

Connecting Nuclear Physics to QCD with the lattice

André Walker-Loud

Bogoliubov Laboratory of Theoretical Physics, Joint Institute for Nuclear Research

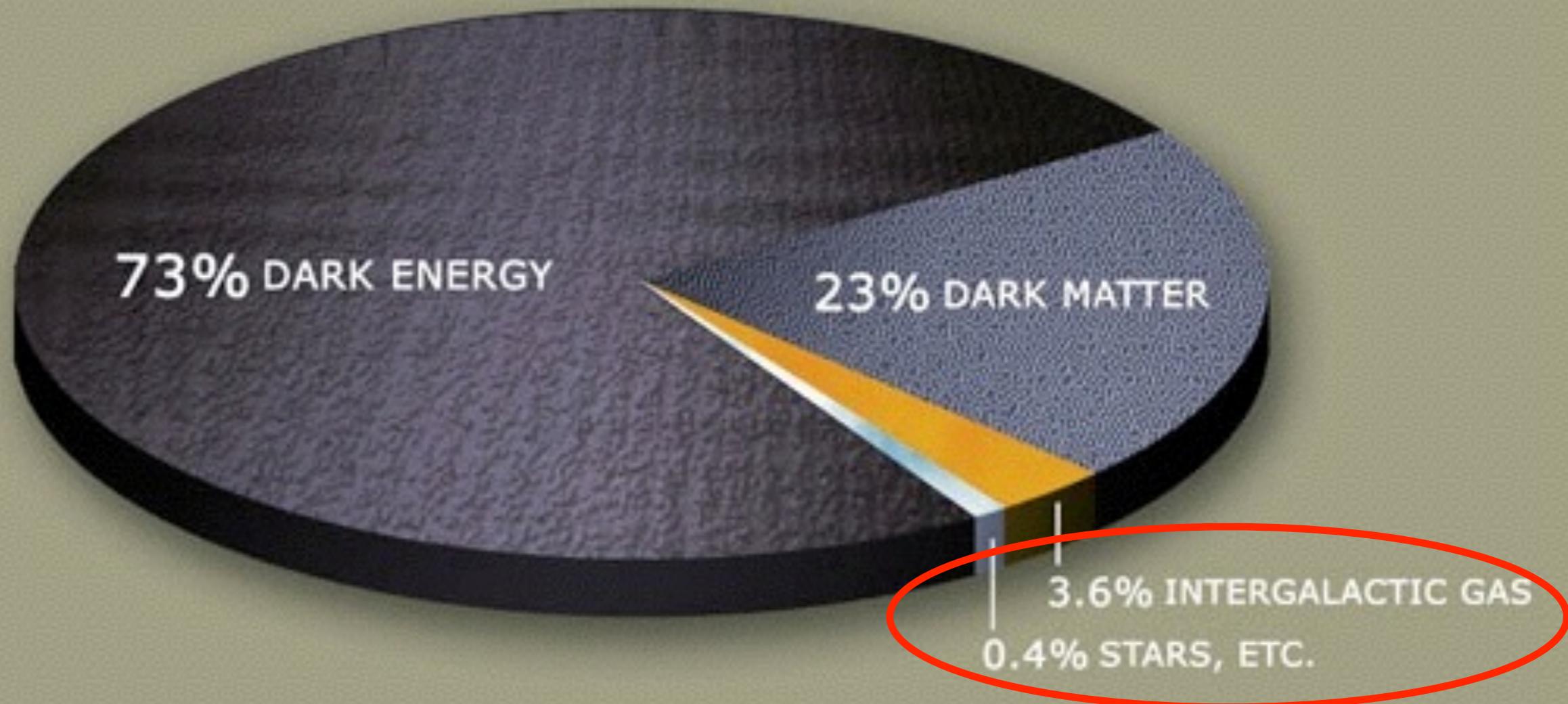
Dubna International Advanced School of Theoretical Physics
Helmholtz International Summer School

Dubna, Russia, **August 25-September 6 , 2014**

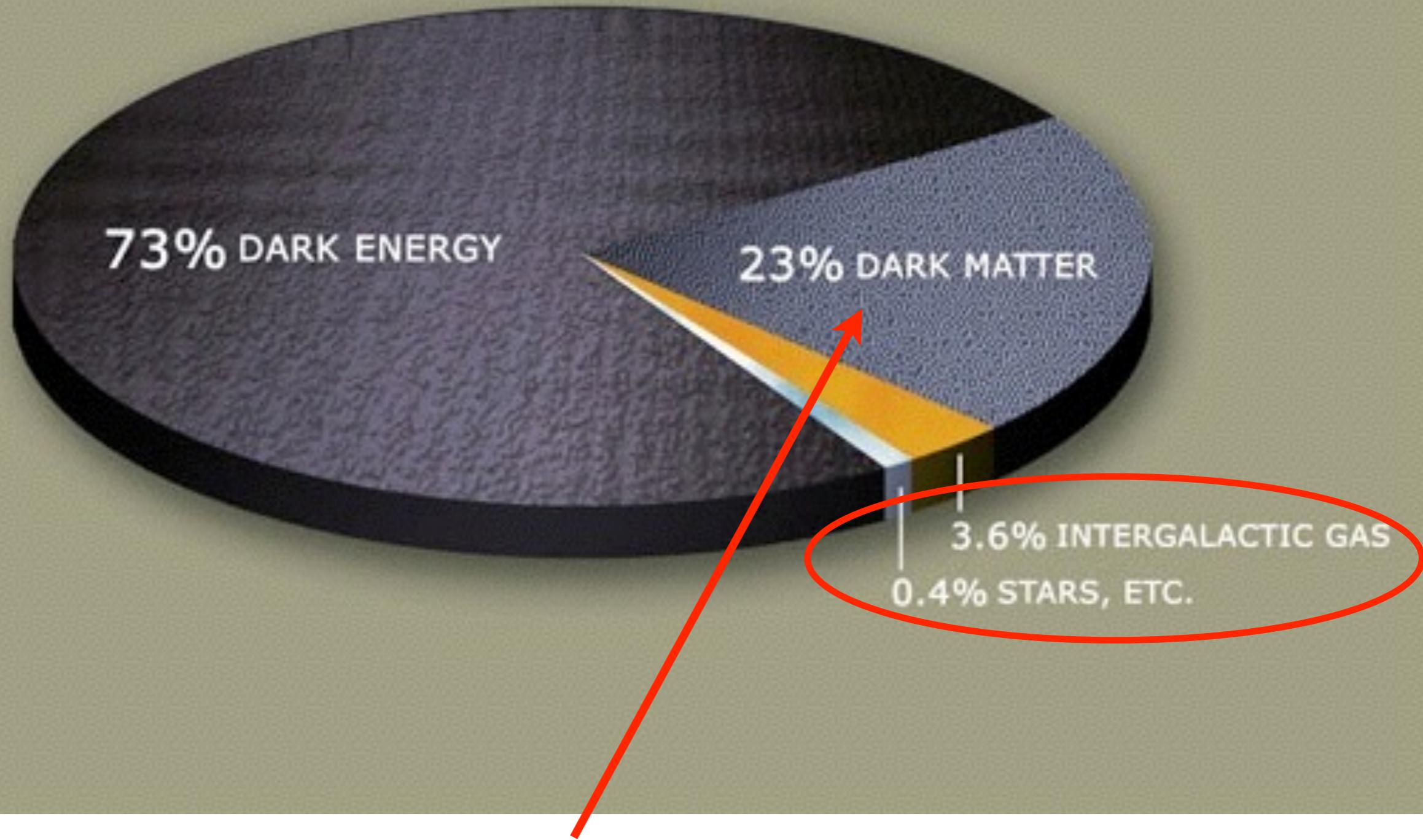
**Lattice QCD,
Hadron Structure
and Hadronic Matter**



Energy Budget of the Universe



Energy Budget of the Universe



quark mass dependence of nucleon mass helps us understand sensitivities of direct dark matter detection

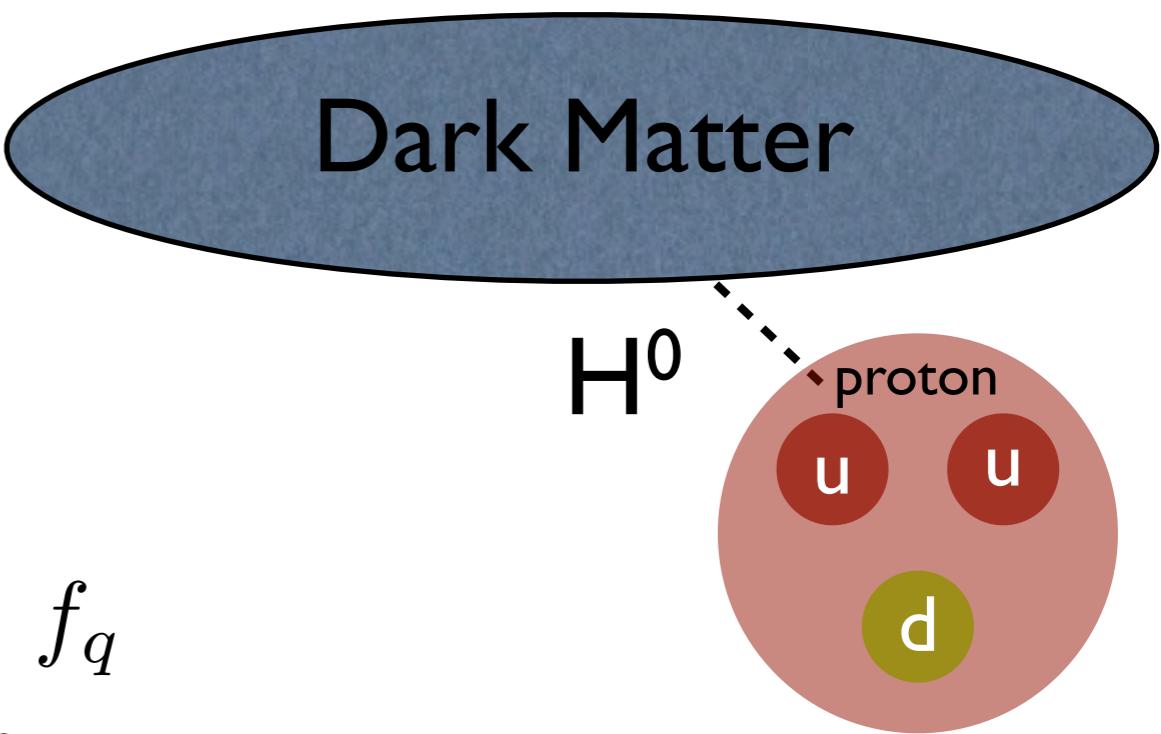
If Dark Matter couples to the scalar current of the nucleon (eg via Higgs) Spin Independent cross section

$$\sigma \propto |f|^2 \quad f = \frac{2}{9} + \frac{7}{9} \sum_{q=u,d,s} f_q$$

$$f_q \equiv \frac{\langle N | m_q \bar{q} q | N \rangle}{m_N}$$

see eg. Cheung, Hall, Pinner, Ruderman
arXiv:1211.4873

with enhancement of A^2 for nucleus (Xenon)



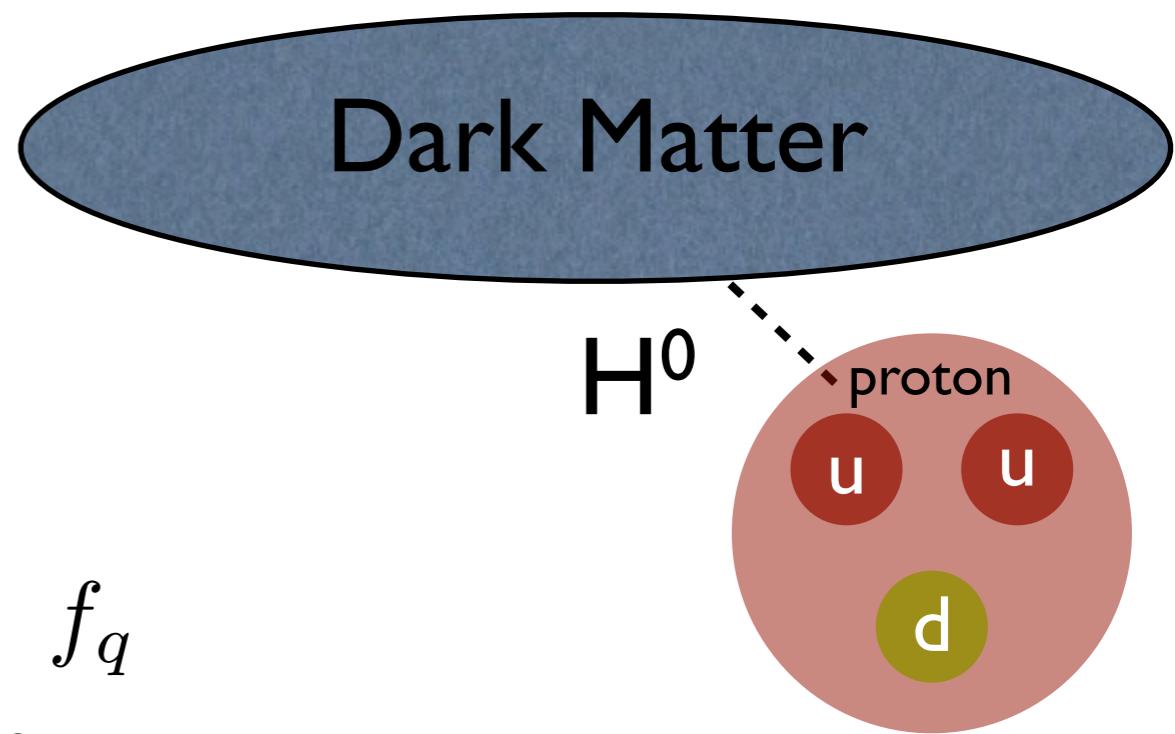
If Dark Matter couples to the scalar current of the nucleon (eg via Higgs) Spin Independent cross section

$$\sigma \propto |f|^2 \quad f = \frac{2}{9} + \frac{7}{9} \sum_{q=u,d,s} f_q$$

$$f_q \equiv \frac{\langle N | m_q \bar{q} q | N \rangle}{m_N}$$

see eg. Cheung, Hall, Pinner, Ruderman
arXiv:1211.4873

with enhancement of A^2 for nucleus (Xenon)



scalar current difficult to measure experimentally

$f_{u,d}$ estimated from pion-nucleon scattering

f_s uncertainty dominates estimates of cross section

Ellis, Olive, Savage
Phys.Rev. D77 (2008)

If Dark Matter couples to the scalar current of the nucleon (eg via Higgs) Spin Independent cross section

$$\sigma \propto |f|^2 \quad f = \frac{2}{9} + \frac{7}{9} \sum_{q=u,d,s} f_q$$

$$f_q \equiv \frac{\langle N | m_q \bar{q} q | N \rangle}{m_N}$$

see eg. Cheung, Hall, Pinner, Ruderman
arXiv:1211.4873

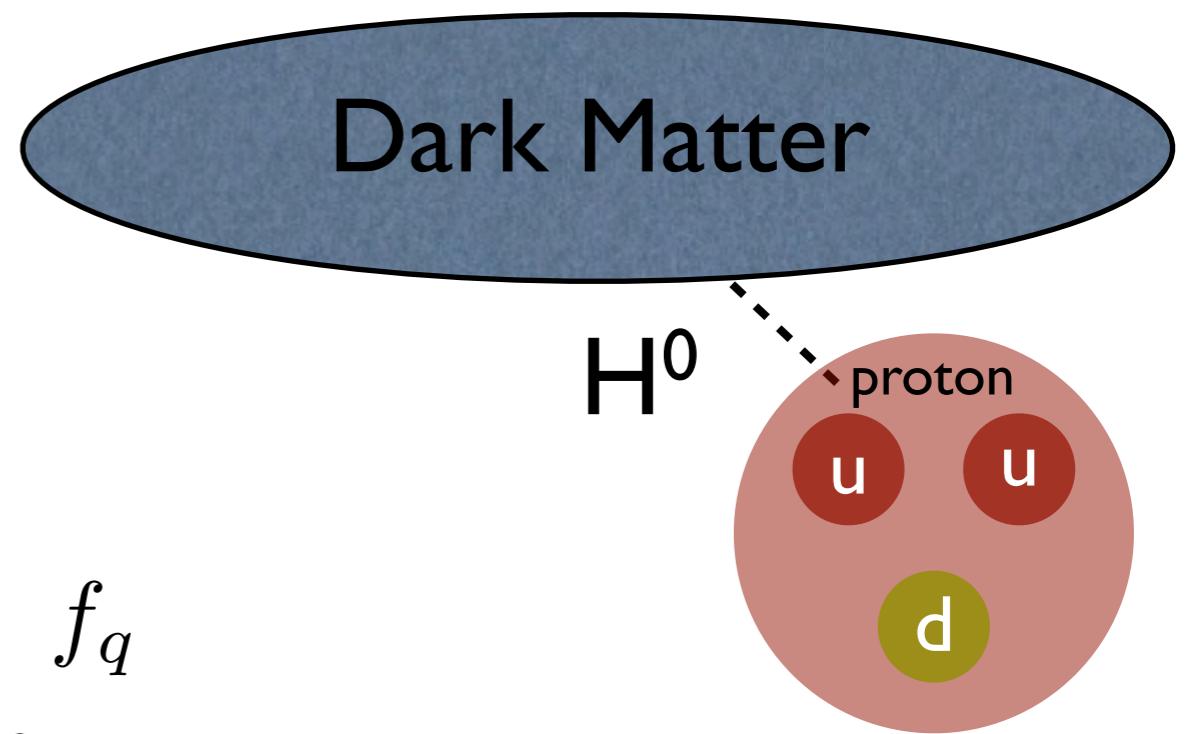
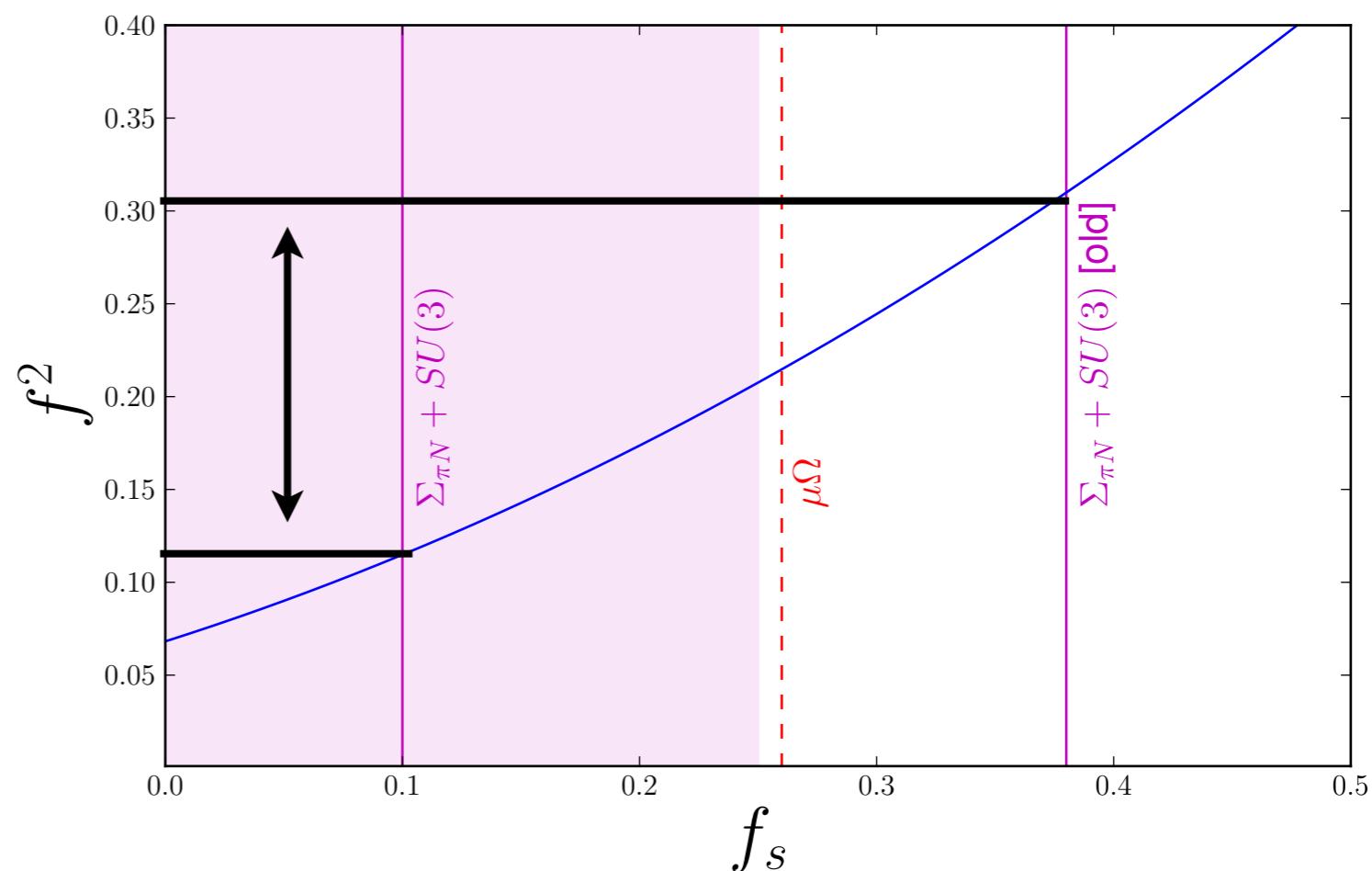


figure adapted from arXiv:1211.4873
thanks to J. Ruderman and collaborators



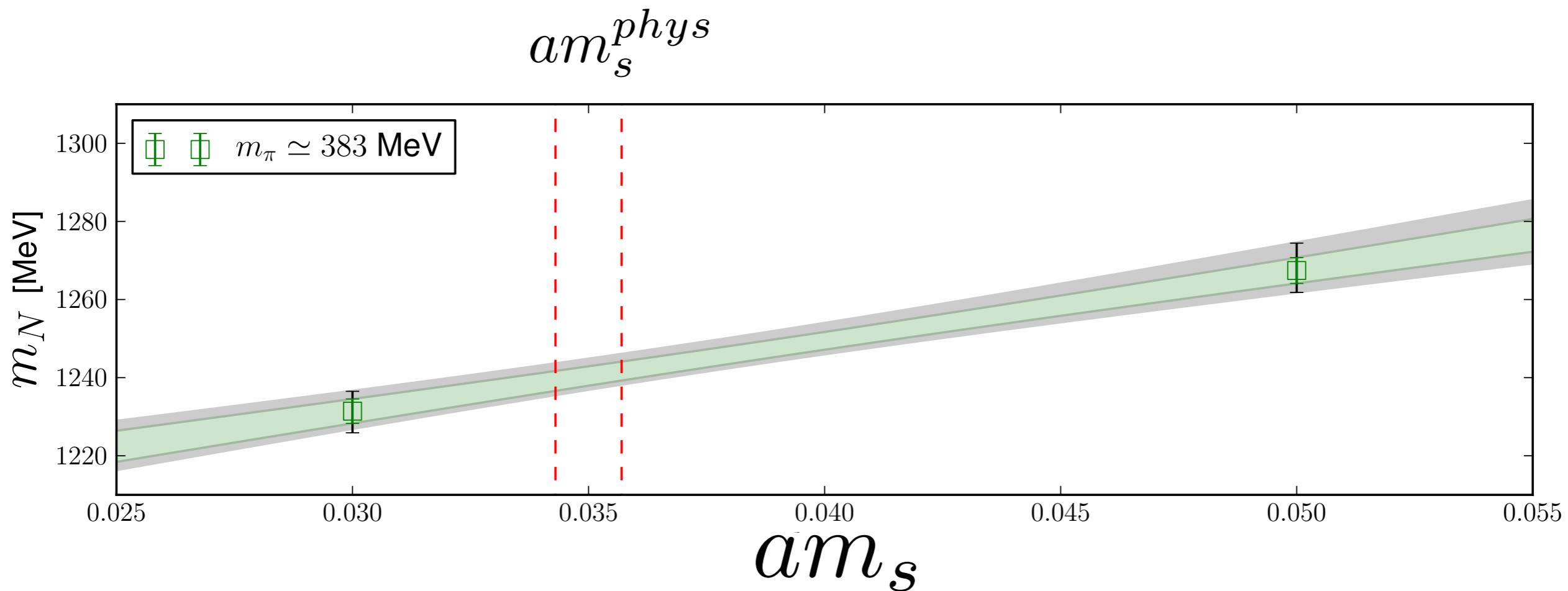
strange content of the nucleon

P.Junnarkar and AWL
arXiv:1301.1114

Lattice QCD perfect tool to compute strange content of nucleon $m_s \langle N | \bar{s}s | N \rangle$

Feynman-Hellmann Theorem

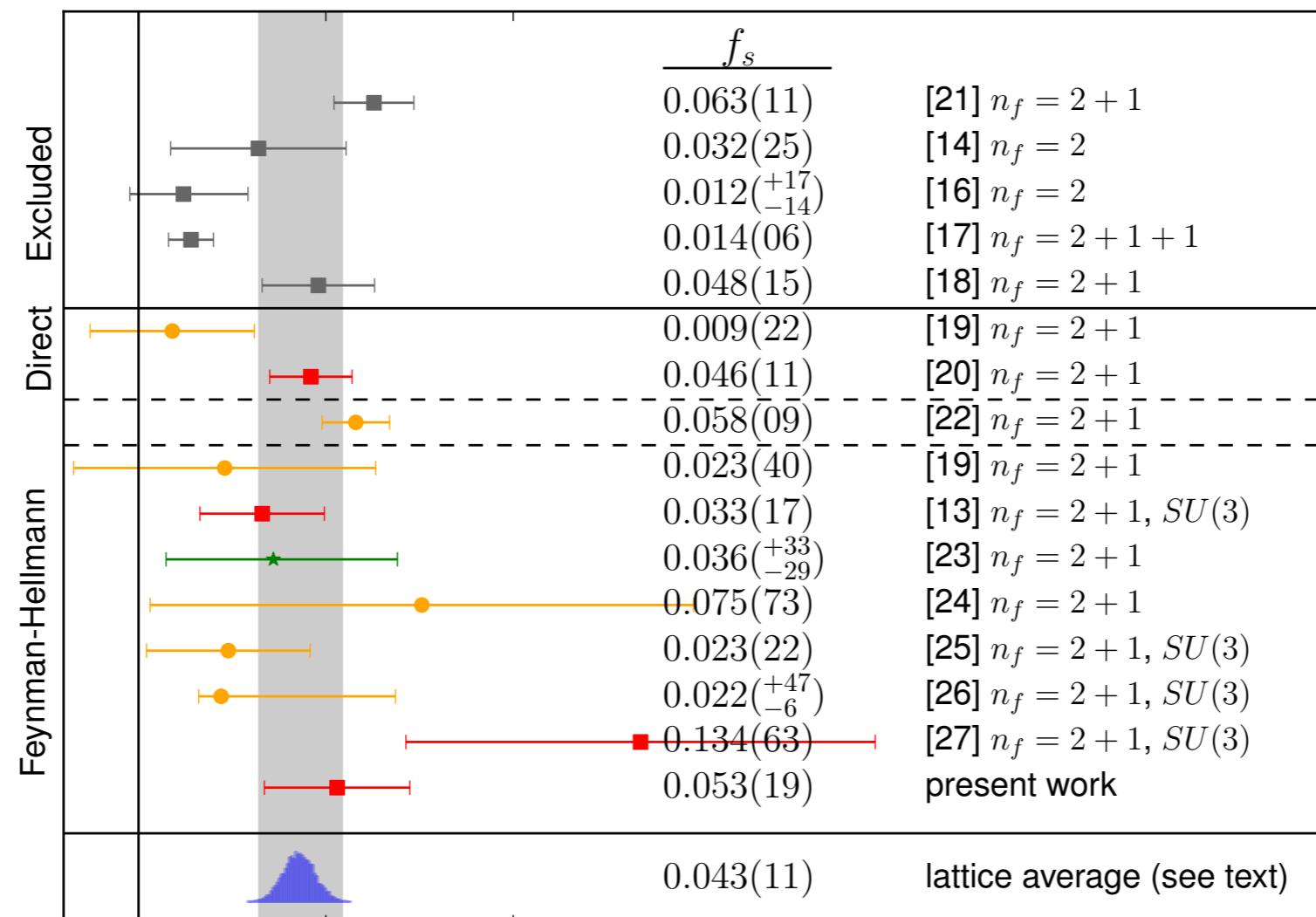
$$m_q \langle N | \bar{q}q | N \rangle = m_q \frac{\partial}{\partial m_q} m_N$$



Lattice QCD perfect tool to compute strange content of nucleon $m_s \langle N | \bar{s}s | N \rangle$

Feynman-Hellmann Theorem

$$m_q \langle N | \bar{q}q | N \rangle = m_q \frac{\partial}{\partial m_q} m_N$$



$$\hat{f}_s = m_s \langle N | \bar{s}s | N \rangle / m_N$$

If Dark Matter couples to the scalar current of the nucleon (eg via Higgs) Spin Independent cross section

$$\sigma \propto |f|^2 \quad f = \frac{2}{9} + \frac{7}{9} \sum_{q=u,d,s} f_q$$

$$f_q \equiv \frac{\langle N | m_q \bar{q} q | N \rangle}{m_N}$$

see eg. Cheung, Hall, Pinner, Ruderman
arXiv:1211.4873

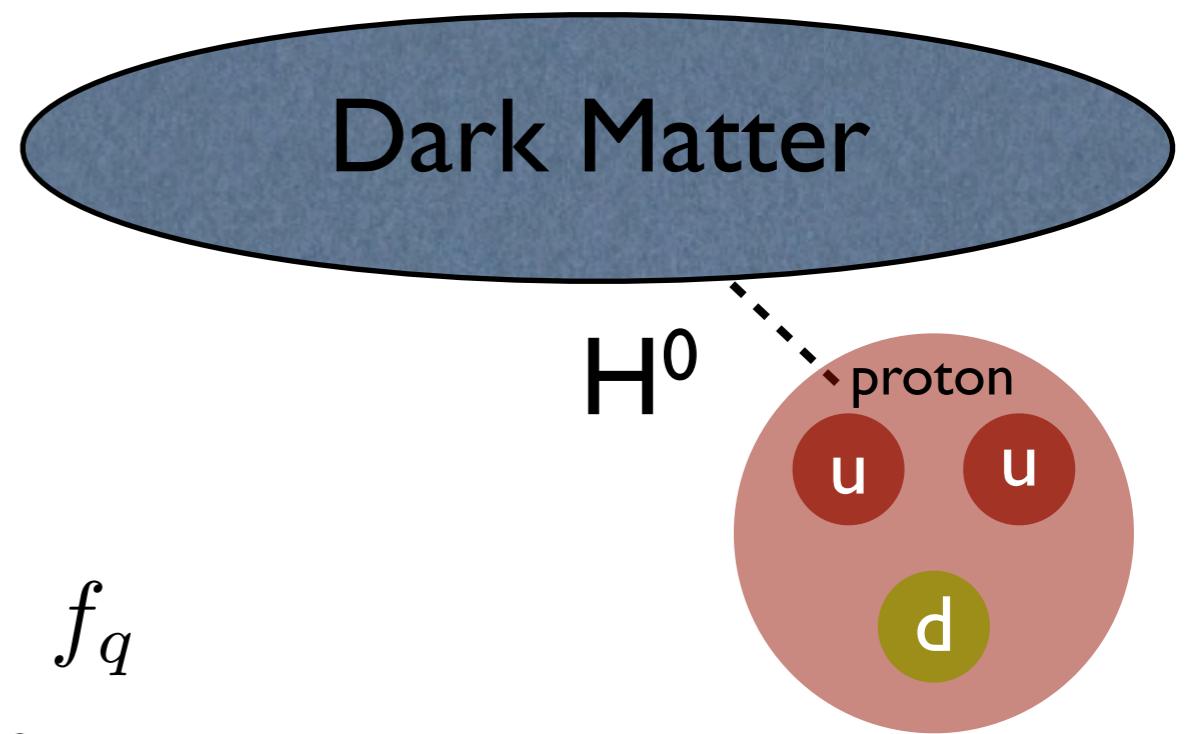
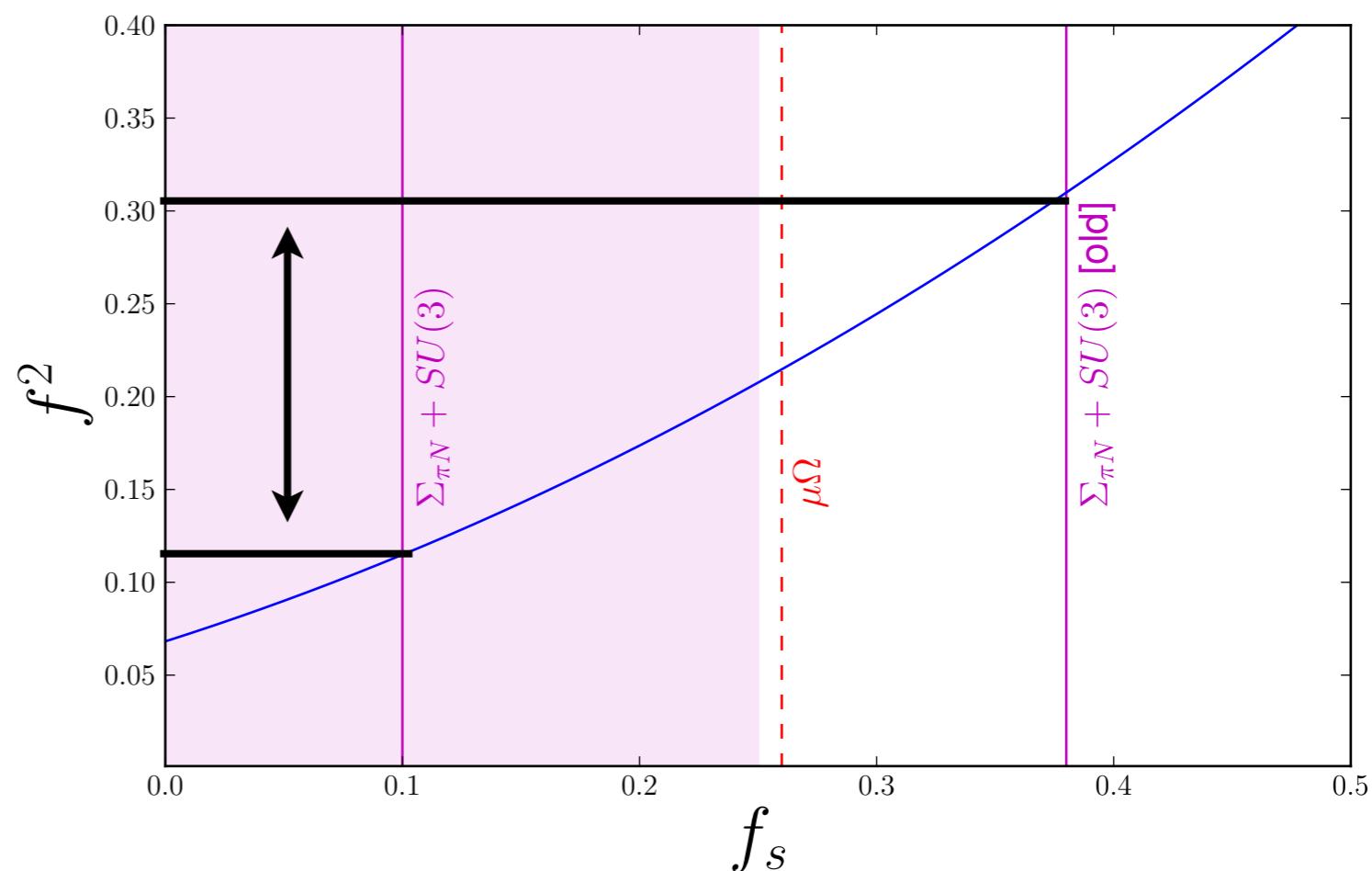


figure adapted from arXiv:1211.4873
thanks to J. Ruderman and collaborators



If Dark Matter couples to the scalar current of the nucleon (eg via Higgs) Spin Independent cross section

$$\sigma \propto |f|^2 \quad f = \frac{2}{9} + \frac{7}{9} \sum_{q=u,d,s} f_q$$

$$f_q \equiv \frac{\langle N | m_q \bar{q} q | N \rangle}{m_N}$$

see eg. Cheung, Hall, Pinner, Ruderman
arXiv:1211.4873

dramatic reduction in uncertainty of cross section

now $f_{u,d}$ gives larger uncertainty - but harder

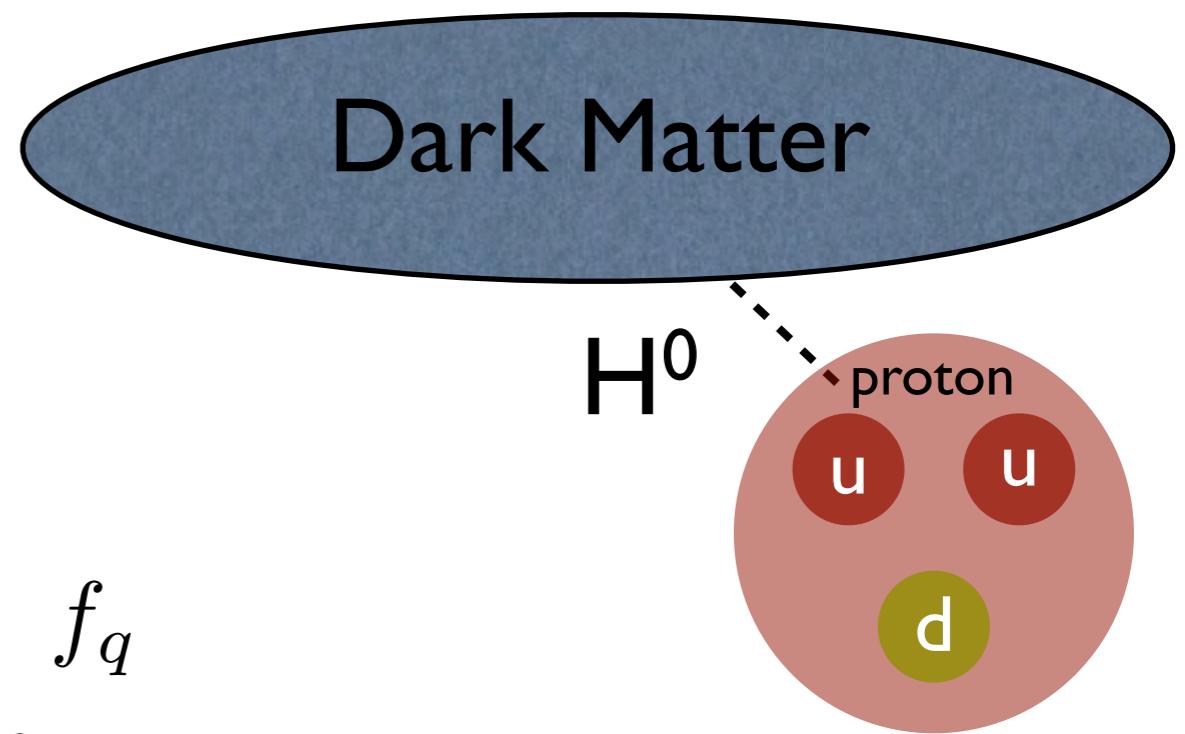
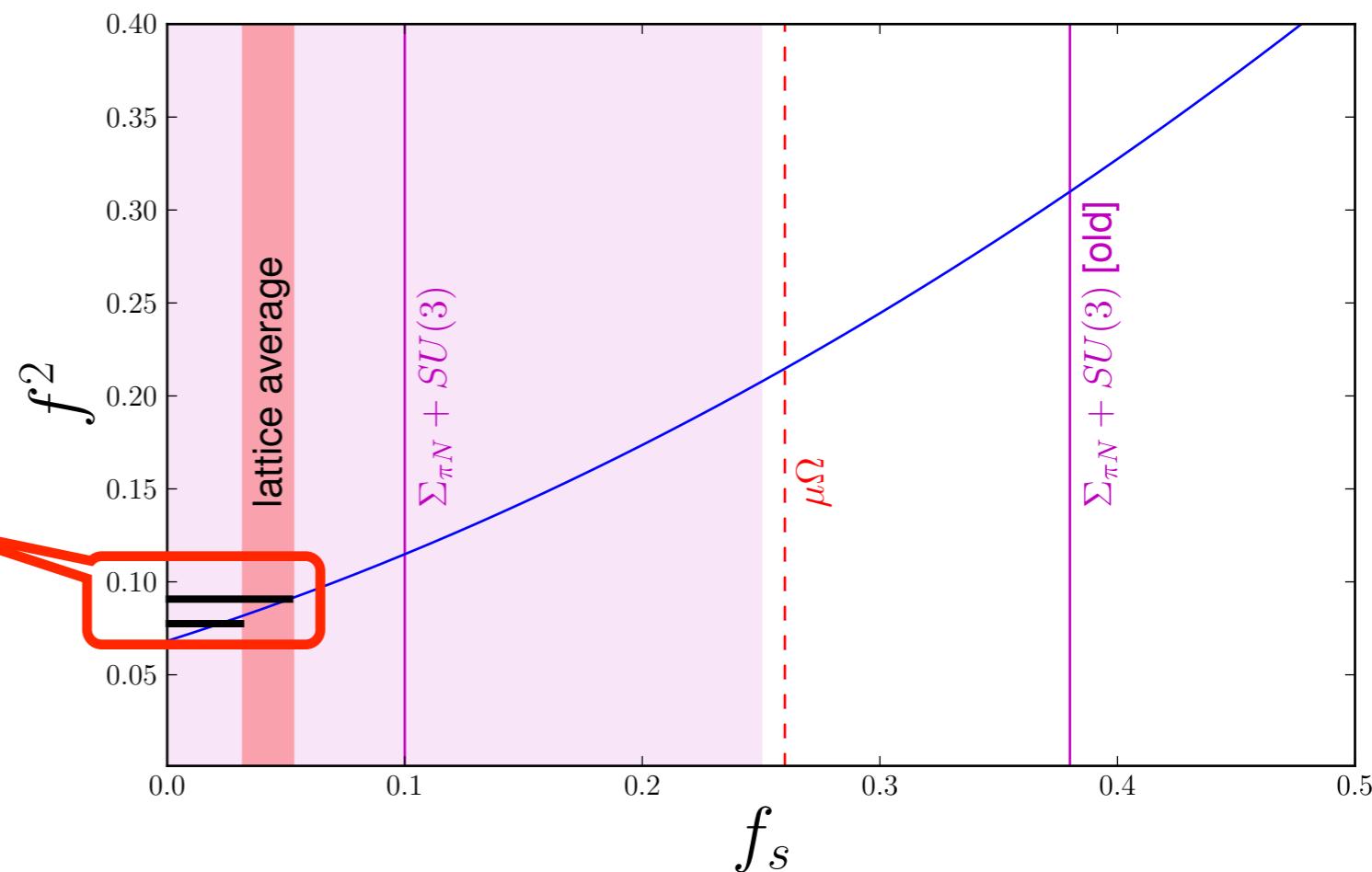


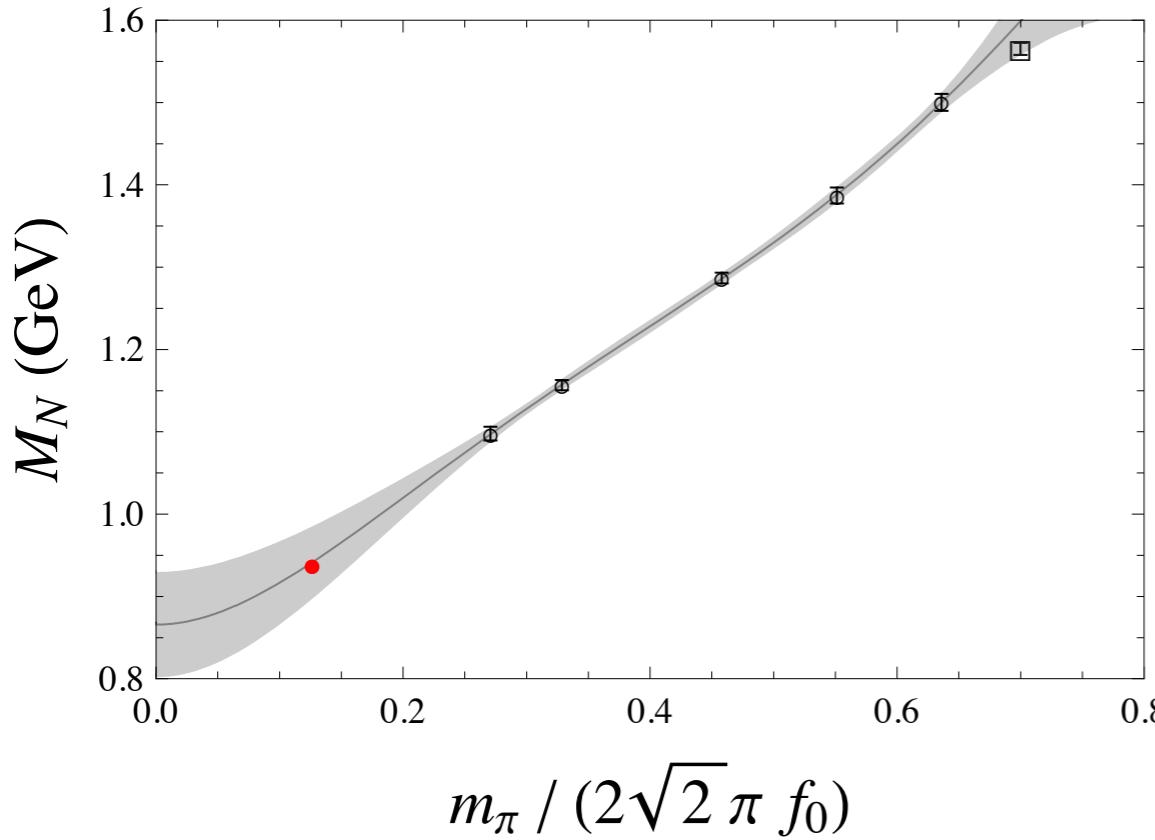
figure adapted from arXiv:1211.4873
thanks to J. Ruderman and collaborators



Baryons in lattice QCD

Light quark mass dependence of M_N

NNLO - m_π^4 , with $g_A=1.2(1)$, $g_{\Delta N}=1.5(3)$



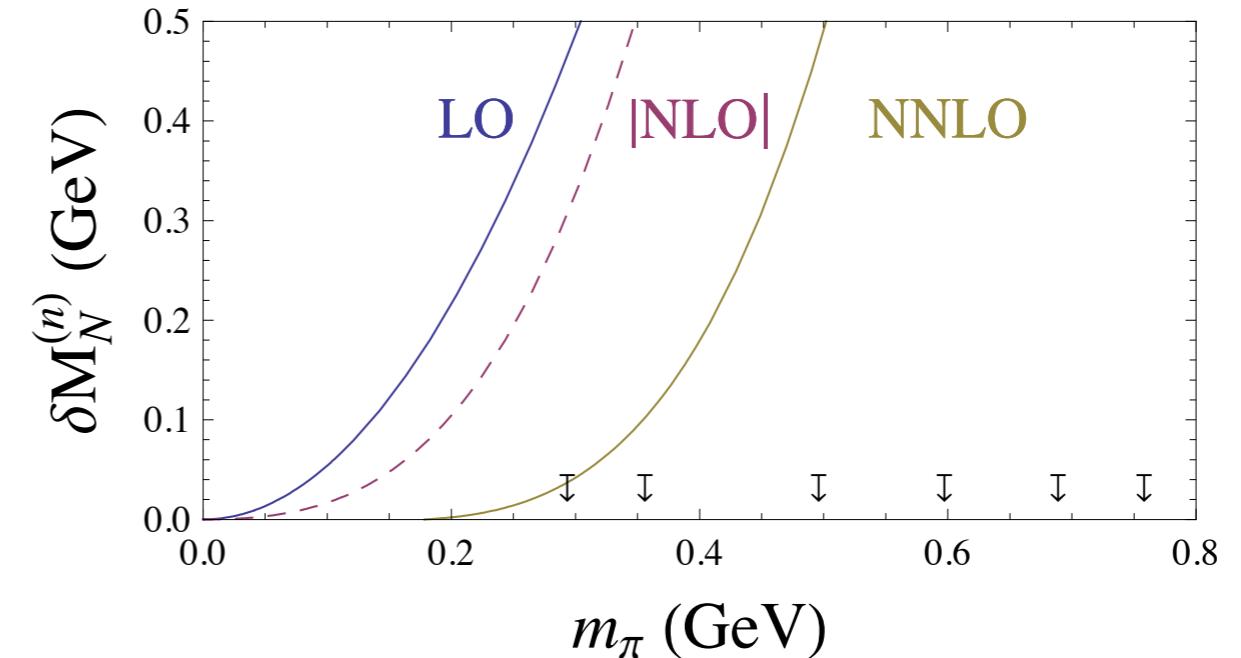
NNLO Heavy Baryon Fit

$$M_N = 954 \pm 42 \pm 20 \text{ MeV}$$

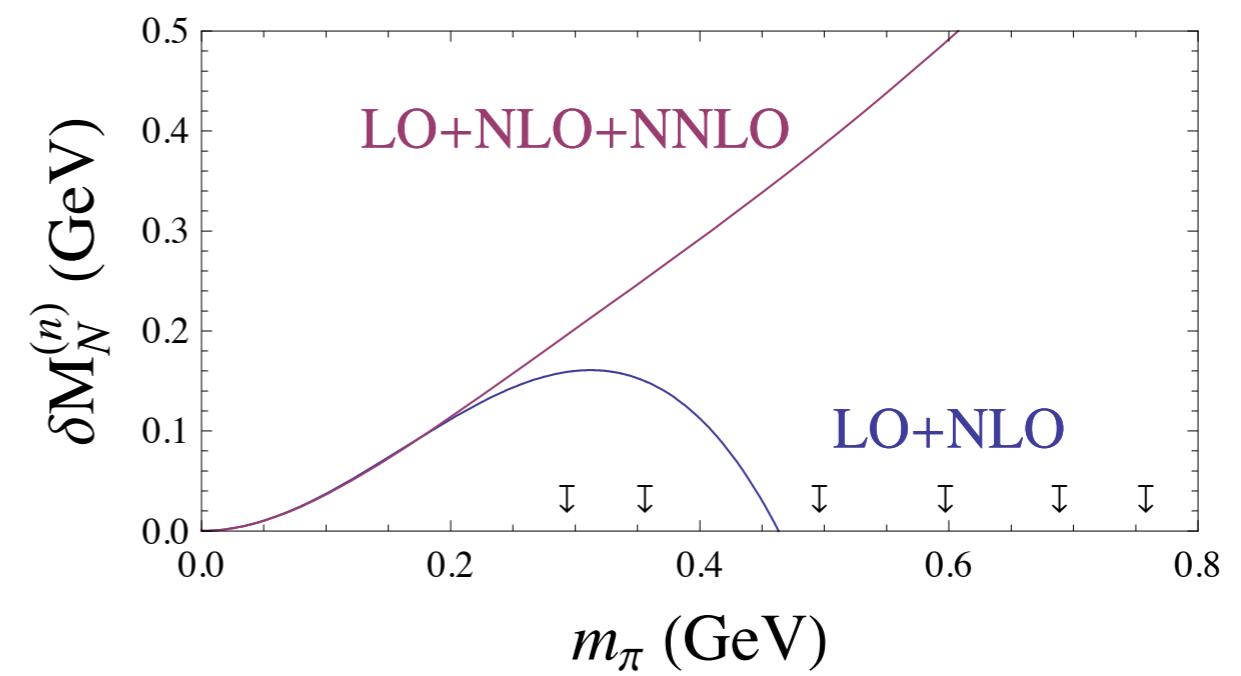
statistical

varying inputs

$g_A=1.2(1)$, $g_{\Delta N}=1.5(3)$



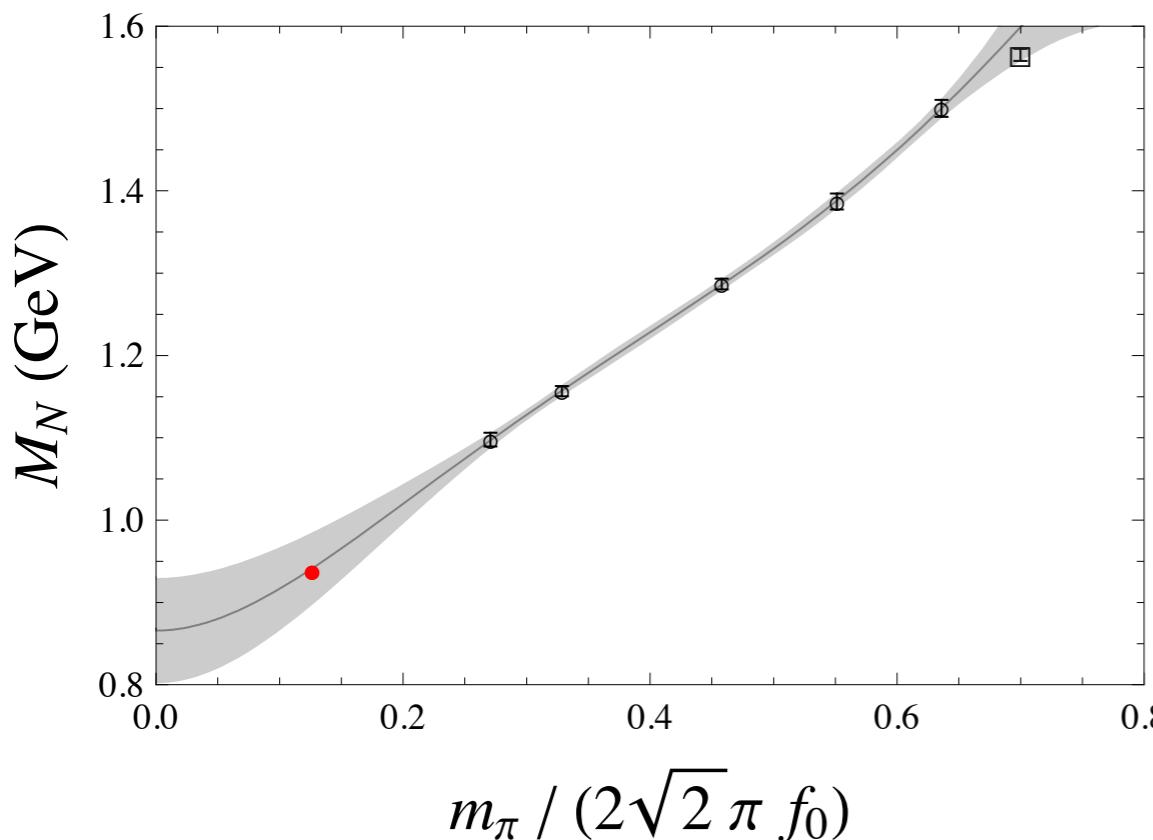
$g_A=1.2(1)$, $g_{\Delta N}=1.5(3)$



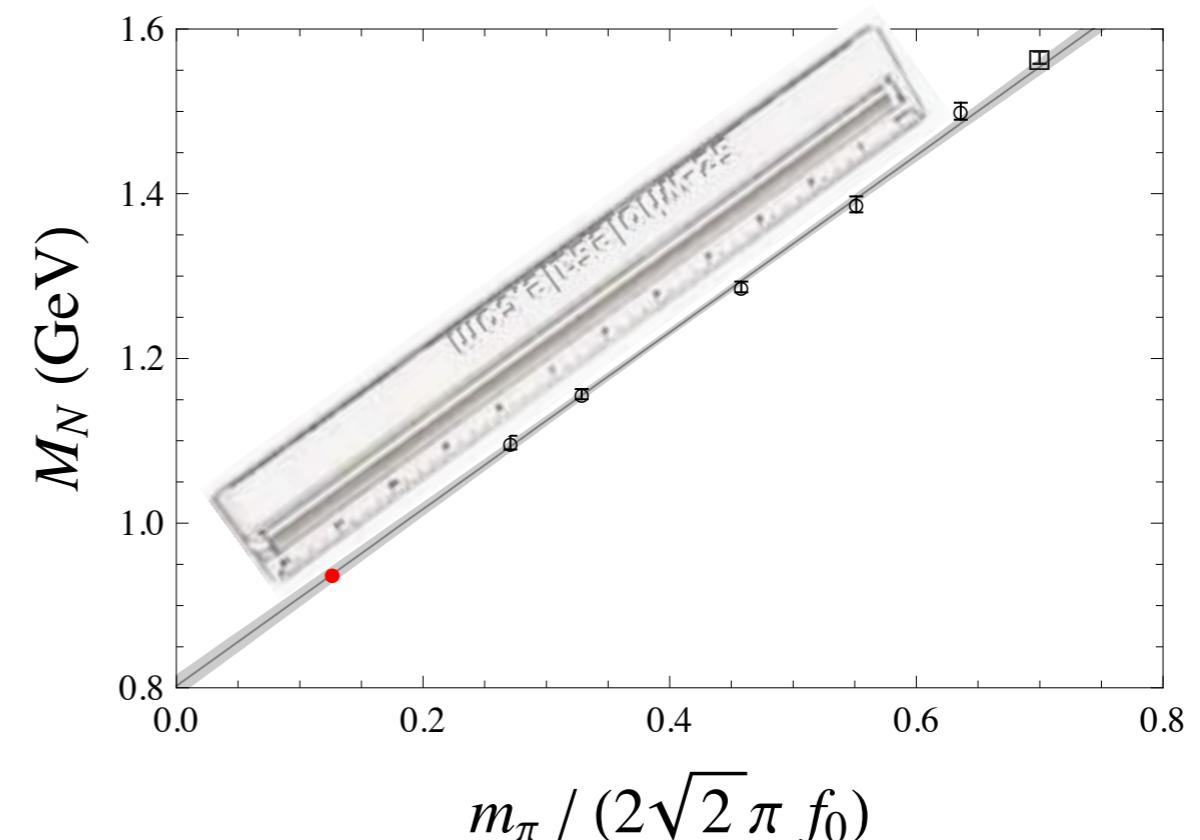
Baryons in lattice QCD

Light quark mass dependence of M_N

NNLO – m_π^4 , with $g_A=1.2(1)$, $g_{\Delta N}=1.5(3)$



$M_N = \alpha_0^N + \alpha_1^N m_\pi$



NNLO Heavy Baryon Fit

$$M_N = 954 \pm 42 \pm 20 \text{ MeV}$$

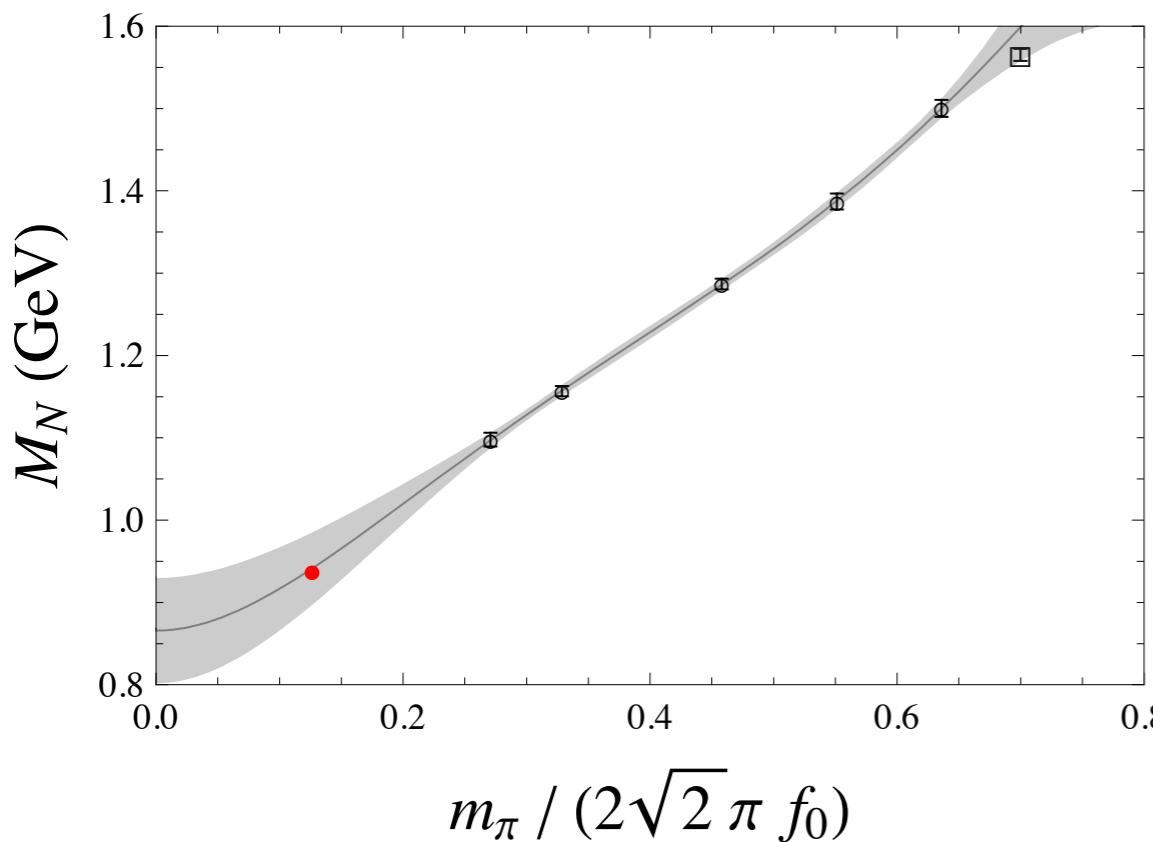
Ruler Approximation

$$M_N = \alpha_0^N + \alpha_1^N m_\pi$$

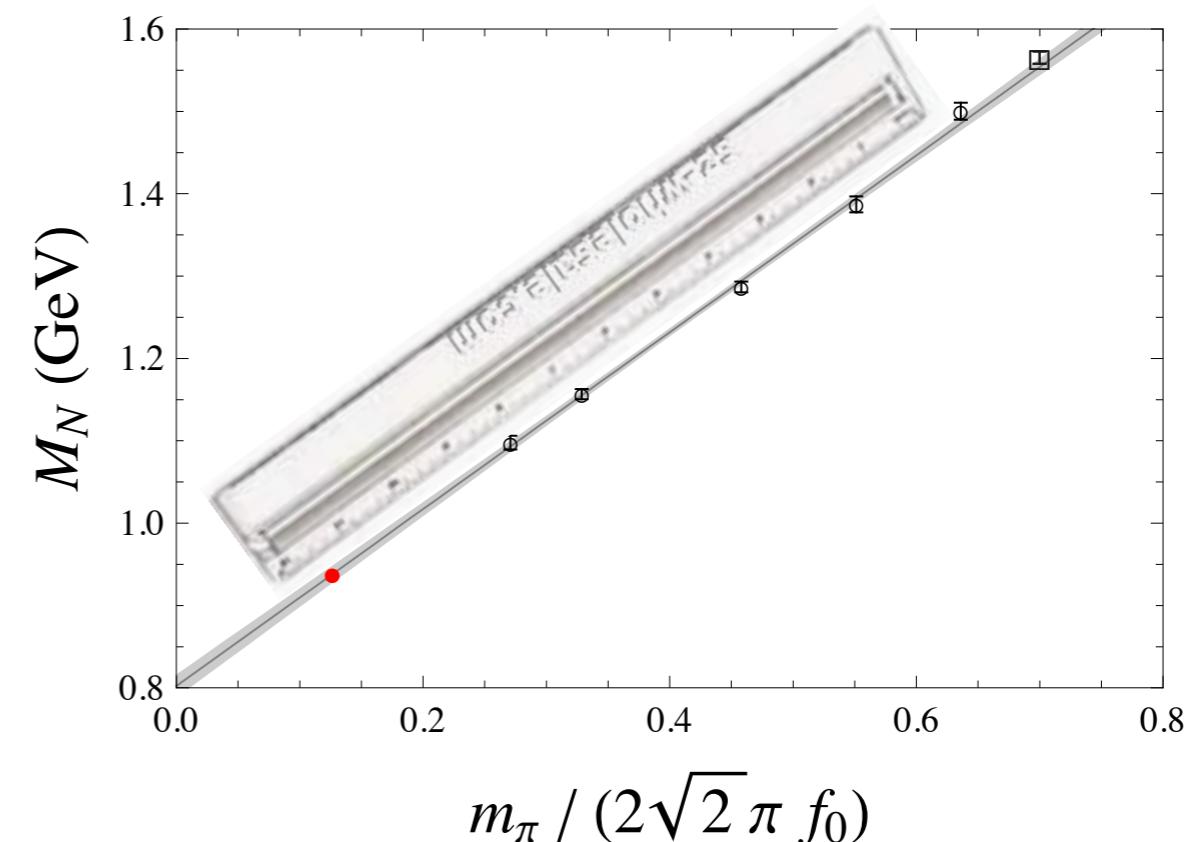
Baryons in lattice QCD

Light quark mass dependence of M_N

NNLO – m_π^4 , with $g_A=1.2(1)$, $g_{\Delta N}=1.5(3)$



$M_N = \alpha_0^N + \alpha_1^N m_\pi$



NNLO Heavy Baryon Fit

$$M_N = 954 \pm 42 \pm 20 \text{ MeV}$$

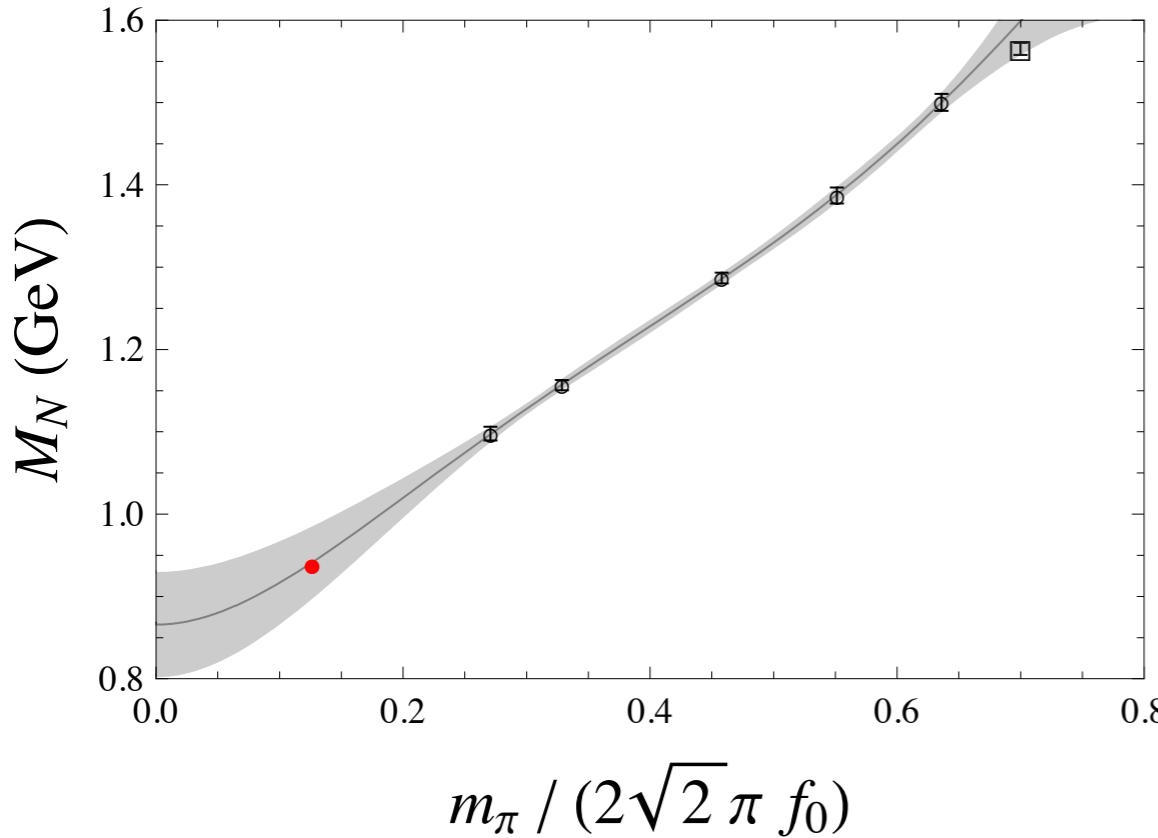
Ruler Approximation

$$\begin{aligned} M_N &= \alpha_0^N + \alpha_1^N m_\pi \\ &= 938 \pm 9 \text{ MeV} \end{aligned}$$

Baryons in lattice QCD

Light quark mass dependence of M_N

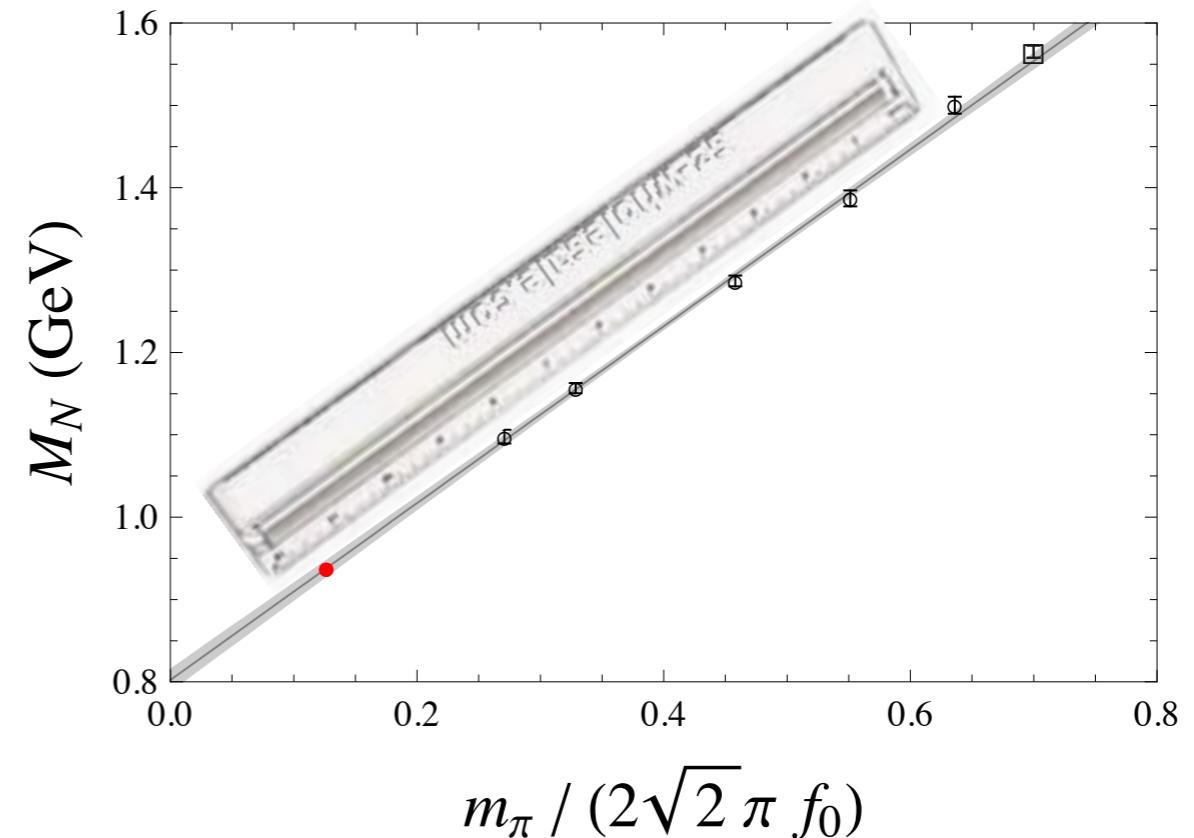
NNLO – m_π^4 , with $g_A=1.2(1)$, $g_{\Delta N}=1.5(3)$



NNLO Heavy Baryon Fit

$$M_N = 954 \pm 42 \pm 20 \text{ MeV}$$

$M_N = \alpha_0^N + \alpha_1^N m_\pi$



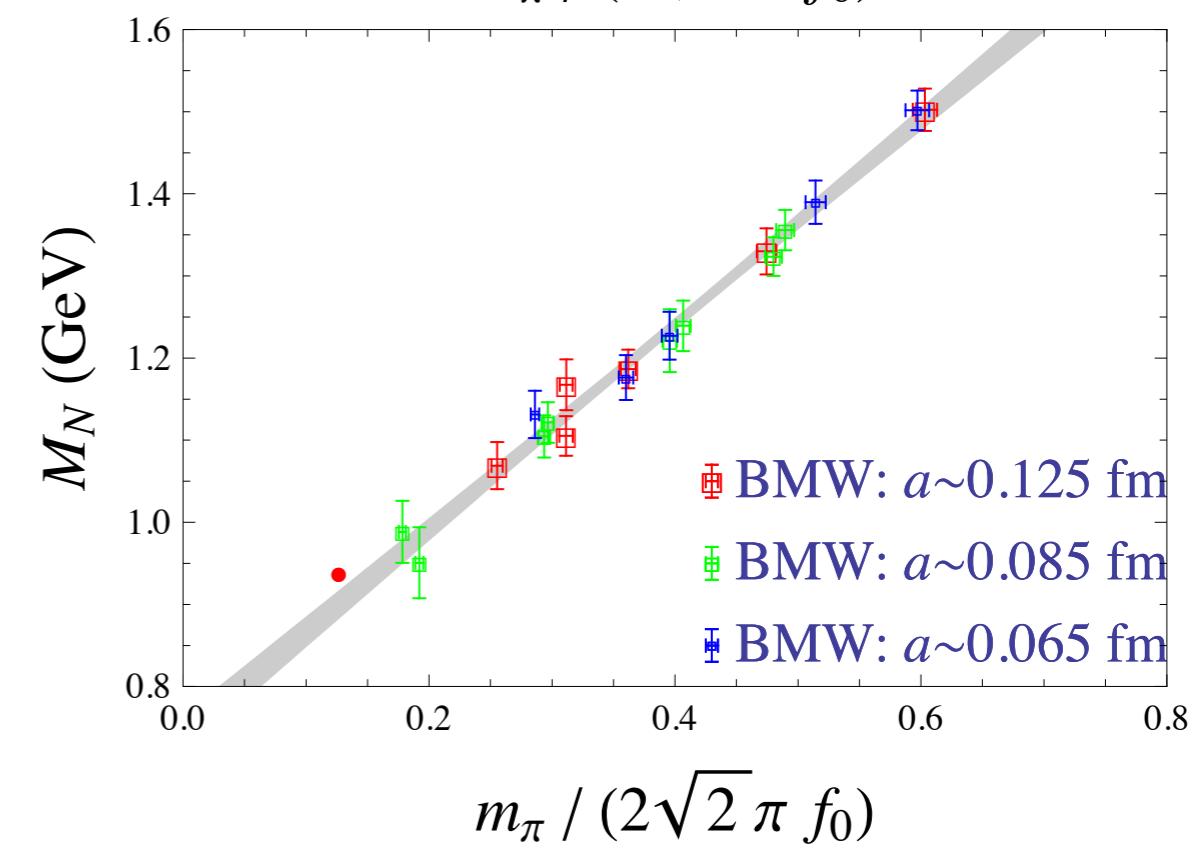
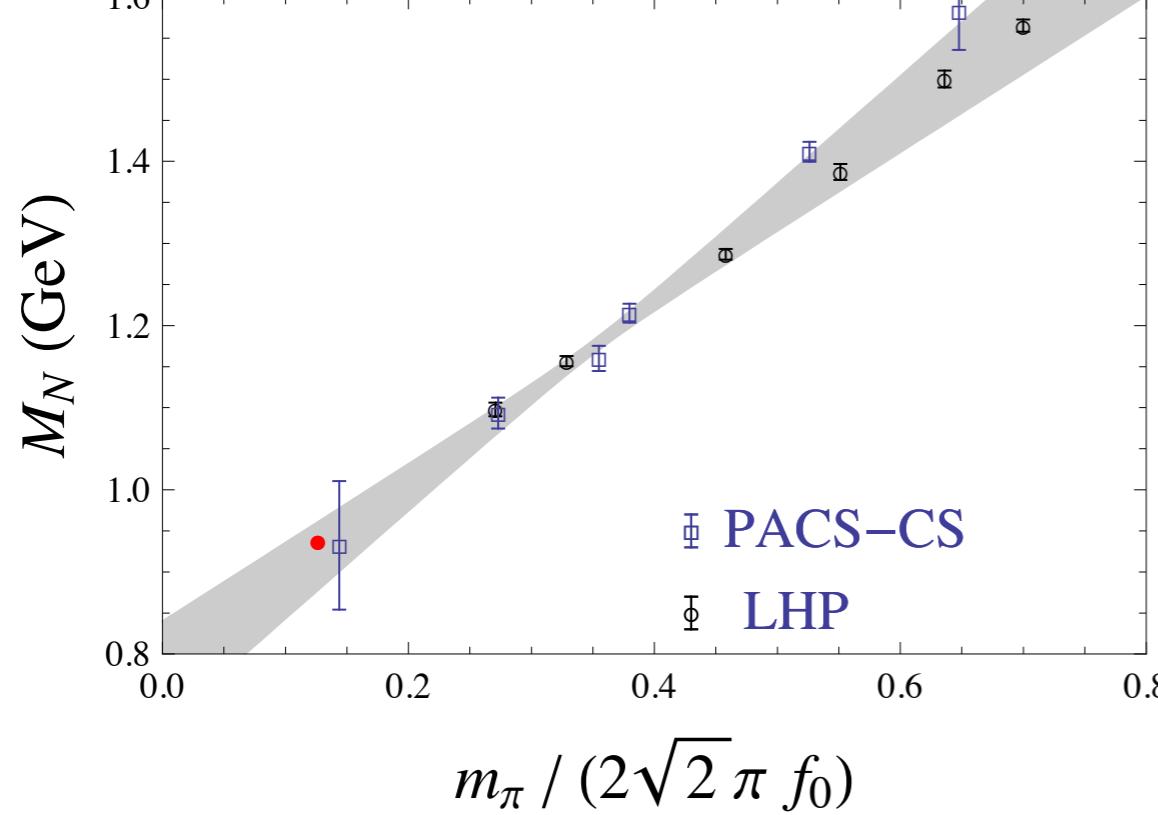
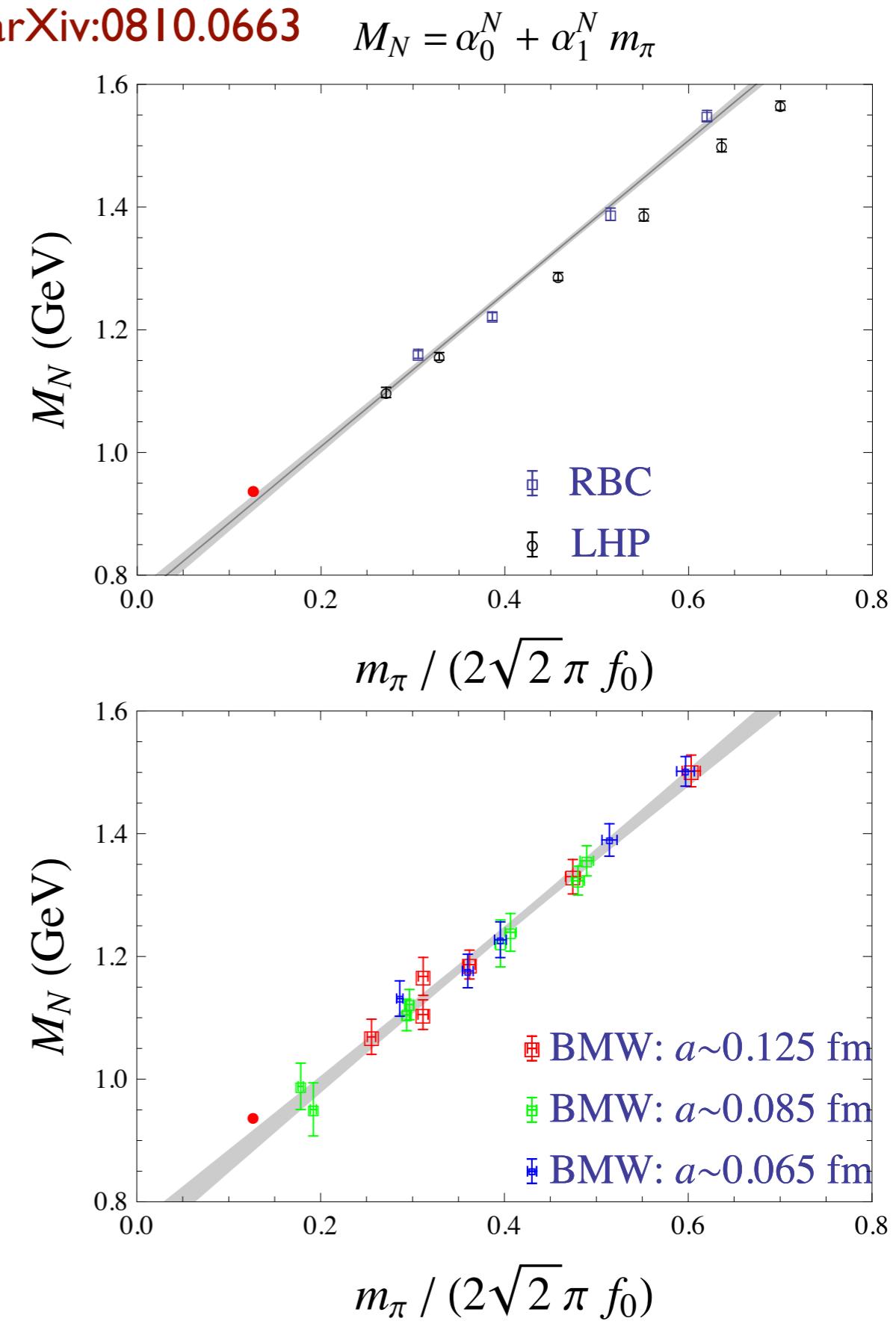
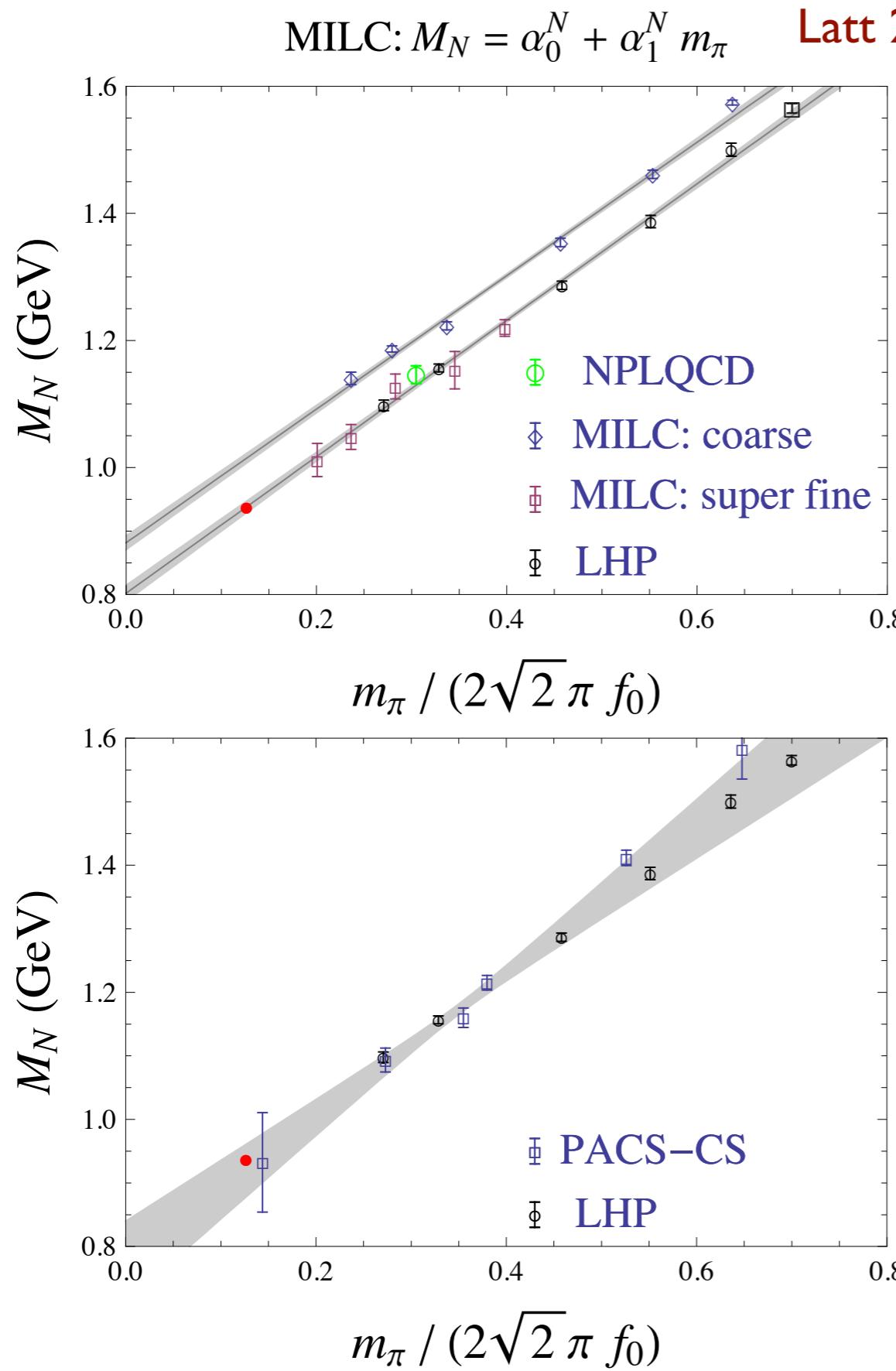
Ruler Approximation

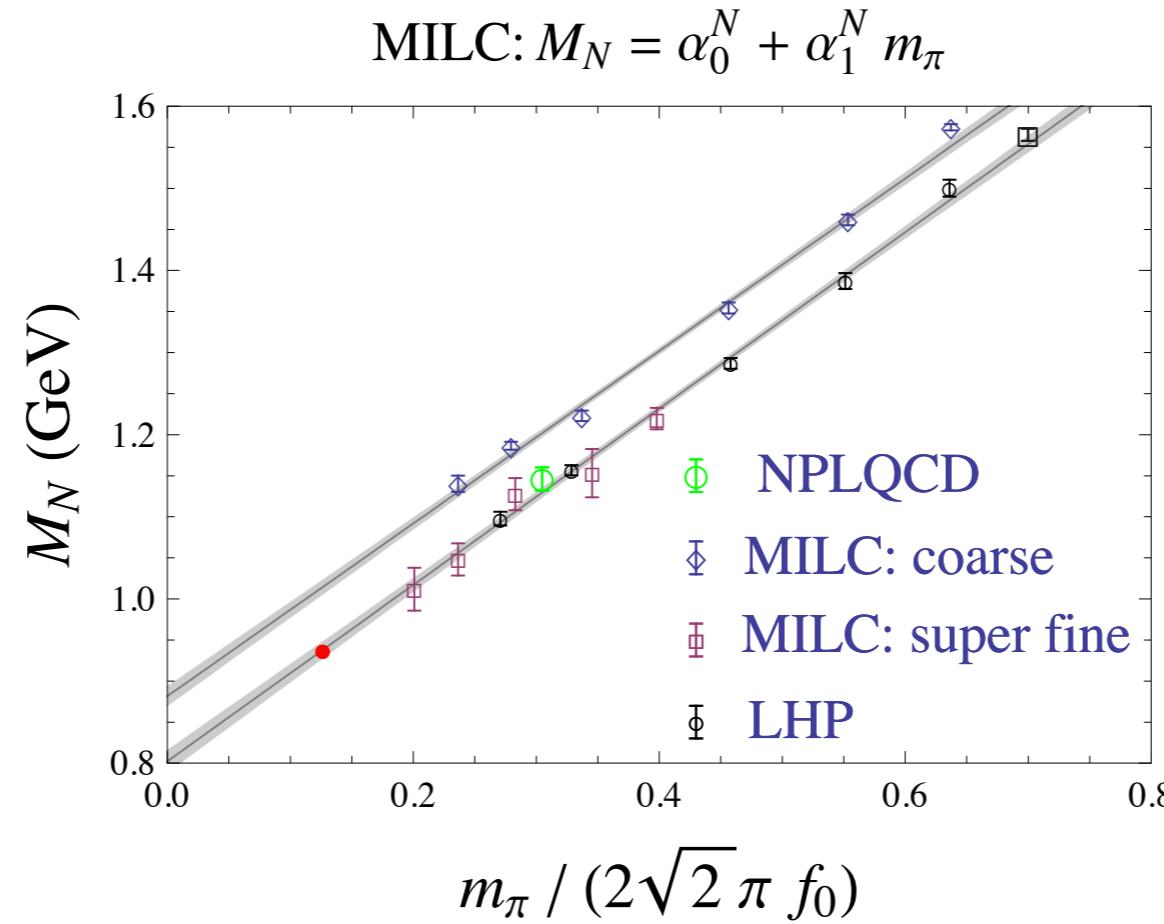
$$\begin{aligned} M_N &= \alpha_0^N + \alpha_1^N m_\pi \\ &= 938 \pm 9 \text{ MeV} \end{aligned}$$

I am not advocating this as
a good model for QCD!

Baryons in lattice QCD

Light quark mass dependence of M_N



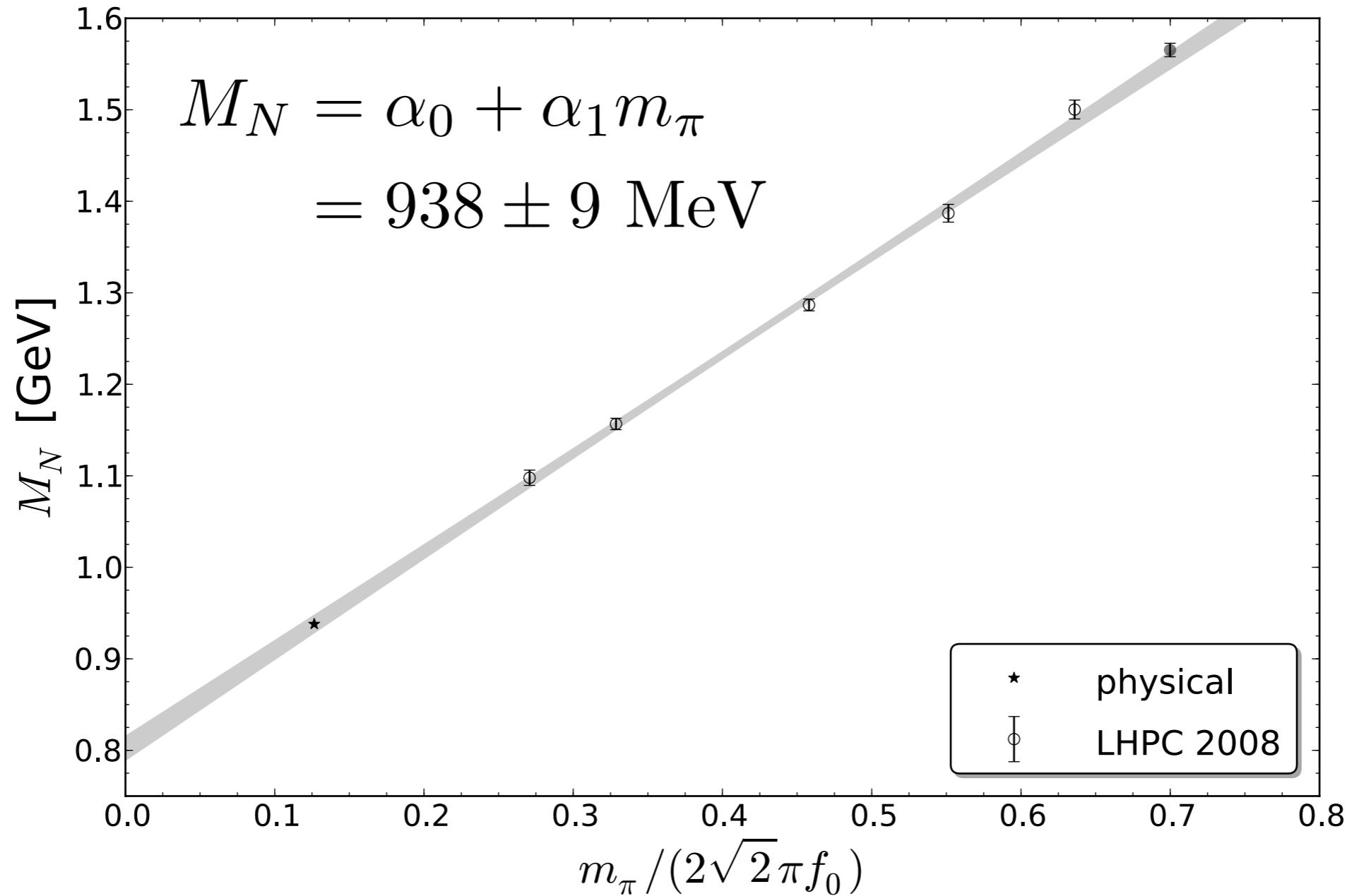


What does this teach us?

For these pion masses, there is a strong cancellation between LO, NLO and NNLO χ PT contributions

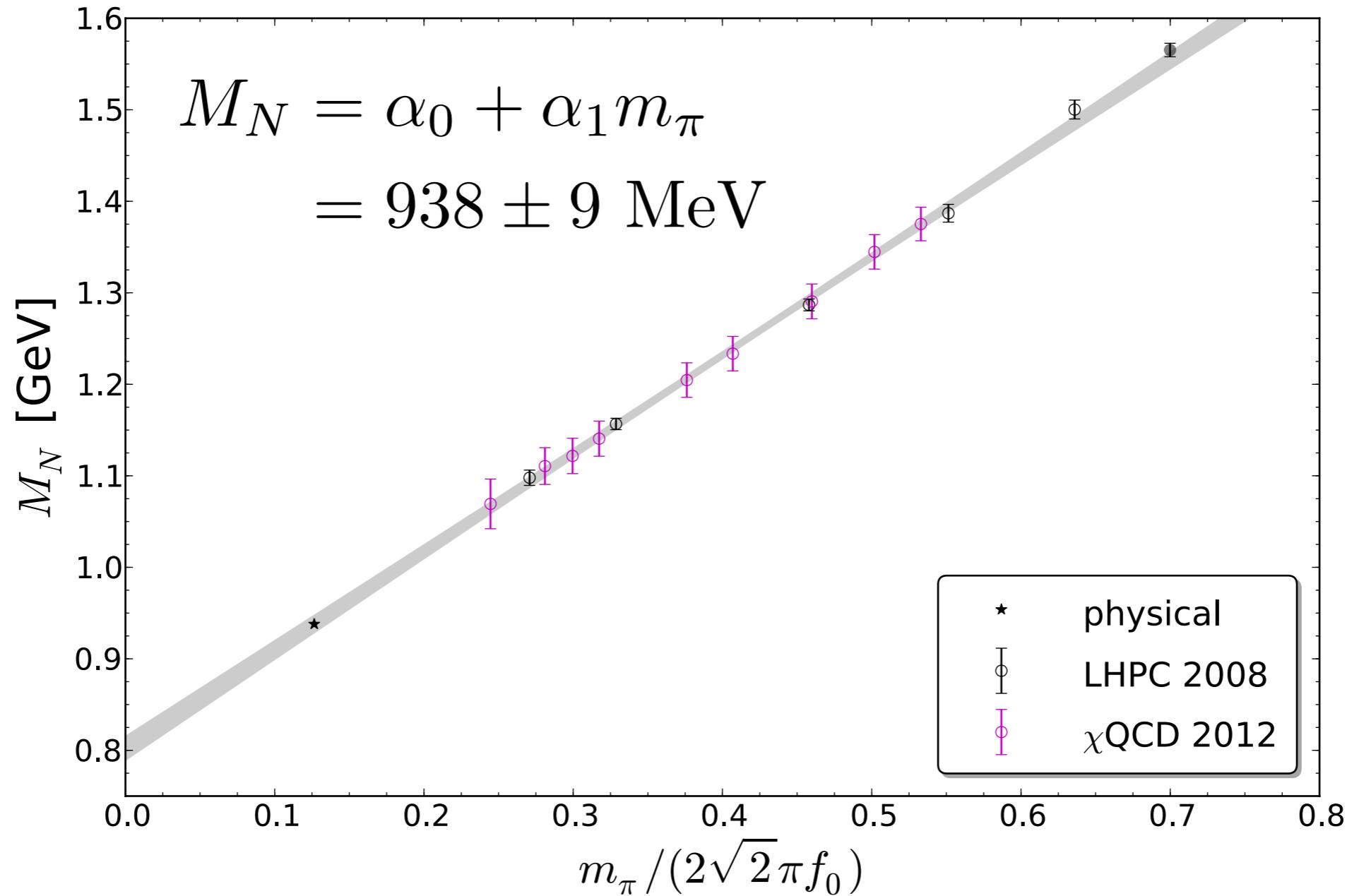
perhaps should have been expected given poor convergence (but just not a straight line!!!)

What is the status now (2012)?



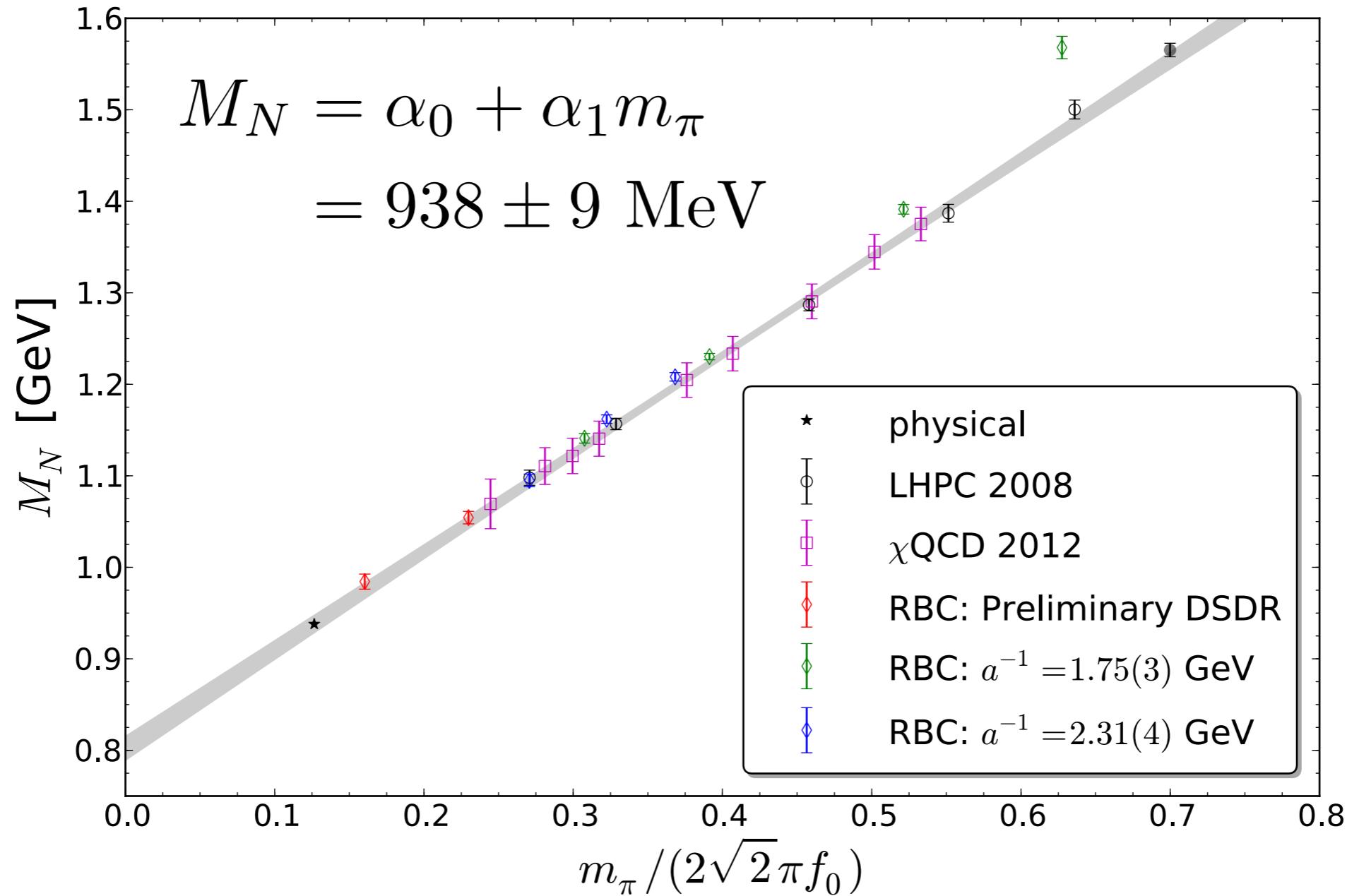
Physical point NOT included in fit

What is the status now (2012)?



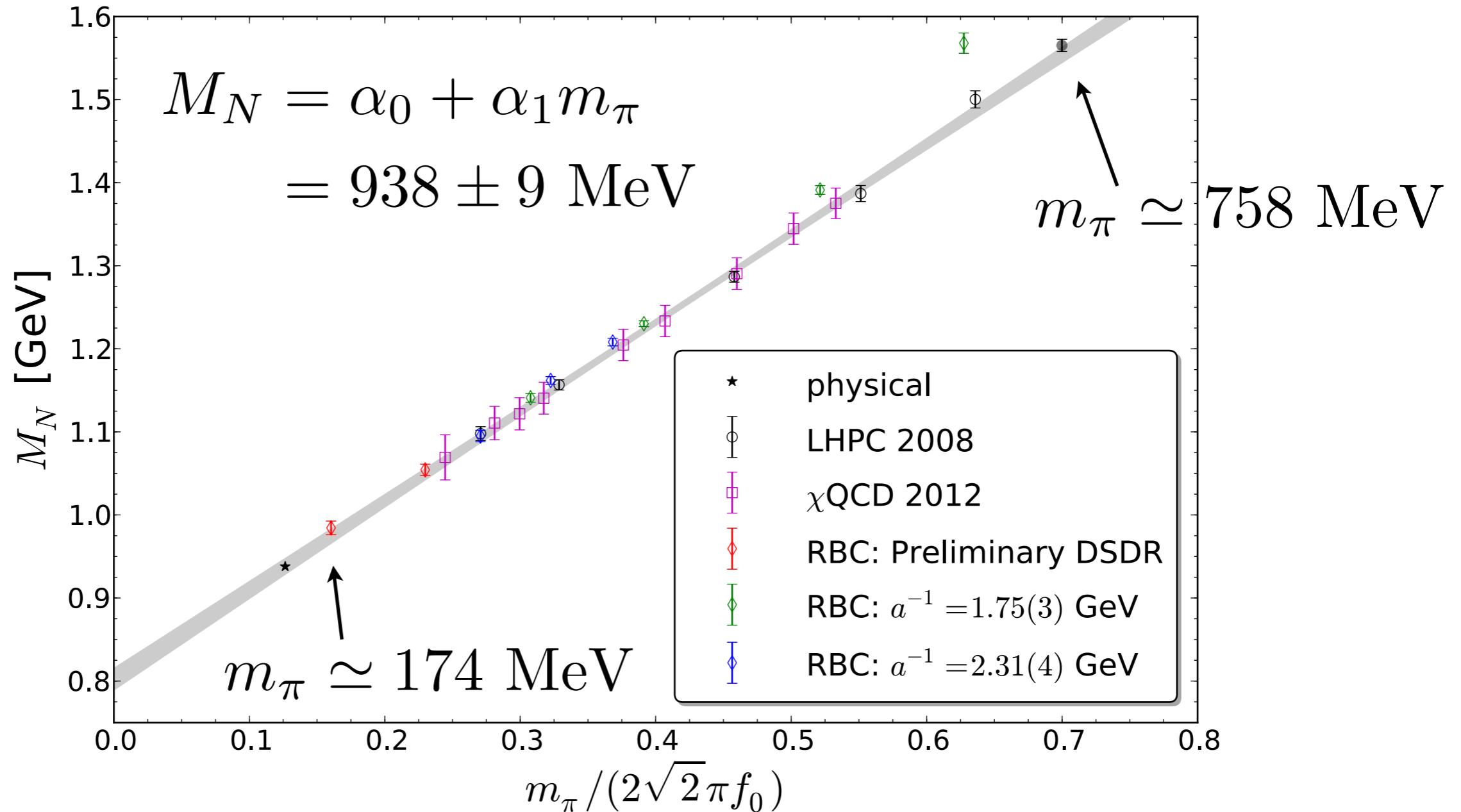
χQCD Collaboration uses Overlap Valence fermions on Domain-Wall (RBC-UKQCD) sea fermions

What is the status now (2012)?



RBC-UKQCD Collaboration uses Domain-Wall valence
and sea fermions

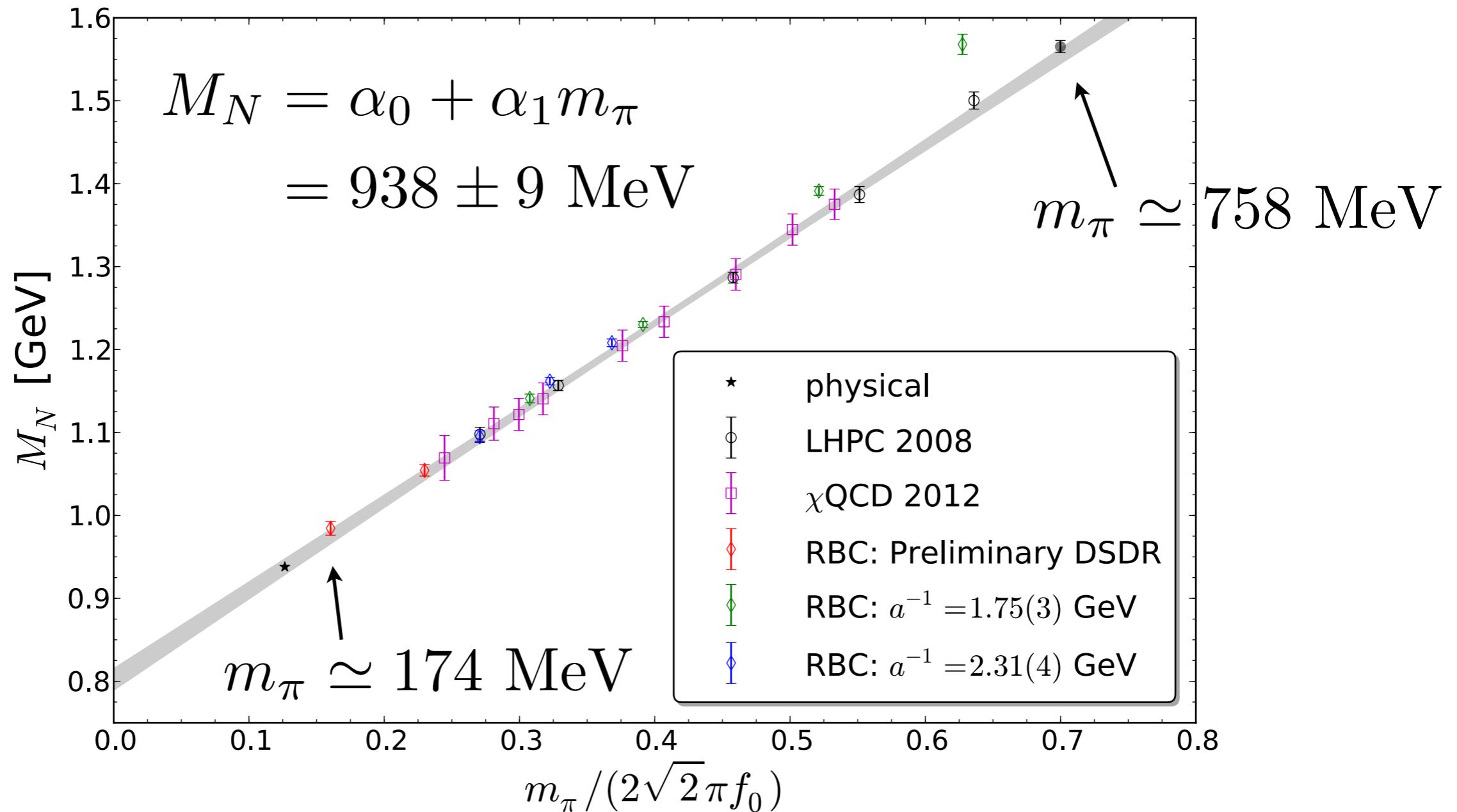
What is the status now (2012)?



Taking this seriously yields

$$\sigma_{\pi N} = 67 \pm 4 \text{ MeV}$$

What is the status now (2012)?



Taking this seriously yields
 $\sigma_{\pi N} = 67 \pm 4 \text{ MeV}$

I am not advocating this as
a good model for QCD!

Thank You