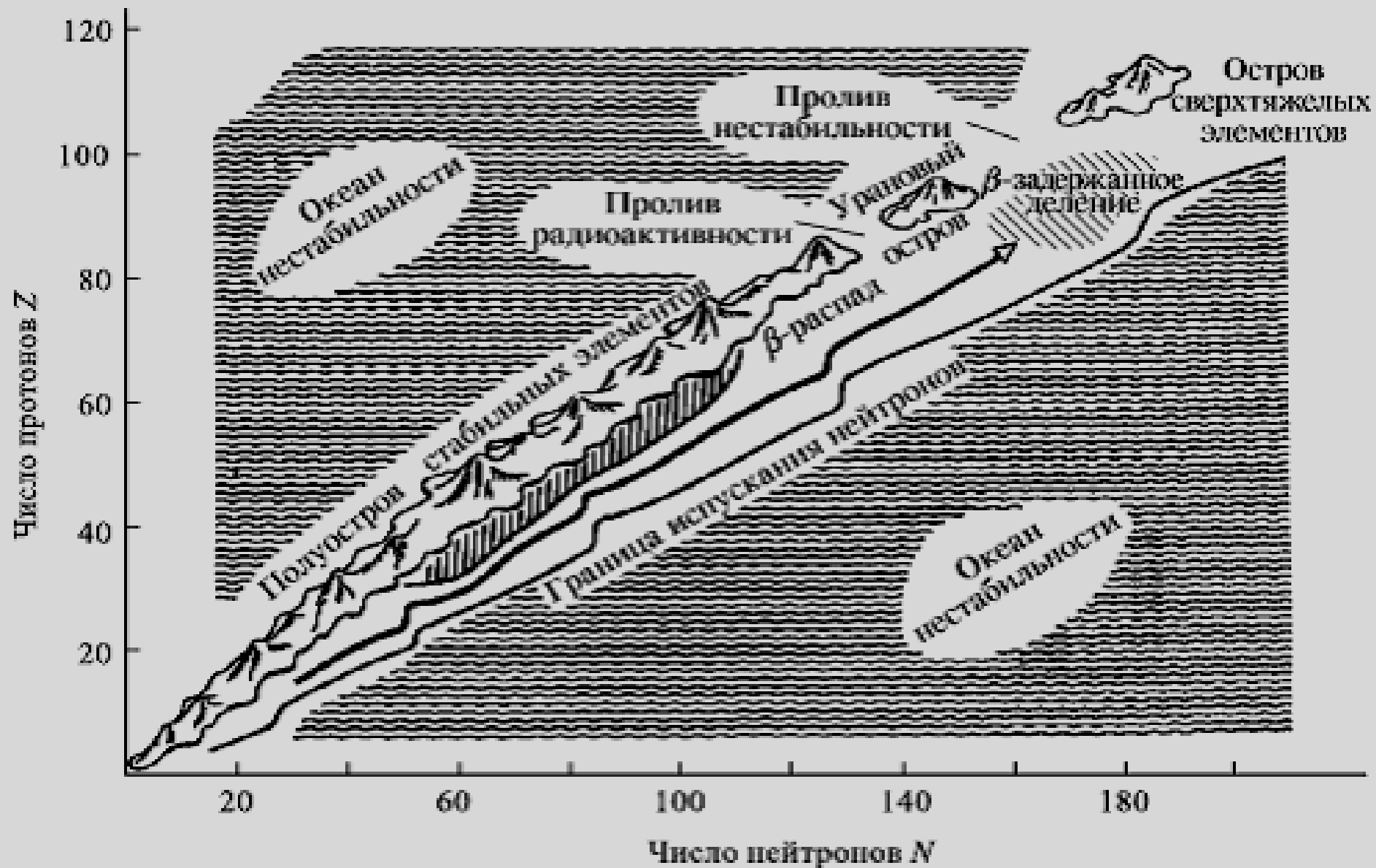
## Fusion



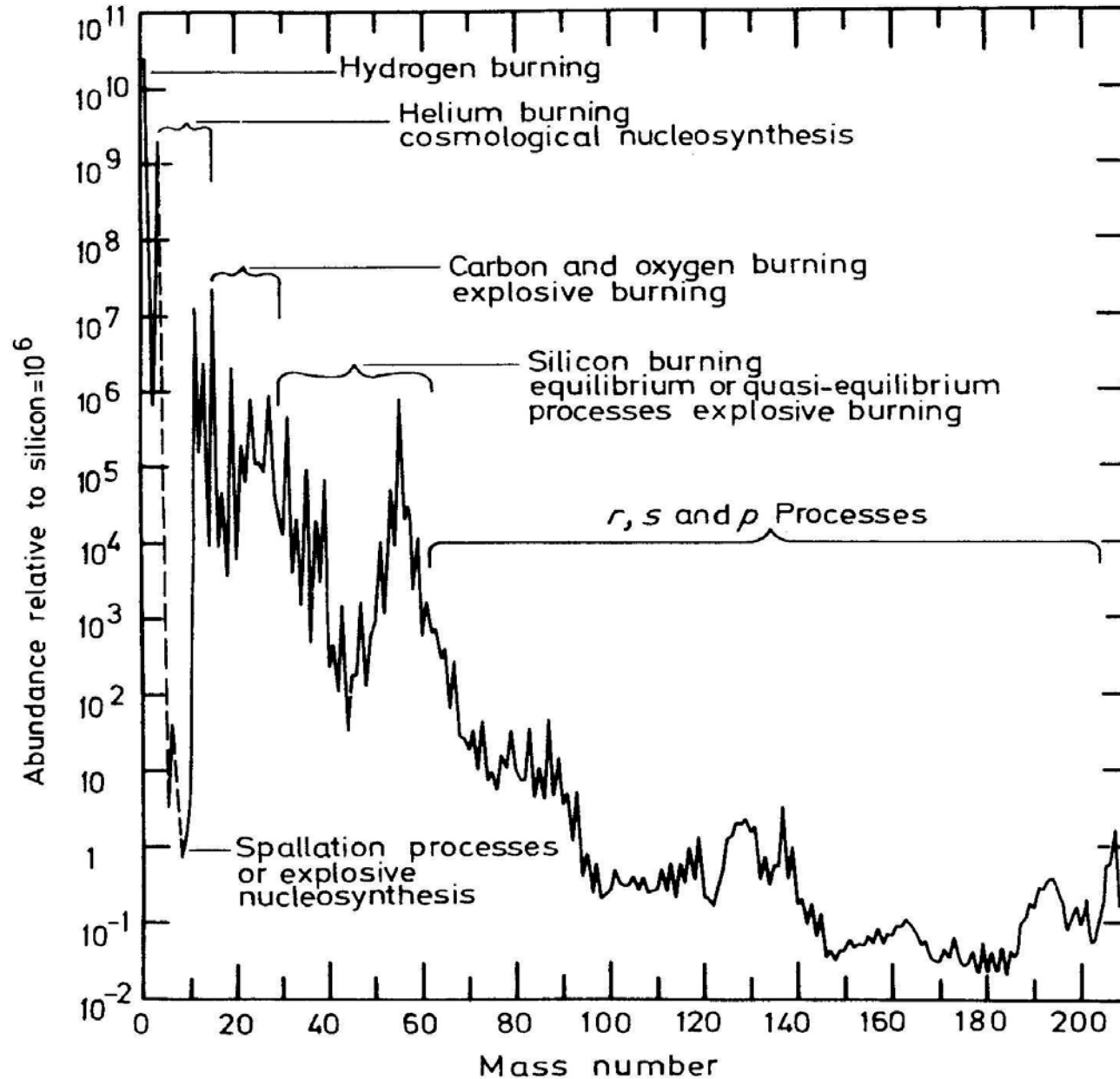
## Fission



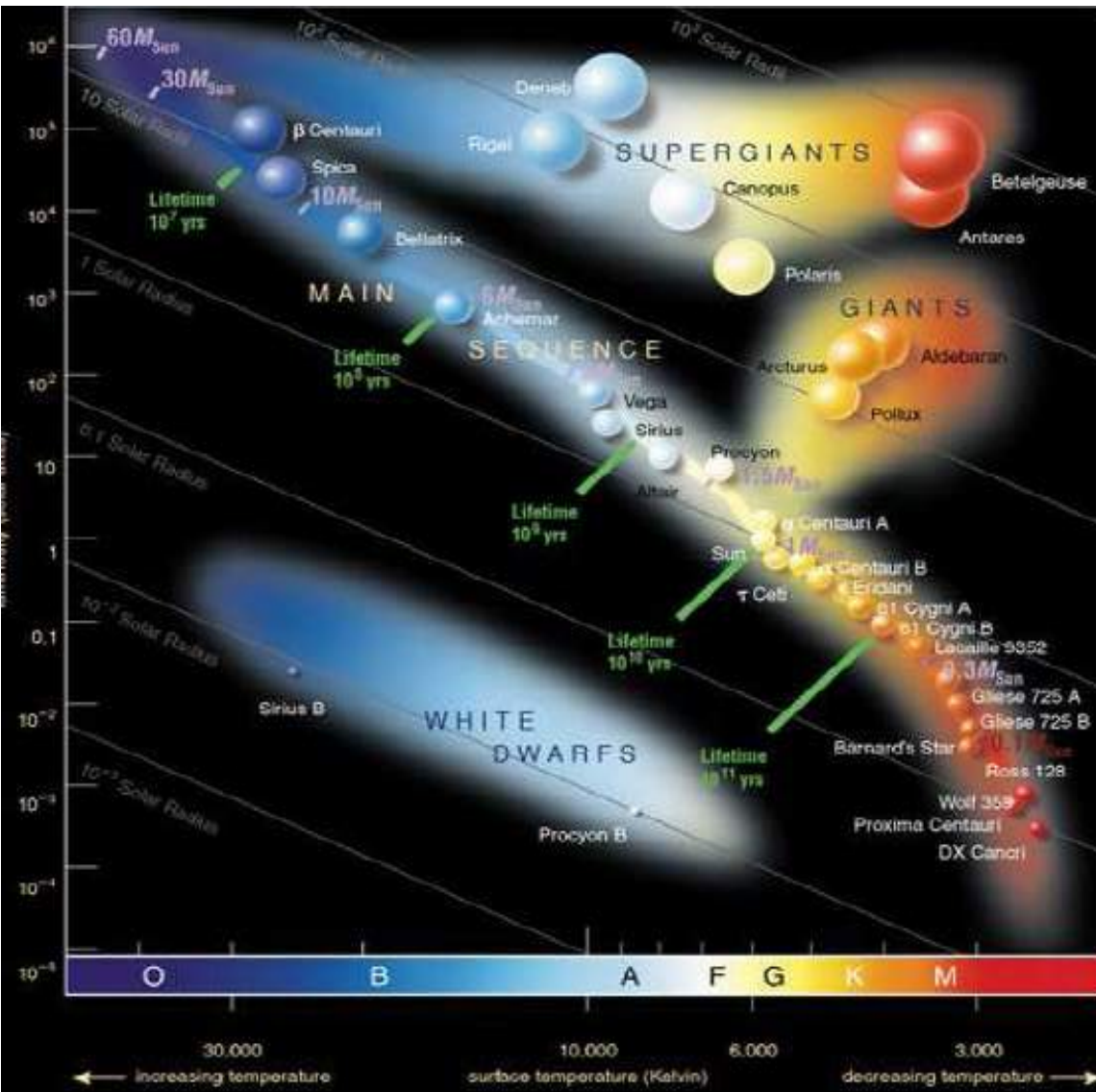
# Nuclear Binding Energy Map



# Solar Abundances



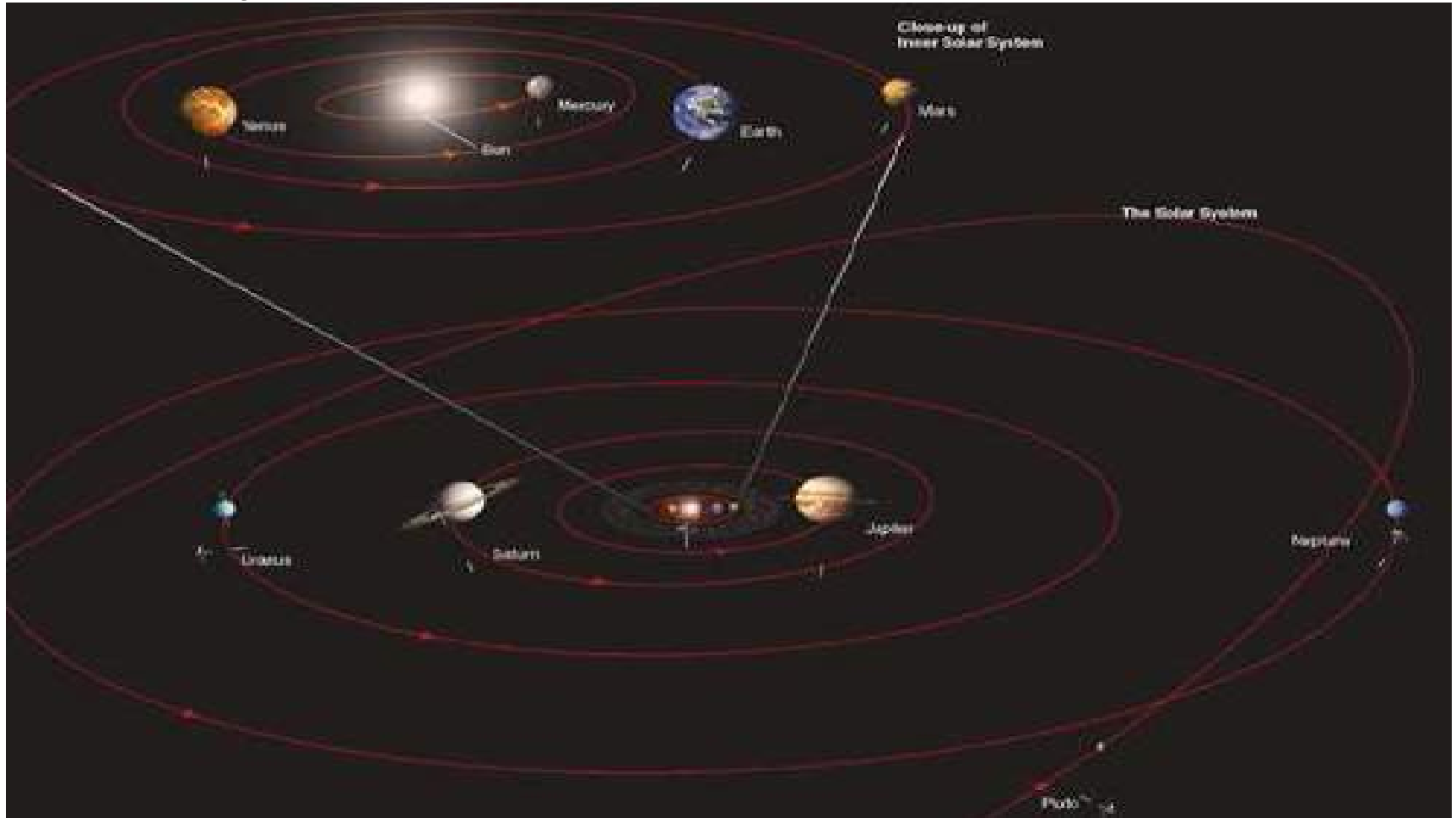
# Hertzprung- Russell (H-R) diagram



Stefan-Boltzmann  
Law for flux

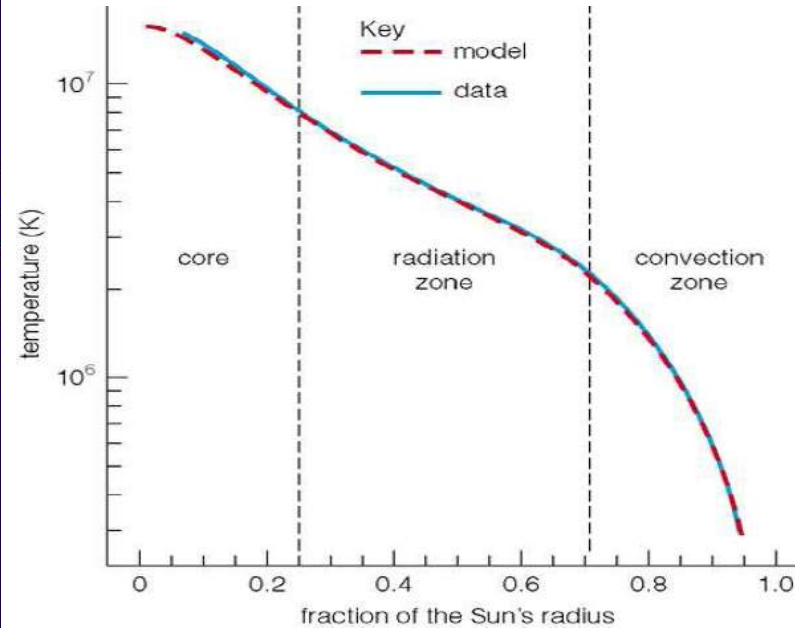
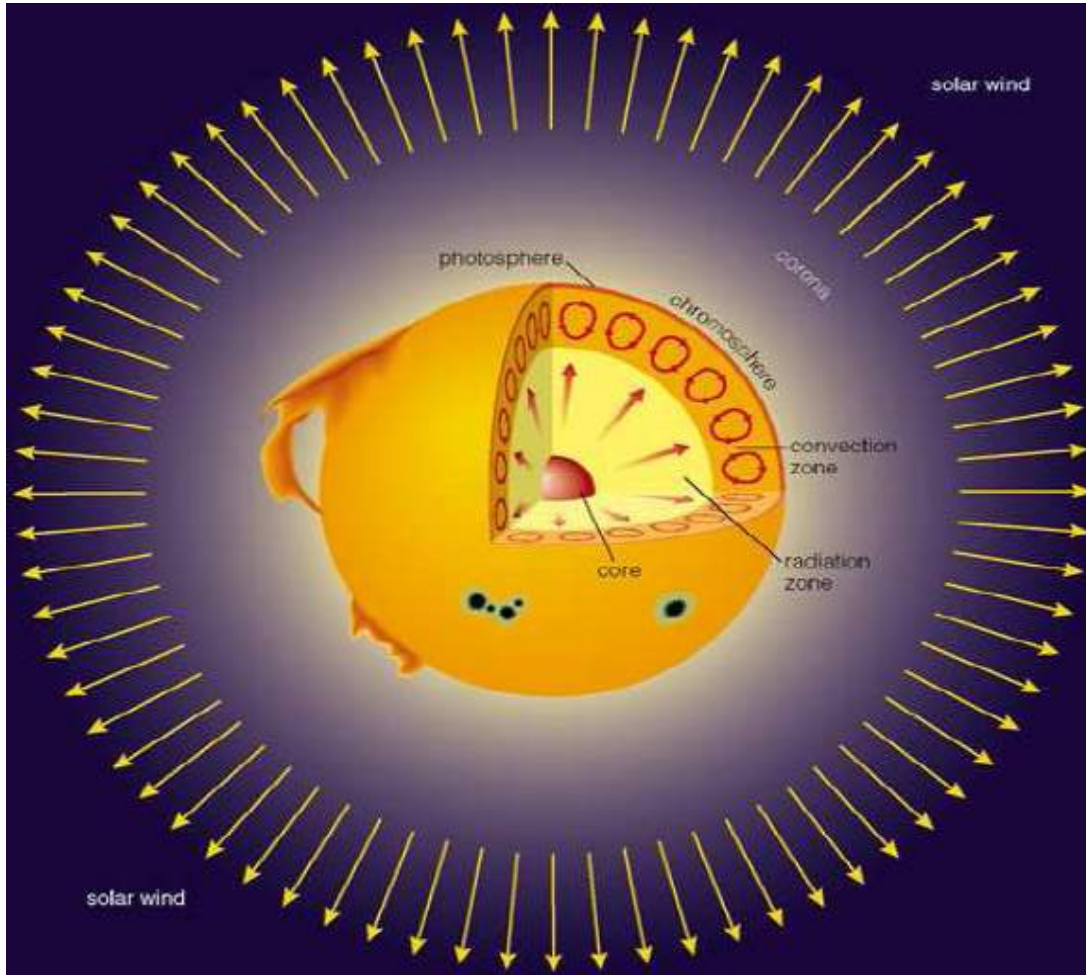
luminosity  $L$   
of a star with  
**radius  $R$  &**  
**surface temperature  $T$**   
 $L \sim (\text{Surface})T^4 \sim R^2T^4$

# Solar System



- Sun(star) + 9 planets - ***M, V, E, Mars, J S N U***
  - (P= dwarf planet)
- Distance Earth to Sun=  $1.5 \times 10^{11}$  m= 1AU
  - Pluto-Sun= 39.5 AU

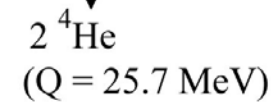
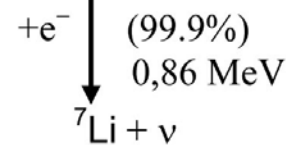
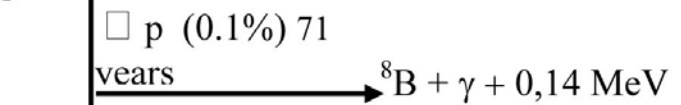
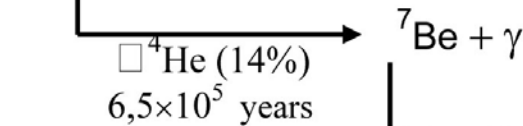
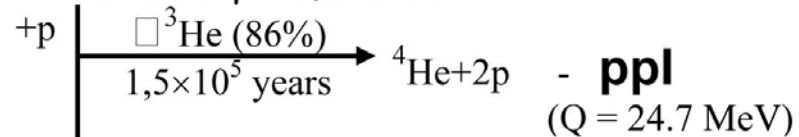
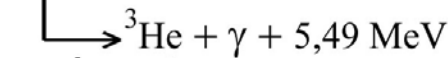
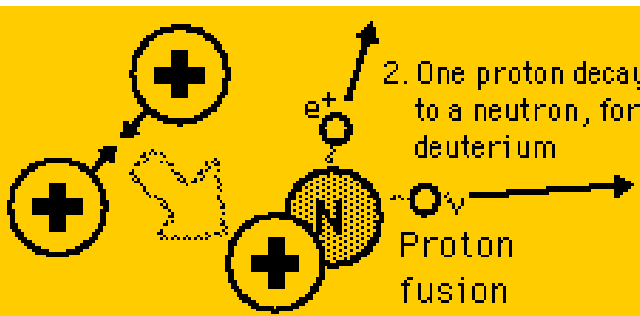
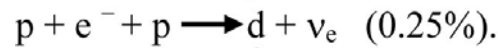
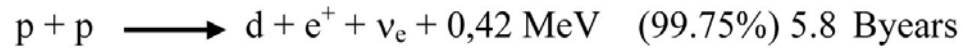
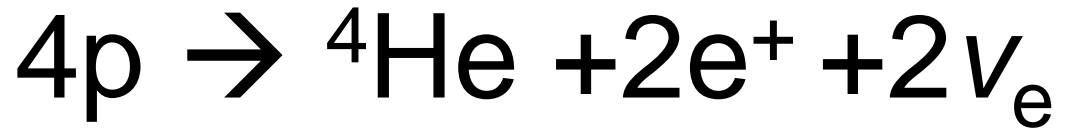
# Structure of low mass star, e.g. the Sun



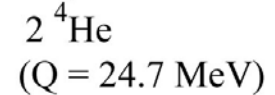
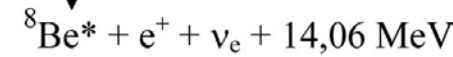
- Nuclear fusion in core with high temperature & pressure
- released energy that is transported to photosphere & emitted



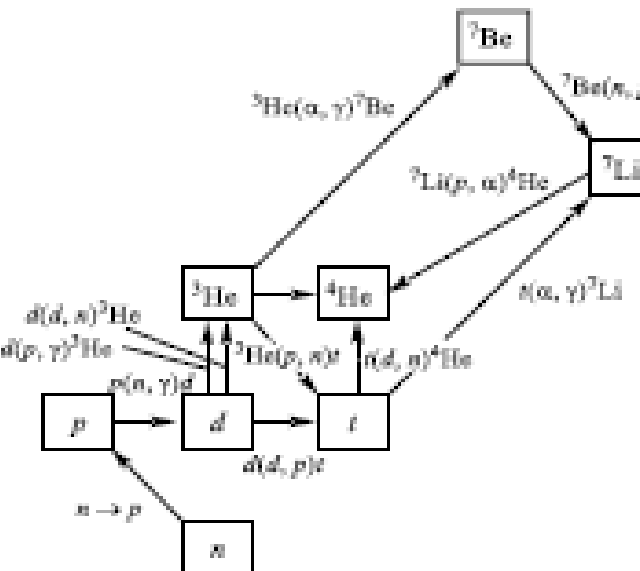
# pp chains



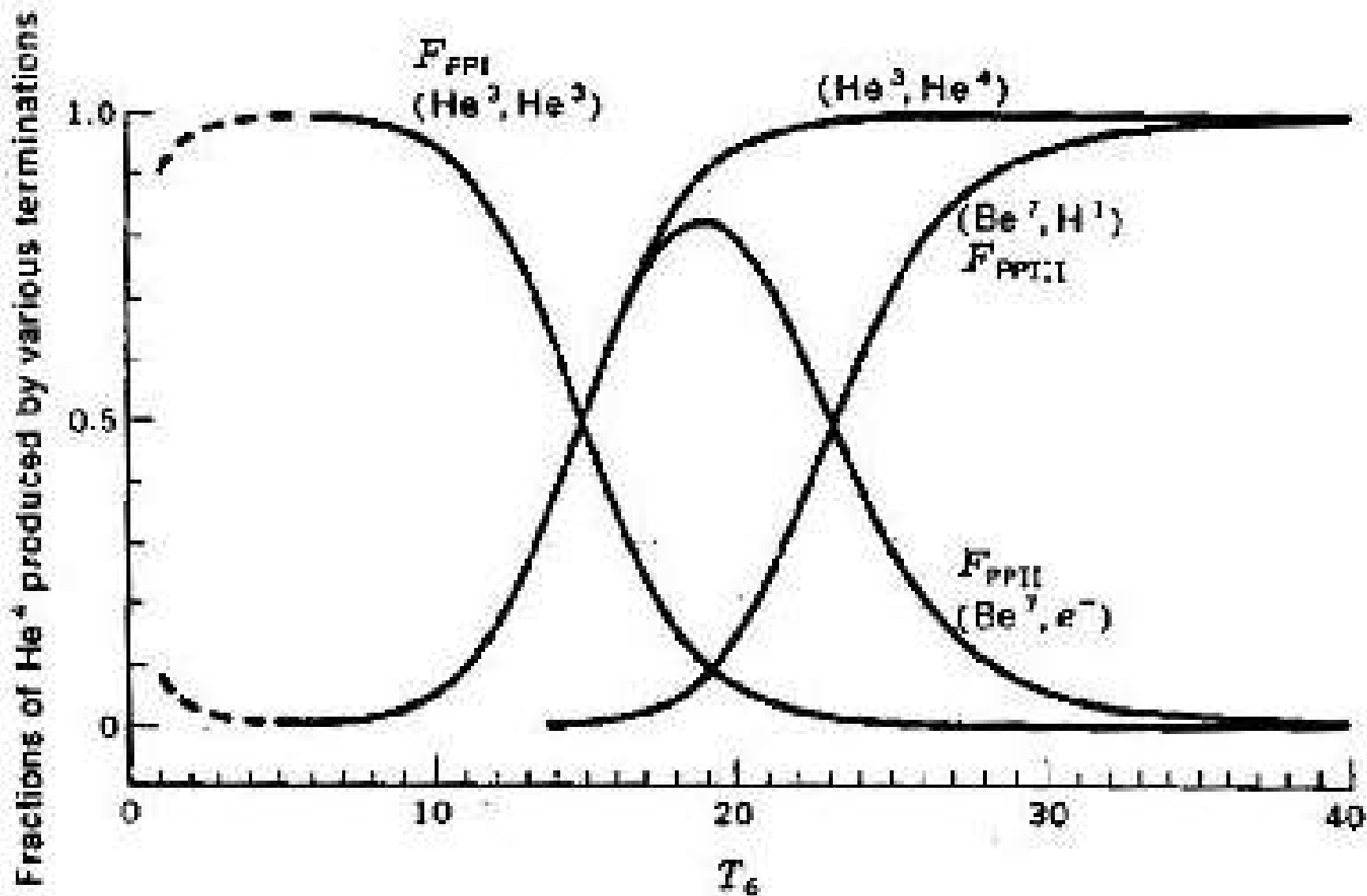
**ppII**



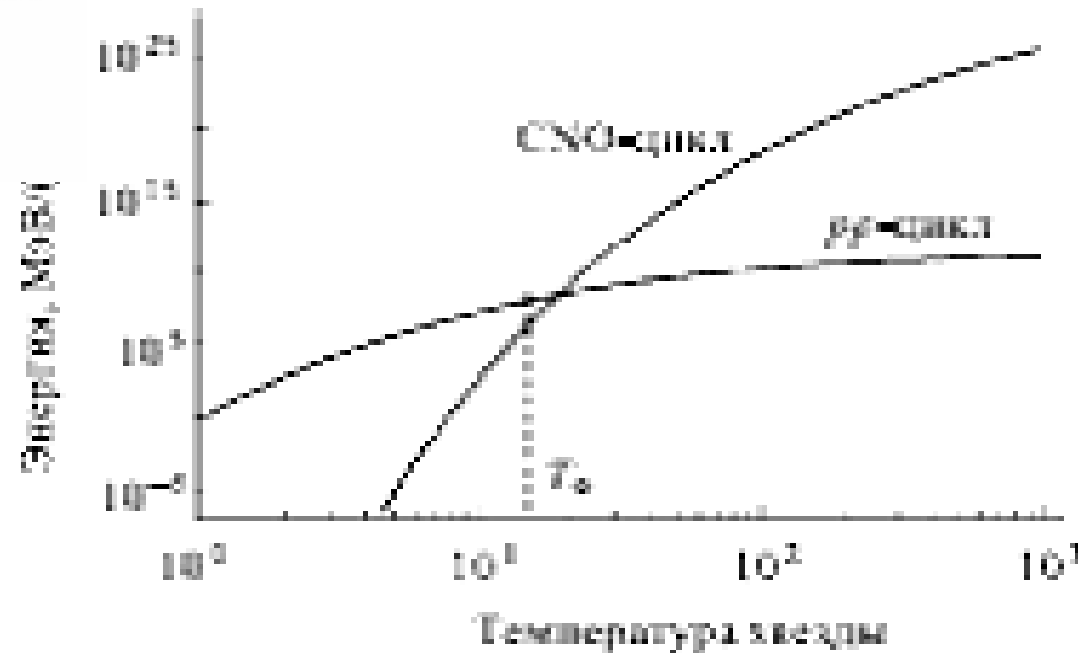
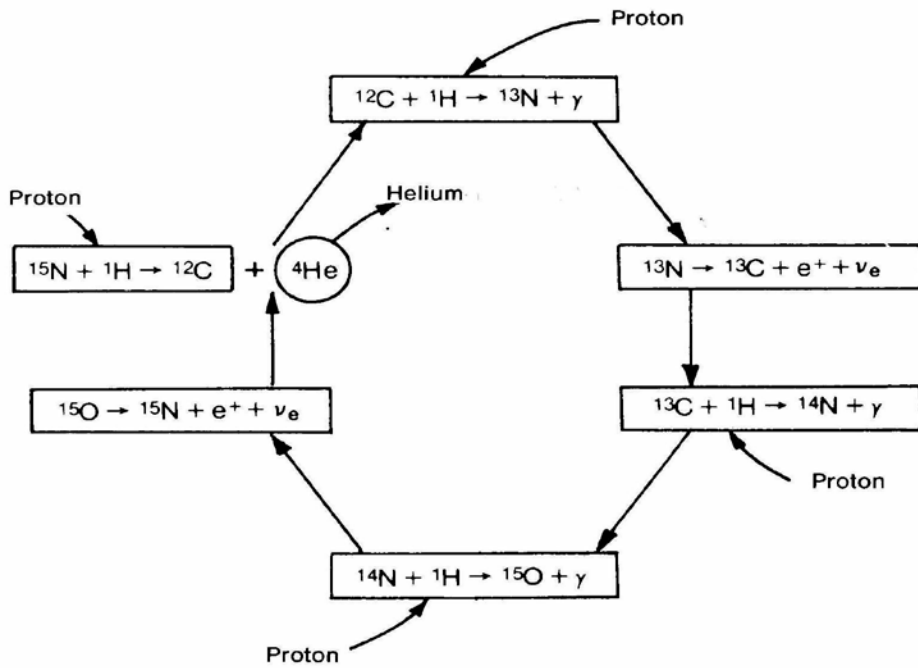
**ppIII**

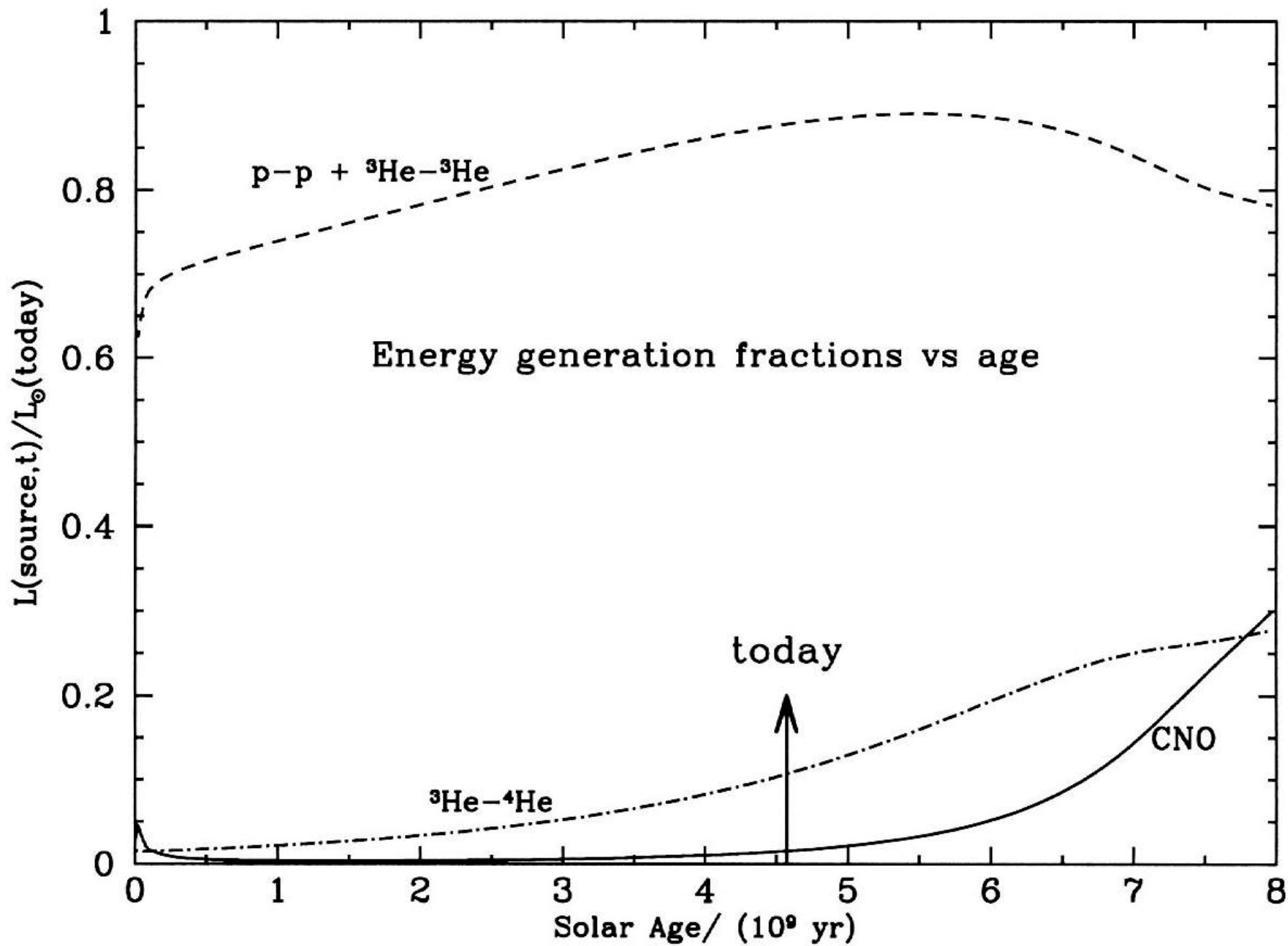


# Competition pp chains vs temperature

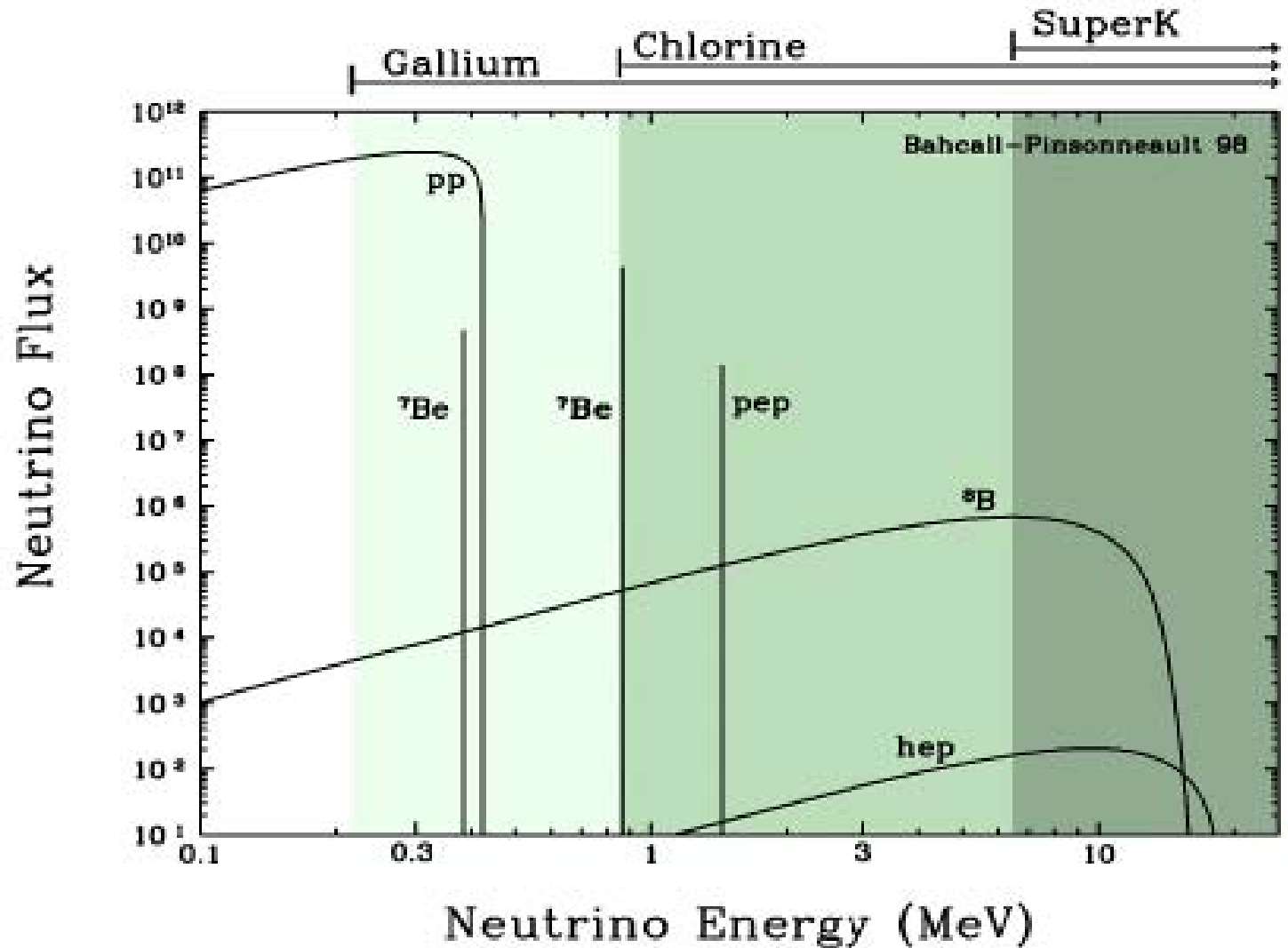


# pp & CNO cycles vs temperature





# solar neutrino spectrum



# Solar (& other) Neutrino

different fluxes



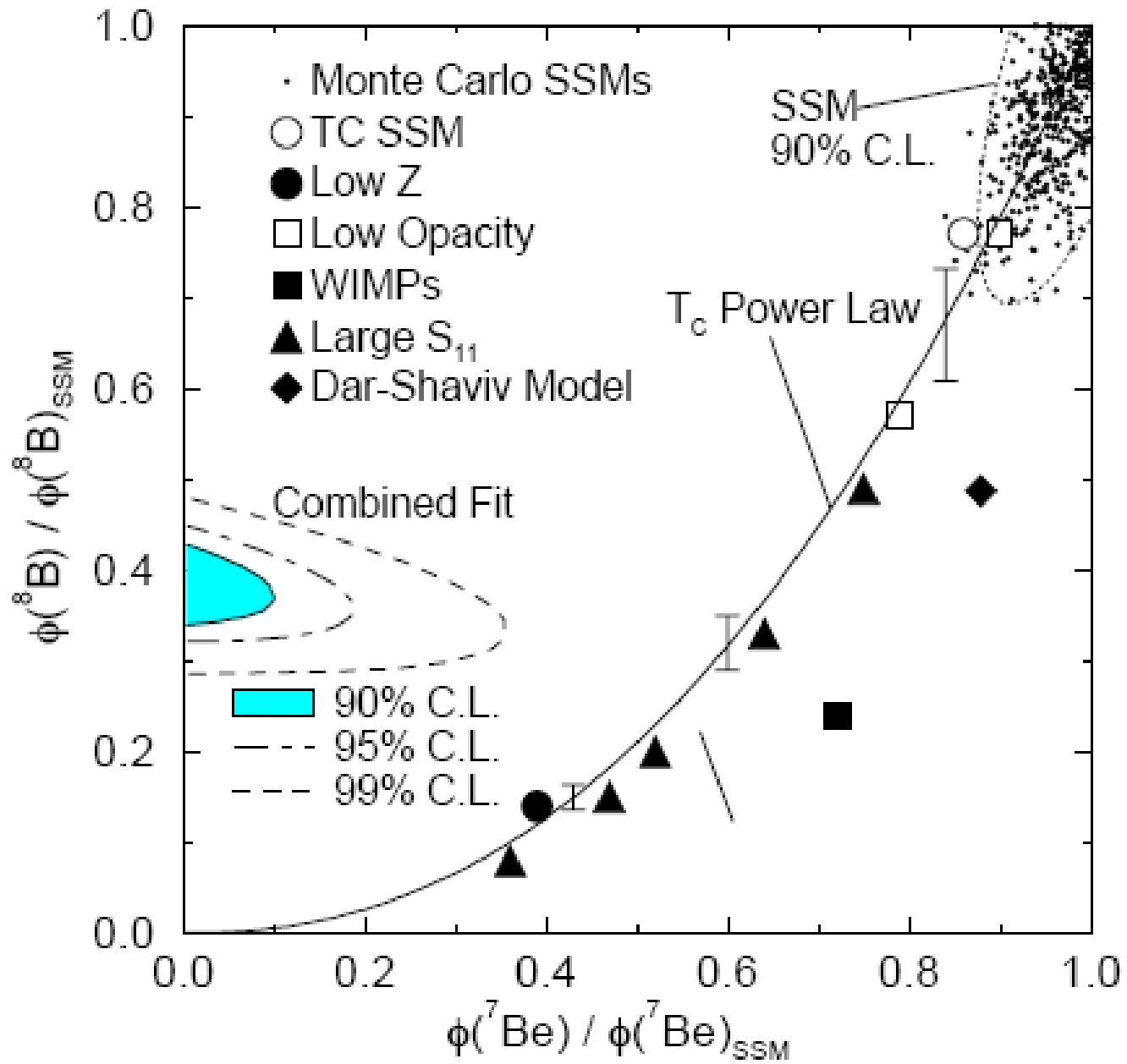
**observations**



**expectations from**  
standard solar model (**SSM**)

## Fluxes

$$\phi({}^7\text{Be}) \approx 0 \quad \phi({}^8\text{B}) \approx 0.43\phi^{\text{SSM}}({}^8\text{B})$$



# Neutrino masses & *vacuum* neutrino oscillations

weak interaction eigenstates

NEUTRINO  $\nu_e \leftrightarrow$  positron

$\nu_\mu \leftrightarrow$  muon

MASS eigenstates  $|\nu_1\rangle$  &  $|\nu_2\rangle$  with masses  
 $m_1$  &  $m_2$

$$|\nu_e\rangle = \cos\theta_\nu |\nu_1\rangle + \sin\theta_\nu |\nu_2\rangle$$

$$|\nu_\mu\rangle = \sin\theta_\nu |\nu_1\rangle + \cos\theta_\nu |\nu_2\rangle$$

(vacuum) mixing angle  $\theta_\nu$ .



time  $t=0 \rightarrow |\nu(t=0)\rangle = |\nu_e\rangle = \cos\theta_\nu |\nu_1\rangle + \sin\theta_\nu |\nu_2\rangle$

Each eigenstate propagates with a phase  $\blacktriangledown$

$$\exp\left\{i\left(\vec{k}\vec{x} - \omega t\right)\right\} = \exp\left\{i\left(\vec{k}\vec{x} - t\sqrt{m_i^2 + k^2}\right)\right\}$$

neutrino mass  $\ll$  momentum  $\blacktriangleleft \sqrt{m_i^2 + k^2} \approx k \left(1 + \frac{m_i^2}{2k^2}\right)$

$$|\nu(t)\rangle = \exp\left\{i\left(\vec{k}\vec{x} - kt - \frac{m_1^2 + m_2^2}{2k}t\right)\right\}$$

$$\left[ \cos\theta_\nu |\nu_1\rangle \exp\left\{i\delta m^2 t / 4k\right\} + \sin\theta_\nu |\nu_2\rangle \exp\left\{-i\delta m^2 t / 4k\right\} \right]$$

**BEAT PHASE**  $\delta m^2 = m_2^2 - m_1^2$

# PROBABILITY for neutrino state

to remain  $|\nu_e\rangle$  at time  $t$

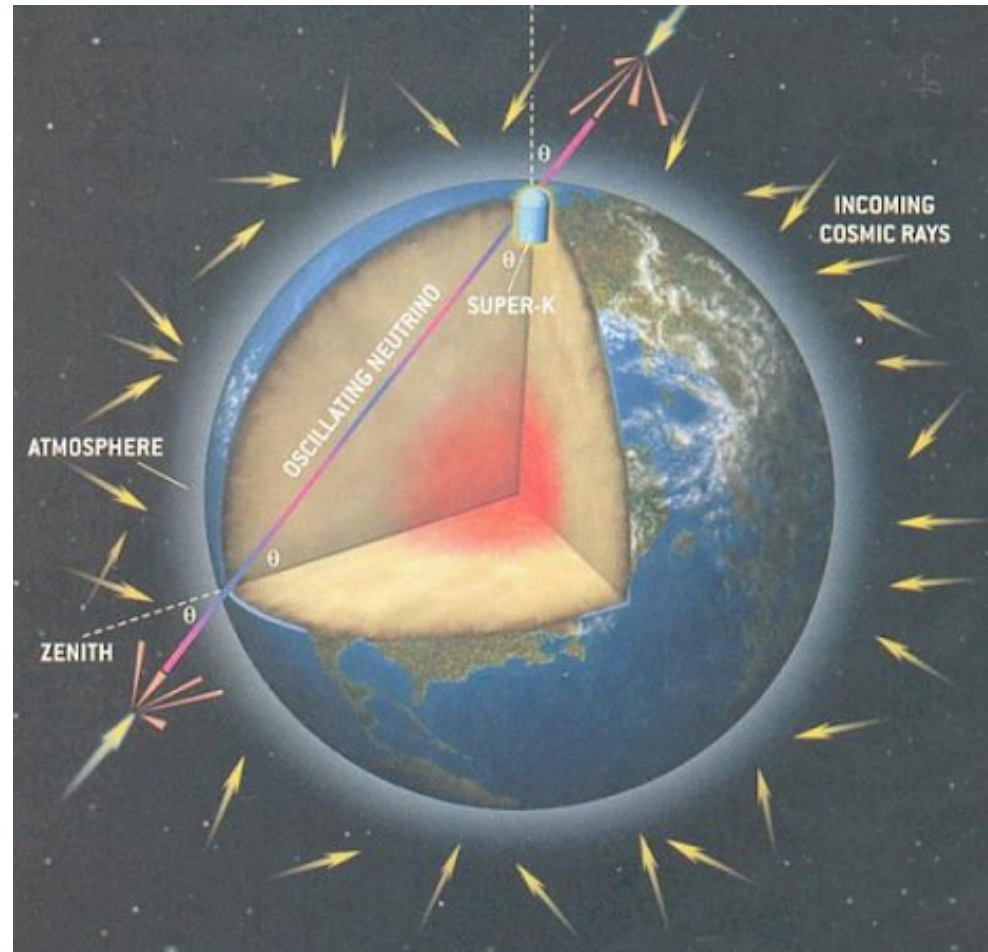
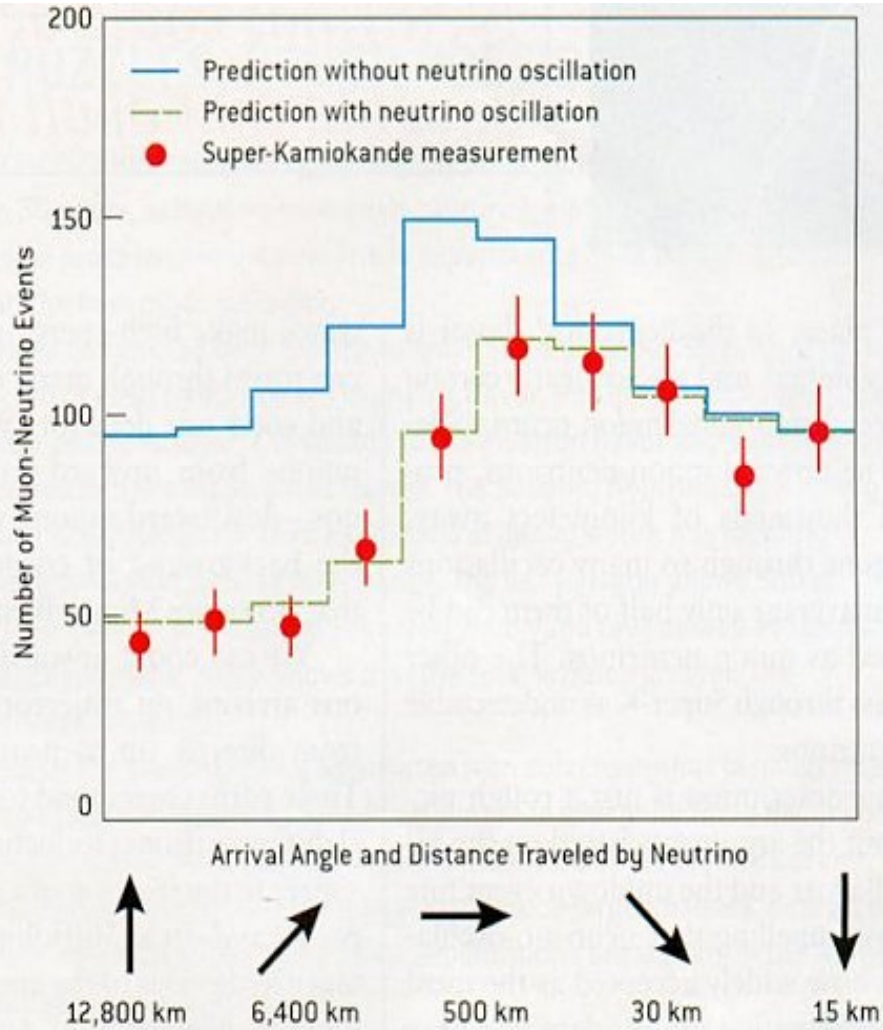
$$P_{\nu_e}(t) = \left| \langle \nu_e | \nu(t) \rangle \right|^2$$
$$= 1 - \sin^2 2\theta_\nu \sin^2 \left( \frac{\delta m^2 t}{4k} \right) \rightarrow 1 - \sin^2 2\theta_\nu$$

$$m \ll E \sim k \rightarrow P_\nu(x) = 1 - \sin^2 2\theta_\nu \sin^2 \left( \frac{\delta m^2 c^4 x}{4\hbar c E} \right)$$

oscillation length  $L_o = \frac{4\pi\hbar c E}{\delta m^2 c^4}$

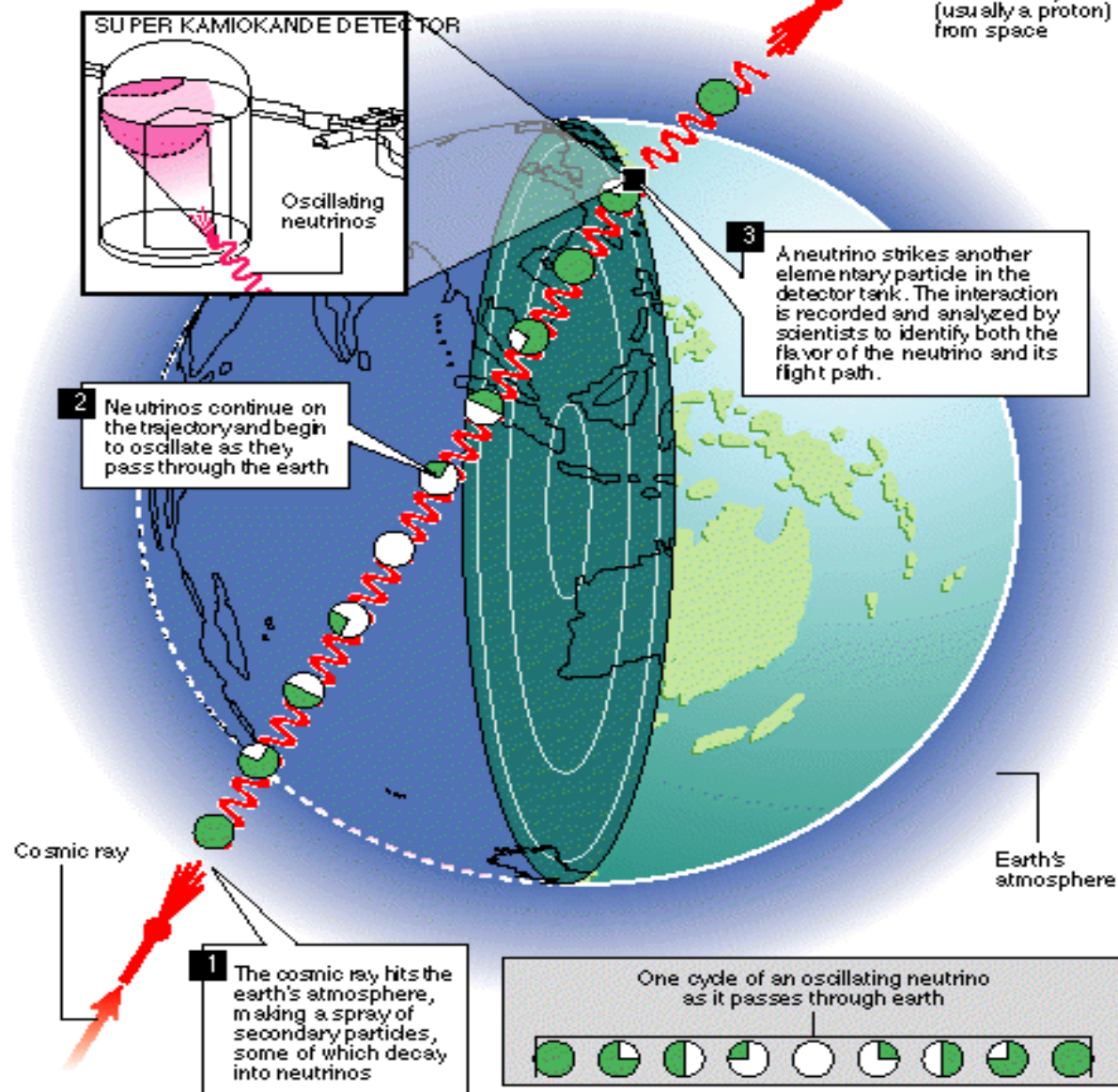
$E \sim 1 \text{ MeV}$ : sensitivity to  $\delta m_\nu^2 \geq 10^{-12} \text{ eV}^2$

# SuperKamiokande

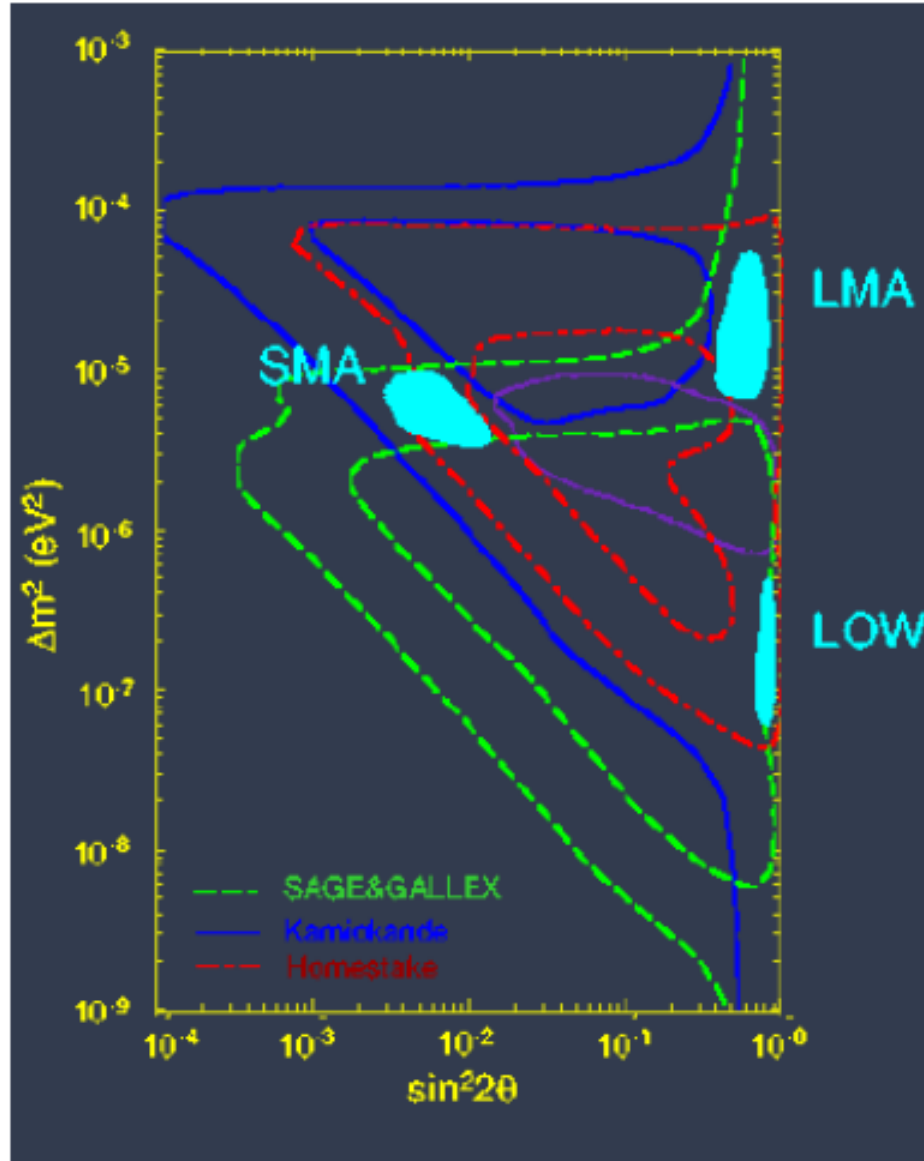


## Discovering Mass

The farther neutrinos travel, the more time they have to oscillate. By comparing the ratio of flavors of neutrinos coming "up" through the Earth to those coming from overhead, physicists determined that neutrinos oscillate, which neutrinos can only do if they have mass.



# Parameters from Cl/Ga/Kamiokande/Super-Kamiokande neutrino data



**LIGHT & HEAVY** local mass *eigenstates*

$$|\nu_L\rangle = \cos\theta(x) |\nu_e\rangle - \sin\theta(x) |\nu_\mu\rangle$$

$$|\nu_H\rangle = \sin\theta(x) |\nu_e\rangle + \cos\theta(x) |\nu_\mu\rangle$$

**local mixing angle**

$$\sin 2\theta(x) = \frac{\sin 2\theta_\nu}{\sqrt{X^2(x) + \sin^2 2\theta_\nu}}$$

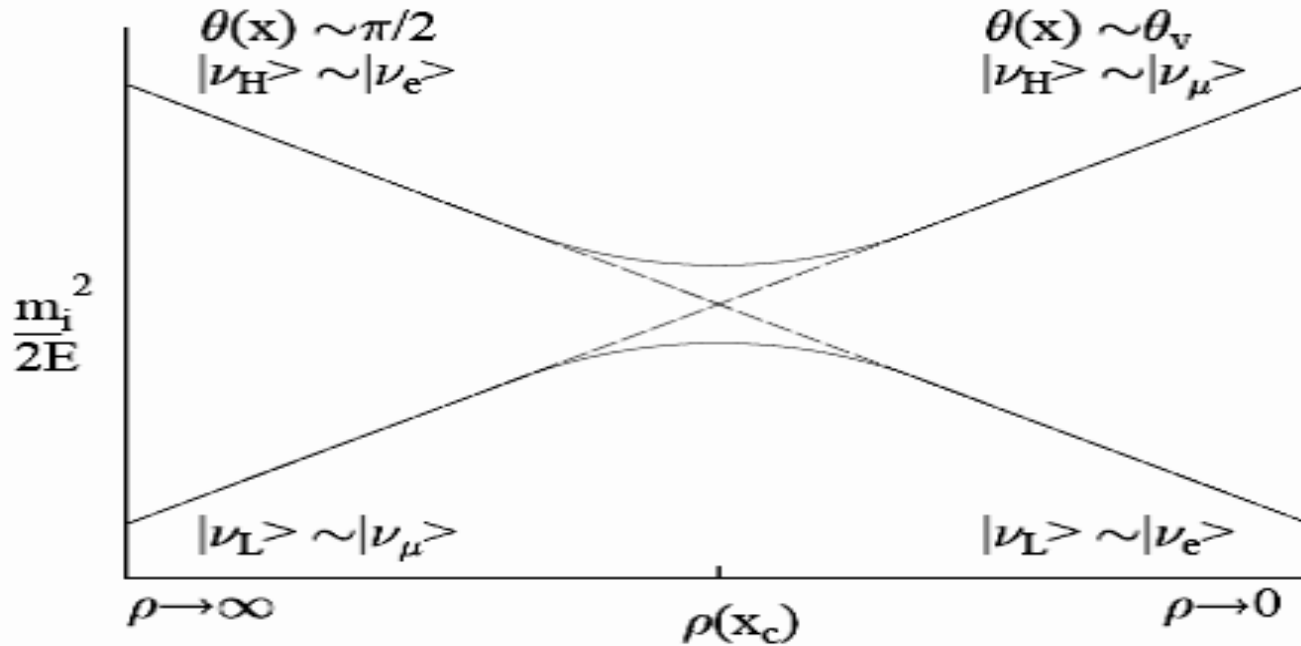
$$\cos 2\theta(x) = \frac{-X(x)}{\sqrt{X^2(x) + \sin^2 2\theta_\nu}}$$

$$X(x) = 2E\sqrt{2}G_F\rho(x)E / \delta m^2 - \cos 2\theta_\nu :$$

$$\theta(x) \in \theta_\nu \rightarrow \pi/2 \text{ as density } \rho(x) \ 0 \rightarrow \infty .$$

# Mikheyev-Smirnov-Wolfenstein

## *avoided level crossing*



**splitting minimum value:** --  $\sin 2\theta_v \delta m^2 / 2E$ ,

at **critical density**  $\rho_c = \rho(x_c) \leftarrow 2\sqrt{2}EG_F \rho_c = \delta m^2 \cos 2\theta_v$

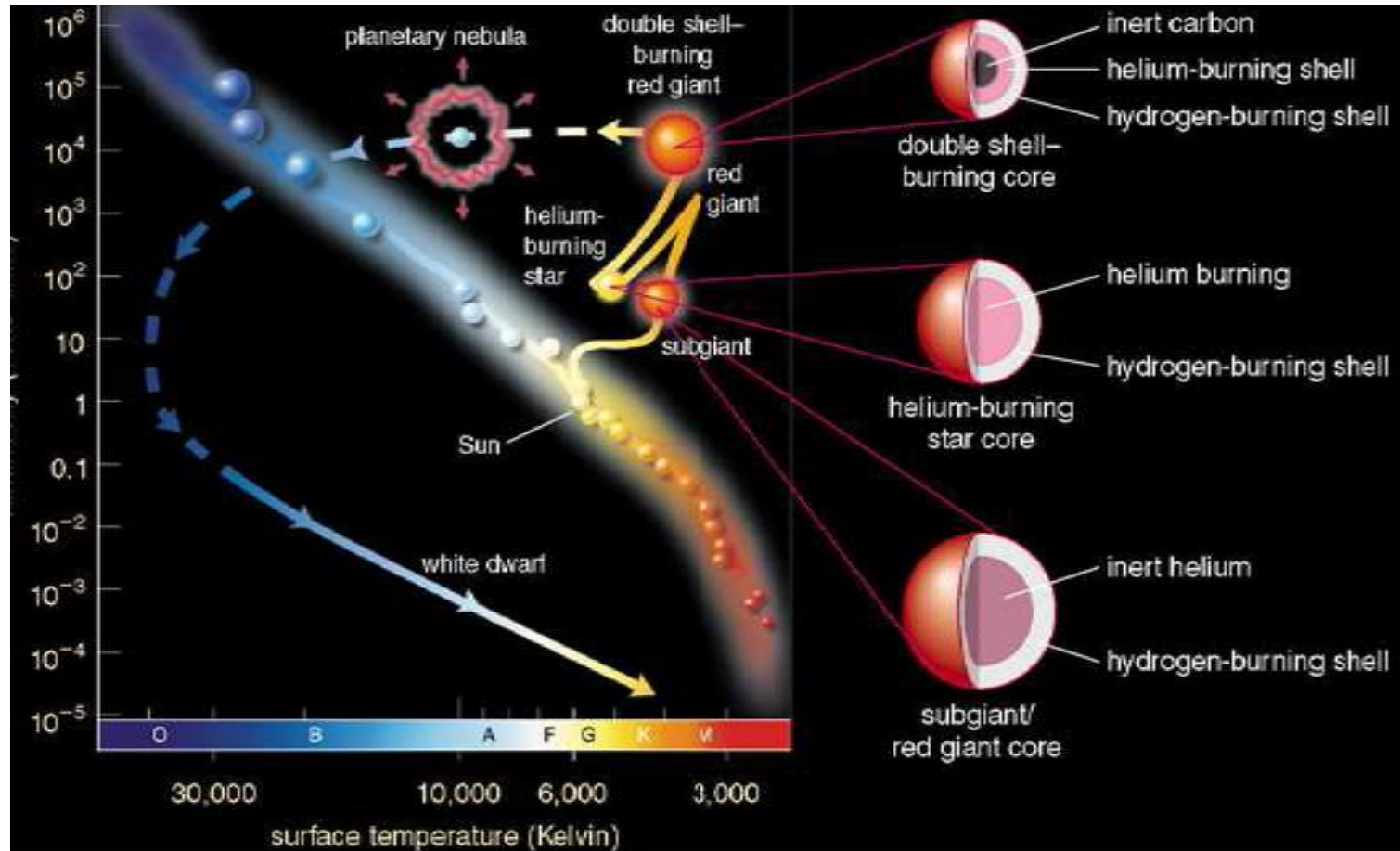
cross point for diagonal elements of original flavor matrix





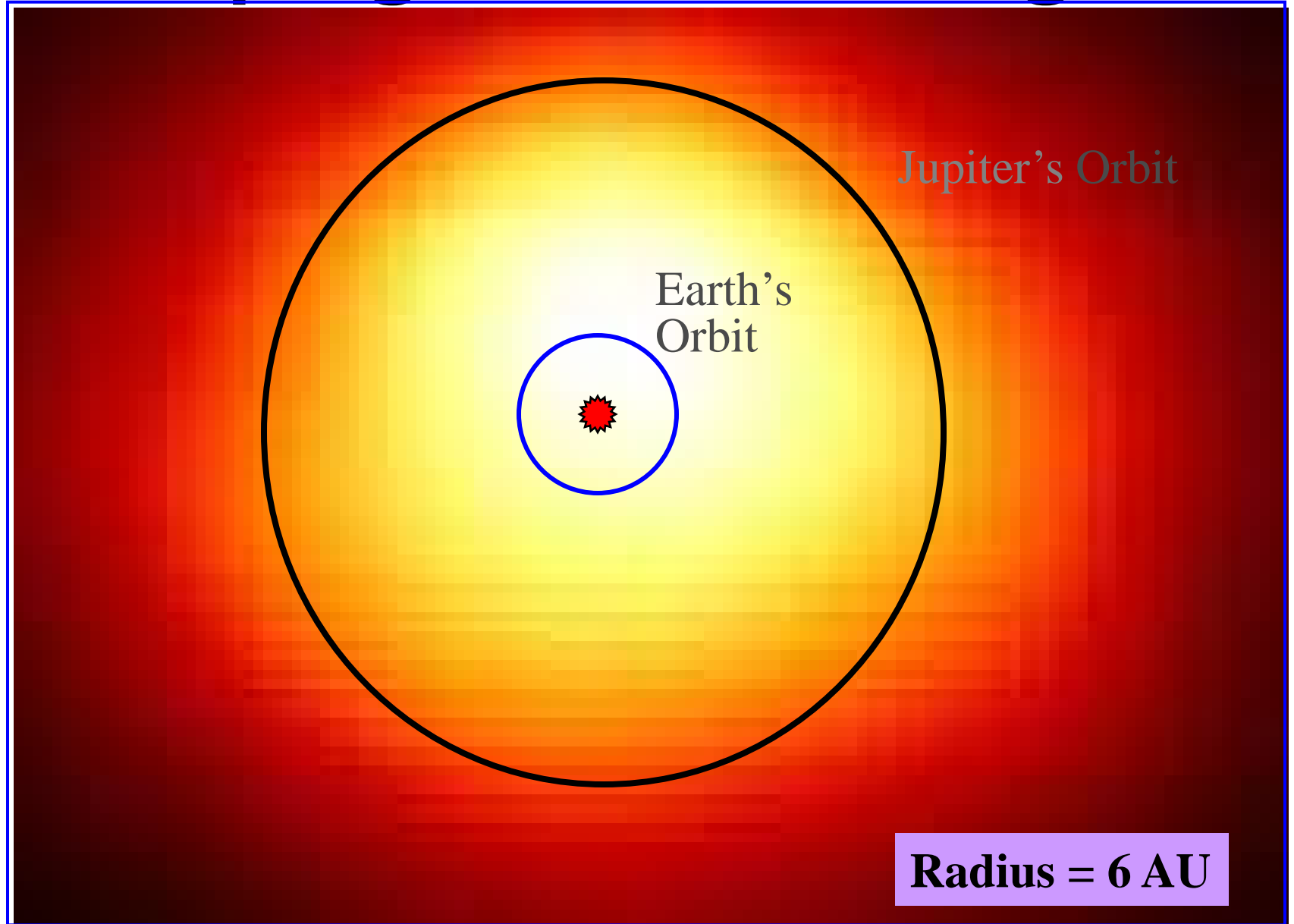
# HR diagr evolution: low-mass star at exhausted H in core

- Main sequence star CORE/reactor fusing H to He at  $T > 10^7 \text{ }^\circ\text{K}$



- Subgiant /Red Giant star (inert He core and Hydrogen- fusing shell)
- All the core H converted to He. To fuse He $\rightarrow$ Carbon  $T_c > 10^8 \text{ K}$
- inert He core shrinks, heats up H-Shell bringing fusion of H to Helium starts
- star expands & enters subgiant/red giant phase: large and luminous

# The Supergiant Star Betelgeuse

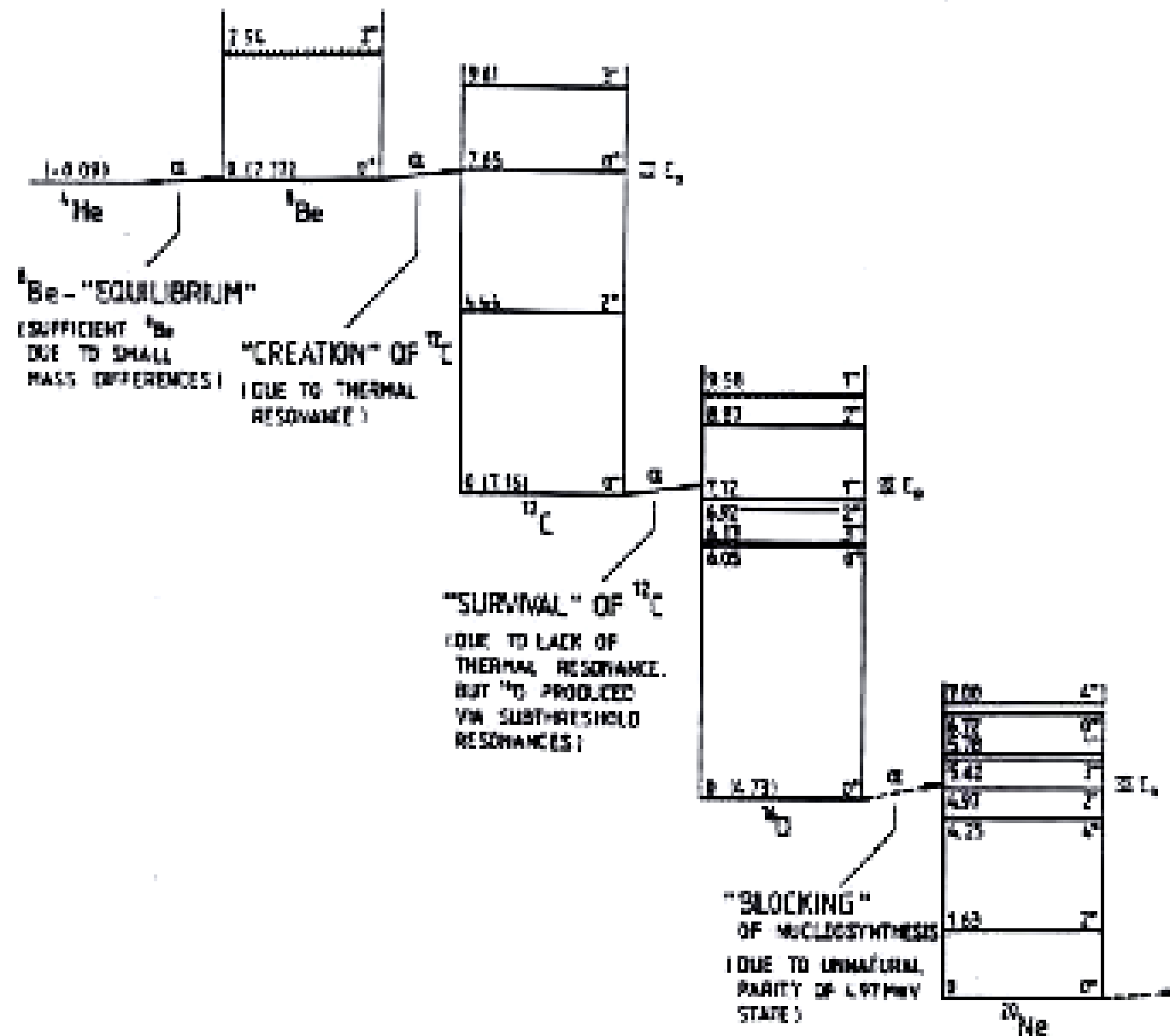


# Red giants and helium burning

- helium core densities  $\sim 10^6 \text{ g/cm}^3$
- core densities and temperatures allow --> helium burning
- $4\text{He} + 4\text{He} + 4\text{He} \rightarrow 12\text{C} + \Gamma\text{-ray}$
- red giants evolve along
  - the “horizontal branch”
- Form **Planetary nebulae & white dwarf**

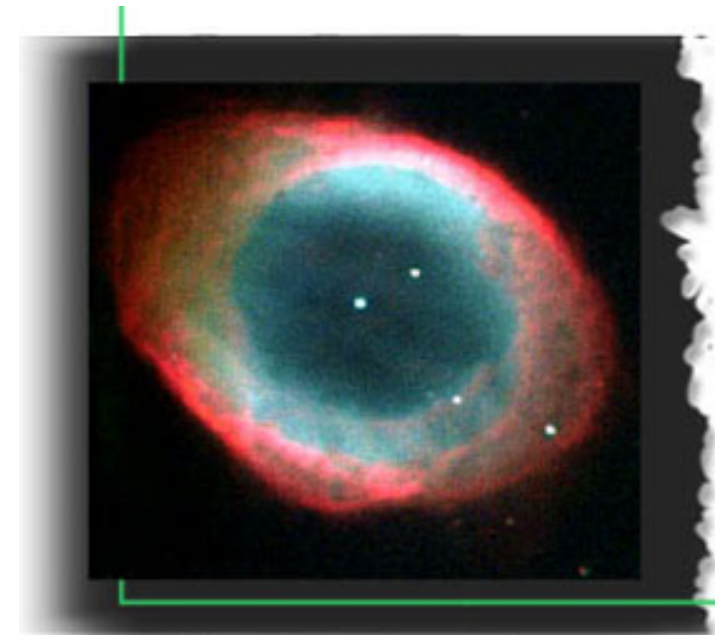
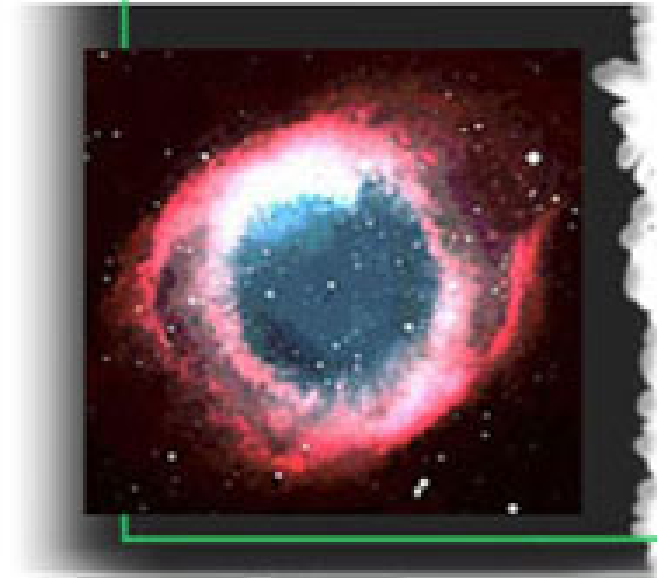
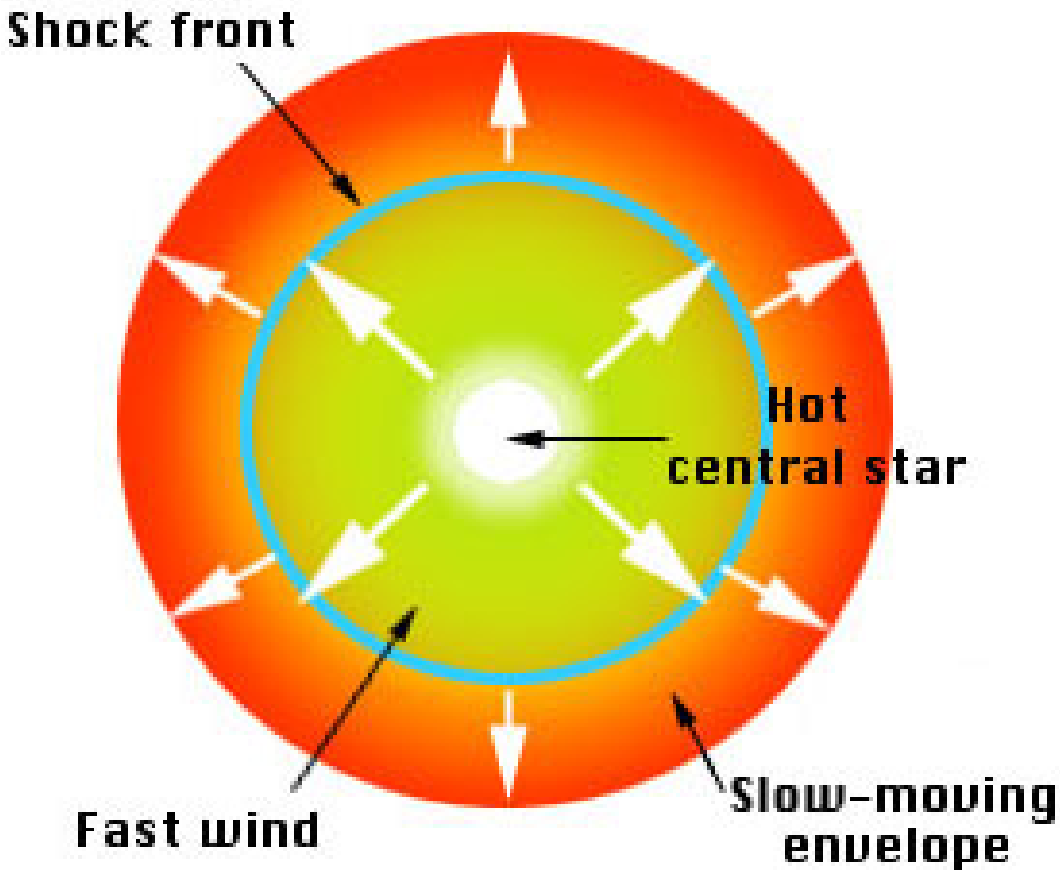
# HELIUM-BURNING REACTIONS IN RED GIANTS

(THE START OF NUCLEOSYNTHESIS OF HEAVY ELEMENTS)

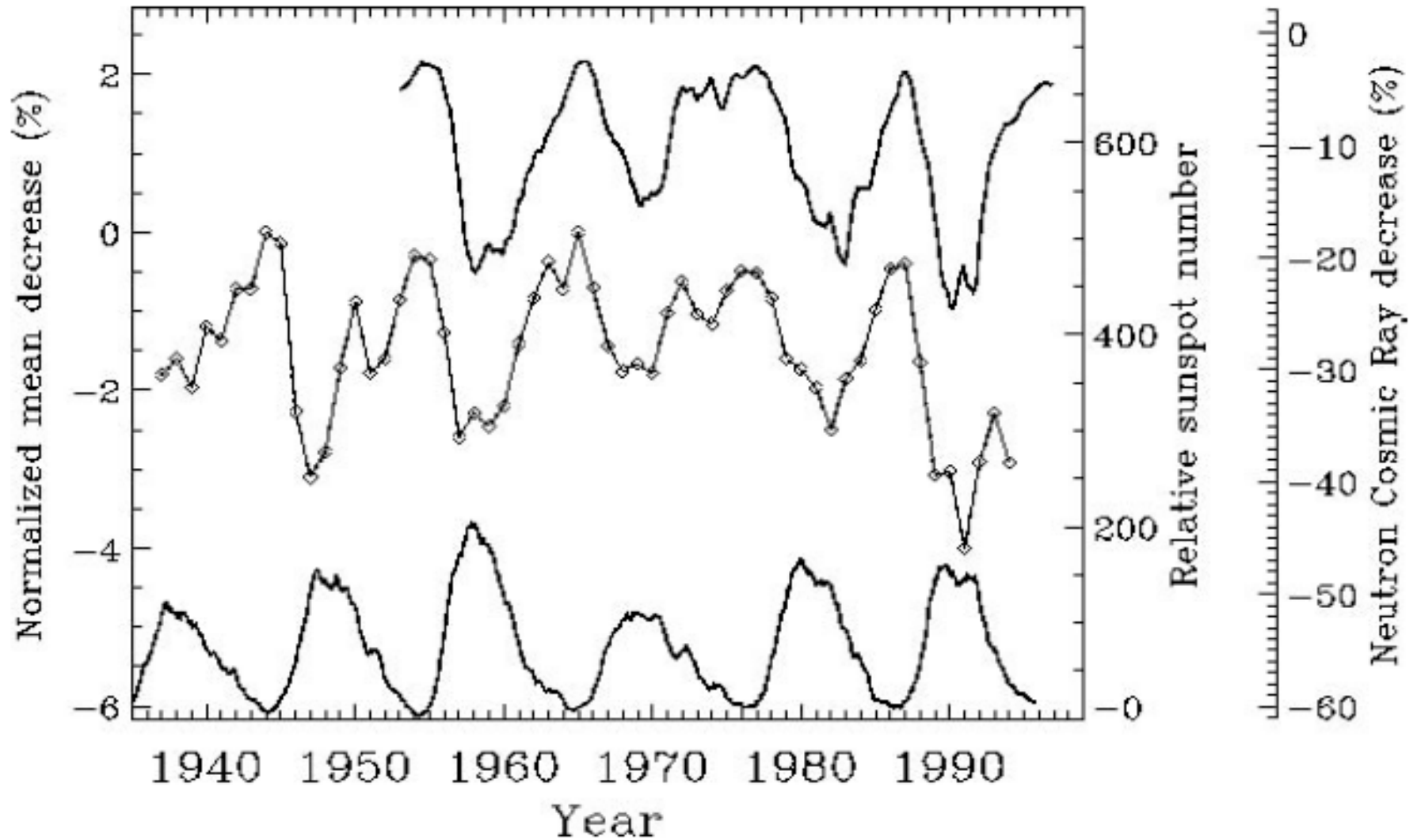


# Planetary Nebula Formation

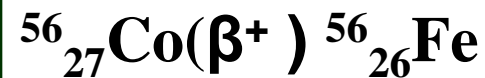
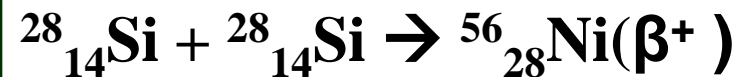
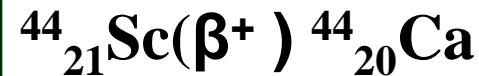
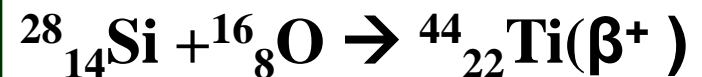
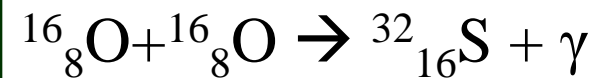
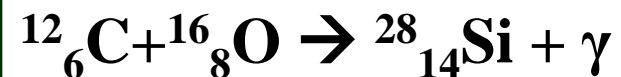
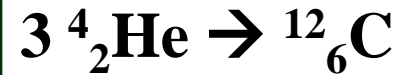
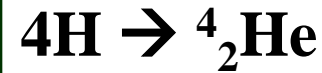
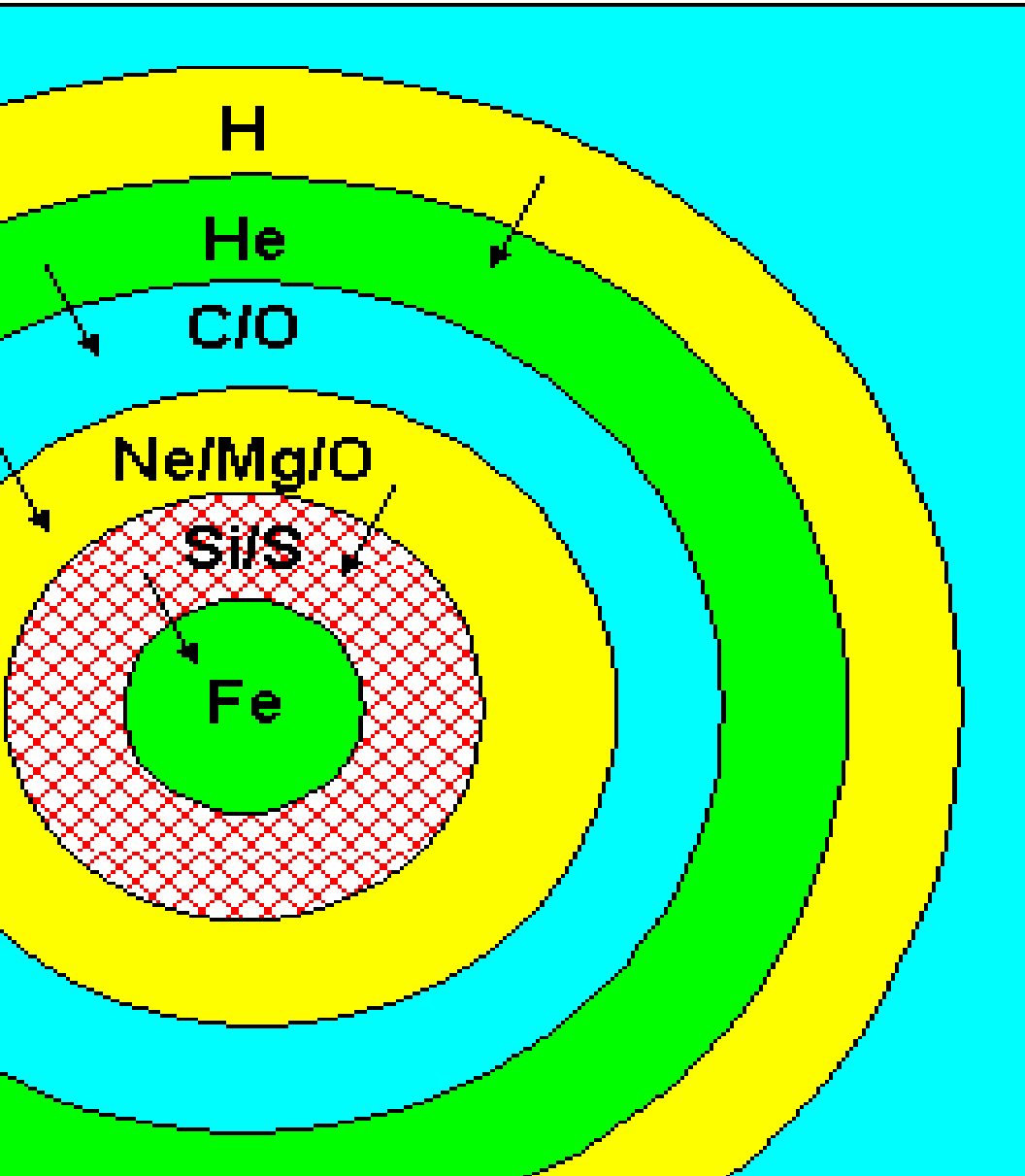
## Helix

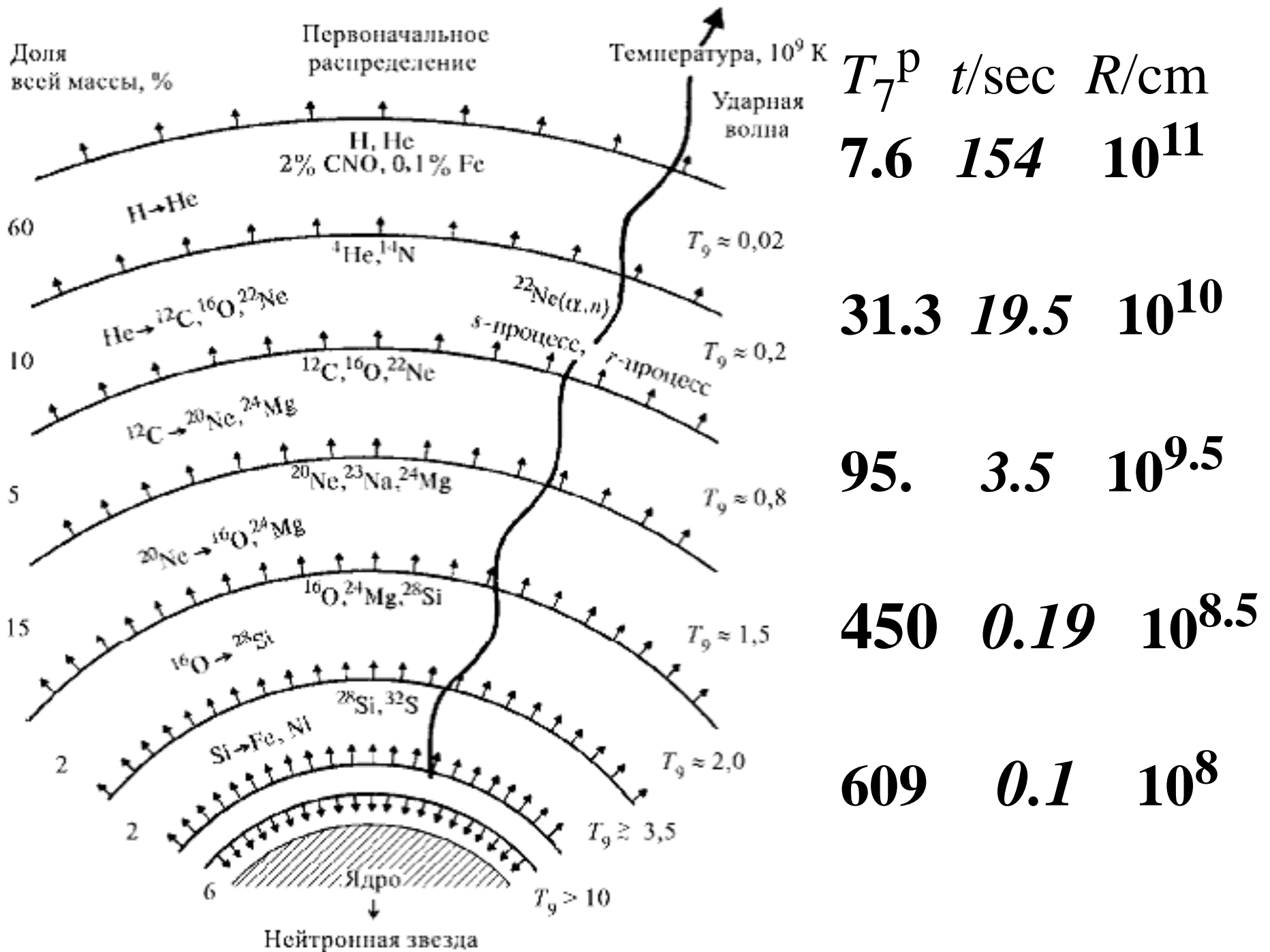


# Cosmic Rays flux vs solar activity



# Massive Star, $M > 8 M_{\odot}$ *onion*







# Crab



Supernova remnant called Crab Nebula; VLT/Optical



X-ray (blue) and optical (red) radiation from Crab Nebula's core region. A pulsar near the center propels particles.



# Горение атомных ядер: ВЗРЫВНОЕ ГОРЕНИЕ

Дубна.2014

**В.Н. Кондратьев**

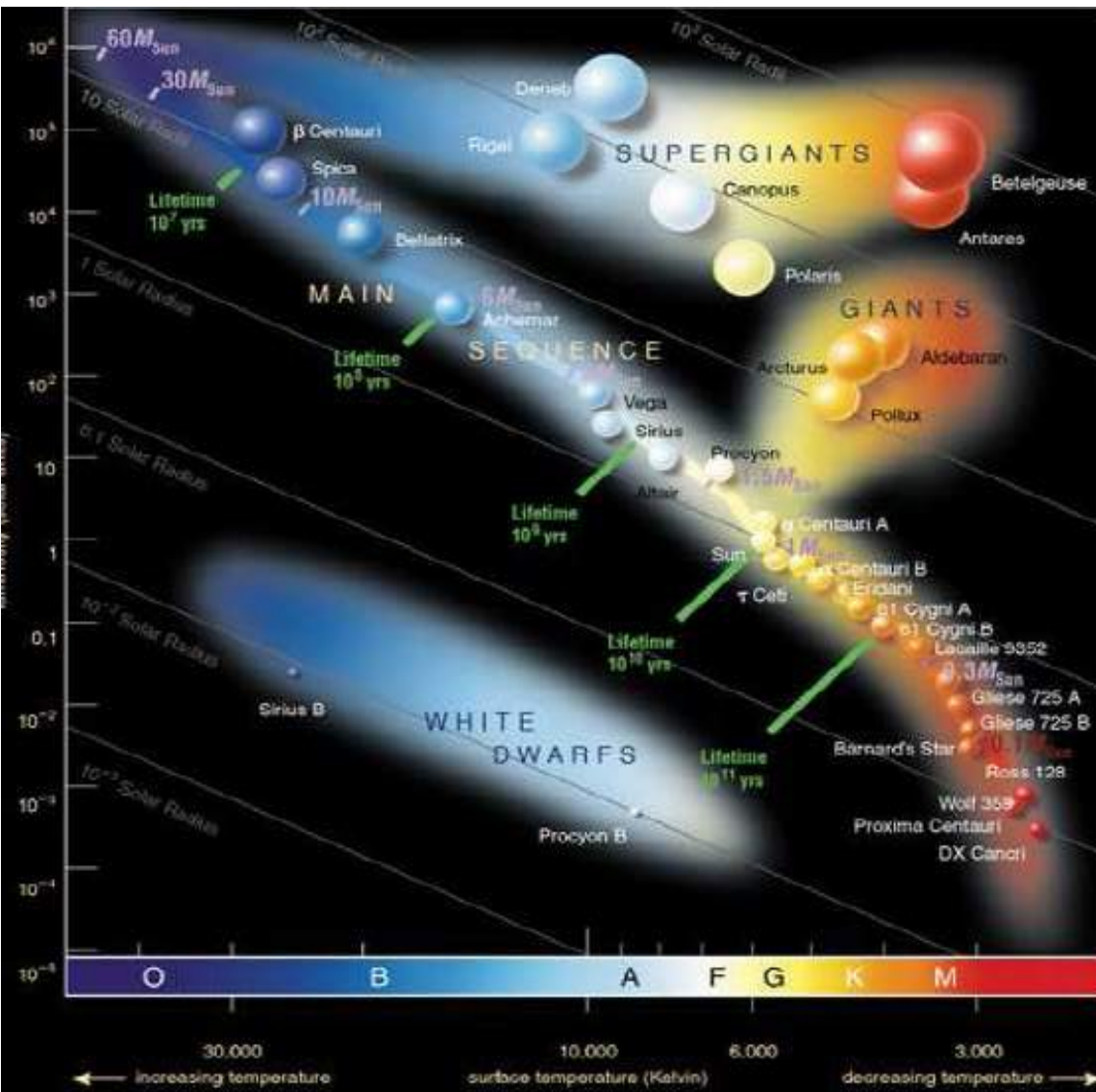
ЛТФ, ОИЯИ, 141980, Дубна Россия

*Киевский национальный университет имени  
Тараса Шевченко, UA-03022 Киев, Украина*

# Before and after pictures of SN1987a



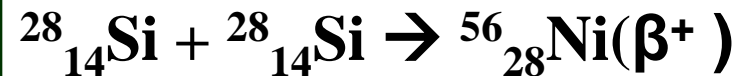
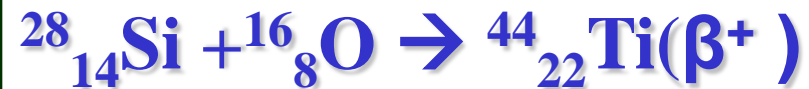
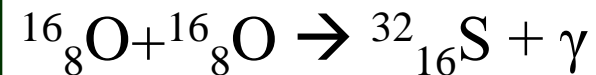
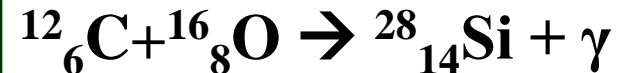
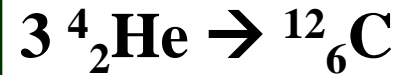
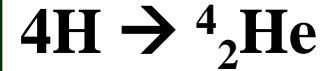
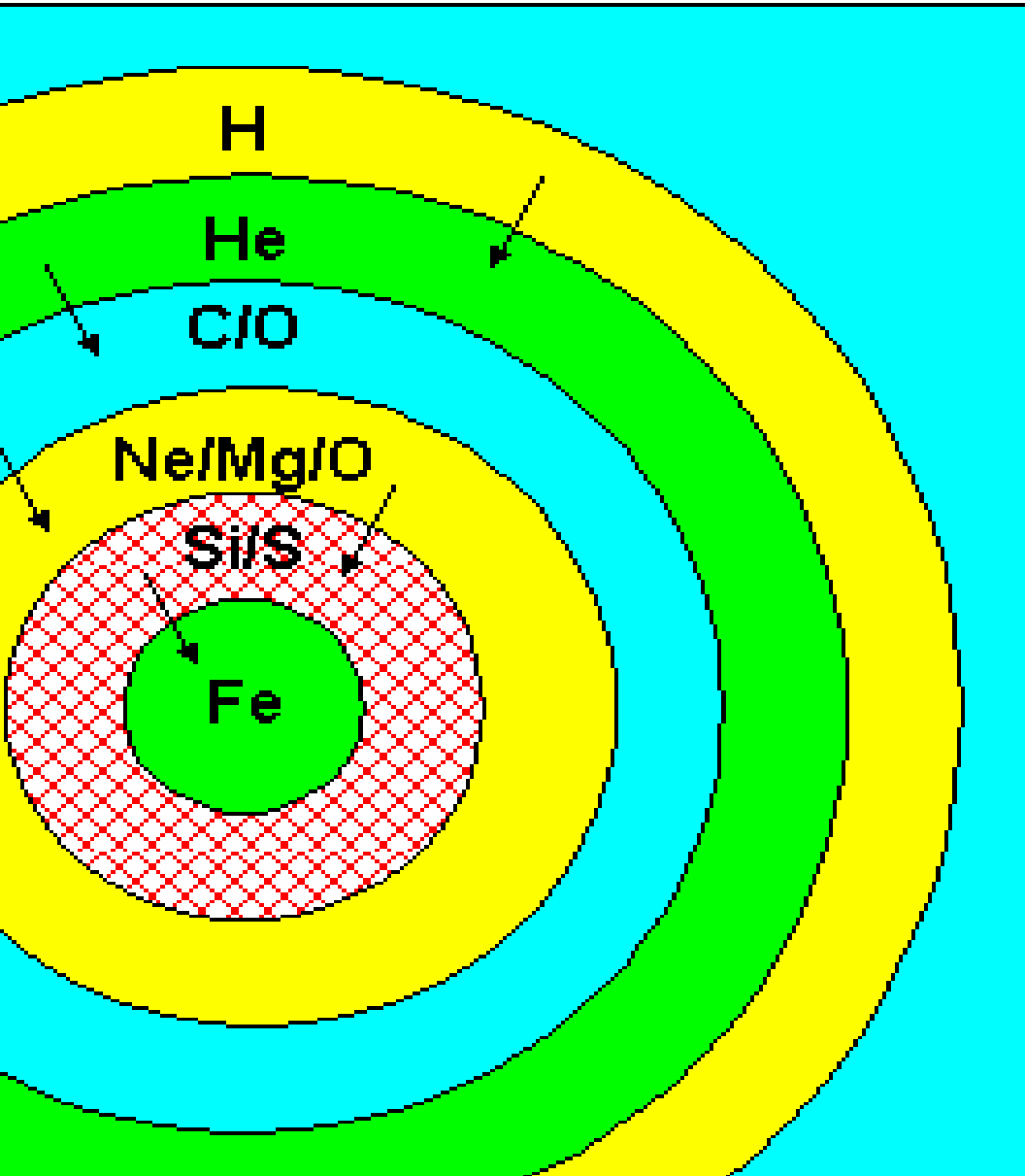
# Hertzsprung- Russell (H-R) diagram



Stefan-Boltzmann  
Law for flux

luminosity  $L$   
of a star with  
*radius  $R$  &*  
*surface temperature  $T$*   
 $L \sim (\text{Surface}) T^4 \sim R^2 T^4$

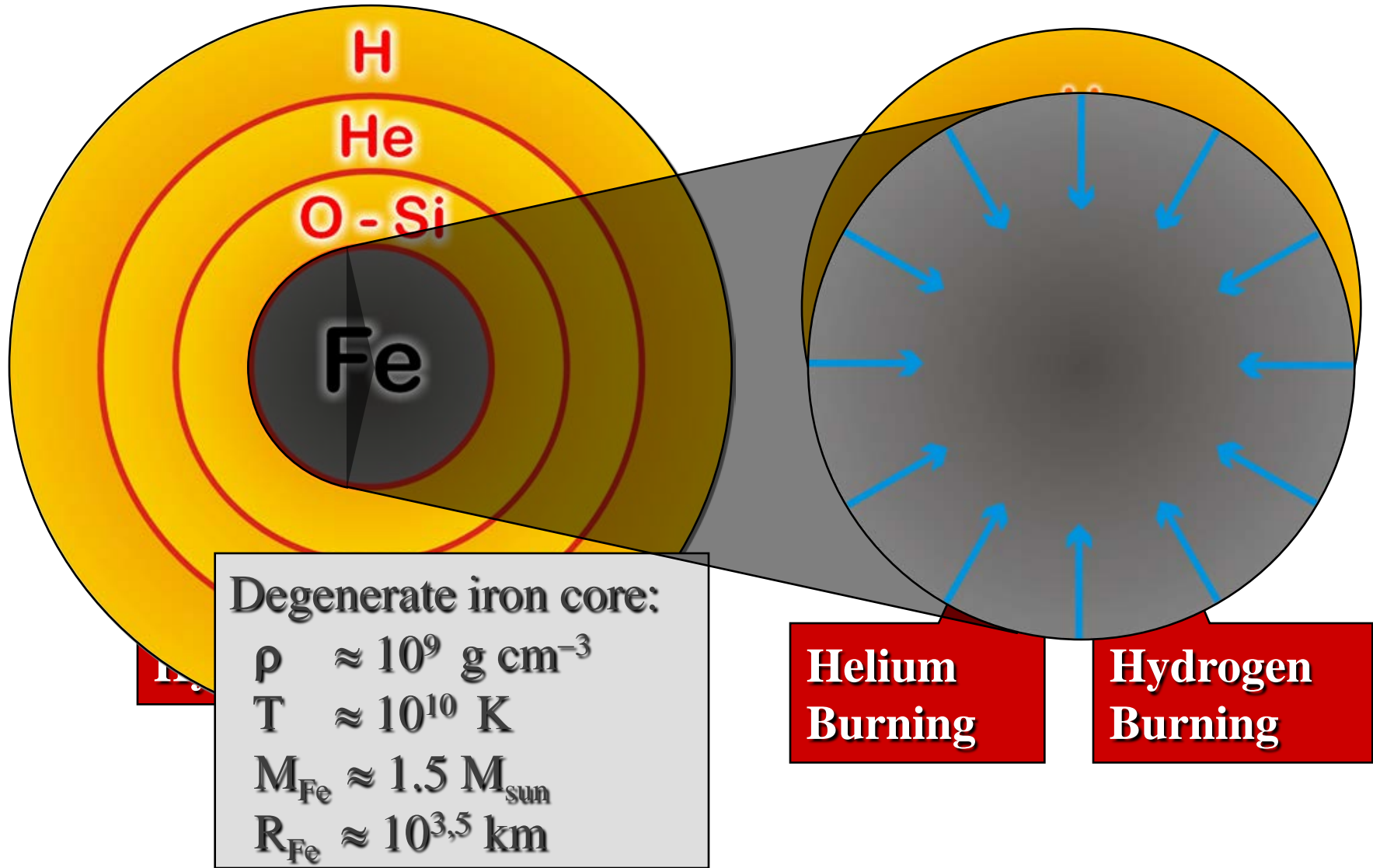
# Massive Star, $M > 8 M_{\odot}$ *onion*



# Stellar Collapse and Supernova Explosion

Onion structure

Collapse (implosion)



Degenerate iron core:

$$\rho \approx 10^9 \text{ g cm}^{-3}$$

$$T \approx 10^{10} \text{ K}$$

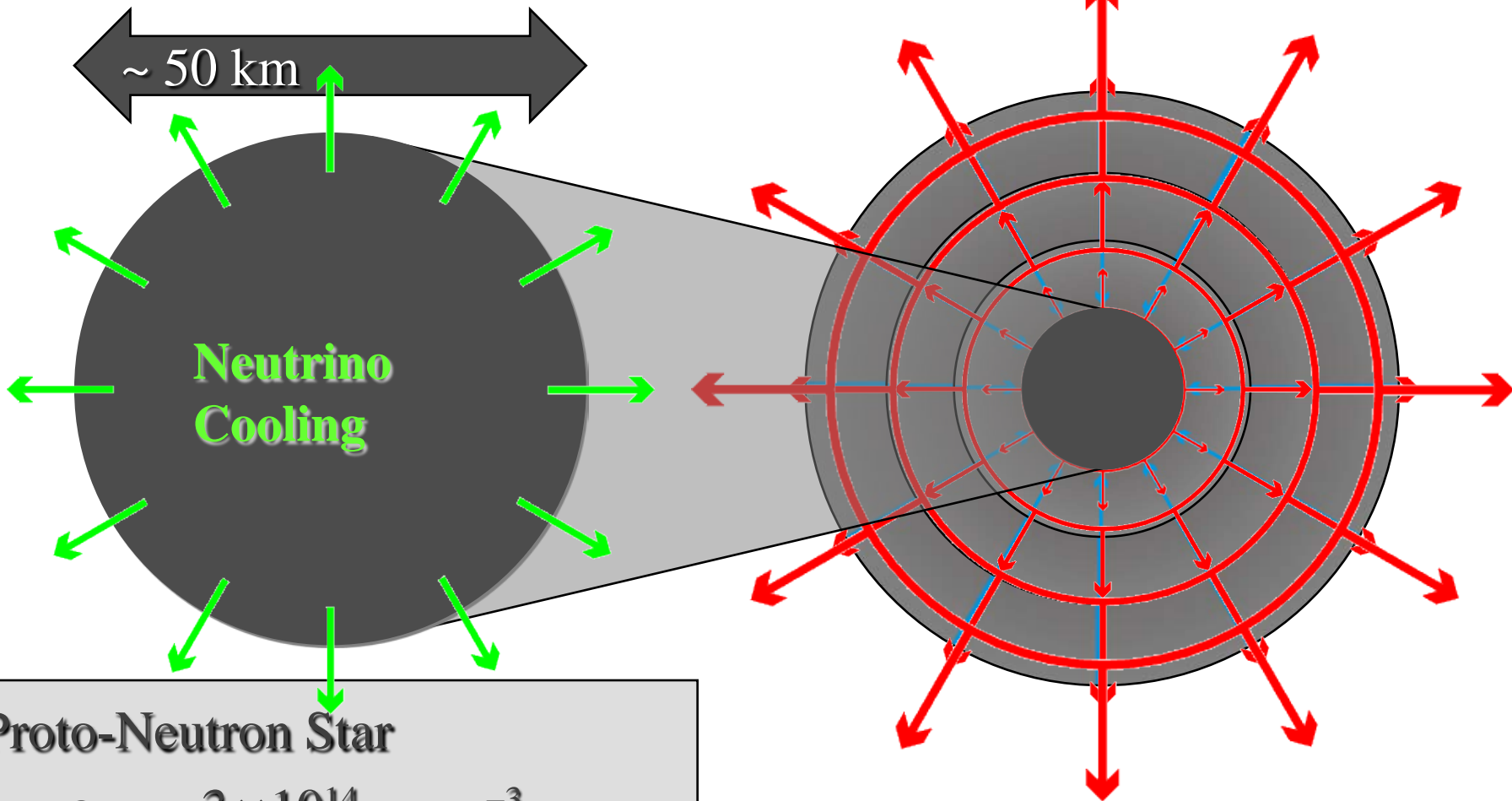
$$M_{\text{Fe}} \approx 1.5 M_{\text{sun}}$$

$$R_{\text{Fe}} \approx 10^{3,5} \text{ km}$$

# Stellar Collapse and Supernova Explosion

**Newborn Neutron Star**

**Explosion**



**Proto-Neutron Star**

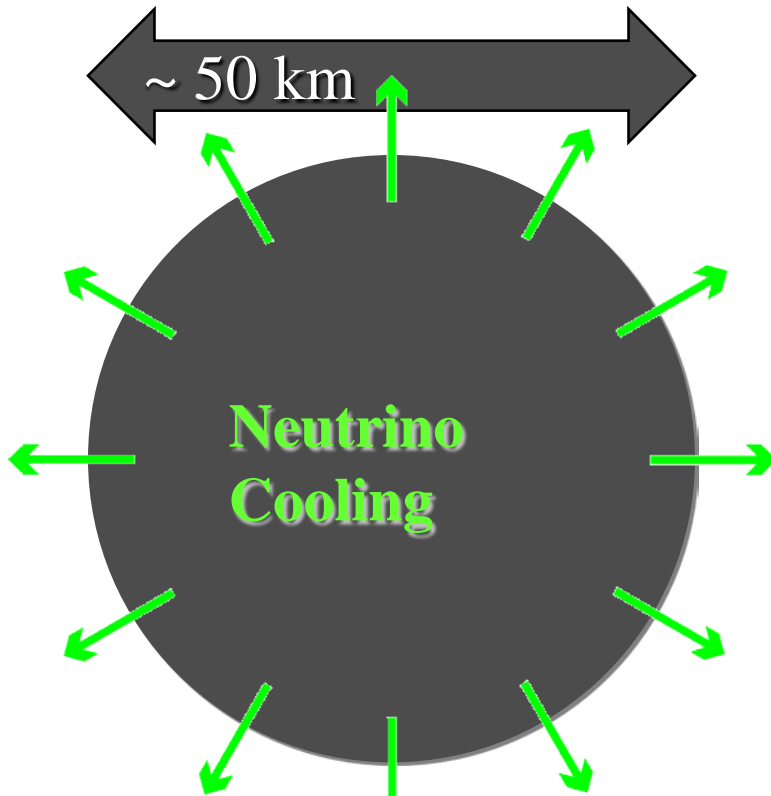
$$\rho \approx \rho_{\text{nuc}} = 3 \times 10^{14} \text{ g cm}^{-3}$$

$$T \approx 30 \text{ MeV}$$



# Stellar Collapse and Supernova Explosion

Newborn Neutron Star



Proto-Neutron Star

$$\rho \approx \rho_{\text{nuc}} = 3 \times 10^{14} \text{ g cm}^{-3}$$

$$T \approx 10 \text{ MeV}$$

Gravitational binding energy

$$E_b \approx 10^{53,5} \text{ erg} \approx 20\% M_{\text{SUN}} c^2$$

This is distributed as

99% Neutrinos

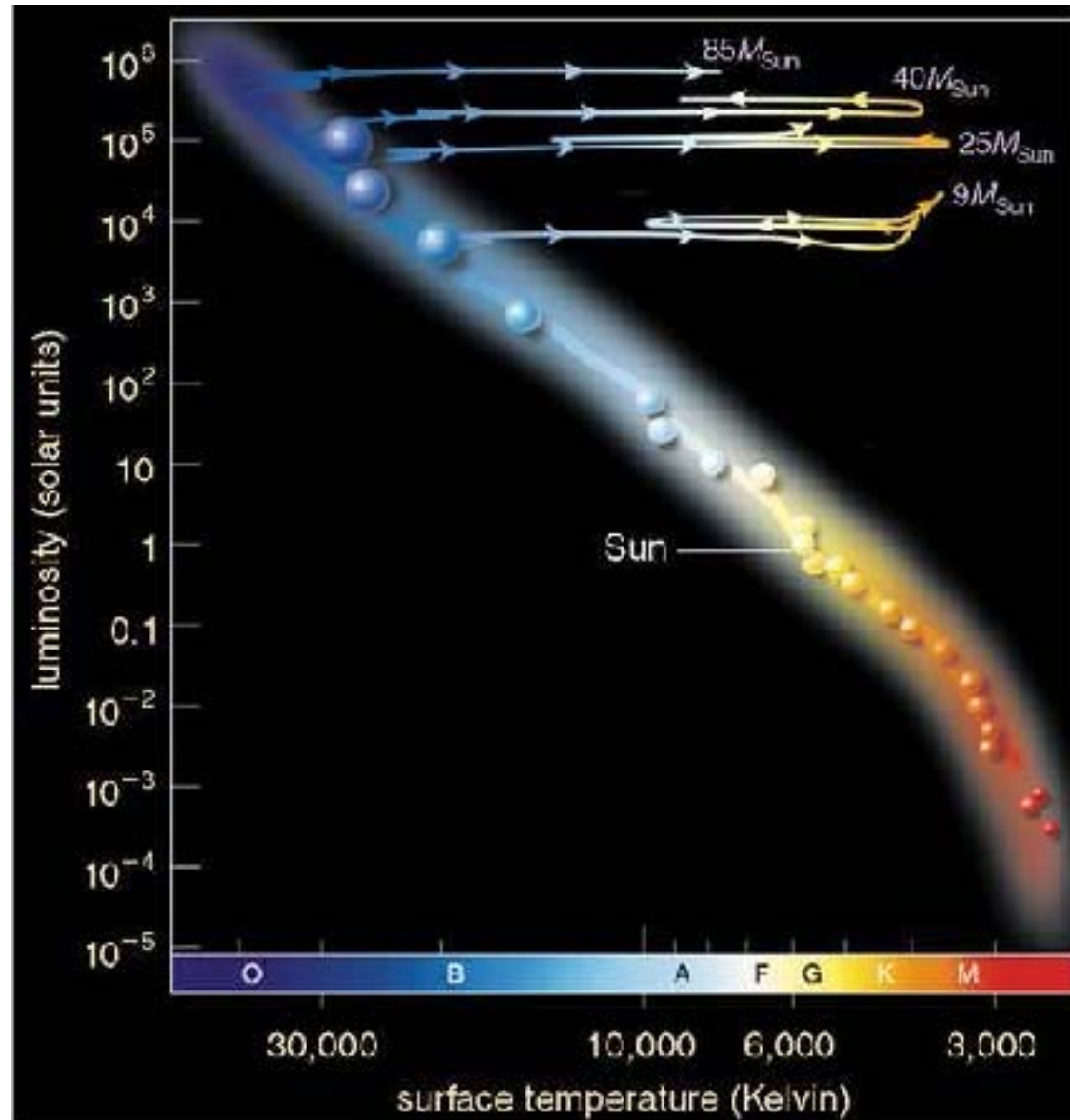
1% Kinetic energy of explosion  
(1% of this into cosmic rays)

0.01% Photons, outshine host galaxy

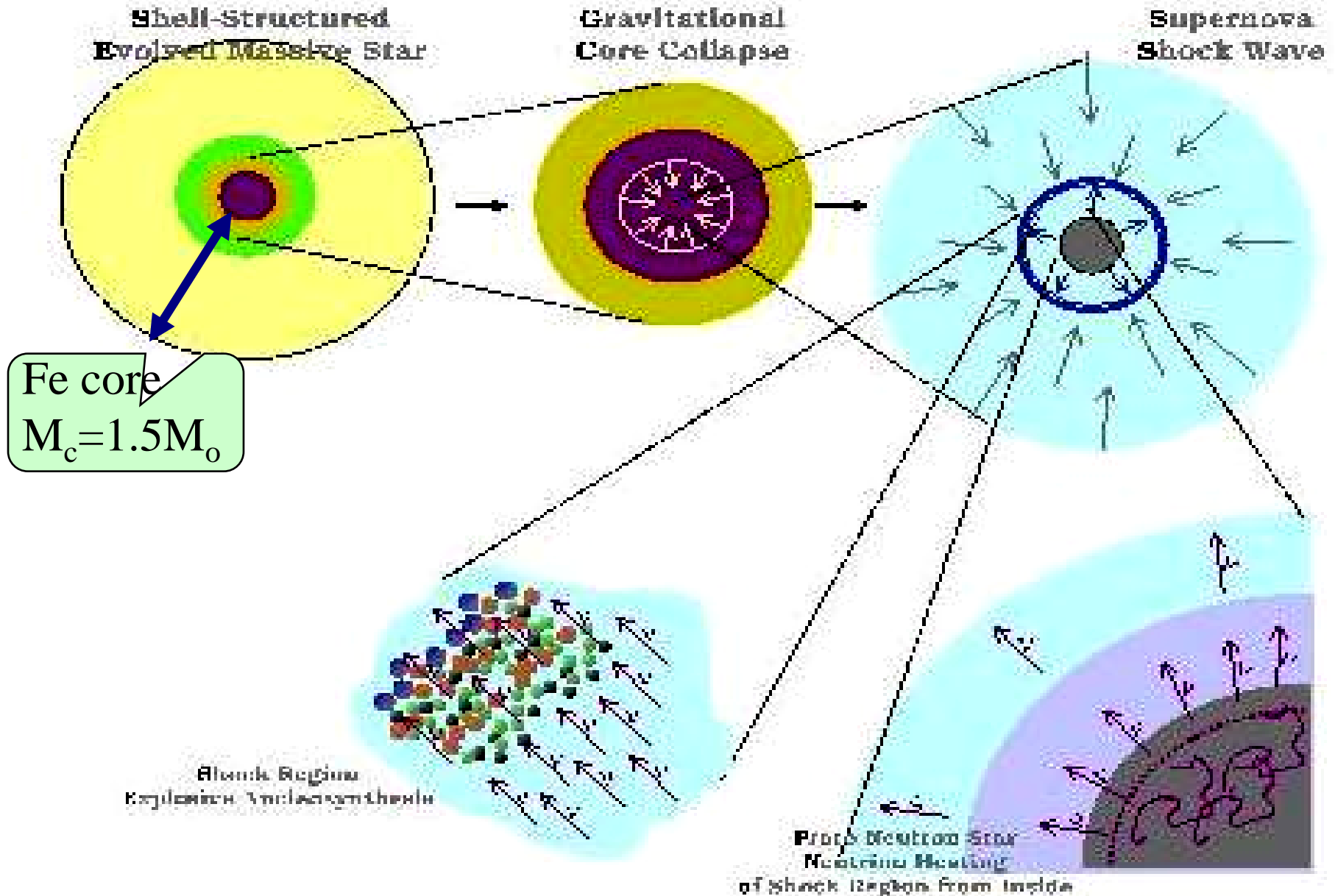
# Core-collapse supernova

*high-mass ( $M > 8M_{\odot}$ ) star  
Evolution on HR diagram*

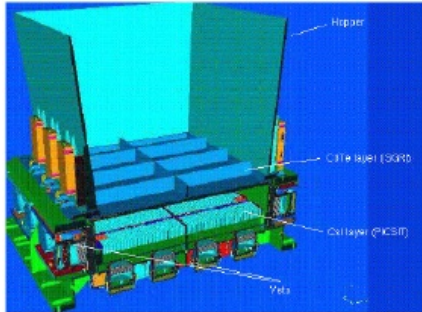
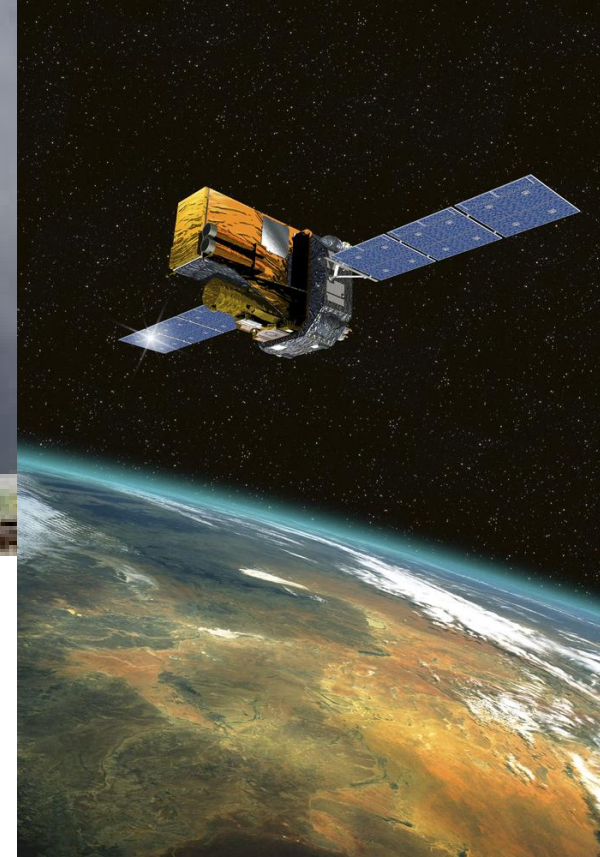
explosive  
nucleosynthesis  
origin of  
Heavy Nuclides



# Explosive nucleosynthesis



# INTEGRAL VIRGO.UA



## IBIS/ISGRI

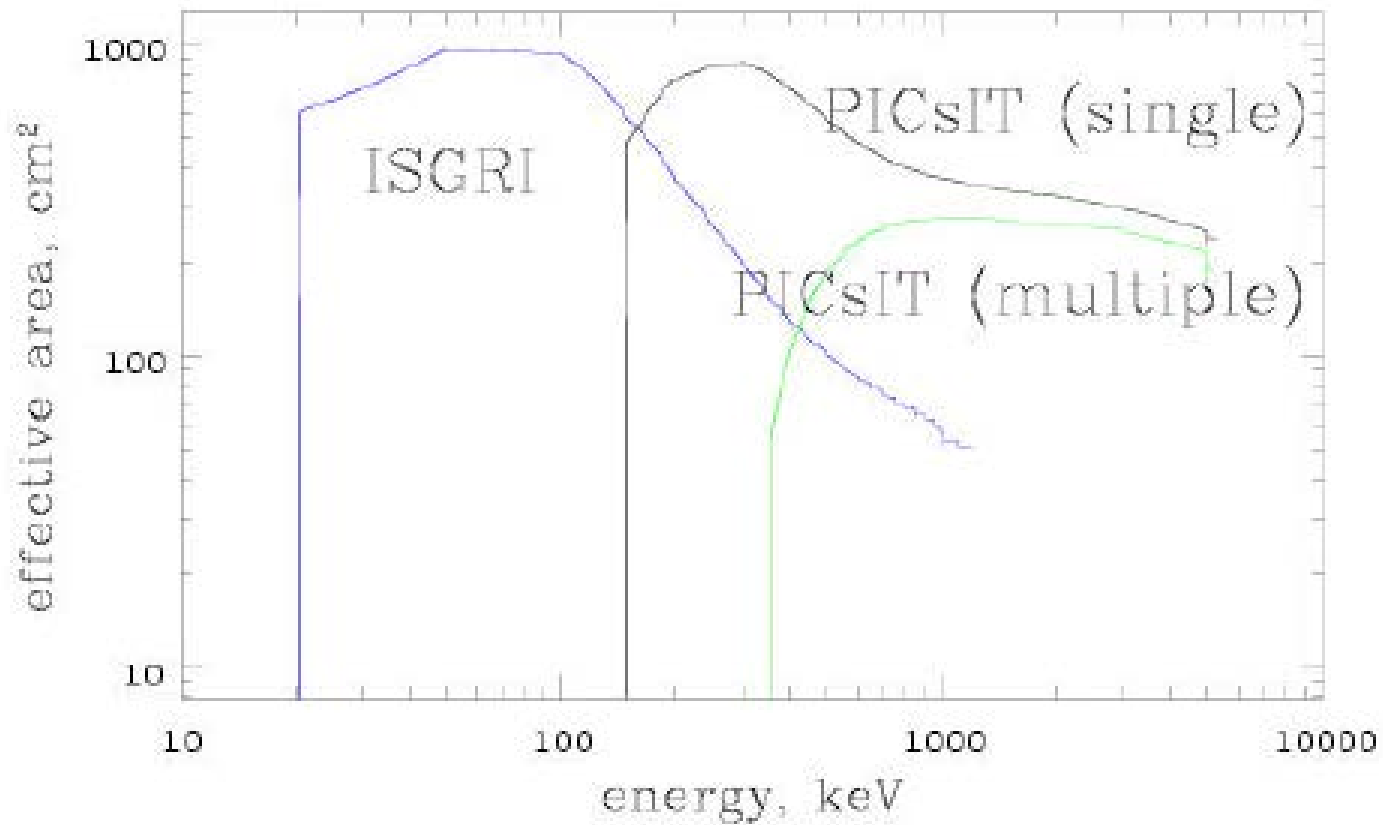
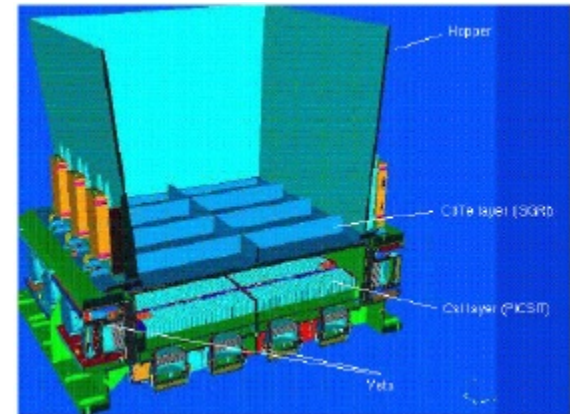
Energy range	20 keV – 1 MeV
Energy resolution (FWHM)	7% at 100 keV
Detector area	960 cm <sup>2</sup> at 50 keV

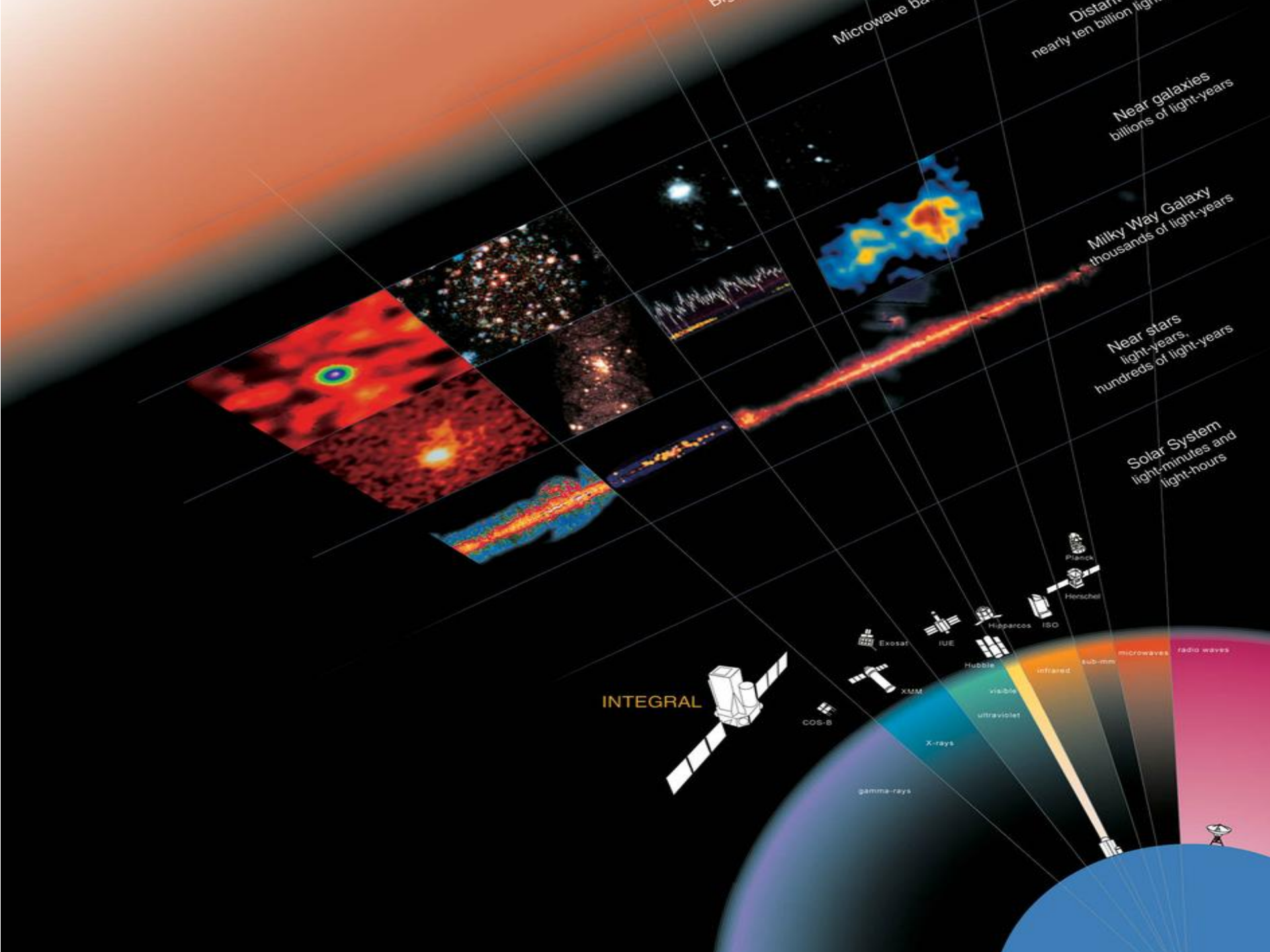


## SPI

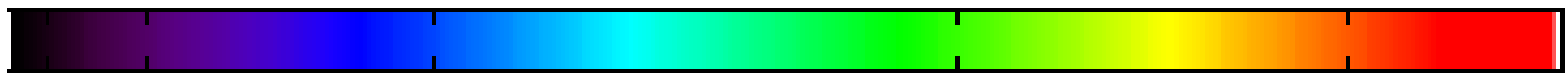
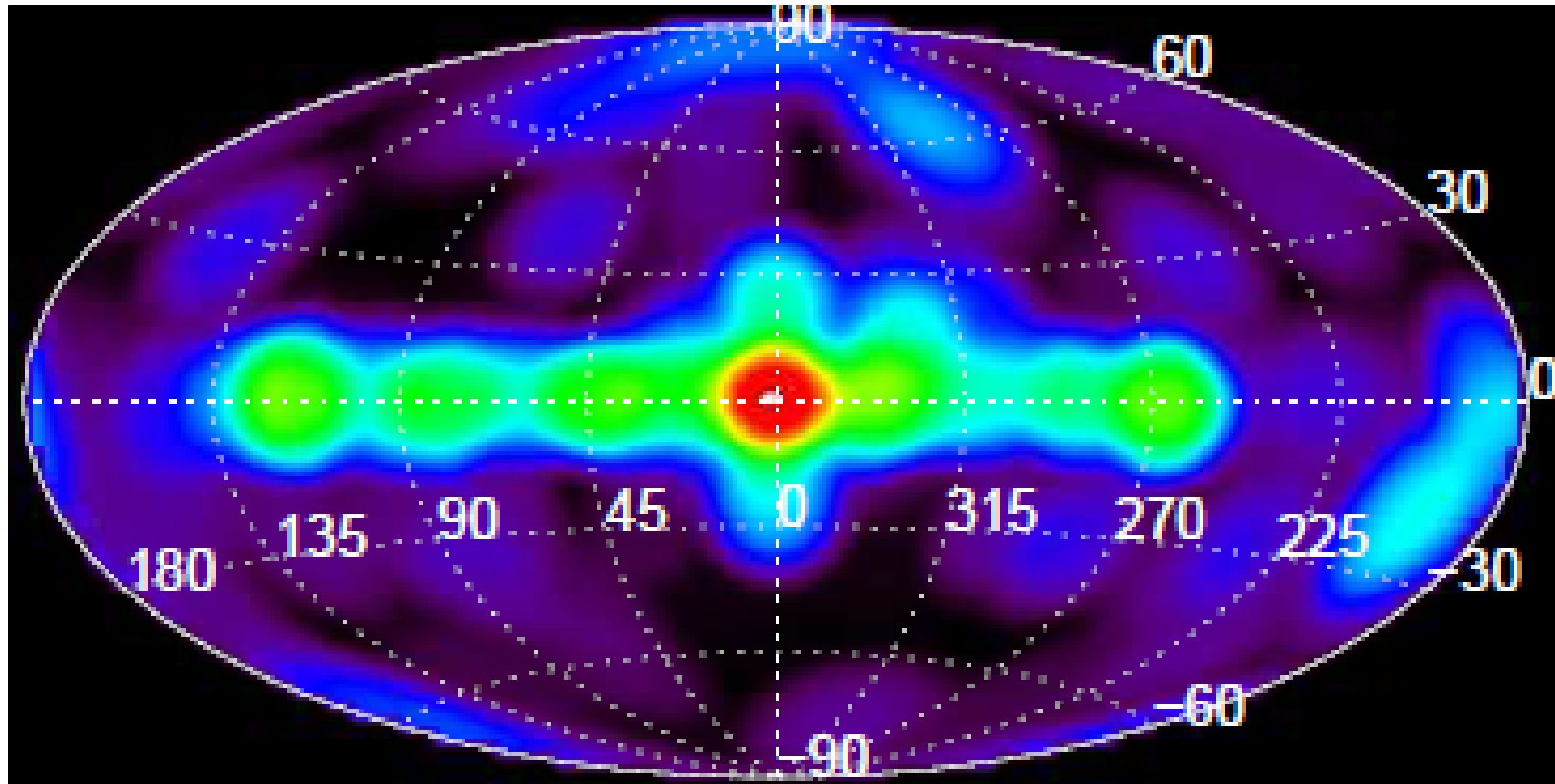
Energy range	20 keV – 8 MeV
Energy resolution (FWHM)	2.35 keV at 1.33 MeV
Detector area	~500 cm <sup>2</sup>

# Integral IBIS/ISGRI



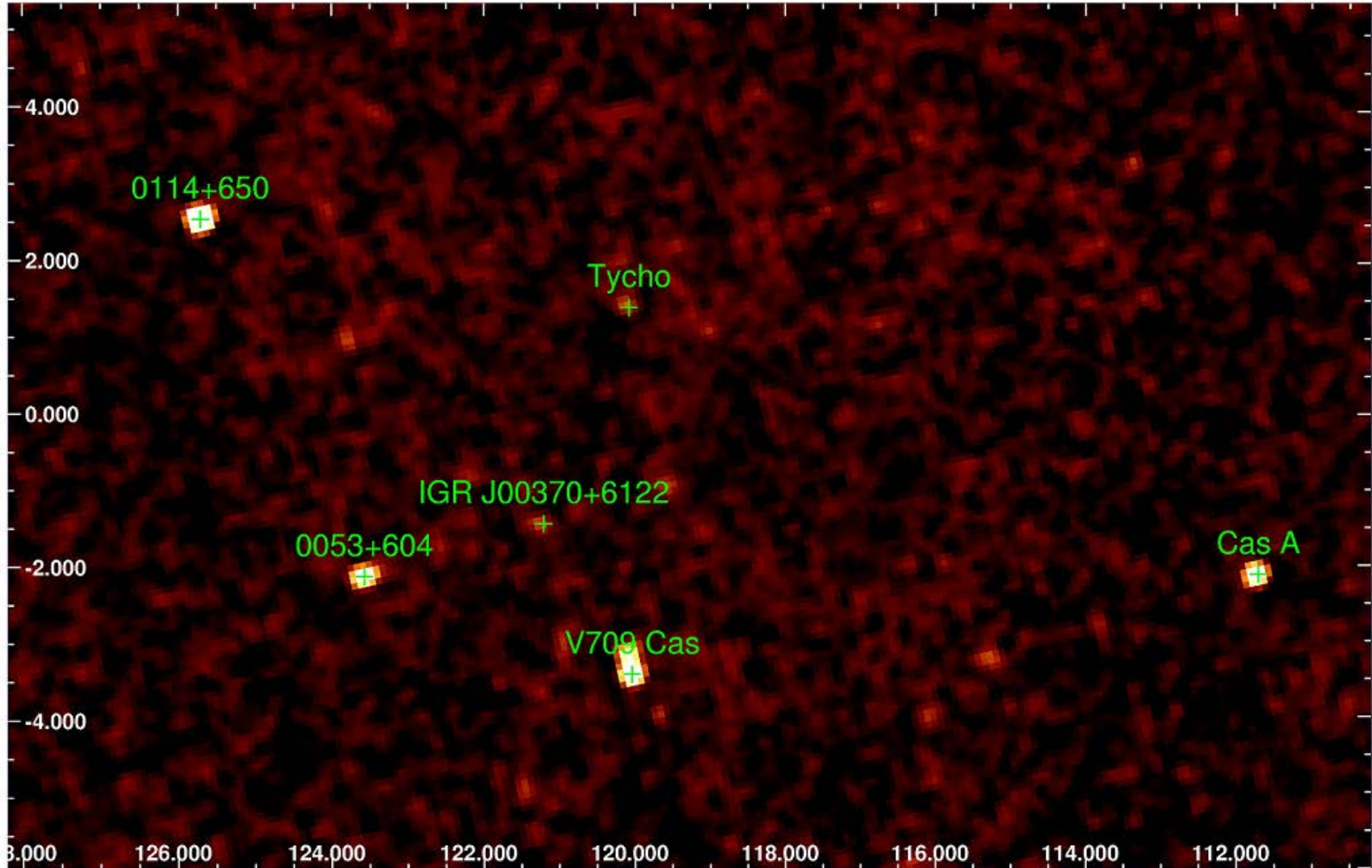


# INTEGRAL Exposure of sky in Galactic coordinates (ksec)



0  $10^2$   $10^3$   $5 \times 10^3$   $10^4$

# CAS A( $3.4_{+0.3-0.1}$ )kpc, TYCHO( $2.2_{\pm 0.3}$ )kpc



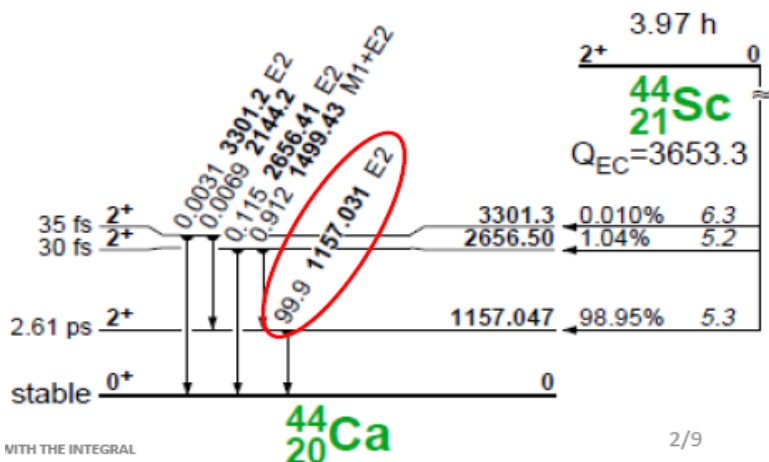
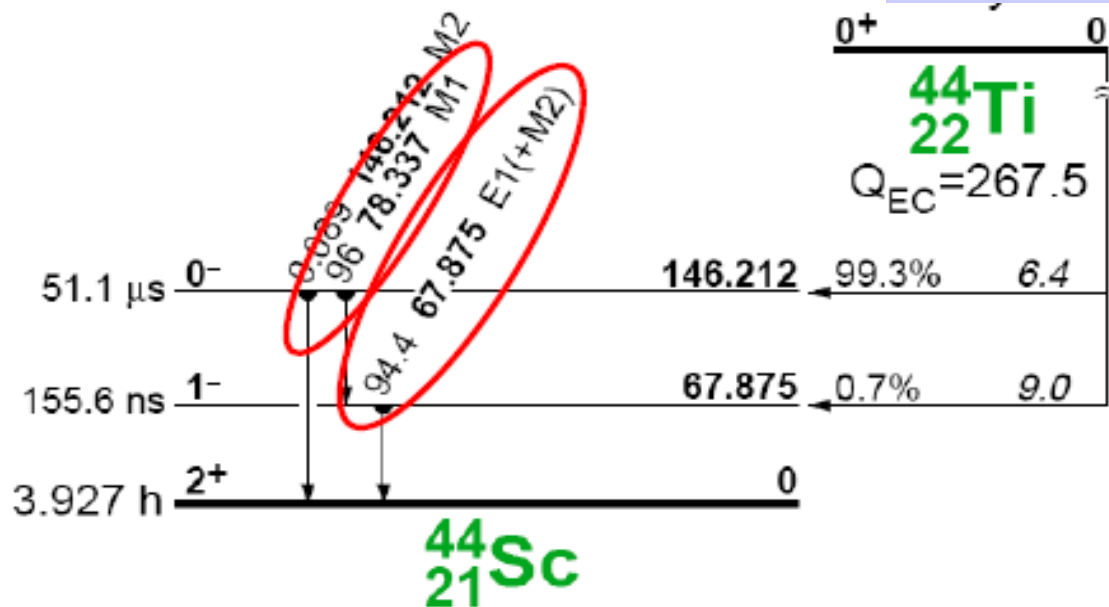


# Investigated Supernovas

Name	Coordinate	Age
SN 1987A	- 279.7 -31.9,	<b>20 y</b>
Cas A	- 111.7 -2.1,	<b>330 y</b>
TYCHO	- 120.1 +1.4,	<b>440 y</b>
Vela Junior	- 266.3 -1.2	<b>?</b>

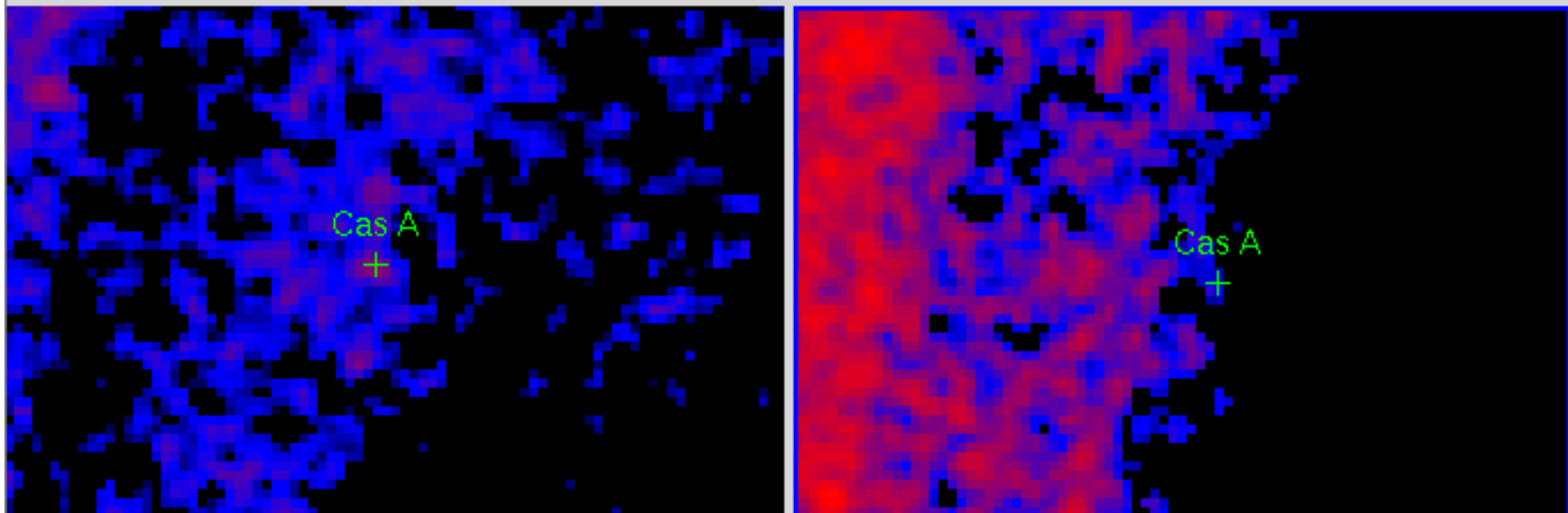
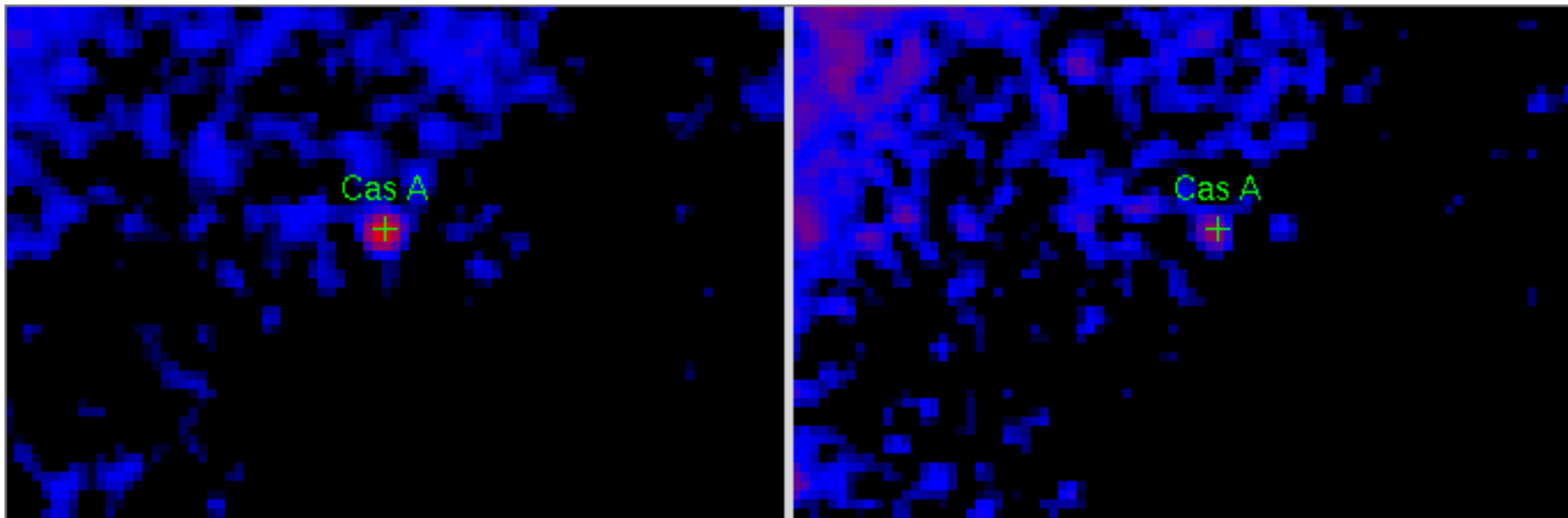
# The scheme of the $^{44}\text{Ti}$ decay *Earth environment*

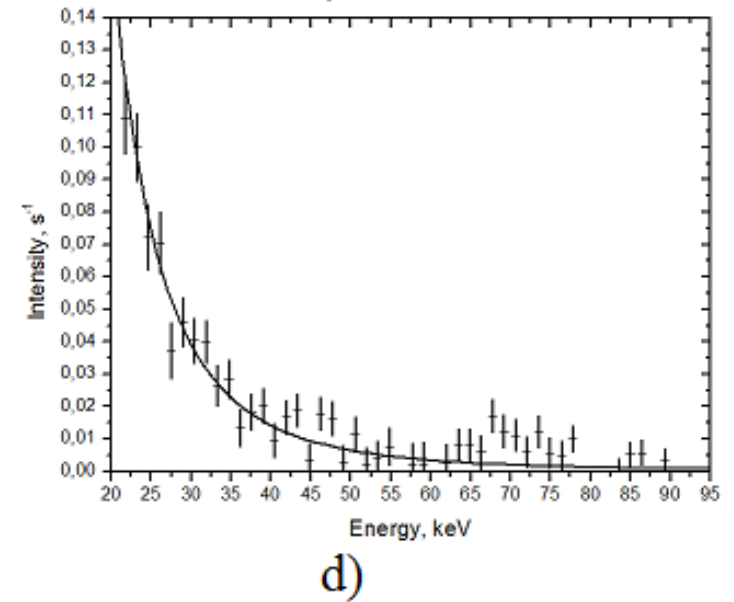
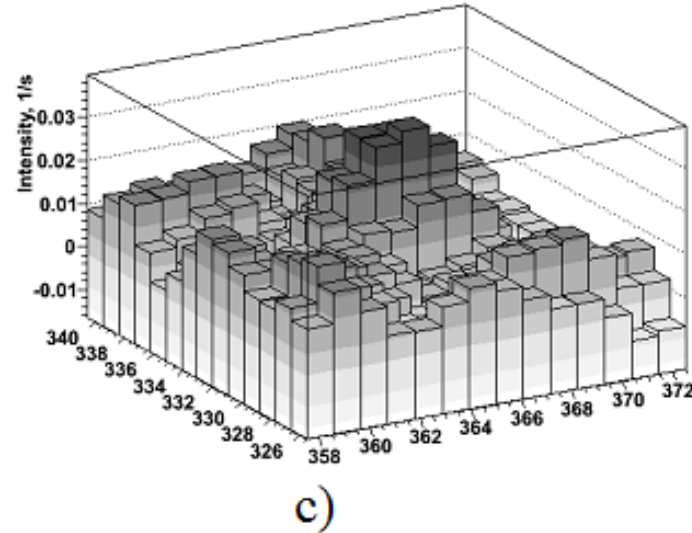
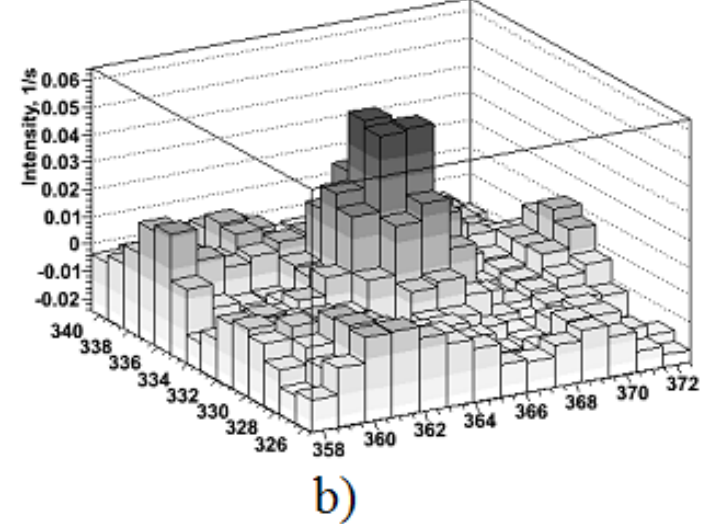
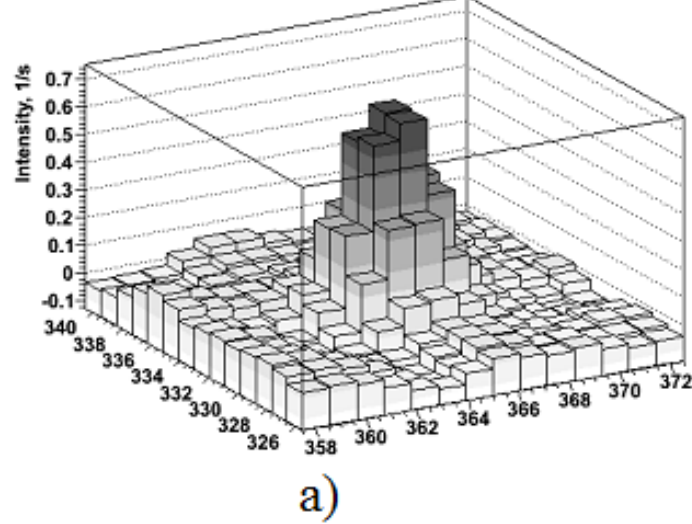
59 y



# Cassiopeia A $(3.4+0.3-0.1)\text{kpc}$

Energy range (keV): **20-62-72-82-100**



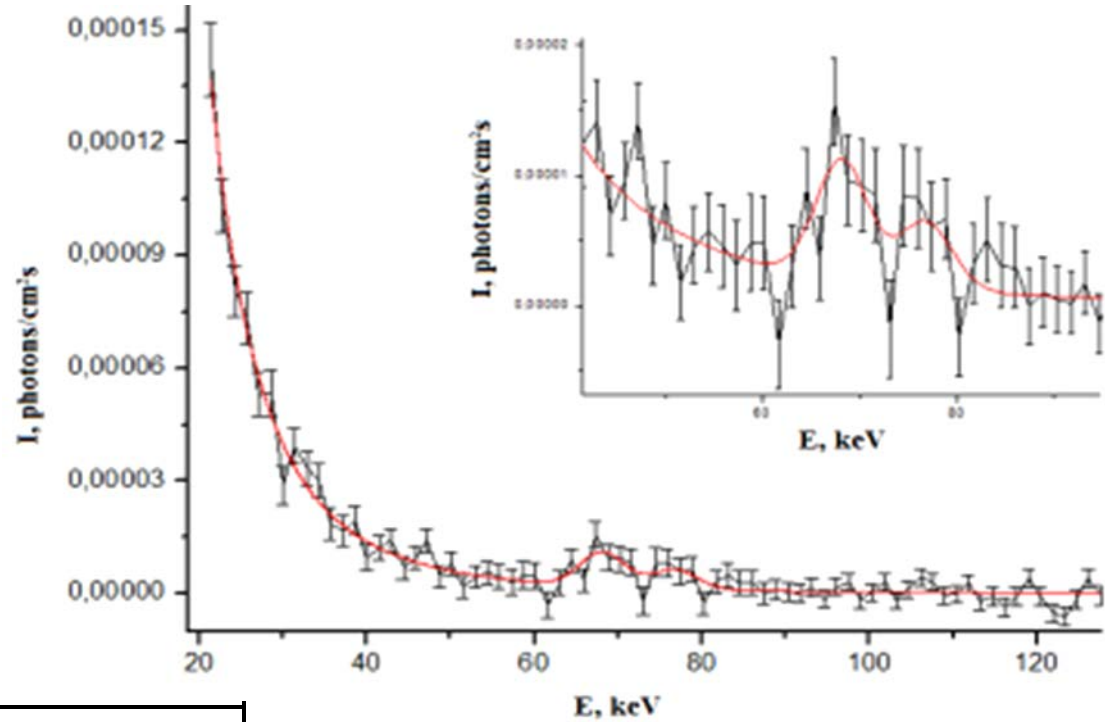


F The direction (i.e.. pixel number) dependence of the registered gamma-ray flux at different energy ranges: **20–62 keV - a**, **62–72 keV - b**, **72–82 keV – c**; for the angle region containing the Cassiopeia A SN remnant. The right bottom panel (d) represents the spectrum from the Cassiopeia A in the energy range 20–95 keV, the solid line shows the fit with the power law energy  $E$  dependence, .

# Background estimate + Lines

$K \cdot E^a$

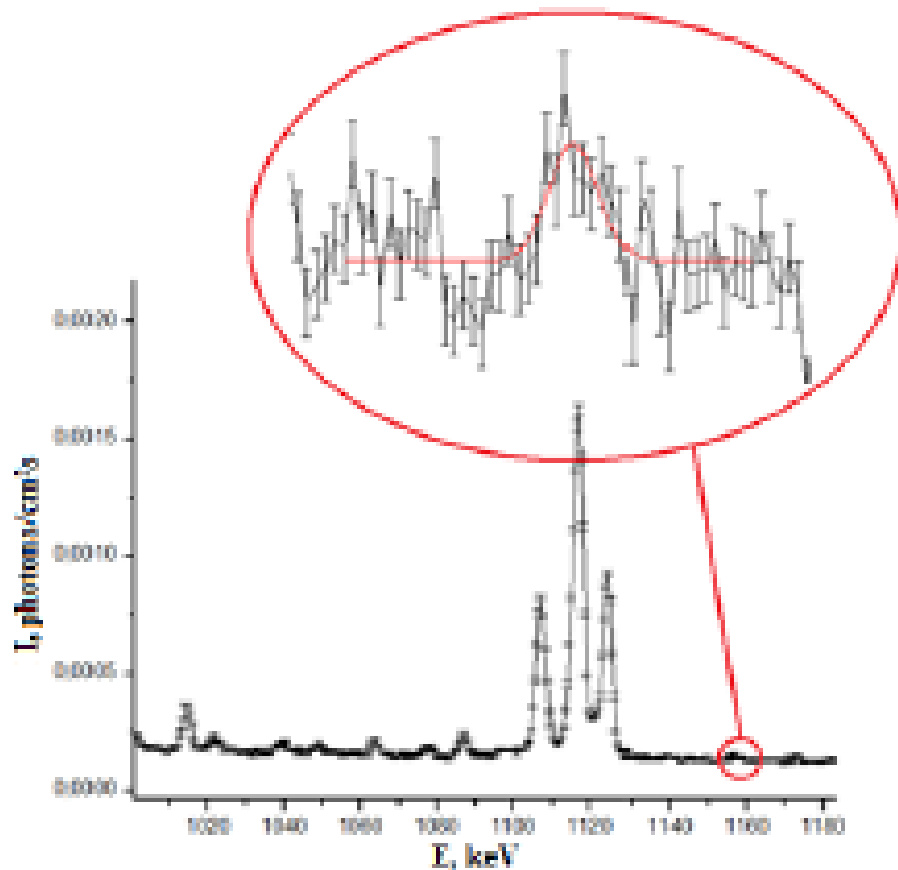
$$a = -3.64 \pm 0.09$$



$E, \text{keV}$	$I \pm \Delta I, \text{photons}/(\text{cm}^2\text{s})$
67.9	$(6.0 \pm 1.0) \cdot 10^{-5}$
78.3	$(4.0 \pm 1.0) \cdot 10^{-5}$

$T = 1.5 \text{Ms}$

# SPI detector data



*Fit results*

$$E = 1157.20 \pm 0.26 \text{ keV}$$

$$FWHM = 3.1 \pm 0.7 \text{ keV}$$

$$Flux = (5.1 \pm 1.0) \cdot 10^{-5} \text{ ph/cm}^2 \text{ s}$$

T=1.5 Ms

# mass of $^{44}\text{Ti}$ synthesized at Cassiopeia A explosion [VNK et al 'Nucleus2004'; PhAN (2009)]

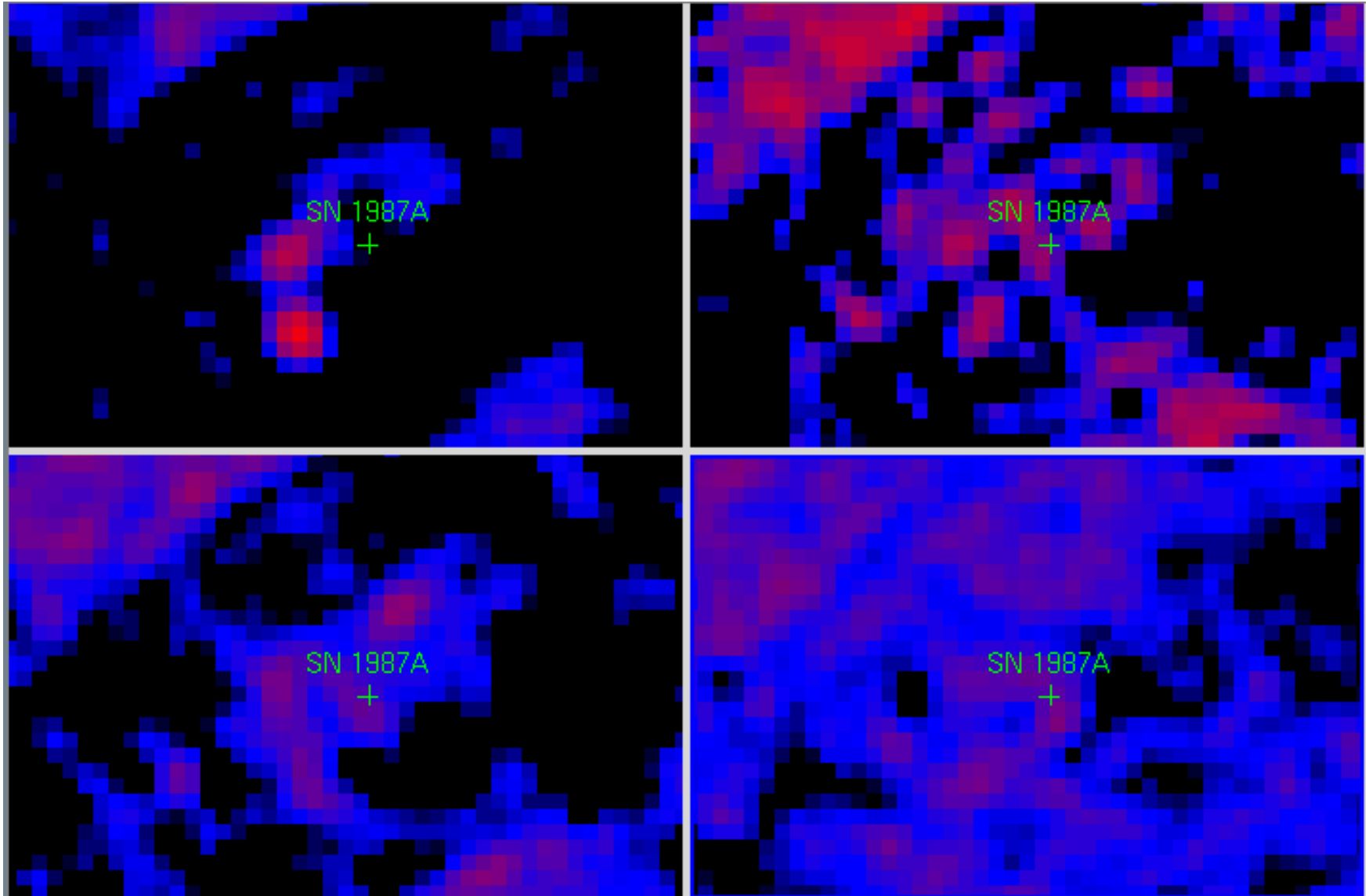
$$m = \frac{4\pi R^2 \cdot T_{1/2} \cdot M \cdot I}{\ln 2 \cdot N_a \cdot p} \cdot 2^{-\frac{t}{T_{1/2}}}$$

R—distance to the object,  $T_{1/2}$ —the element half-life,  $N_a$ —the Avogadro constant,  $M$ —molar mass,  $I$ — $\gamma$ -quanta flux,  $p$ —quantum yield,  $t$ —remnant age.

E, keV	$I \pm \Delta I,$ $10^{-5}$ photons/cm <sup>2</sup> s	$m \pm \Delta m, 10^{-4} M_{\odot}$
67.9 keV	$6.0 \pm 1.0$	$4.0 \pm 0.7$
78.3 keV	$4.0 \pm 1.0$	$2.6 \pm 0.7$
1157.1 keV	<b>5.1</b> $\pm 1.1$	<b>3.3</b> $\pm 0.7$

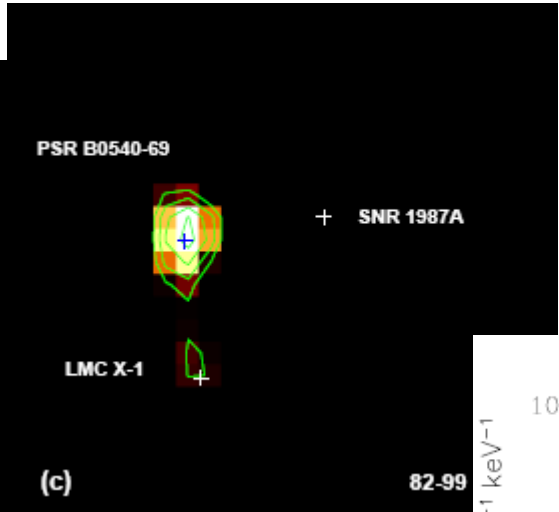
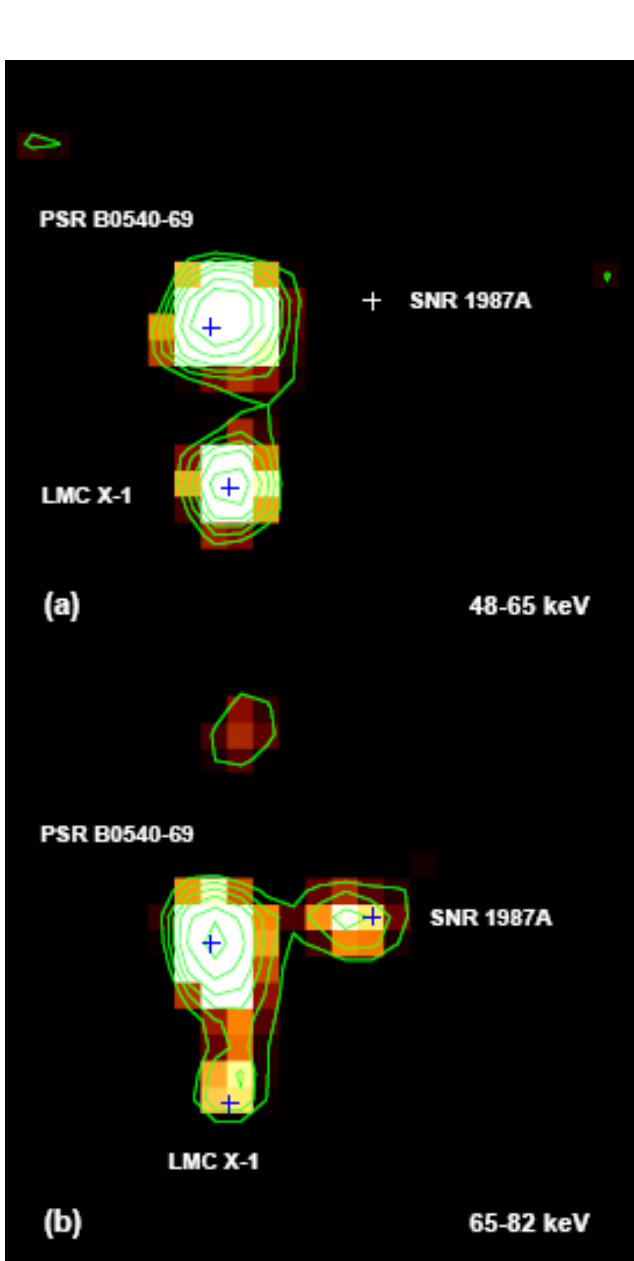
# SN1987 A (50kpc)

Energy range (keV): **20-62-72-82-100**

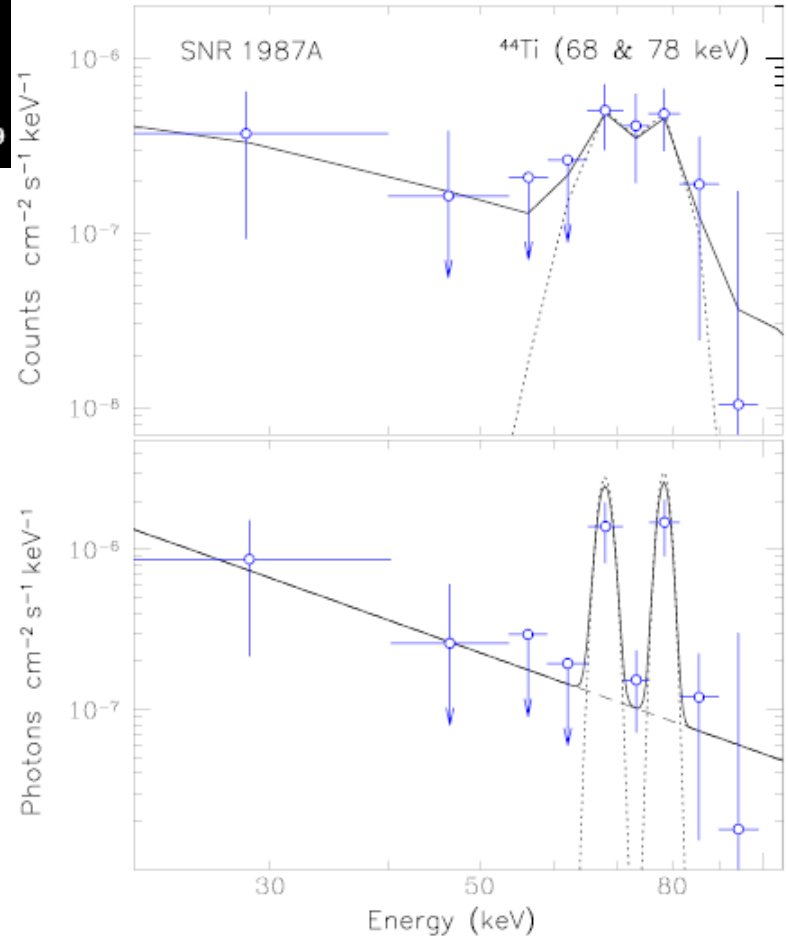




# S.A.Grebenev et al *Nature*, 490, 373-375 (2012).



$$(3.1 \pm 0.8) \times 10^{-4} M_{\odot}$$



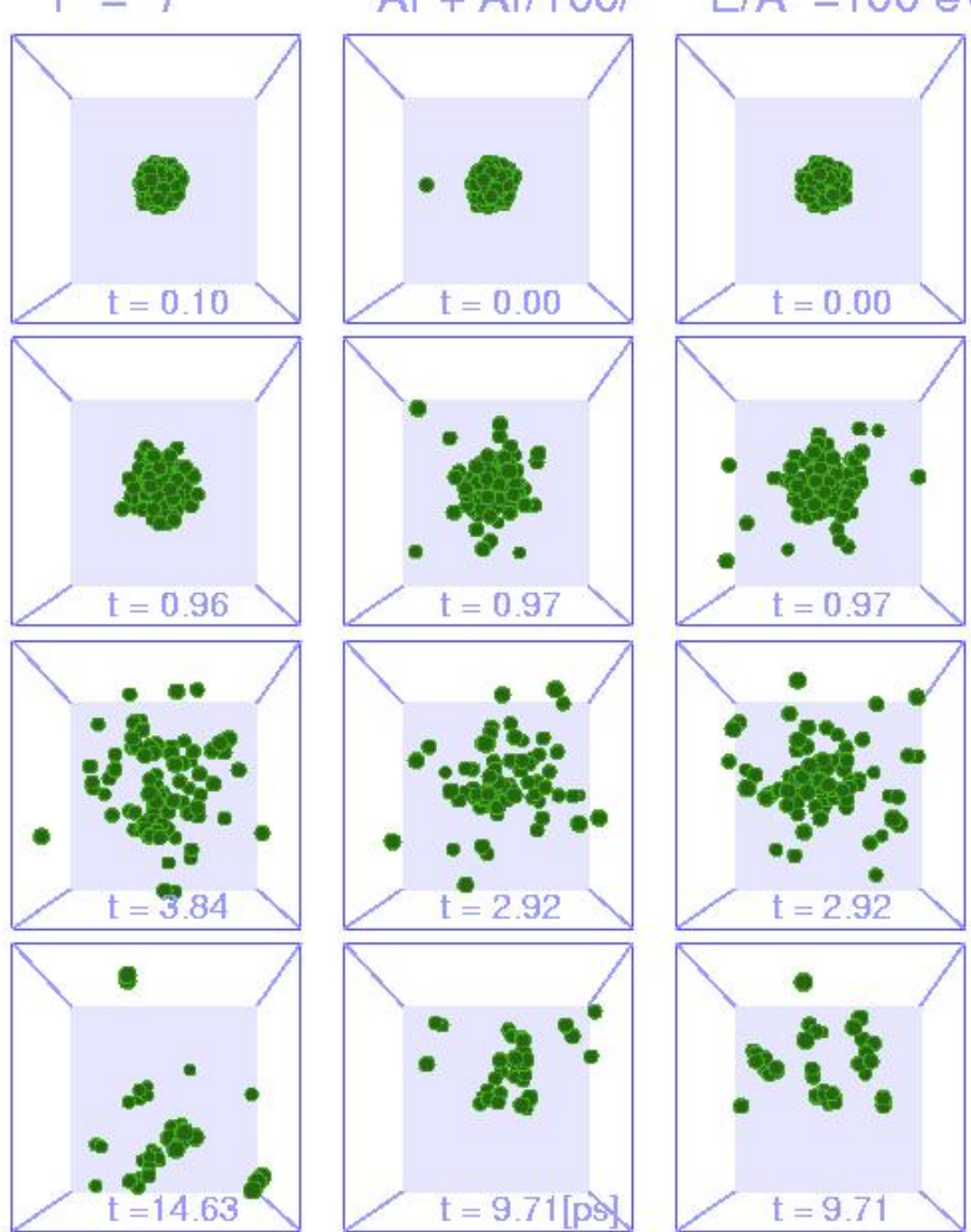
# THEORY

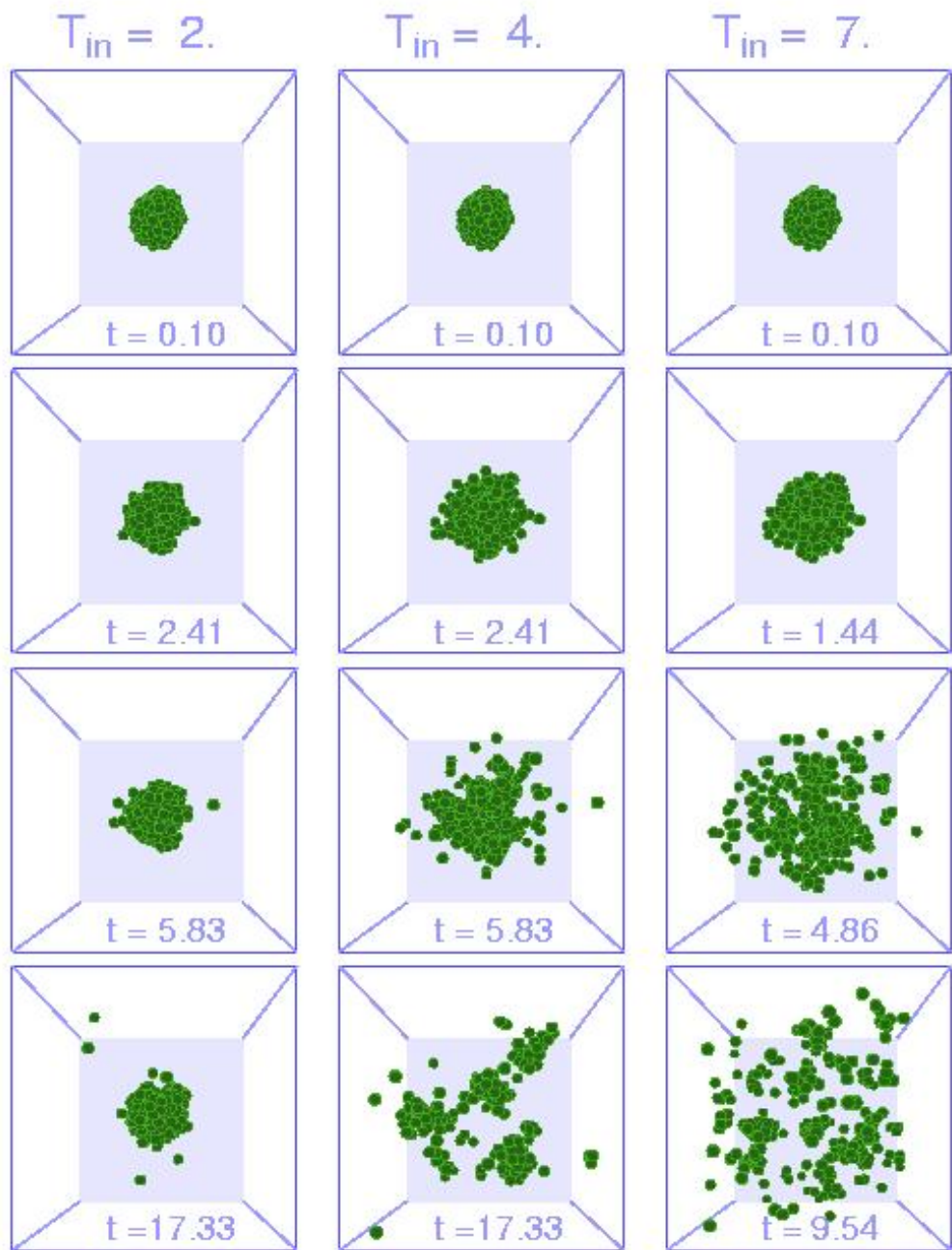
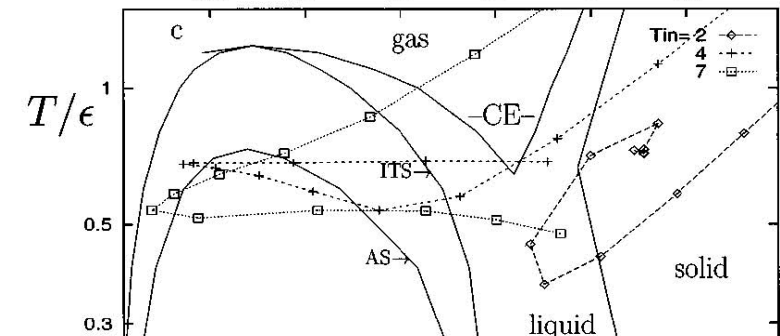
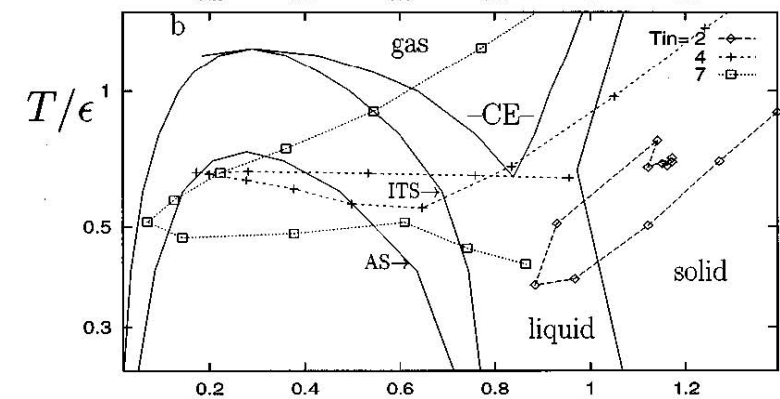
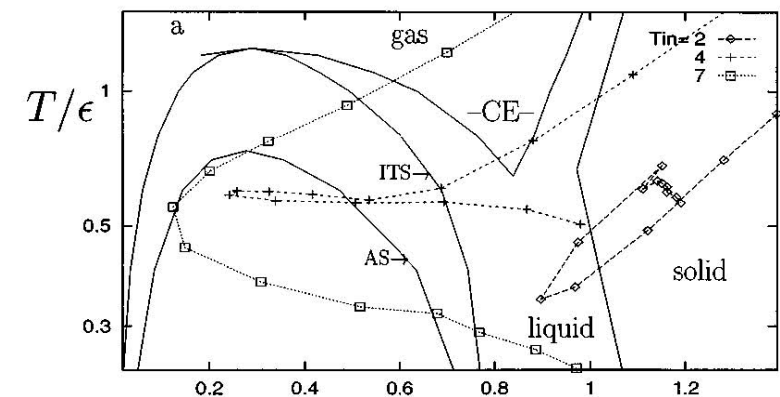
$$(0.02 - 2.5) \times 10^{-4} M_{\odot}$$

Thielemann, F.-K., Hashimoto, M. & Nomoto, K. Explosive nucleosynthesis in SN1987A. II. Composition, radioactivities, and the neutron star mass. *Astrophys. J.* **349**, 222–240 (1990)

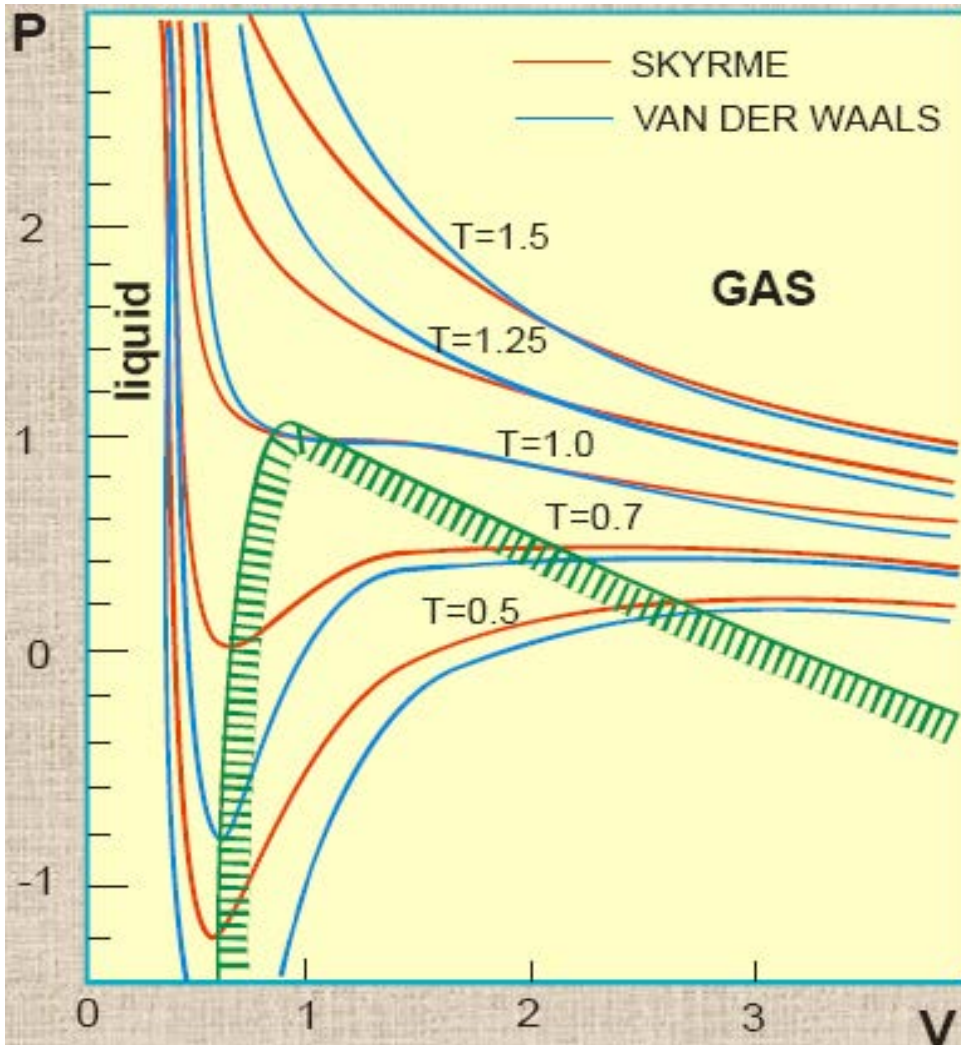
Woosley, S.E., & Hoffman, R. D.  $^{57}\text{Co}$  and  $^{44}\text{Ti}$  production in SN1987A. *Astrophys. J.*, **368**, L31-L34 (1991).







# *State Equation (SE): Nuclear Matter vs Regular Liquid (Noble Gas)*

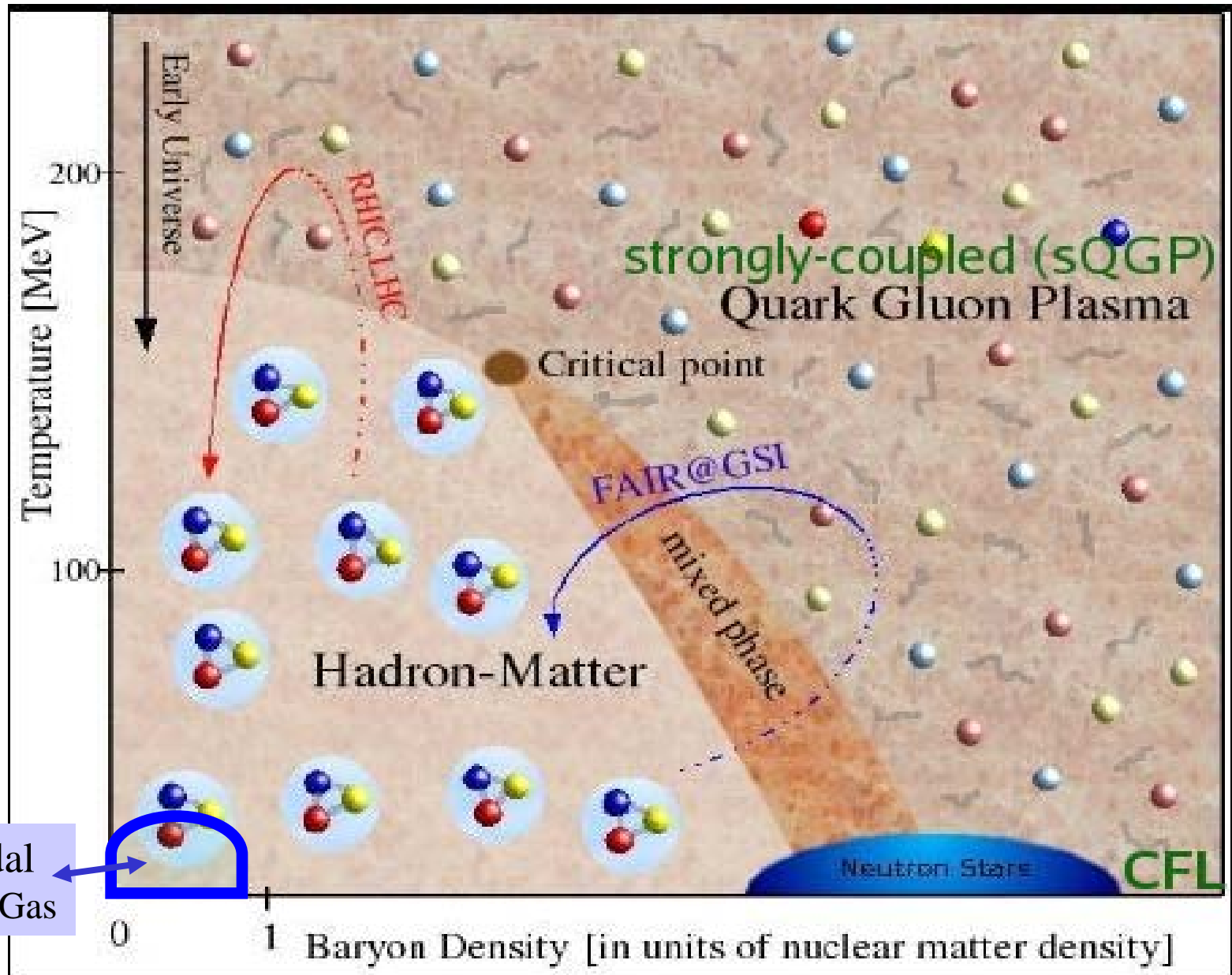


**Van der Waals SE (1875)  
& Nuclear Matter SE  
(1983)**(relative units:  
 $V/V_c$ ;  $p/p_c$  ;  $T/T_c$  )

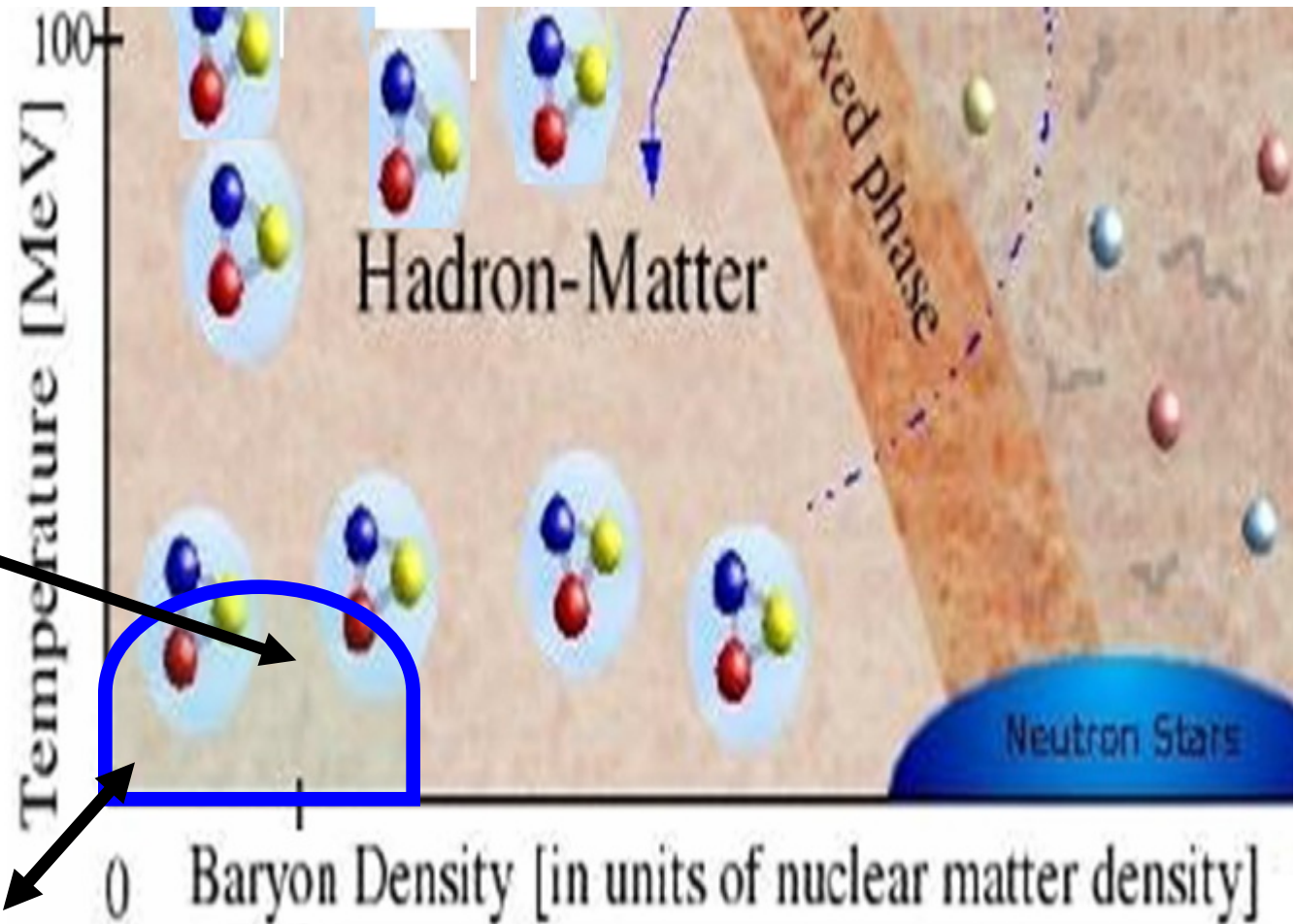
Instability region

$$\left. \frac{\partial p}{\partial V} \right|_T > 0$$

# equation of state vs phase diagram



# equation of state vs phase diagram



Spinodal  
Liquid-Gas

$D=10^9 \text{ g/cm}^3$   
 $T=0.5 \text{ MeV}$

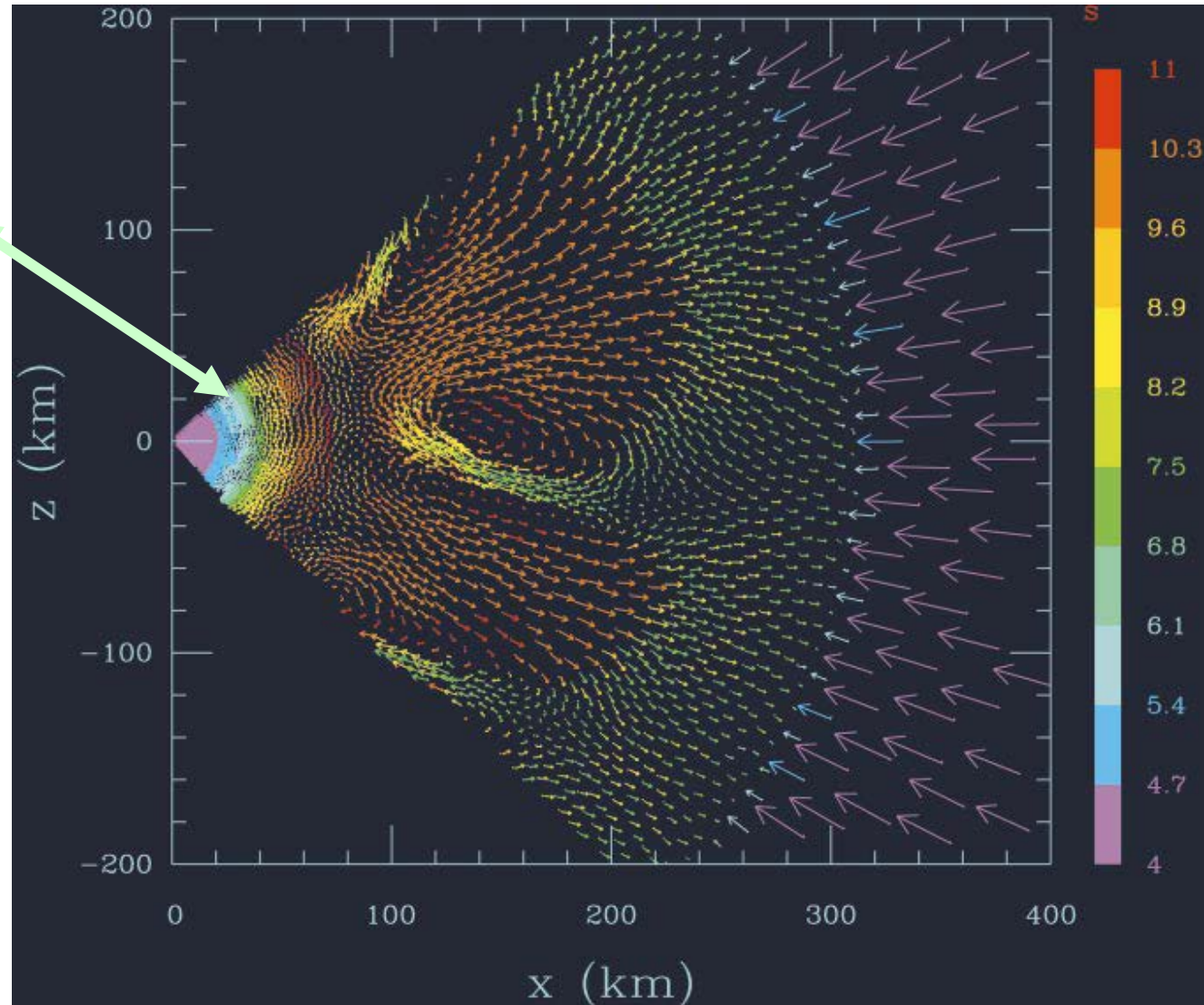


# explosion proceeds through convection processes

$\mathbf{V}$ -sphere

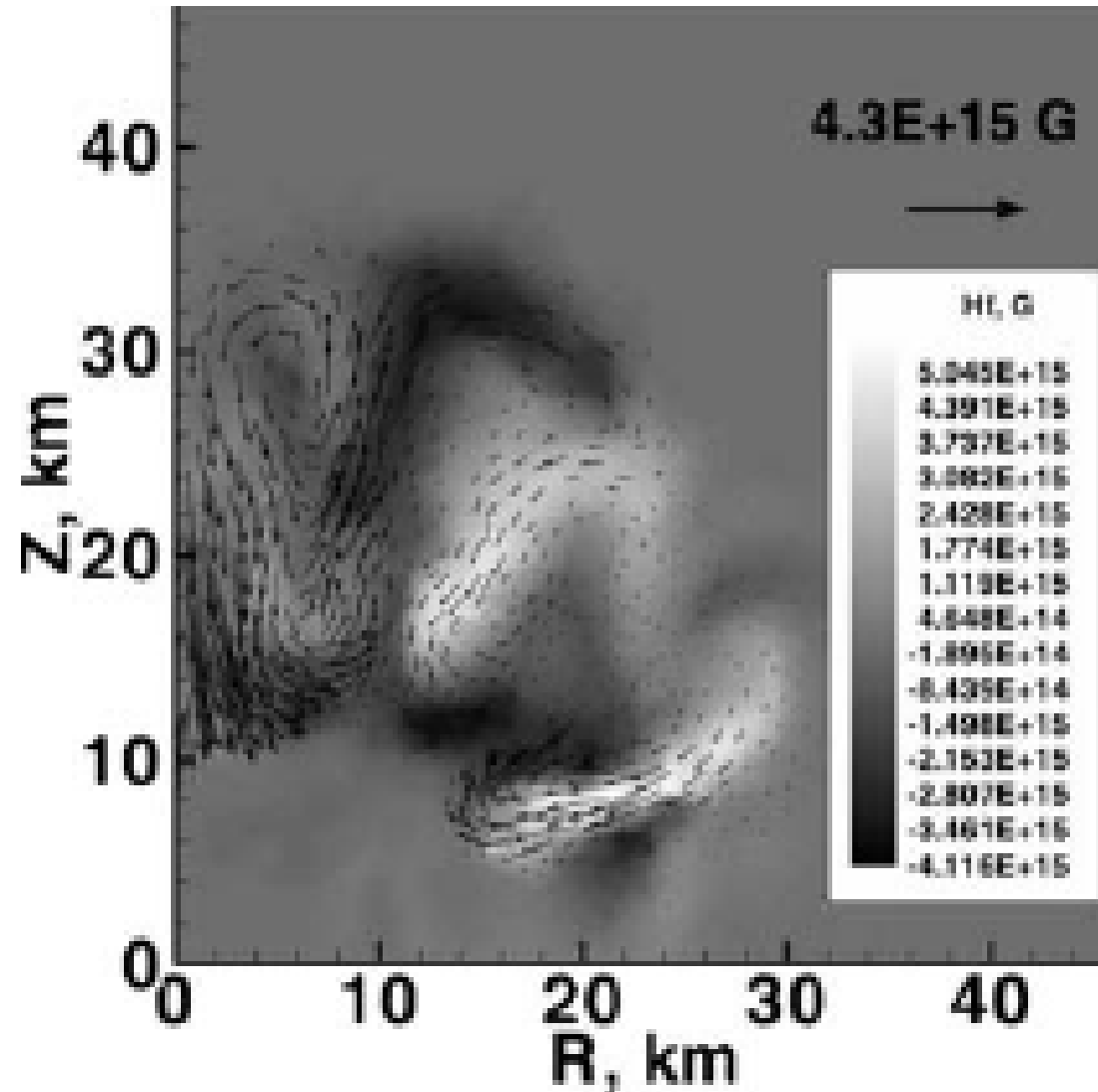
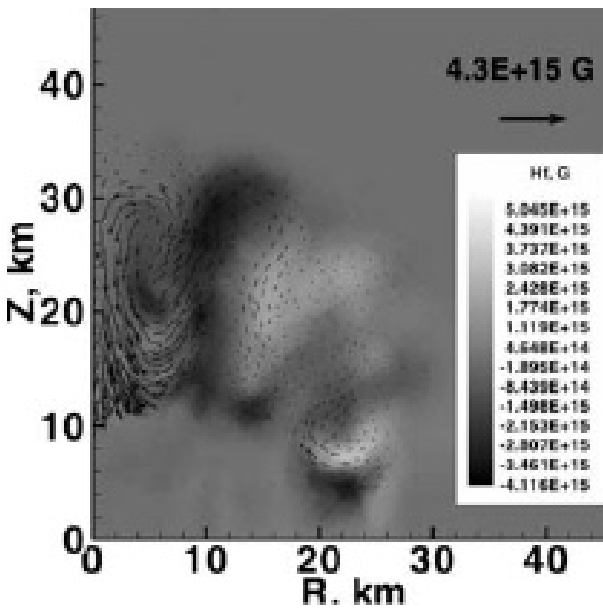
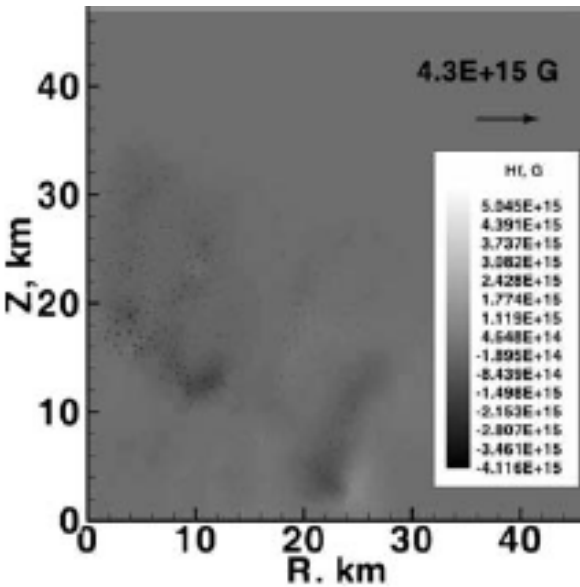
magneto-rotational  
instabilities &  
dynamo-action  
→ amplifying

Magnetic fields  
up to strengths  
hundred *tera-tesla*



# The magnetic field evaluation

(S.G.Moiseenko, G.S.Bisnovatyj-Kogan, N.V.Ardeljan, MNRAS 370 (2006) 501)



# Magnetic field estimates

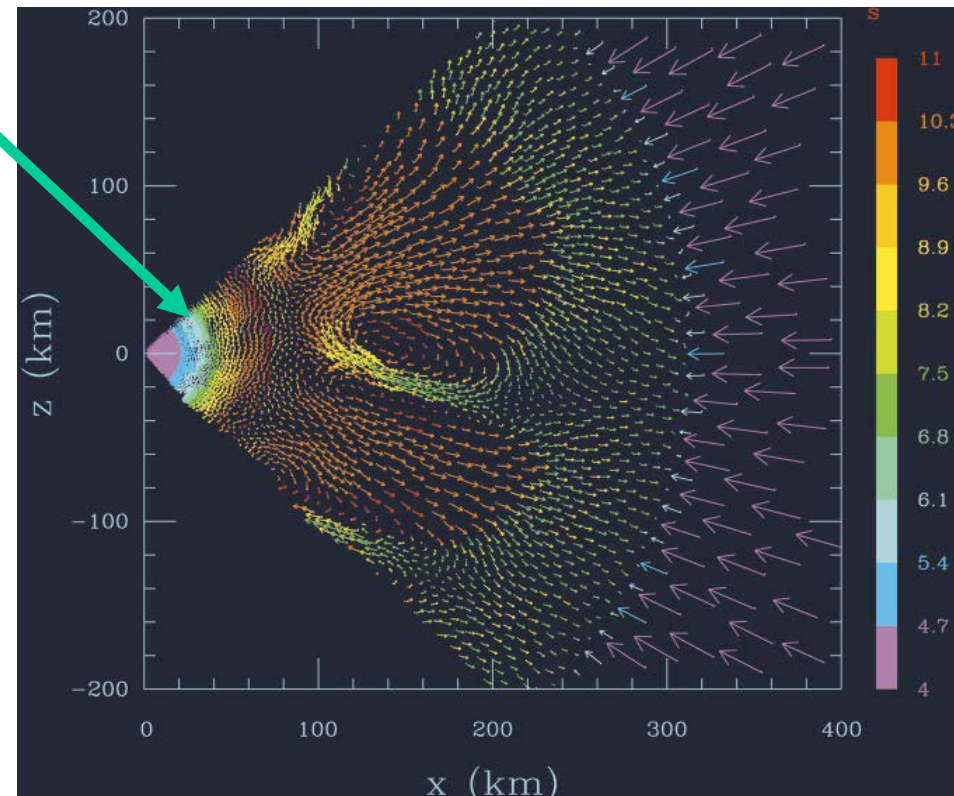
predominant energy component of shock wave  $E_S$   
originates from the magnetic pressure

$$\langle B_v^2 \rangle R_v^2 \Delta R \sim 2E_S \sim 10^{51,5} \text{ ergs}$$

$$R_v \sim 40 \text{ km}; \Delta R \sim 1 \text{ km}$$

$$B_v \sim 10^1 - 10^2 \text{ TeraTesla}$$

$$B(R) \sim B_v \Delta R_v / R$$



# Magnetic field estimates

Magnetic and gravitational forces

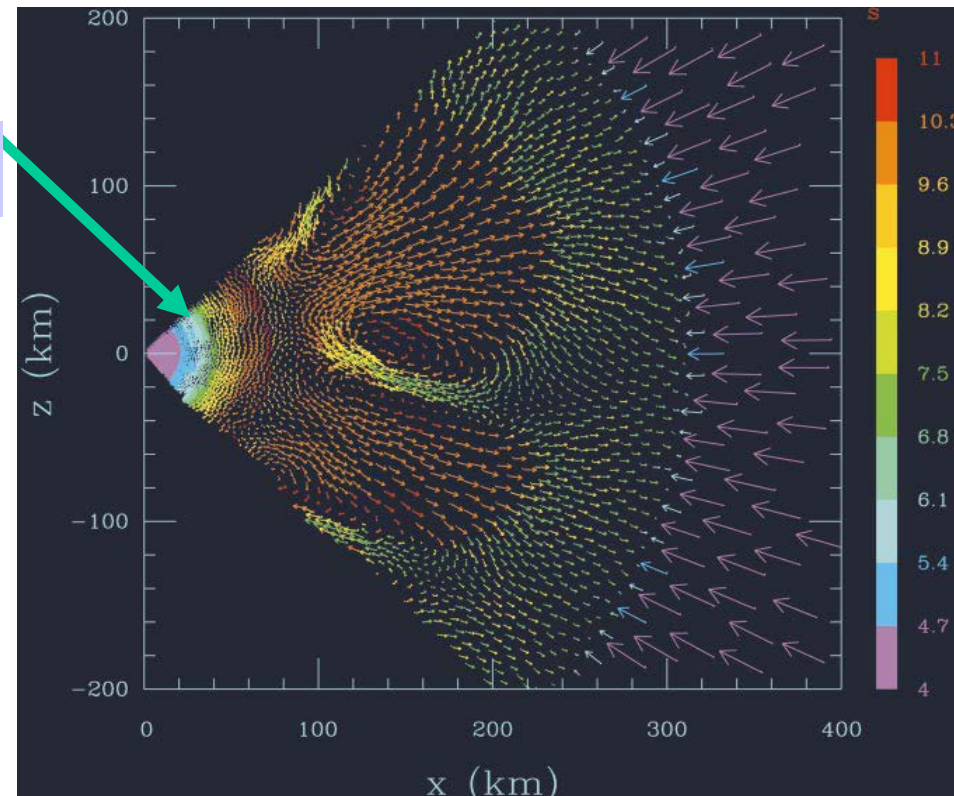
$$dB_v^2/dR \sim 8\pi GM n(R)/R^2$$

$$4\pi R^2 n(R) = dM/dR$$

$$B \sim 10^{1.5} \text{TeraTesla} (M/M_\odot)(10\text{km}/R)^2$$

$$R_v \sim 40\text{Km}; \pi R \sim 1\text{Km}$$

$$B_v \sim 10^1 - 10^2 \text{TeraTesla}$$



# NUCLEAR STATISTICAL EQUILIBRIUM in Ultra-Strong Magnetic Fields

Entropy  $S$  extremum  $\rightarrow$   $TdS = \sum_i \lambda_i dY_i = 0$

Nuclear composition at temperature  $T$

$$Y \propto \frac{G_A}{G_n^N G_p^Z} \exp\{-B/kT\}$$

Binding Energy  $B$

spin-magnetic part in partition function

$$G_i = \sum_M \exp\{g_i M \omega_L / kT\} / (2I_i + 1) \quad \omega_L = \mu_N H$$

# Nuclear Shell Effects at *Ultra-Strong Magnetic Fields*

V.N.K. //PRL 2002. V.88, 221101 // J.Nucl.Sci.Technol. V.1 Sup.2. P.550 //  
J.Nucl.Radiochem.Sci. 2002. V.3. P.205 // PRC 2004 V.69, 038801/ЯФ. 2012

$$N = \int_{-\infty}^{\varepsilon_F} d\varepsilon \cdot \rho(\varepsilon)$$

*Nucleons*

Binding Energy

$$B = \int_{-\infty}^{\varepsilon_F} d\varepsilon \varepsilon \rho(\varepsilon) = B^{sm} + B^{sh}$$

Level density

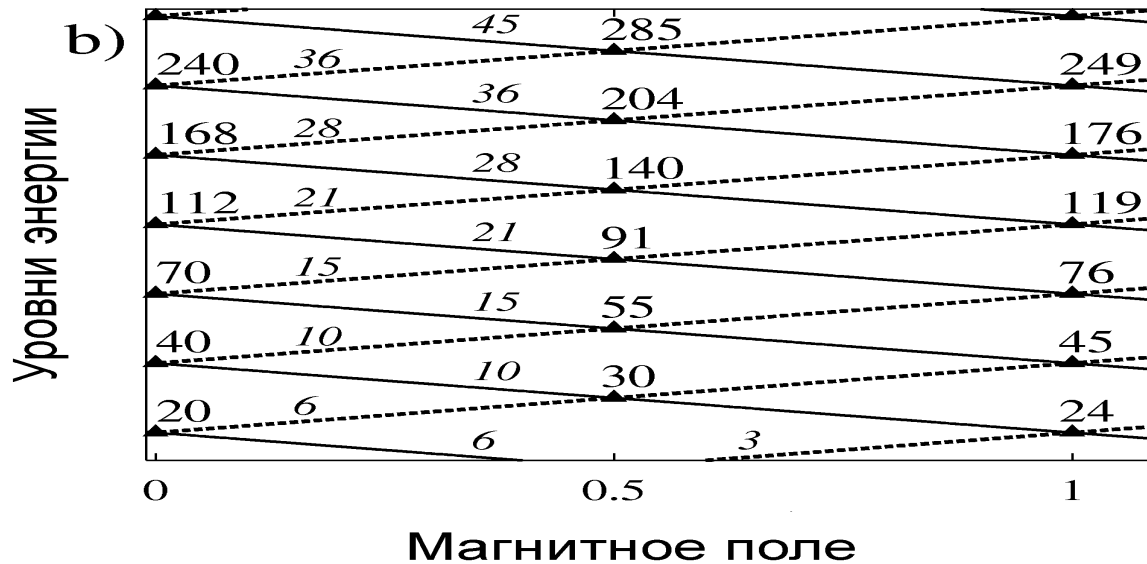
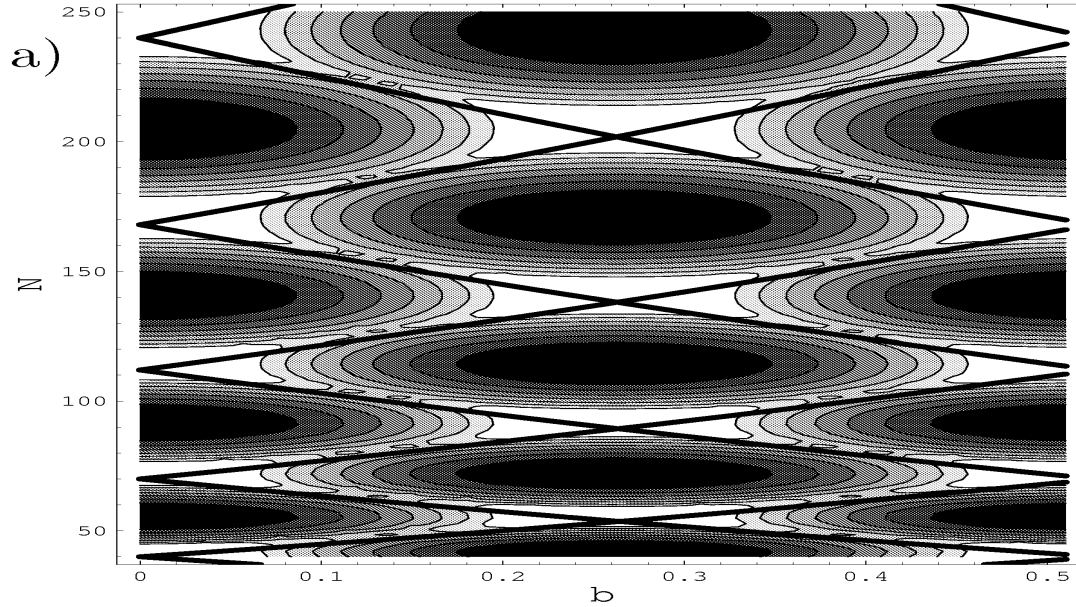
$$\rho = \sum_n \delta(\varepsilon - \varepsilon_n) = \rho^{sm} + \rho^{sh}$$

With Single particle levels  $\varepsilon_n$  *filled up to the Fermi energy*

# the Hartree self-consistent mean field approach in magnetic field : $\hbar$

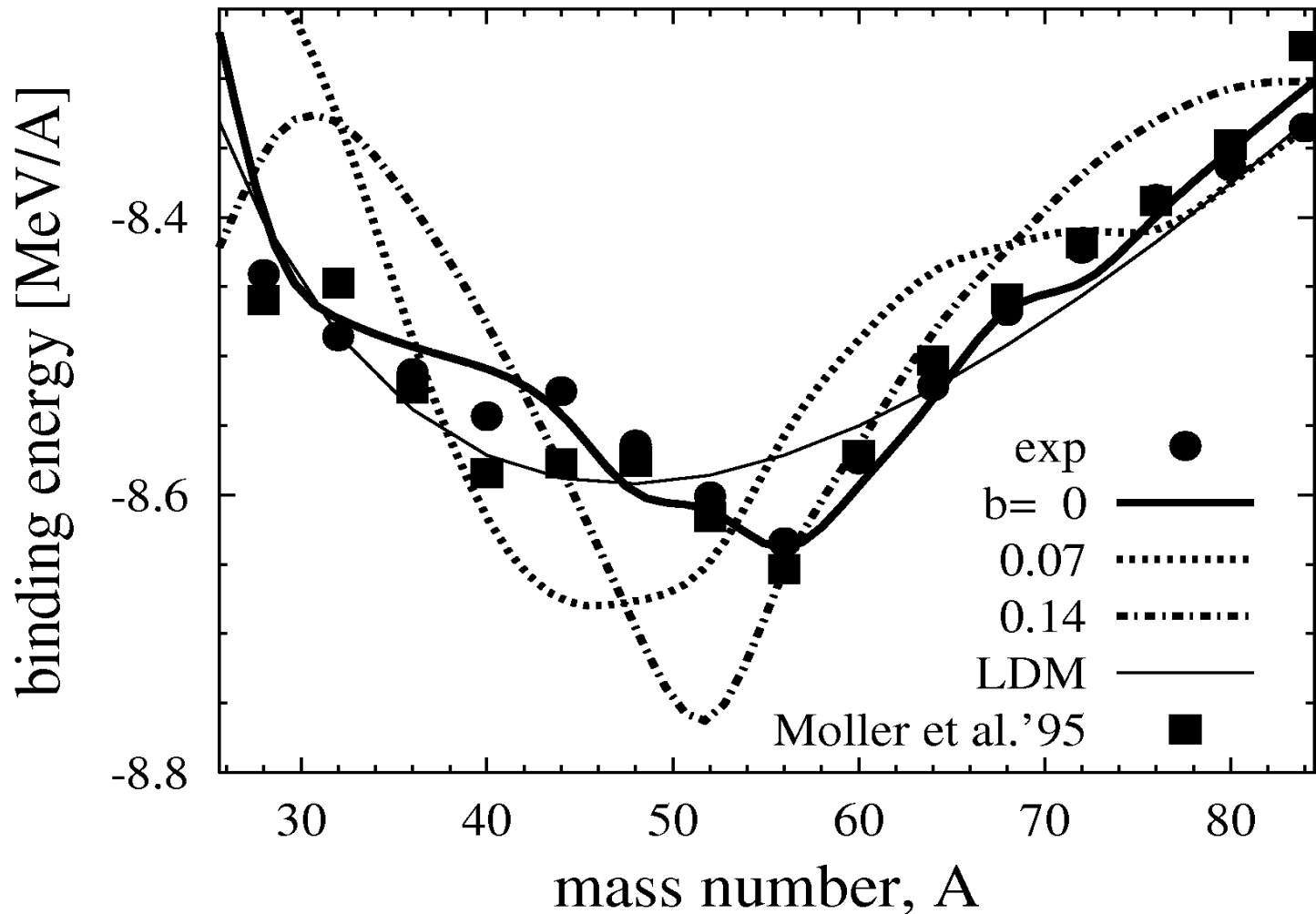
- Single particle Hamiltonian
- $\mathbf{H} = \mathbf{H}_{\text{MF}} + (\text{so})^*(\mathbf{l}\mathbf{S}) + (\text{Magnetic terms})$
- Pauli–spin ( $S$ )  $\rightarrow -\mathbf{M}(\hbar\mathbf{S})$
- Landau–orbital ( $l$ )  $\rightarrow -\mathbf{M}(\hbar\mathbf{l})$  :protons

# Neutrons



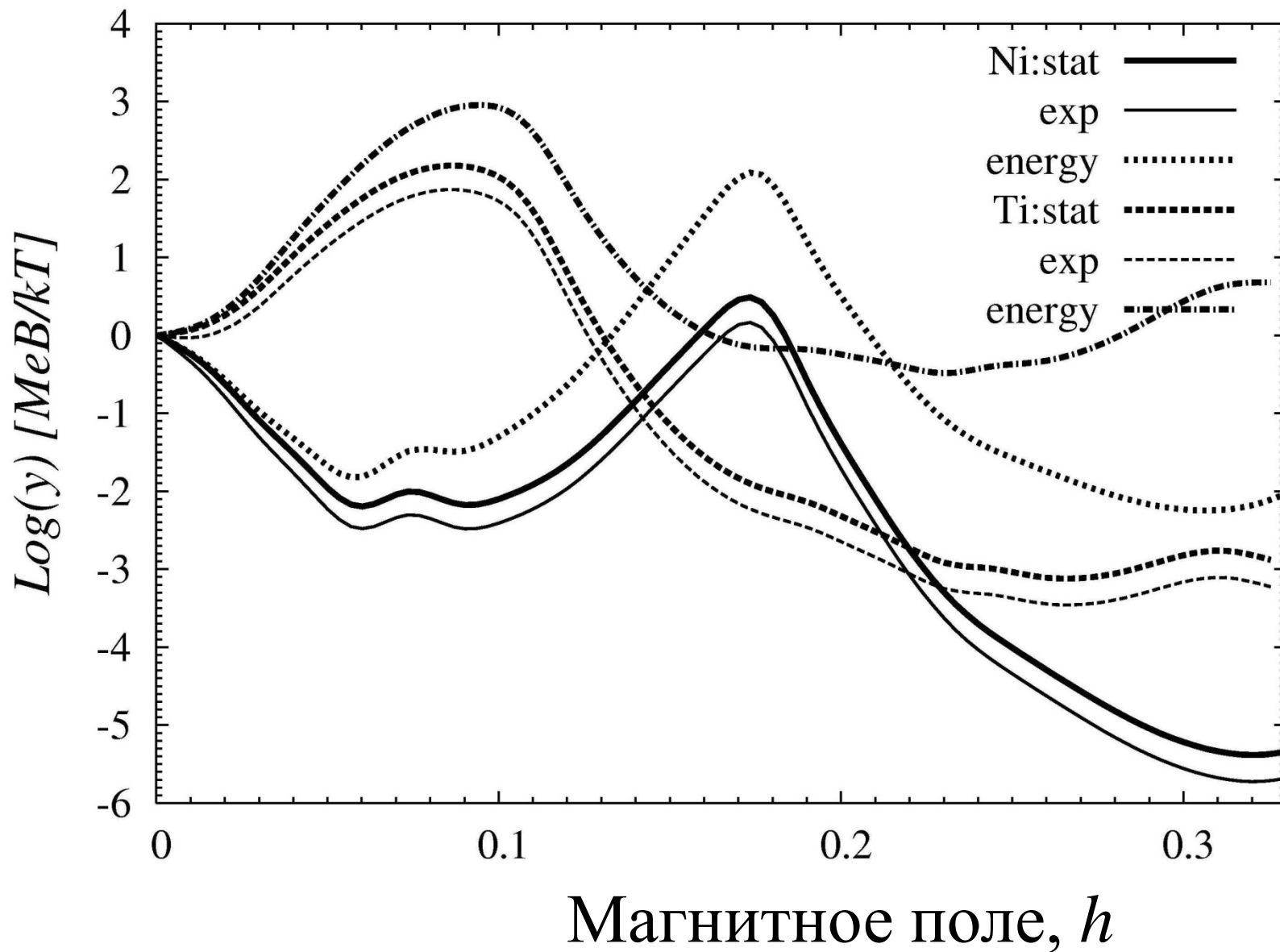


# Binding energy of even-even symmetric nuclei *at magnetic fields $h$*



relative yield  $y = Y(H)/Y(0)$

$^{56}\text{Ni}$ (solid) i  $^{44}\text{Ti}$ (dashed line)



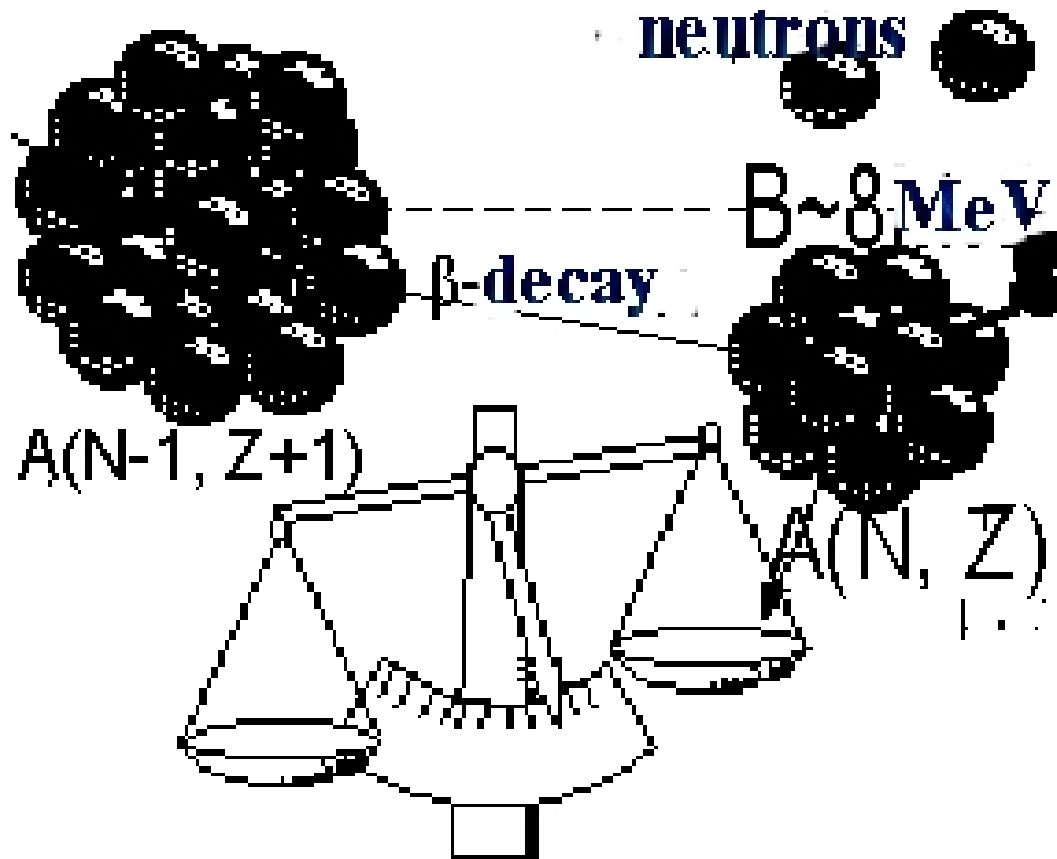


# Explosive nucleosynthesis: **s-process**

- The **s-process** (slow-process) proceeds via neutron-capture ( $n, \gamma$ ) reaction chains at low neutron flux.
- Process definition:  $T_{\text{beta-decay}} < T_{(n, \gamma)}$  --  
**new nucleus decays to stable nuclide before further neutron is captured:** (sec vs  $\sim 10^3$  yrs)
- s-process works through chain of stable nuclides ('beta stability valley') & n-capture much **slower** decay times
- Famous s-element Technetium ( $Z = 43$ ) has no stable isotopes - the longest lived  $\sim 10^6$  years -- must be produced recently in stars where observed.
- Tc was first observed by Merrill (1952) and provided the first real proof that nucleosynthesis was indeed happening in stars!

# slow- or s-process.

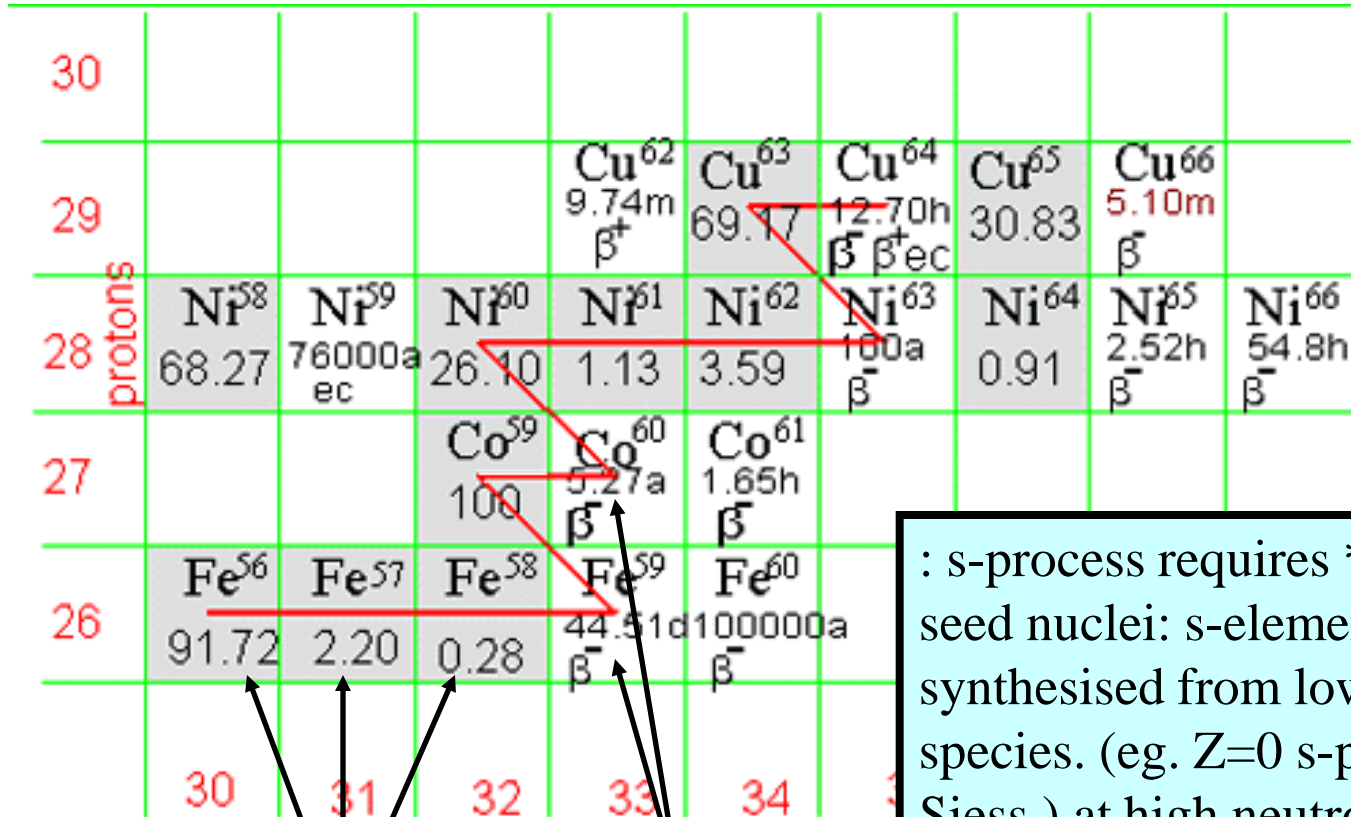
Beta-decay faster n-capture rate  
weak interactions keep the proton&neutron  
Fermi levels at equilibrium



# Sample of the s-process:

neutron capture: 

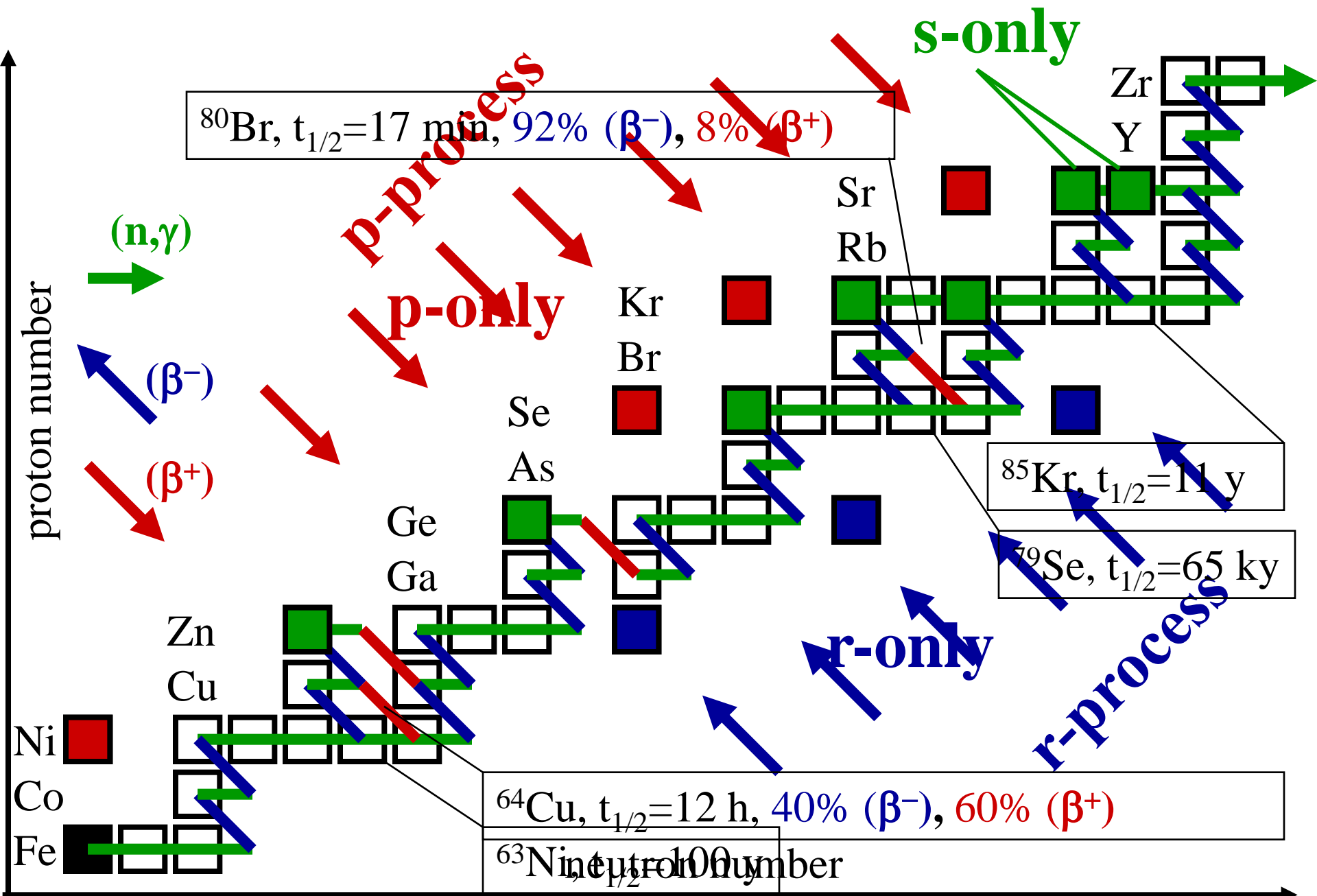
$\beta^+$  decay: 



: s-process requires \*no\* Fe group seed nuclei: s-elements can be synthesised from lower-mass species. (eg. Z=0 s-process, Siess.) at high neutron fluxes.

Stable

Unstable ( $\beta^-$  - decay)



# Nuclear Magics in s-process:

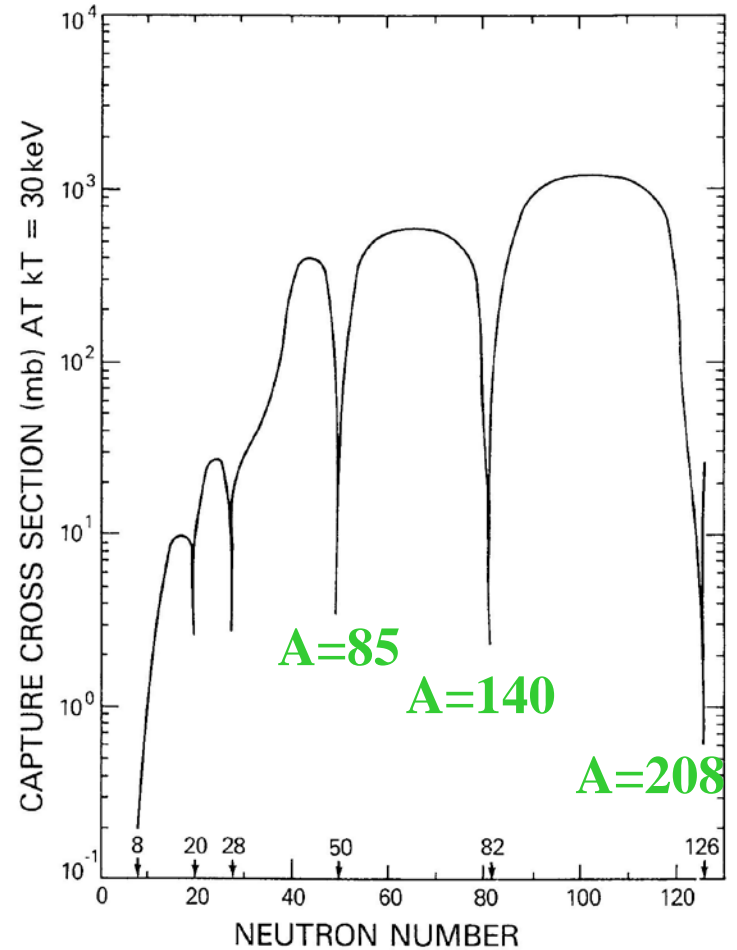
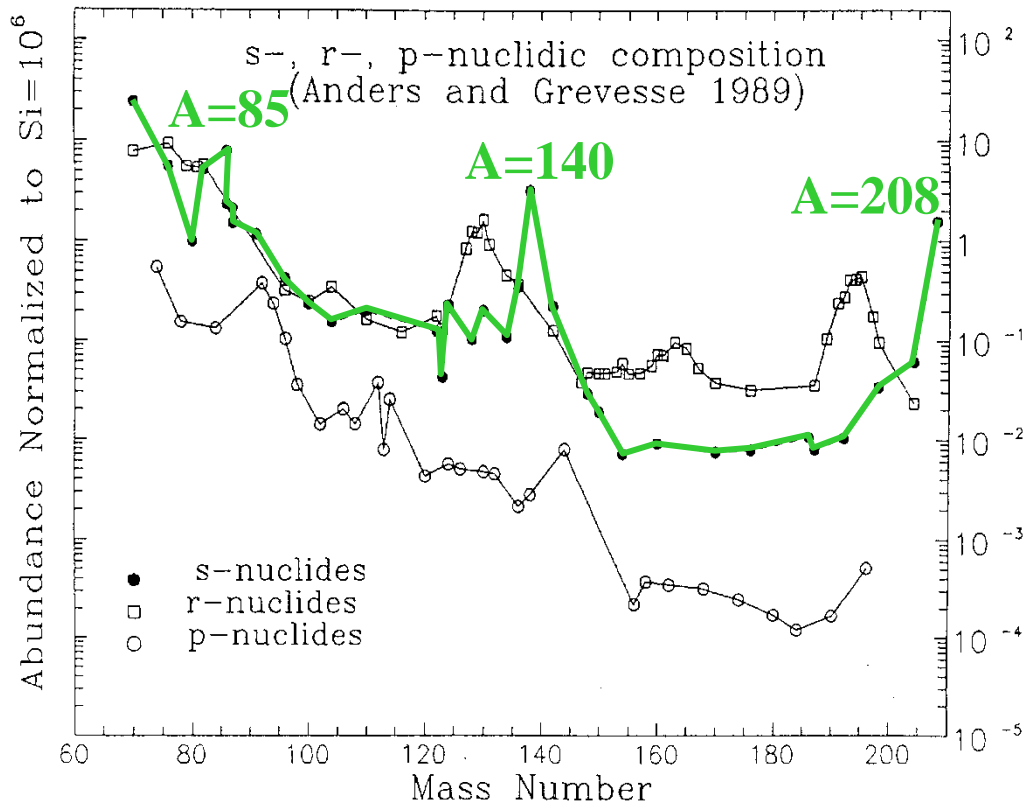
- For certain neutron numbers --  $N = 28, 50, 82, 126$  -- neutron capture cross-sections are much smaller than others.
- When one of these "magic" numbers is reached, the nucleus hardly captures more neutrons.
- significant number of nuclei will build up at magic numbers
- more abundant Elements corresponding to neutron "magic" numbers produce peaks in the observations.
- Solar System as abundance peaks at  $^{88}\text{Sr}$ ,  $^{138}\text{Ba}$ , &  $^{208}\text{Pb}$ .



$$N_A \propto \frac{1}{\langle \sigma \rangle_A}$$

small capture cross sections at neutron magic numbers

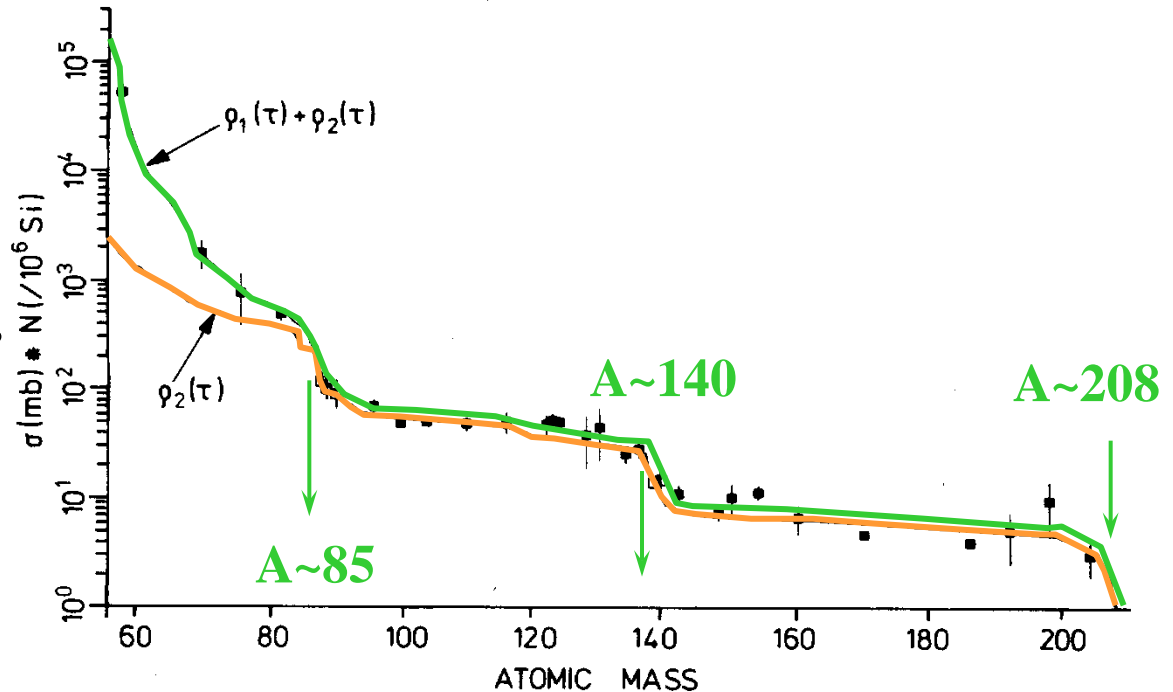
⇔ pronounced abundance peaks



$$\langle \sigma \rangle_A N_A = \text{constant}$$

condition fulfilled between  
magic numbers of neutrons

sudden drops observed at  
neutron magic numbers



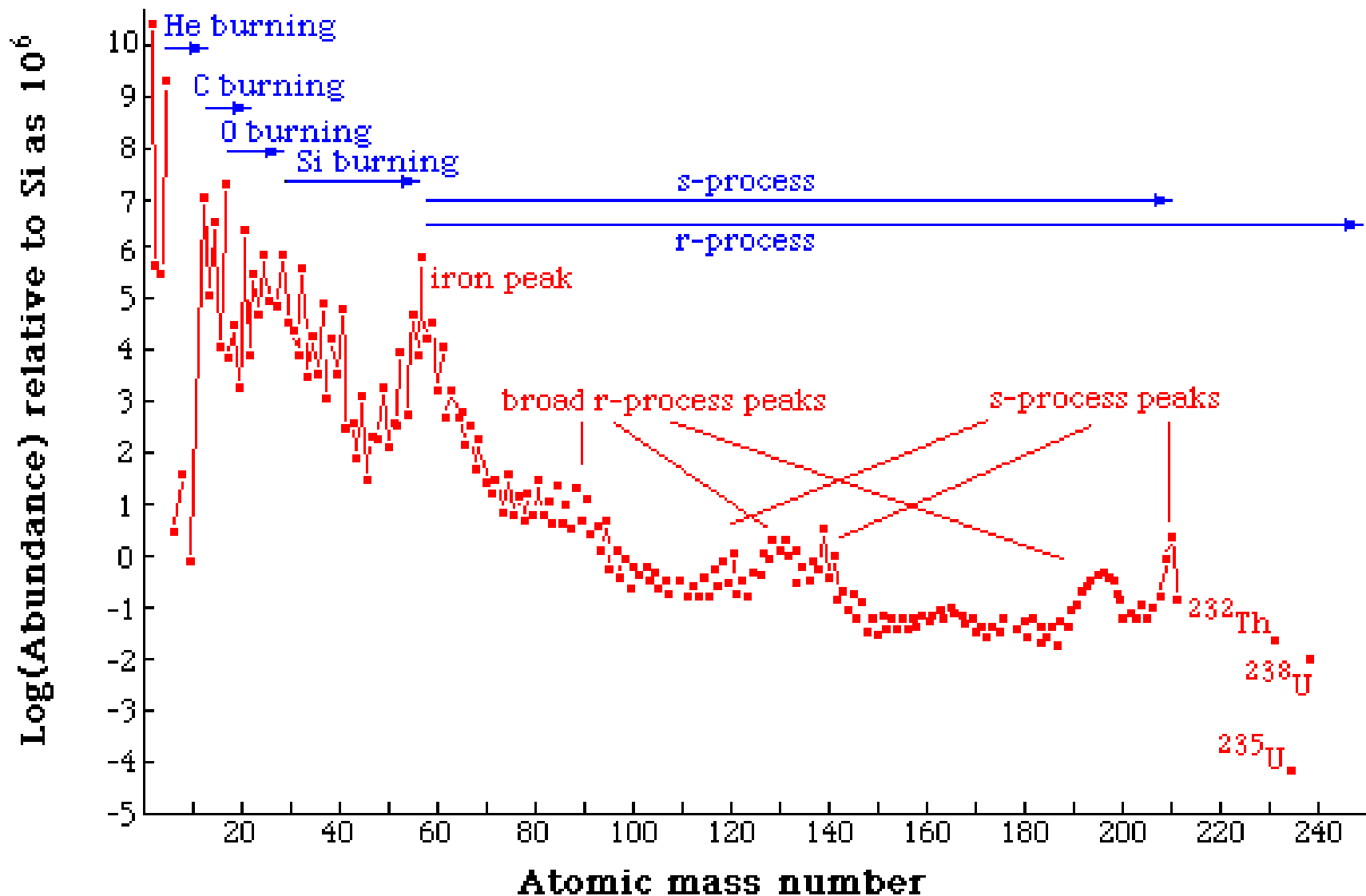
### NOTE

a superposition of many neutron irradiations is needed to correctly  
reproduce the abundance curve

- **main component** ( $A=88-209$ )
- **weak component** ( $A < 90$ )

s-process: best understood nucleosynthesis process from nuclear point of view

The s (& r) processes explain the production of higher mass nuclei (beyond the Fe peak).

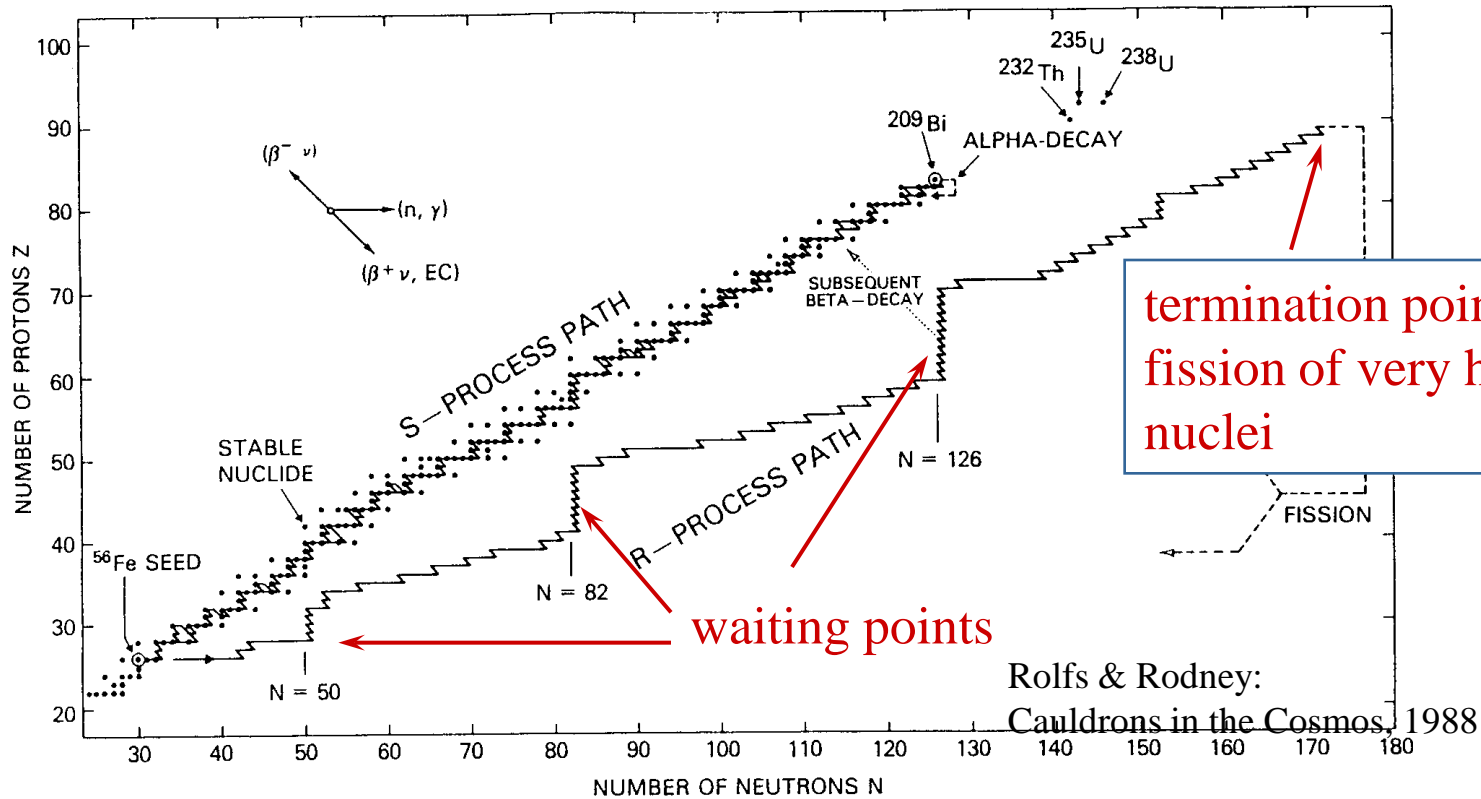


# r-process (r = rapid) neutron capture process)

unstable nucleus reacts before the decay



$$\tau_n \ll \tau_\beta$$



termination point:  
fission of very heavy  
nuclei

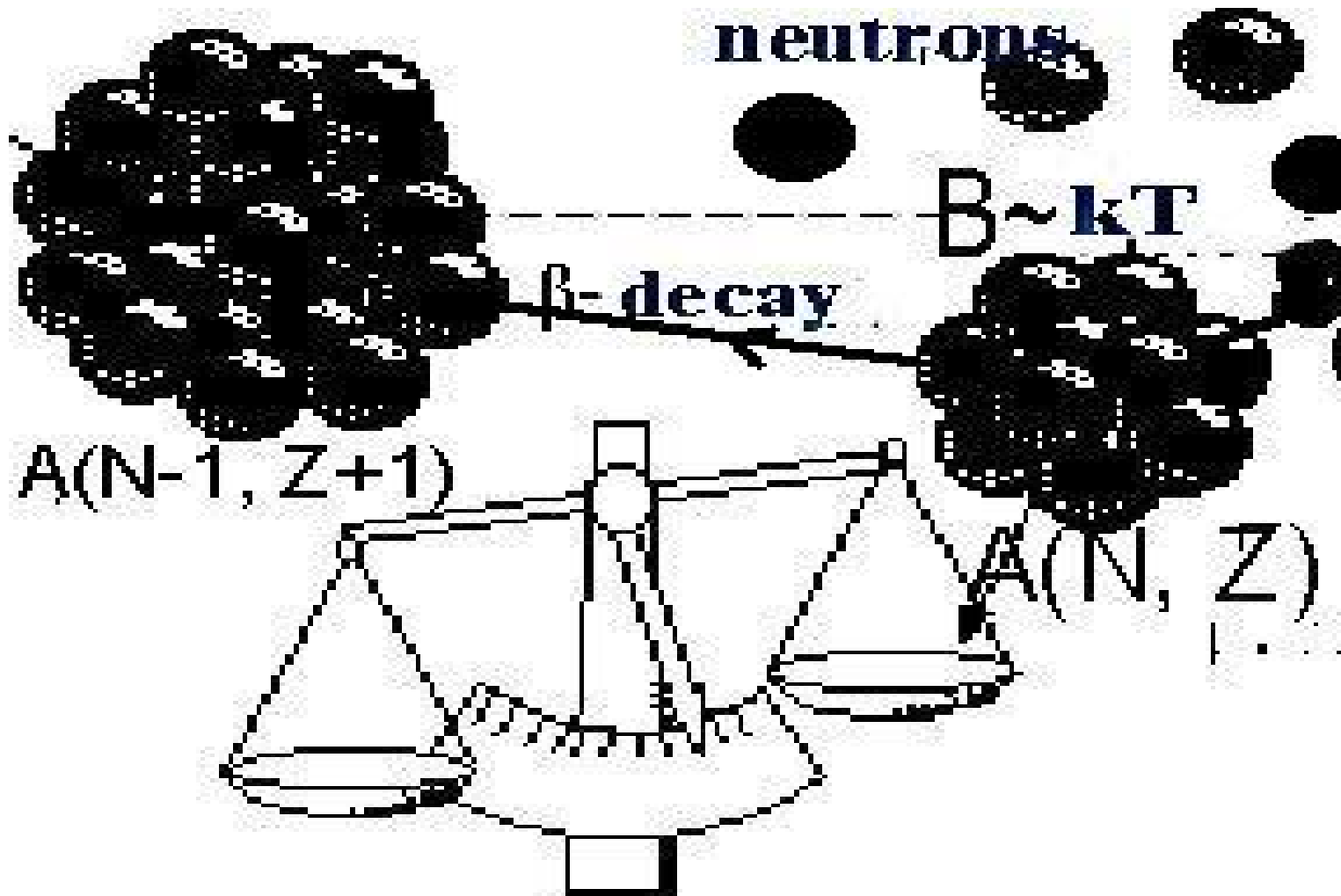
waiting points

typical lifetimes for unstable nuclei far from the valley of  $\beta$  stability:  $10^{-6} - 10^{-2}$  s

requiring:

$$\tau_n \sim 10^{-4} \text{ s} \quad \Leftrightarrow \quad N_n \sim 10^{20} \text{ n/cm}^3$$

explosive scenarios needed to account for such high neutron fluxes



R-process ← | exceptionally explosive conditions  
 $N(n) \sim 10^{20}/\text{cc} : T \sim 10^9\text{K} : t \sim 1\text{sec}$

# Transmutations of ultramagnetized atomic nuclei under intensive neutron

nuclear reaction rate

$$r = \int \sigma v dD_x dD_X$$

radiative *neutron*-capture ( $n, \gamma$ ) reaction

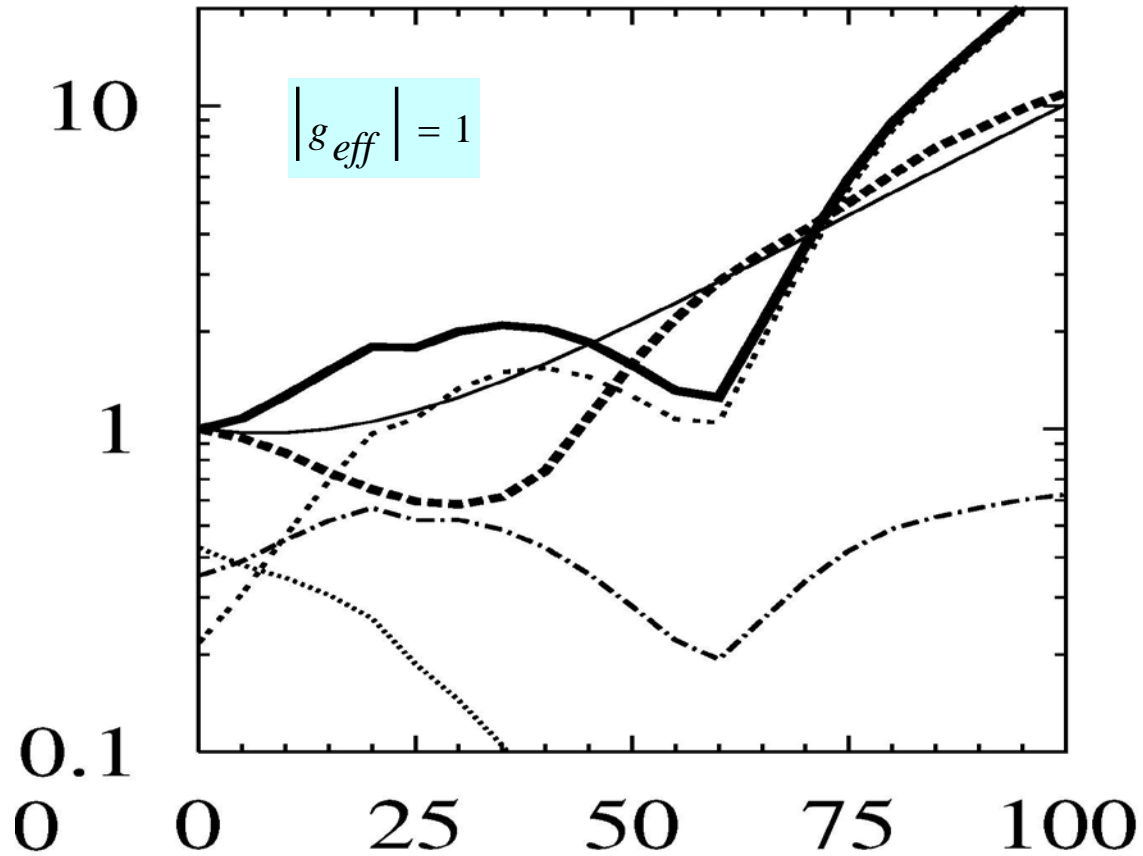
Hauser-Feshbach statistical approach

$$\sigma_{(n,\gamma)} \propto \Gamma / \kappa^2 \propto 1 / v_n$$

Relative rate

$$r(H) / r(0) = \sigma(H) / \sigma(0)$$

# Neutron-capture rate



# SUMMARY

- We analyze the synthesis and decay of  $n$  nuclides in SNRs
- Obtained the spectra and flux for the specific lines of decay products



# SUMMARY

- Magnetism of Atomic Nuclei  
Thermodynamic formalism
- Magnetic field effect on Nuclear Shell Structure  
Phase-shift in shell oscillations of Nuclear Masses
- Pauli type Paramagnetic Response
- Landau-type Orbital Magnetism
- Magnetic fields of Tera-Tesla shift  
Nuclear Magics of Iron region towards  
Smaller Masses approaching Ti-44
- Enhancement in Yield at nucleosynthesis

# SUMMARY

- We analyze the synthesis and decay of  $^{44}\text{Ti}$  in SNRs
- Obtained the spectra and flux for the lines of  $^{44}\text{Sc}$

# Models explaining titanium excess

Motizuki Yuko, Kumagai  
Shiomi,  $^{44}\text{Ti}$  radioactivity in  
young supernova remnants:  
Cas A and SN 1987A., New  
Astr. Rev. 2004. V. 48. P. 69



The radioactivity  
suppression is induced by  
the high level of ionization  
of matter. As a result –  
decreasing of the activity at  
the early phases of the  
evolution of the SNR.

Kondratyev V.N., Kadenko  
I.M., Nucleosynthesis in  
strong magnetic fields at  
statistical equilibrium;  
MNRAS 2005. V.359, P.927



Increasing of the  $\text{Ti-44}$   
synthesis due to astrophysics  
environment

$$1 \text{ kpc} = 3.08567758 \times 10^{19} \text{ метра}$$

