

Beyond the Standard Model with the LHC. Current status

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What is the LHC?

What did we know before the LHC start? What were our expectations?

Whether or not the LHC energy scale is an appropriate one?

What the LHC experiments tell us?

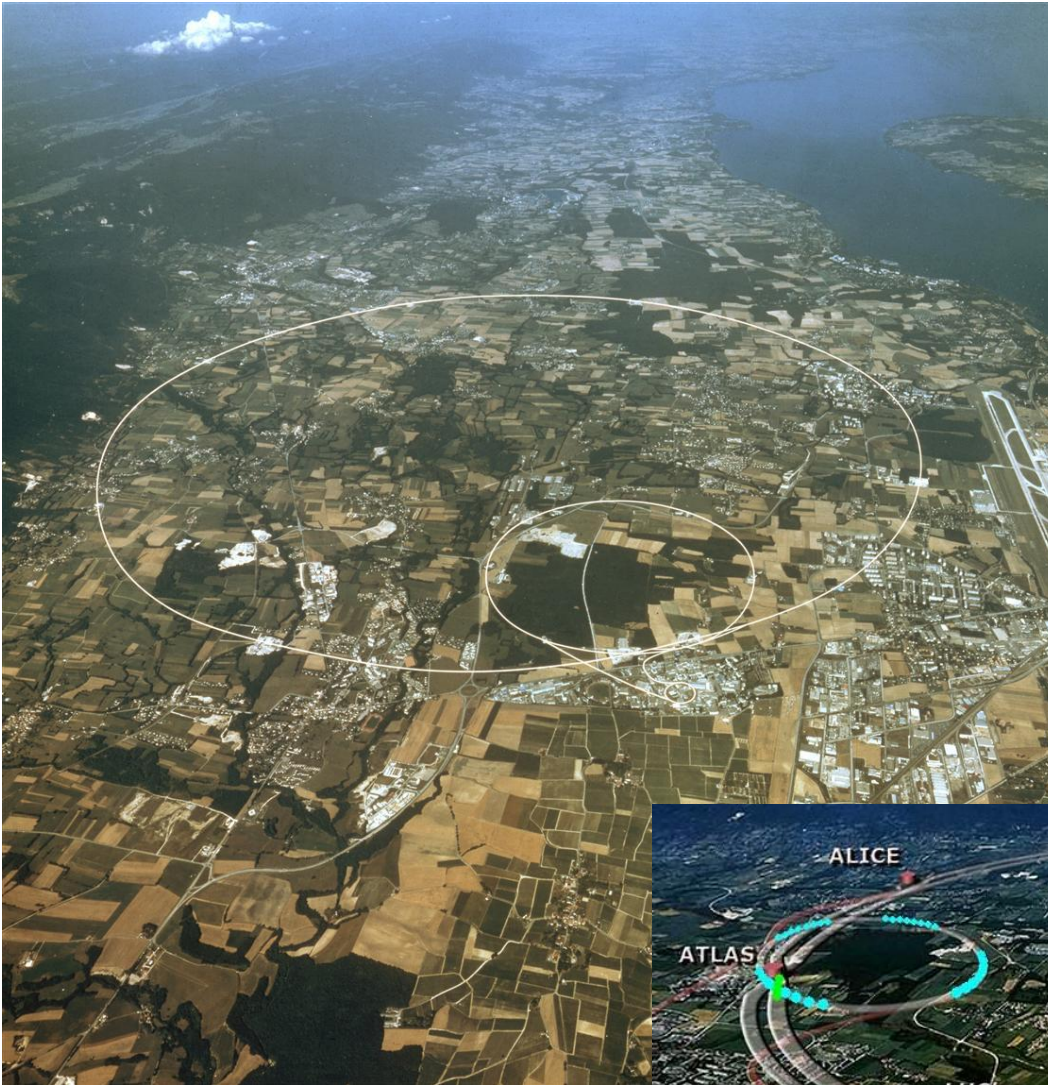
Higgs-like boson is found.
What is found and what does it mean?



BSM searches. And where we are?

What is the LHC?

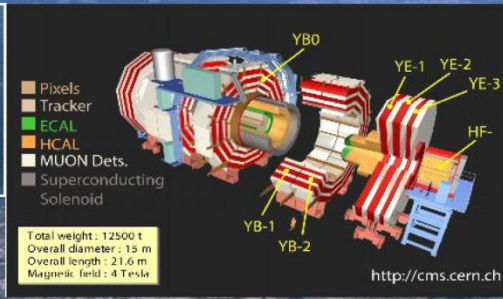
Collider LHC



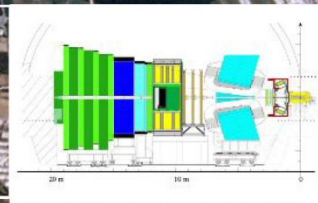
LHC collider (4 detectors: ATLAS, CMS, LHCb, ALICE)

27 km circumference, about 100 m underground

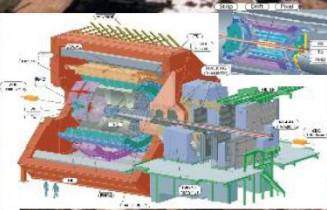
CMS
2900 Physicists
184 Institutions
38 countries



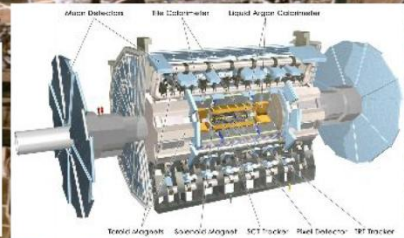
LHCb
700 Physicists
52 Institutions
15 countries



ALICE
1000 Physicists
105 Institutions
30 countries



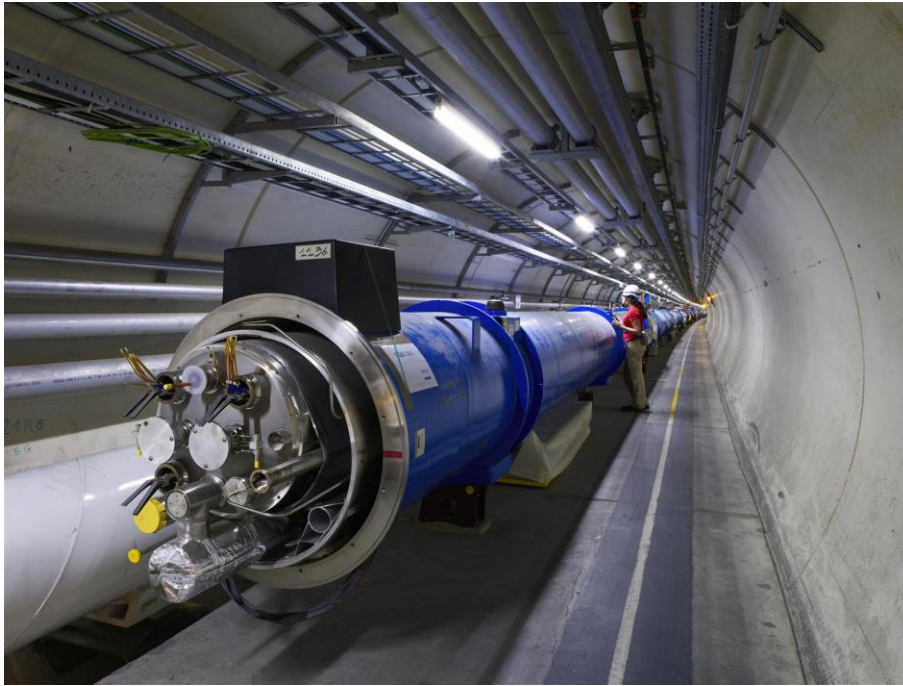
ATLAS
2800 Physicists
169 Institutions
37 countries



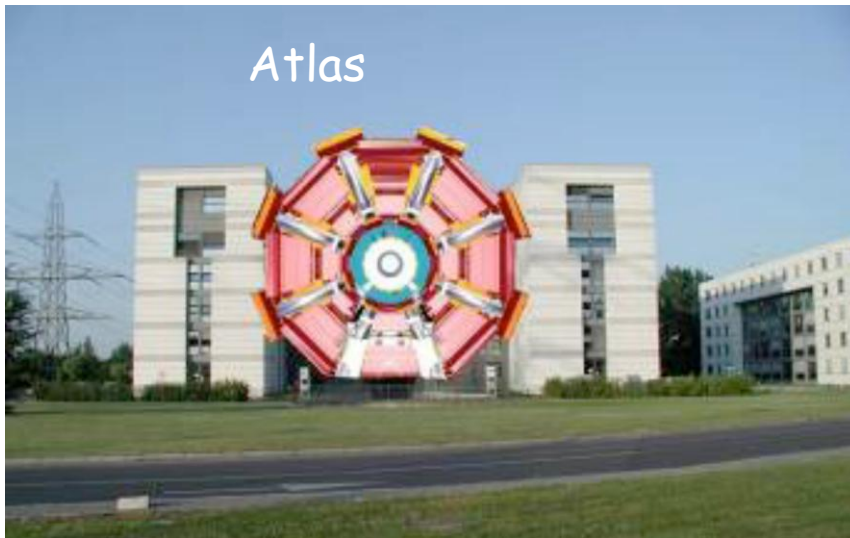
September 19 (2008) - an accident

2012 - run at 8 TeV

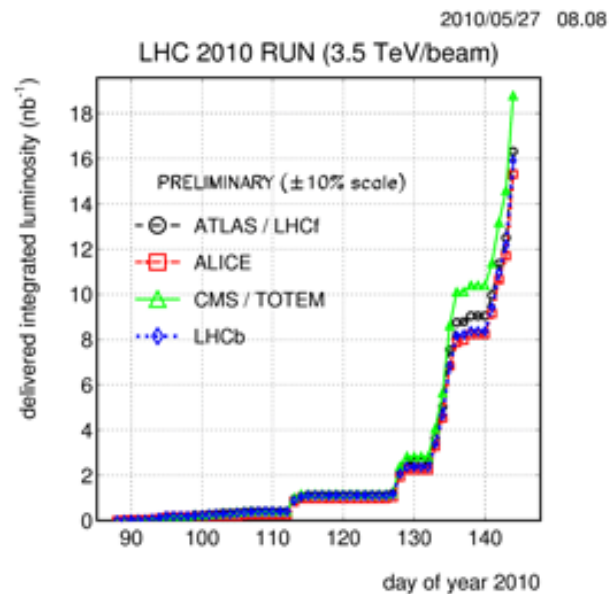
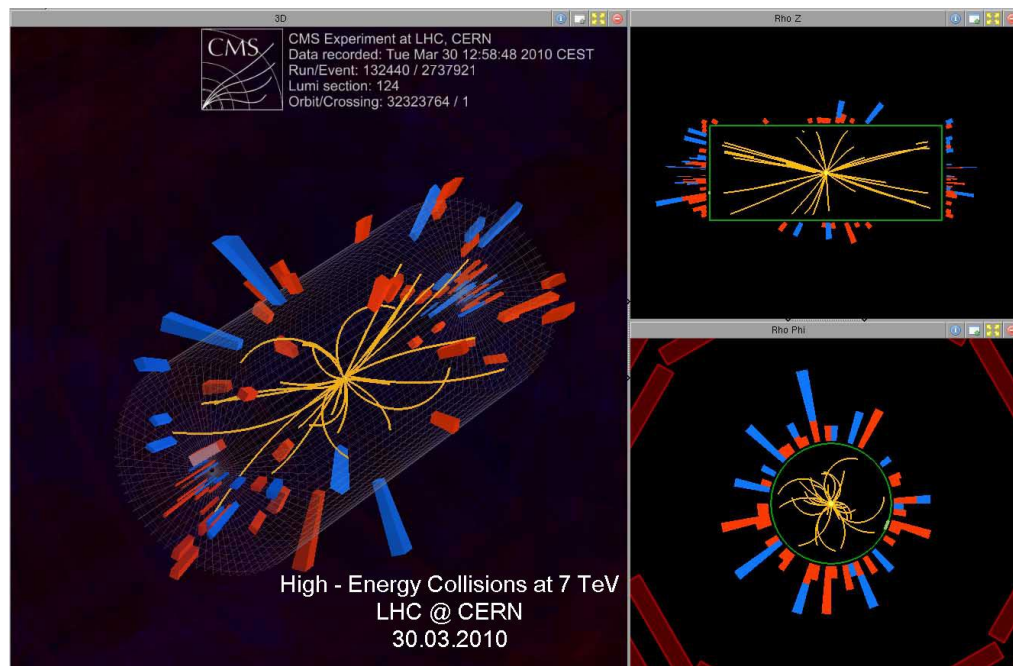
LHC is the most complicated and expensive project in fundamental science

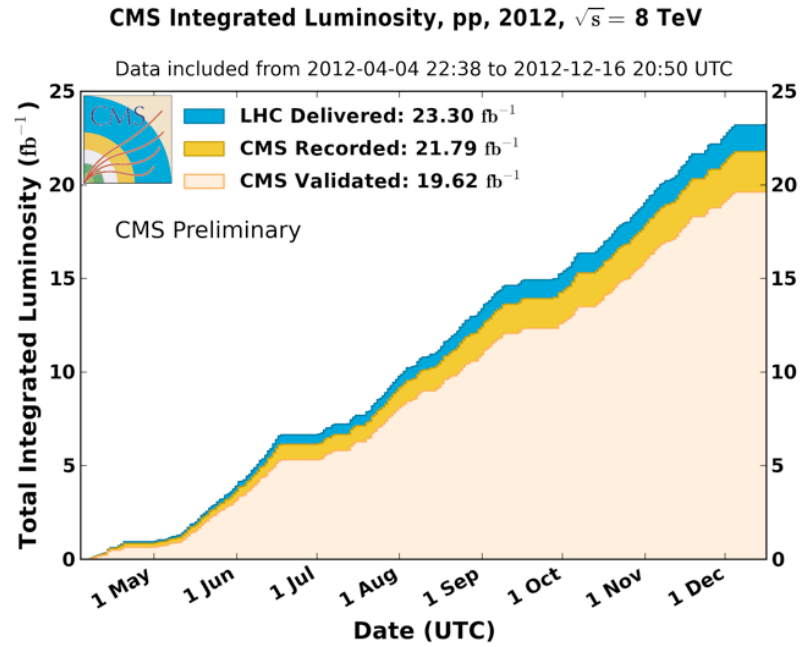
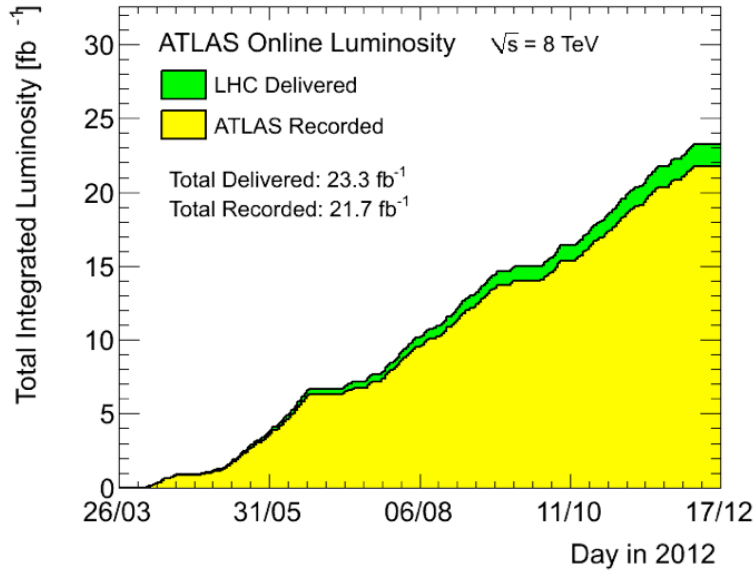


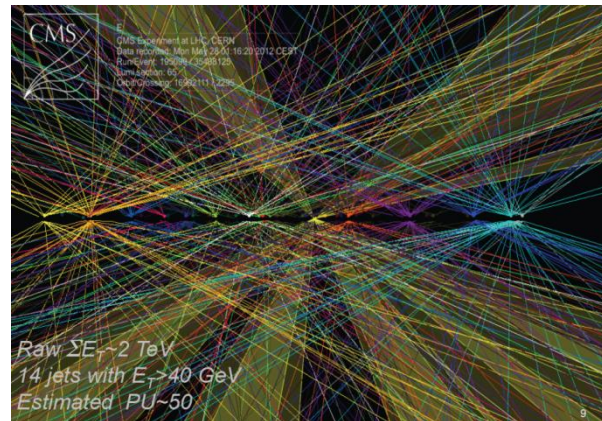
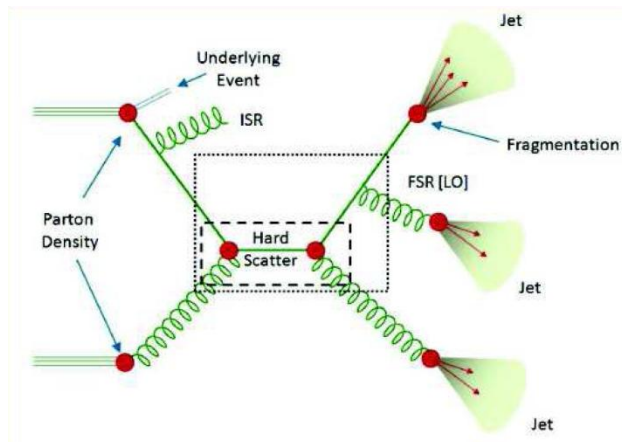
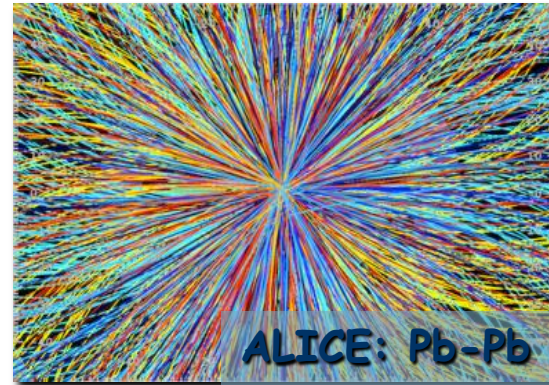
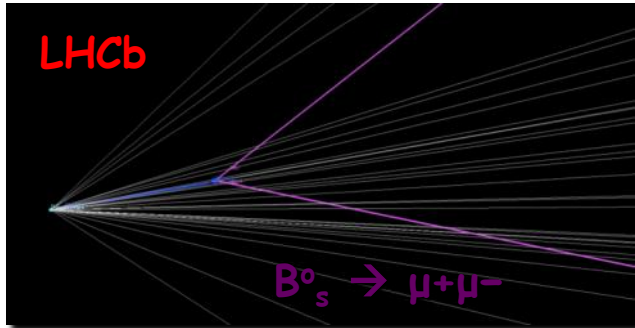
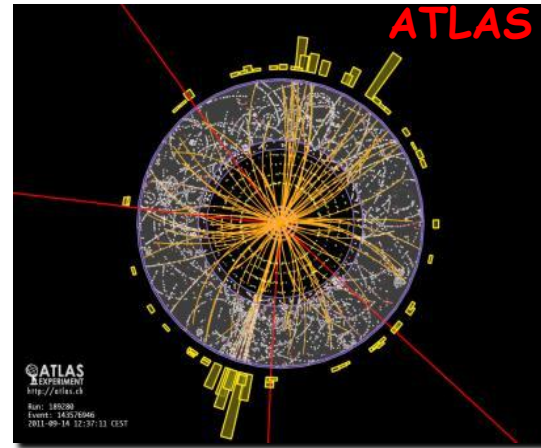
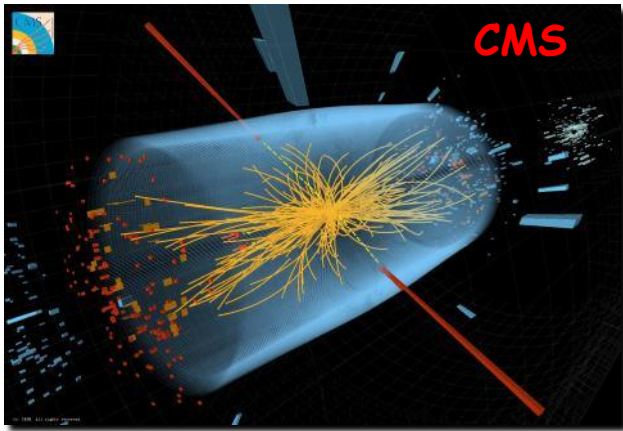
LHC vs Tevatron
Energy: 14 TeV vs 2 TeV
Luminosity: 10^{34} vs 10^{32} cm⁻²s⁻¹



30 March 2010
LHC&7TeV
has started







LHC physics programme

ATLAS and CMS (multipurpose detectors), ALICE and LHCb (dedicated detectors)

Detail studies of various SM processes (including diffraction) and comparisons to NLO (Next to Leading Order), NNLO computations

Search for the Higgs boson in various production and decay modes, measurements the Higgs properties

Search for deviations from SM in top quark production (pair/single) and decays, search for anomalous top properties expected for the heaviest SM particle

Search for best motivated BSM scenarios:
supersymmetry, extra dimensions, new strong dynamics
Model independent searches (Leptoquarks, Leptogluons, Z' , W' , ...)

Search for any other possible exotics (unparticles, hidden valleys...)

Detail studies of b-physics, b-meson oscillations, CP violation, BSM in loops

Detail studies of strongly interacting quark-gluon color medium

**What did we know before
the LHC start?**

SM - the quantum field theory
based on few principles and requirements :

- gauge invariance with lowest dimension (dimension 4) operators;
SM gauge group: $SU(3)_c \times SU(2)_L \times U(1)_y$

- correct electromagnetic neutral currents and
(V-A) charged currents (Fermi):

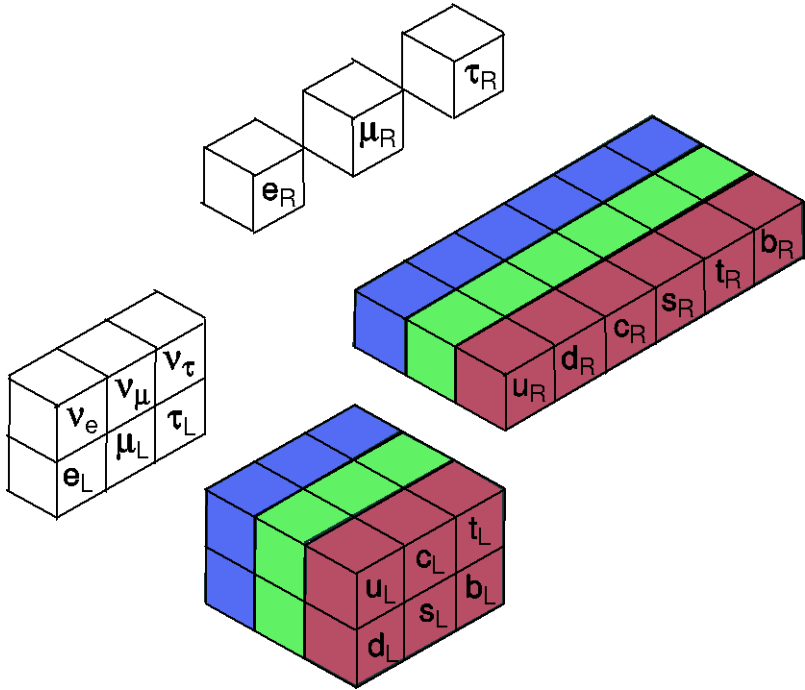
$$\frac{G_F}{\sqrt{2}} \cdot [\bar{\nu}_\mu \cdot \gamma_\alpha (1 - \gamma_5) \cdot \mu] \cdot [\bar{e} \cdot \gamma_\alpha (1 - \gamma_5) \cdot \nu_e] + h.c.$$

- 3 generations without chiral anomalies

- Higgs mechanism of spontaneous symmetry breaking

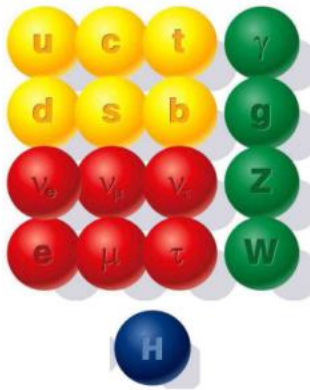
Standard Model - one of the main intellectual achievement for about
last 50 years, a result of many theoretical and experimental studies

Fermions are combined into 3 generations forming left doublets and right singlets with respect to weak isospin



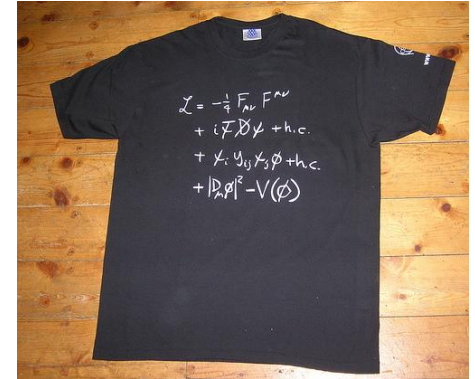
$$f_{L,R} = \frac{1}{2}(1 \mp \gamma_5)f$$

$$I_f^{3L,3R} = \pm \frac{1}{2}, 0 : \begin{aligned} L_1 &= \begin{pmatrix} \nu_e \\ e^- \end{pmatrix}_L, & e_{R1} &= e_R^-, & Q_1 &= \begin{pmatrix} u \\ d \end{pmatrix}_L, & u_{R1} &= u_R, & d_{R1} &= d_R \\ L_2 &= \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix}_L, & e_{R2} &= \mu_R^-, & Q_2 &= \begin{pmatrix} c \\ s \end{pmatrix}_L, & u_{R2} &= c_R, & d_{R2} &= s_R \\ L_3 &= \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix}_L, & e_{R3} &= \tau_R^-, & Q_3 &= \begin{pmatrix} t \\ b \end{pmatrix}_L, & u_{R3} &= t_R, & d_{R3} &= b_R \end{aligned}$$



Standard Model

$$SU(2)_L \times U(1)_Y \times SU(3)_c$$



$$L = -\frac{1}{4} W_{\mu\nu}^i (W^{\mu\nu})^i - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} G_{\mu\nu}^a (G^{\mu\nu})^a + \sum_{f=\ell,q} \bar{\Psi}_L^f (iD_\mu^L \gamma^\mu) \Psi_L^f + \sum_{f=\ell,q} \bar{\Psi}_R^f (iD_\mu^R \gamma^\mu) \Psi_R^f + L_H$$

$$L_H = L_\Phi + L_{Yukawa}$$

$$L_\Phi = D_\mu \Phi^\dagger D^\mu \Phi - \mu^2 \Phi^\dagger \Phi - \lambda (\Phi^\dagger \Phi)^4$$

$$L_{Yukawa} = -\Gamma_d^{ij} \bar{Q}_L^i \Phi d_R^j + h.c. - \Gamma_u^{ij} \bar{Q}_L^i \Phi^C u_R^j + h.c. - \Gamma_e^{ij} \bar{L}_L^i \Phi e_R^j + h.c.$$

$$B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$$

$$G_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g_S f^{abc} A_\mu^b A_\nu^c$$

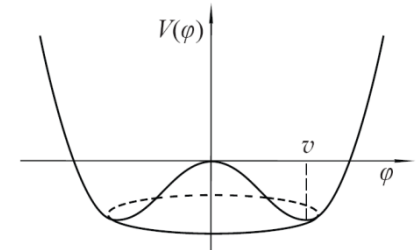
$$W_{\mu\nu}^i = \partial_\mu W_\nu^i - \partial_\nu W_\mu^i + g_2 \varepsilon^{ijk} W_\mu^j W_\nu^k$$

$$D_\mu^L = \partial_\mu - ig_2 W_\mu^i \tau^i - ig_1 B_\mu \left(\frac{Y_L^f}{2} \right) - ig_S A_\mu^a t^a$$

$$D_\mu^R = \partial_\mu - ig_1 B_\mu \left(\frac{Y_R^f}{2} \right) - ig_S A_\mu^a t^a$$

$$i = 1, 2, 3; a = 1, \dots, 8,$$

$$Y_f = 2Q_f - 2I_f^3 \Rightarrow Y_{L_i} = -1, Y_{e_{R_i}} = -2, Y_{Q_i} = \frac{1}{3}, Y_{u_{R_i}} = \frac{4}{3}, Y_{d_{R_i}} = -\frac{2}{3}$$



$$M_V, M_h, M_f \sim v$$

A very elegant theoretical construction!

$$L_{SM} = L_{Gauge} + L_{FG} + L_H$$

**Kinetic terms for the gauge fields;
Interaction terms of the gauge fields**

**Kinetic terms for fermions;
Interactions of fermions with the gauge fields
(NC and CC currents)**

**Kinetic and self-interaction terms for the higgs boson fields;
Higgs - gauge boson interaction terms;
Higgs-fermion interaction terms;
Mass terms for the gauge bosons and fermions;
+ Goldstone bosons and ghosts interactions**

$$L_H = \frac{1}{2}(\partial^\mu h)(\partial_\mu h) + \frac{M_h^2}{2}h^2 - \frac{M_h^2}{2v}h^3 - \frac{M_h^2}{8v^2}h^4 +$$

$$+ (M_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2}M_Z^2 Z_\mu Z^\mu) \left(1 + \frac{h}{v}\right)^2 - \sum_f m_f \bar{f} f \left(1 + \frac{h}{v}\right)$$

$$M_H^2 = 2\lambda v^2$$

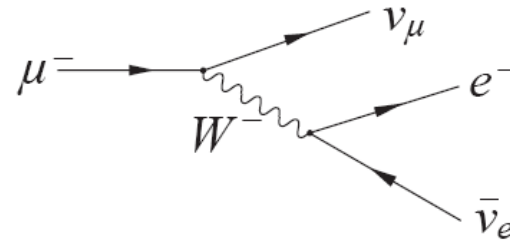
Electroweak Standard Model

The Fermi constant G_F is measured with high precision from muon life time

$$G_F = 1.166\,378\,7(6) \times 10^{-5} \text{ GeV}^{-2}$$

Since the muon mass $m_\mu \ll M_W$ one can neglect the W -boson mass in the propagator and immediately get the following relation

$$\frac{g_2^2}{8M_W^2} = \frac{G_F}{\sqrt{2}}$$



As we have seen the W boson mass is obtained in SM due to the Higgs mechanism and proportional to the Higgs vacuum expectation value v

$$M_W^2 = \frac{1}{4}g_2^2v^2$$

$$v = \frac{1}{\sqrt{\sqrt{2}G_F}} = 246.22 \text{ GeV}$$

From these two relations we obtain

At this point one can see the power of gauge invariance principle, g_2 is the same gauge coupling

The Higgs field expectation value v is determined by the Fermi constant G_F introduced long before the Higgs mechanism appeared!

$$M_W^2 = \frac{1}{4}g_2^2v^2 \quad g_2s_W = e \quad M_W = M_Zc_W$$

$$M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi\alpha_{em}}{\sqrt{2}G_F} \equiv A_0^2$$

$\alpha_{em} = e^2/4\pi$ is the electromagnetic fine structure constant. The low energy value follows mainly from the electron anomalous magnetic measurements

$$\alpha_{em} = (137.035\,999\,074(44))^{-1}$$

One gets A_0 very precisely from low energy measurements

$$A_0 = 37.2804 \text{ GeV}$$

From the other hand one gets A_0 from measured values for the masses of W and Z bosons

$$\begin{array}{l} M_W = 80.385 \pm 0.015 \text{ GeV} \\ M_Z = 91.1876 \pm 0.0021 \text{ GeV} \end{array} \quad \longrightarrow \quad A_0 = 37.95 \text{ GeV}$$

Values are close. The difference is about 1%.

CC and NC interactions of SM fermions, as we know already, have the following structure

$$L_{CC} = \frac{g_2}{2\sqrt{2}} \sum_{ij} V_{ij} \bar{u}_i \gamma_\mu (1 - \gamma_5) d_j = \frac{e}{2\sqrt{2}s_W} \sum_{ij} V_{ij} \bar{u}_i \gamma_\mu (1 - \gamma_5) d_j,$$

$$L_{NC} = e \sum_f Q_f \bar{f} \gamma_\mu f + \frac{e}{4s_W c_W} \sum_f \bar{f} \gamma_\mu (v_f - a_f \gamma_5) f Z^\mu,$$

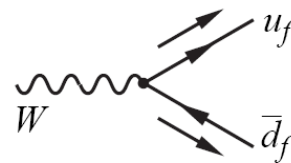
where V_{ij} is the CKM matrix element, $i, j = 1, 2, 3$ - number of fermion generation

$$v_{u_i} = 1 - \frac{8}{3}s_W^2, \quad a_{u_i} = 1; \quad v_{d_i} = -1 + \frac{4}{3}s_W^2, \quad a_{d_i} = -1$$

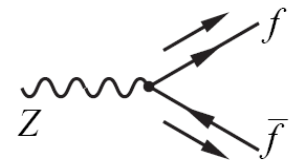
$$v_\ell = -1 + 4s_W^2, \quad a_\ell = -1; \quad v_\nu = 1, \quad a_\nu = 1.$$

$$v_f = 2T_3^f - 4Q_f s_W^2, \quad a_f = 2T_3^f$$

The Feynman rules following from L_{CC} and L_{NC} allow to get tree level formulas for the W and Z boson widths



$$\Gamma(W \rightarrow u_f \bar{d}_f) = |V_{ij}|^2 N_c \frac{\alpha}{12s_W^2} M_W,$$



$$\Gamma(Z \rightarrow f \bar{f}) = N_c \frac{\alpha M_Z}{12 \sin^2(2\theta_W)} [v_f^2 + a_f^2]$$

$N_c = 3$ for quarks, and $N_c = 1$ for leptons

Since CC for all fermions have the same $(V - A)$ structure one can very easily obtain branching fractions for W decay modes

$$\begin{aligned}\sum_q \text{Br}(W \rightarrow q\bar{q}) &= 2N_c \cdot \frac{1}{9} = \frac{2}{3} \\ \sum_\ell \text{Br}(W \rightarrow \ell\nu) &= 3 \cdot \frac{1}{9} = \frac{1}{3}\end{aligned}$$

Measured $\text{Br}(W \rightarrow \ell\nu) = (10.80 \pm 0.09)\%$ is in a reasonable agreement with simple tree level result $1/9 = 11\%$

QCD corrections to $\text{Br}(W \rightarrow q\bar{q})$ improved the agreement

The decay width of the Z -boson to neutrinos, the invisible decay mode, allows to measure the number of light ($m_\nu < M_Z/2$) neutrinos

$$\Gamma_{inv}^Z = \Gamma_{tot}^Z - \Gamma_{had}^Z - \Gamma_{\ell^+\ell^-}^Z$$

$$\Gamma_{tot}^Z = 2.4952 \pm 0.0023 \text{ GeV} \quad \text{is measured from the shape of the } Z\text{-boson resonance}$$

$$\Gamma_{\ell^+\ell^-}^Z = 83.984 \pm 0.086 \text{ MeV}$$

$$\Gamma_{had}^Z = 1744.4 \pm 2.0 \text{ MeV}$$

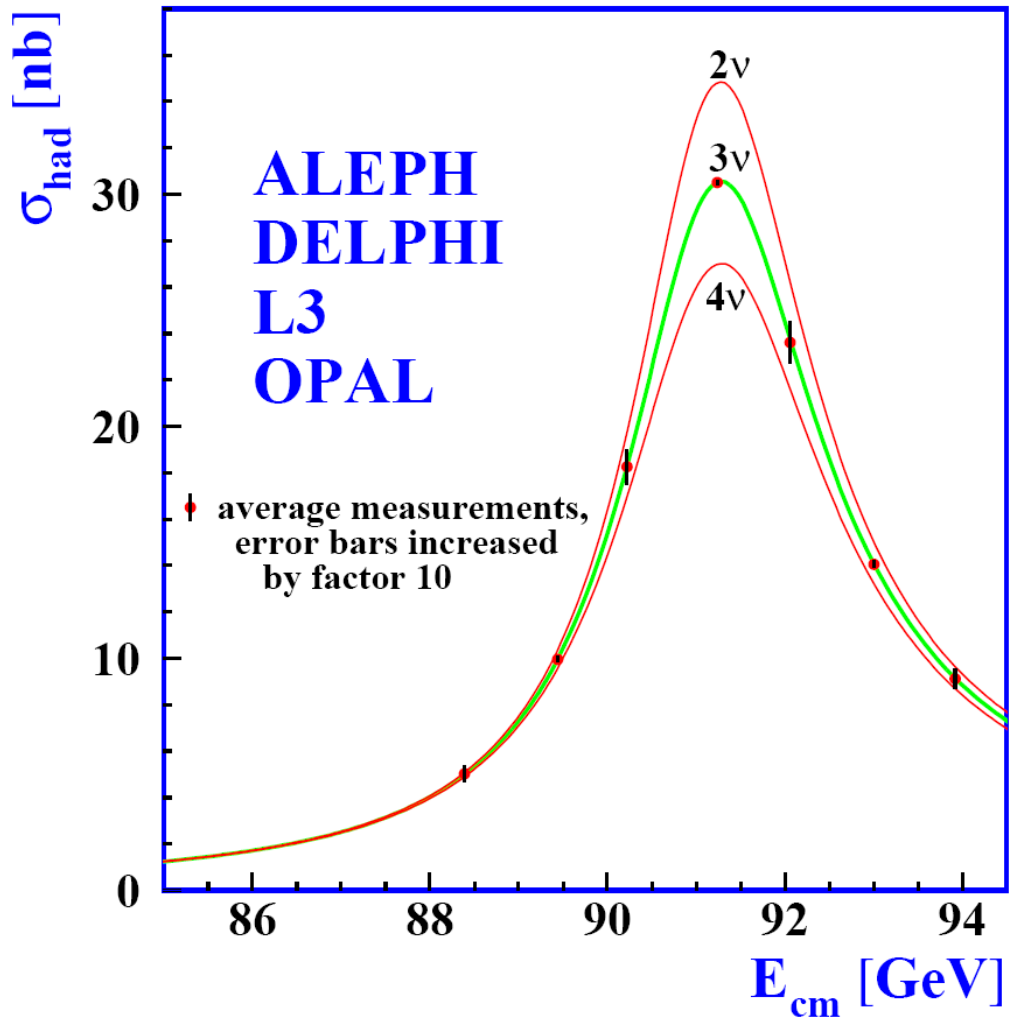
$$\longrightarrow \Gamma_{inv}^Z = 499.0 \pm 1.5 \text{ MeV}$$

$$\text{In SM} \quad \Gamma(Z \rightarrow f\bar{f}) = N_c \frac{\alpha M_Z}{12 \sin^2(2\theta_W)} [v_f^2 + a_f^2] \quad \Longrightarrow \quad \Gamma_{inv}^Z = \Gamma_{\nu\bar{\nu}}^Z = N_\nu \cdot \frac{\alpha M_Z}{12 \sin^2(2\theta_W)} (1 + 1)$$



$$2.9840 \pm 0.0082$$

Confirmation of 3 fermion generations assumed in the SM and observed in nature



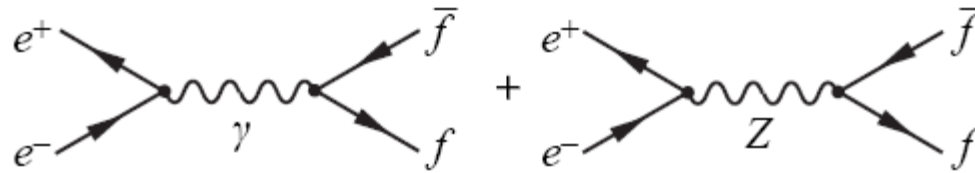
Another way to make this test

$$\frac{\Gamma_{inv}^Z}{\Gamma_{e^+e^-}^Z} = \frac{2N_\nu}{1 + (1 - 4s_W^2)^2}$$

The experimental value **5.942 ± 0.016**

$N_\nu = 3$ gives for the ratio about **5.970** in an agreement with the measured value ($s_W^2 = 0.2324$)

An important part of information about EW fermionic interactions and couplings comes from e^+e^- annihilation to fermion-antifermion pairs



$$\frac{d\sigma}{d\cos\theta} = \frac{2\pi\alpha^2}{4s} N_C \left\{ (1 + \cos^2\theta) \cdot [Q_f^2 - 2\chi_1 v_e v_f Q_f + \chi_2 (a_e^2 + v_e^2)(a_f^2 + v_f^2)] + 2\cos\theta [-2\chi_1 a_e a_f Q_f + 4\chi_2 a_e a_f v_e v_f] \right\}$$

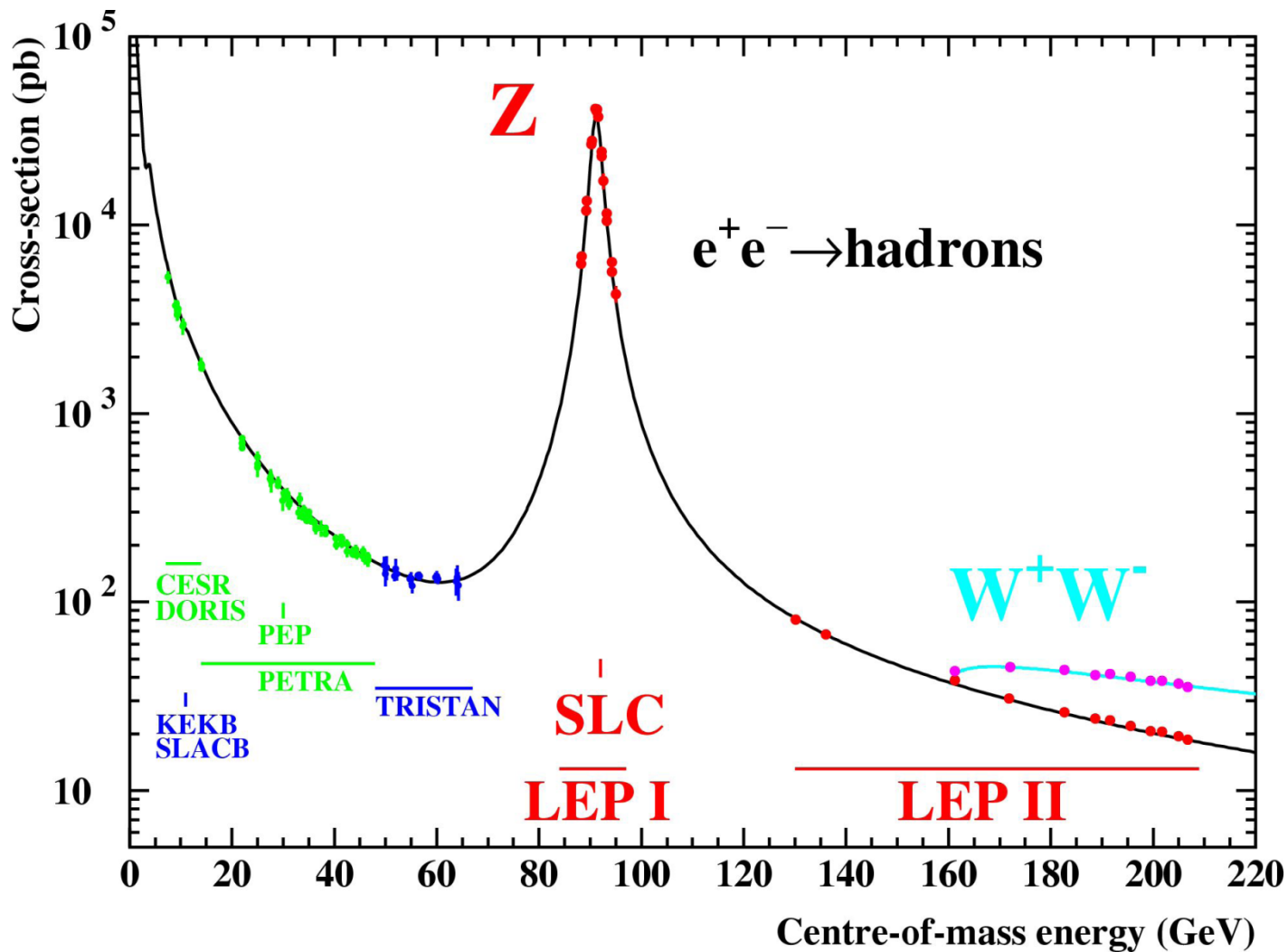
$$\chi_1 = \frac{1}{16s_W^2 c_W^2} \frac{s(s - M_Z^2)}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2},$$

$$\chi_2 = \frac{1}{256s_W^2 c_W^2} \frac{s^2}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2}.$$

In the region much below Z-boson pole one can neglect Z-boson exchange diagram and well known QED formula is restored

$$\frac{d\sigma}{d\cos\theta} = \frac{\pi\alpha^2}{2s} Q_f^2 N_C (1 + \cos^2\theta), \quad \sigma = \frac{4\pi\alpha^2}{3} Q_f^2 N_C$$

$N_c = 3$ for quarks, and $N_c = 1$ for leptons

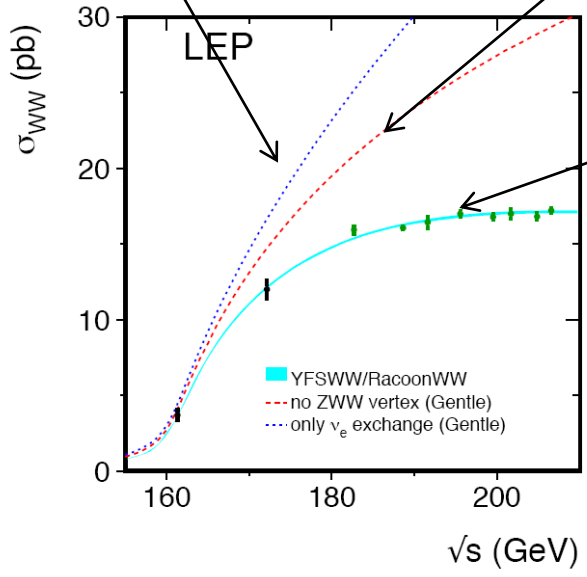
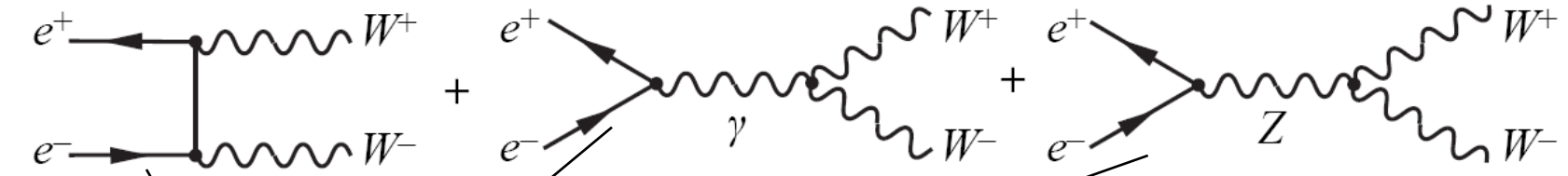


Well known example demonstrating correctness of the Yang-Mills interaction of gauge bosons is **W-boson pair production**. Triple gauge boson vertex $WW\gamma$ and WWZ have been tested at LEP2 ($e^+e^- \rightarrow W^+W^-$) and at the Tevatron ($q\bar{q} \rightarrow W^+W^-, q\bar{q}' \rightarrow W\gamma, q\bar{q}' \rightarrow WZ$).

The triple vertex of Yang-Mills interaction

$$\Gamma_{m_1 m_2 m_3}^{WW\gamma/Z}(p_1 p_2 p_3) = g_{\gamma,Z} [(p_1 - p_2)_{m_3} g_{m_1 m_2} + (p_3 - p_1)_{m_2} g_{m_1 m_3} + (p_2 - p_3)_{m_1} g_{m_2 m_3}]$$

$$g_\gamma = e, g_Z = g_2 c_W = e \frac{c_W}{s_W}$$



Three SM diagrams

The quartic gauge couplings $WW\gamma\gamma$, $WW\gamma Z$, $WWZZ$ have not been tested yet. This is challenging task for the LHC. It will require high luminosity regime at a linear collider

All terms of the SM Lagrangian have dimension 4, and all the coupling constants are dimensionless. So, the SM is the renormalizable theory in the same manner as QED.

The perturbation theory expansion EW parameters α/π with $\alpha_{em} \sim 1/129$ and $\alpha_{weak} \sim 1/30$ are very small

Naively - the EW higher order corrections are not that important

However, the experimental accuracies are in some cases so high, that even 1-loop EW corrections might not be sufficient

M_Z	=	91.1875	\pm	0.0021	GeV	0.002%	
Γ_Z	=	2.4952	\pm	0.0023	GeV	0.09%	$\rho_0 = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W}$
M_W	=	80.385	\pm	0.015	GeV	0.02%	
M_{top}	=	173.2	\pm	0.9	GeV	0.52%	$\rho_l = 1.0050 \pm 0.0010$
							$\sin^2 \theta_{eff}^{lept} = 0.23153 \pm 0.00016$

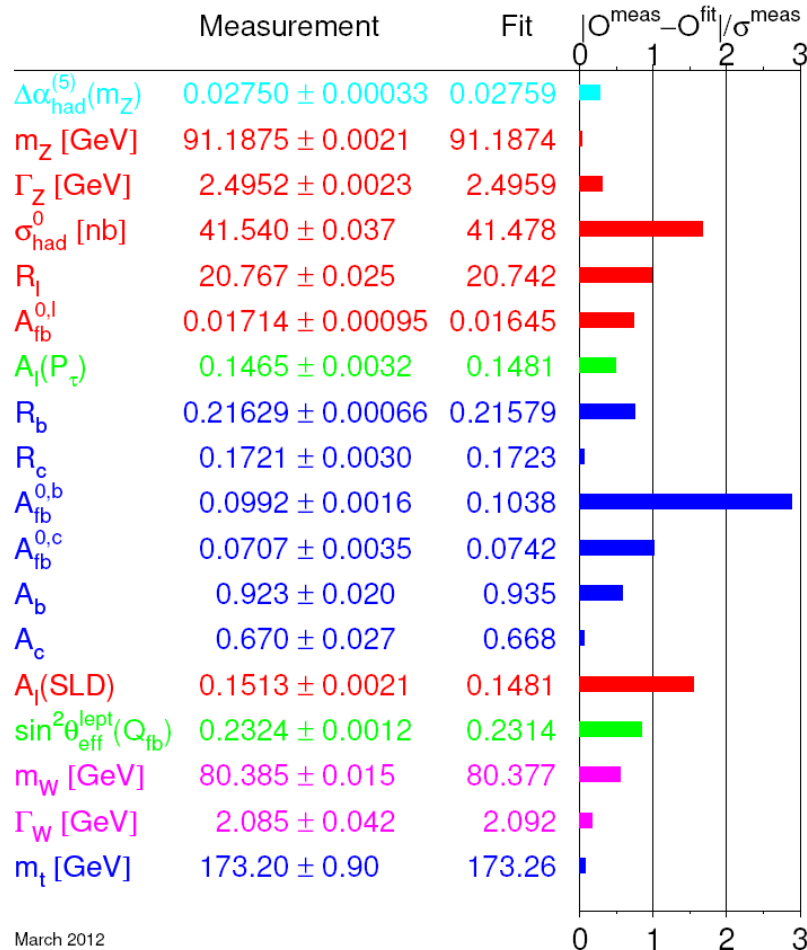
Most important corrections:

Resummation of large logs - $\log (M_{top}^2 / m_e^2) \approx 24.2$;

Corrections proportional to M_{top}^2 / M_W^2 coming from longitudinal modes

Loop corrections lead to the fact that SM parameters (coupling constants, masses, widths) are the running parameters, and **they are nontrivial functions of each other.**

Summary of comparisons of the EW precision measurements at LEP1, LEP2, SLD, and the Tevatron and a global parameter fit



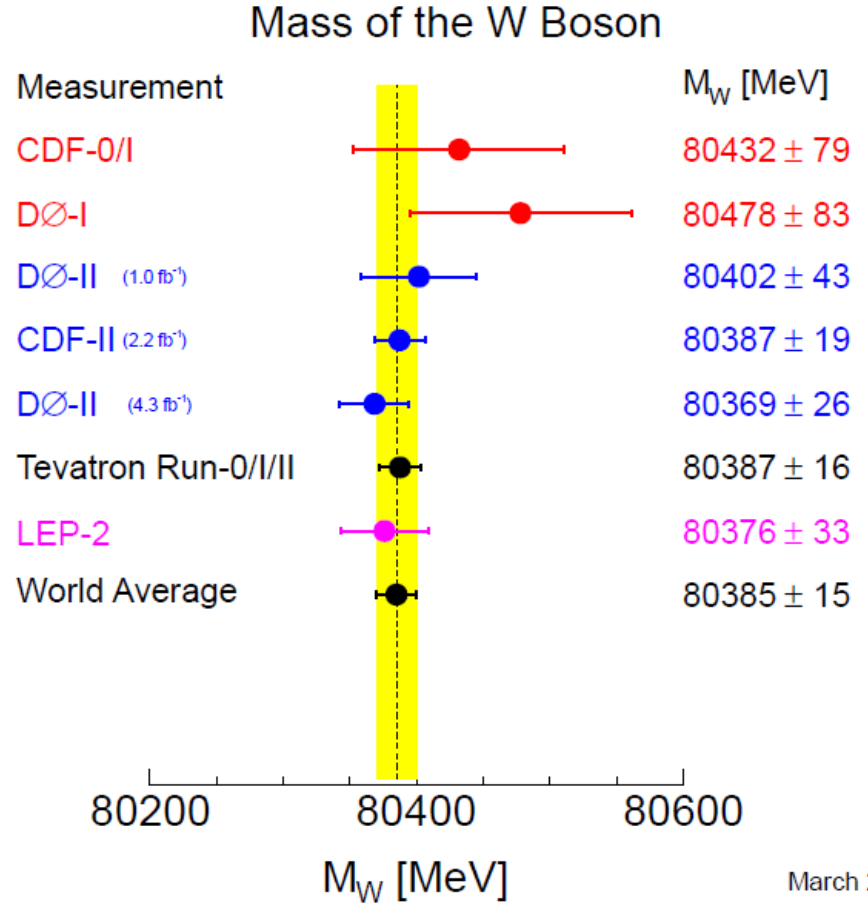
March 2012

$$\sin^2\theta_{\text{eff}}^{\text{lept}} \equiv \frac{1}{4} \left(1 - \frac{v_l}{a_l} \right)$$

CDF ($\int L dt = 2.2 \text{ fb}^{-1}$)
 Electron and Muon
 $M_W = 80387 \pm 19 \text{ MeV}$

Dzero ($\int L dt = 5.2 \text{ fb}^{-1}$)
 Electron only
 $M_W = 80369 \pm 26 \text{ MeV}$

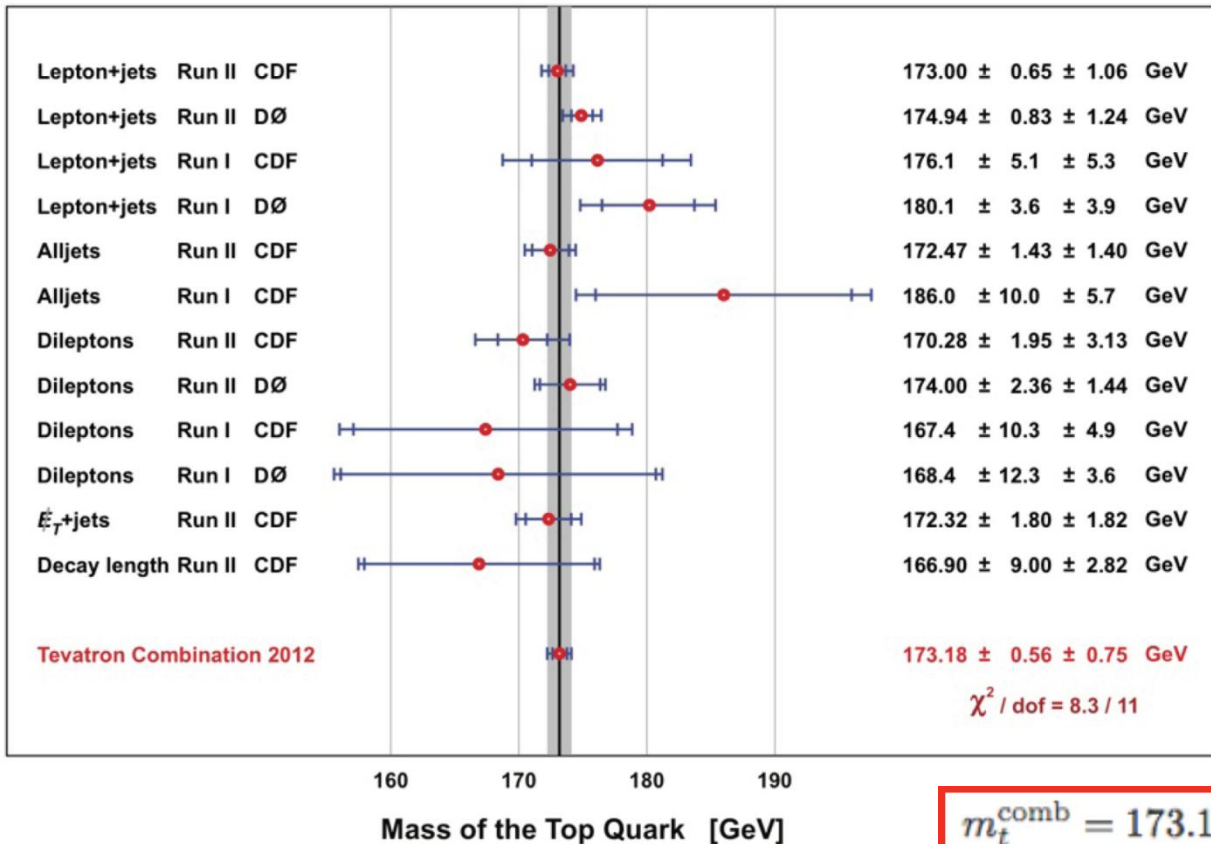
difficult analysis
 Calibration / alignment
 Understanding of recoil



March 2012

Combination : $M_W = 80385 \pm 15 \text{ MeV}$
 0.02%

Top quark mass measurements



Most precisely measured quark mass !

$$m_t^{\text{comb}} = 173.18 \pm 0.56 (\text{stat}) \pm 0.75 (\text{syst}) \text{ GeV}$$

$$= 173.18 \pm 0.94 \text{ GeV}$$

$$\text{LHC: } m_t = 173.3 \pm 0.5 (\text{stat.}) \pm 1.3 (\text{sist.}) \text{ GeV}$$

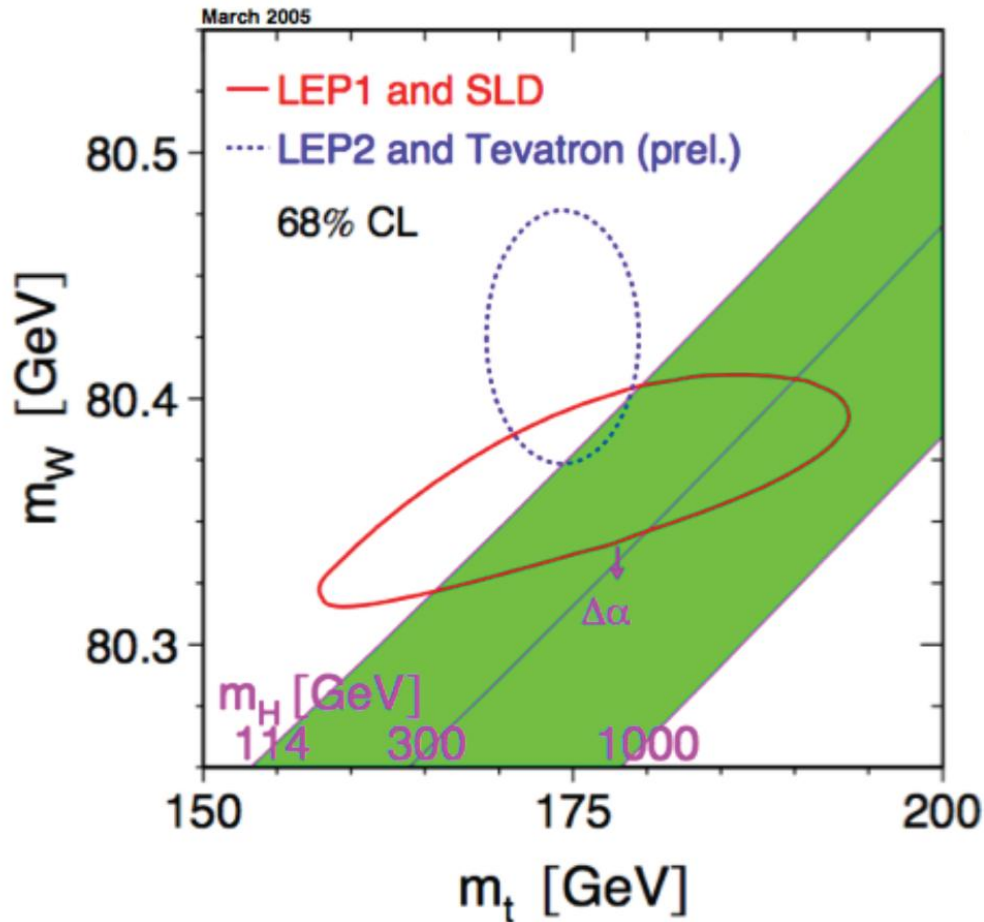
Measurement of top and anti-top mass difference - check of CPT theorem

$$\text{DØ : } = 0.84 \pm 1.87 \text{ GeV}$$

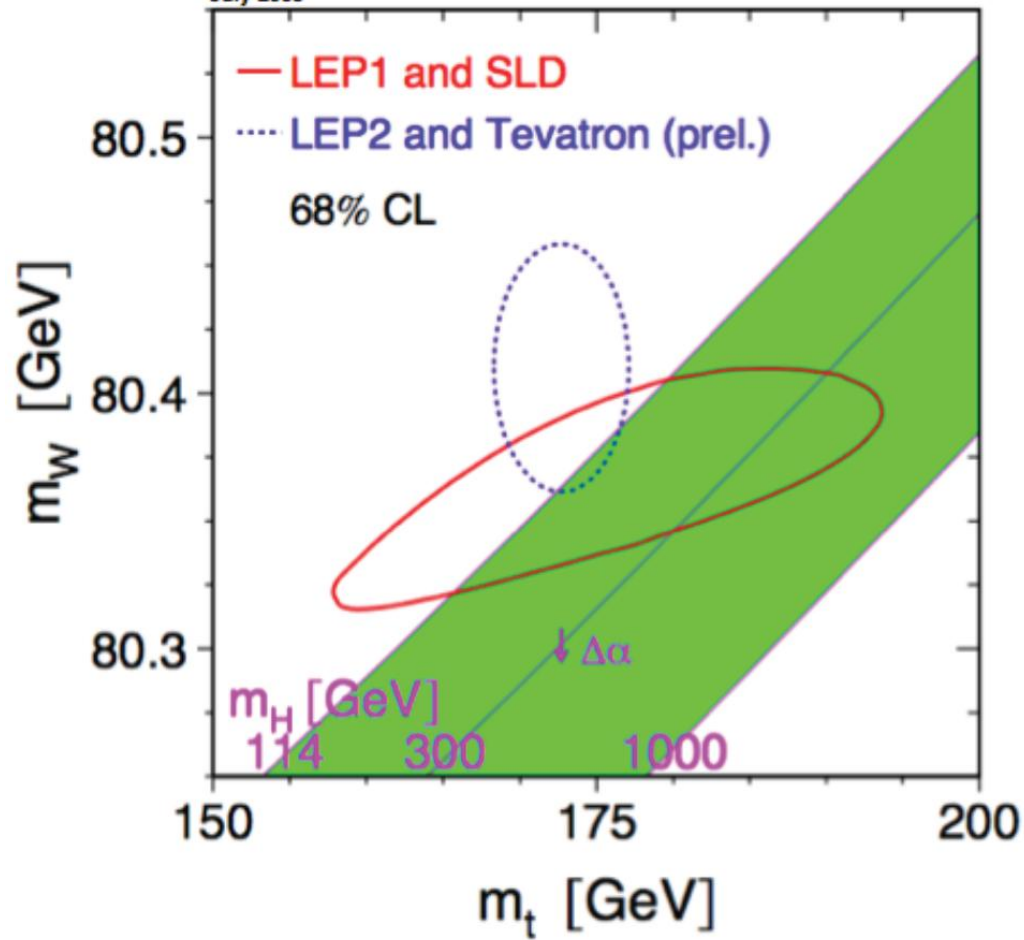
$$\text{CDF : } = -3.3 \pm 1.7 \text{ GeV}$$

$$\text{For comparison CMS : } = -0.44 \pm 0.53 \text{ GeV}/c^2$$

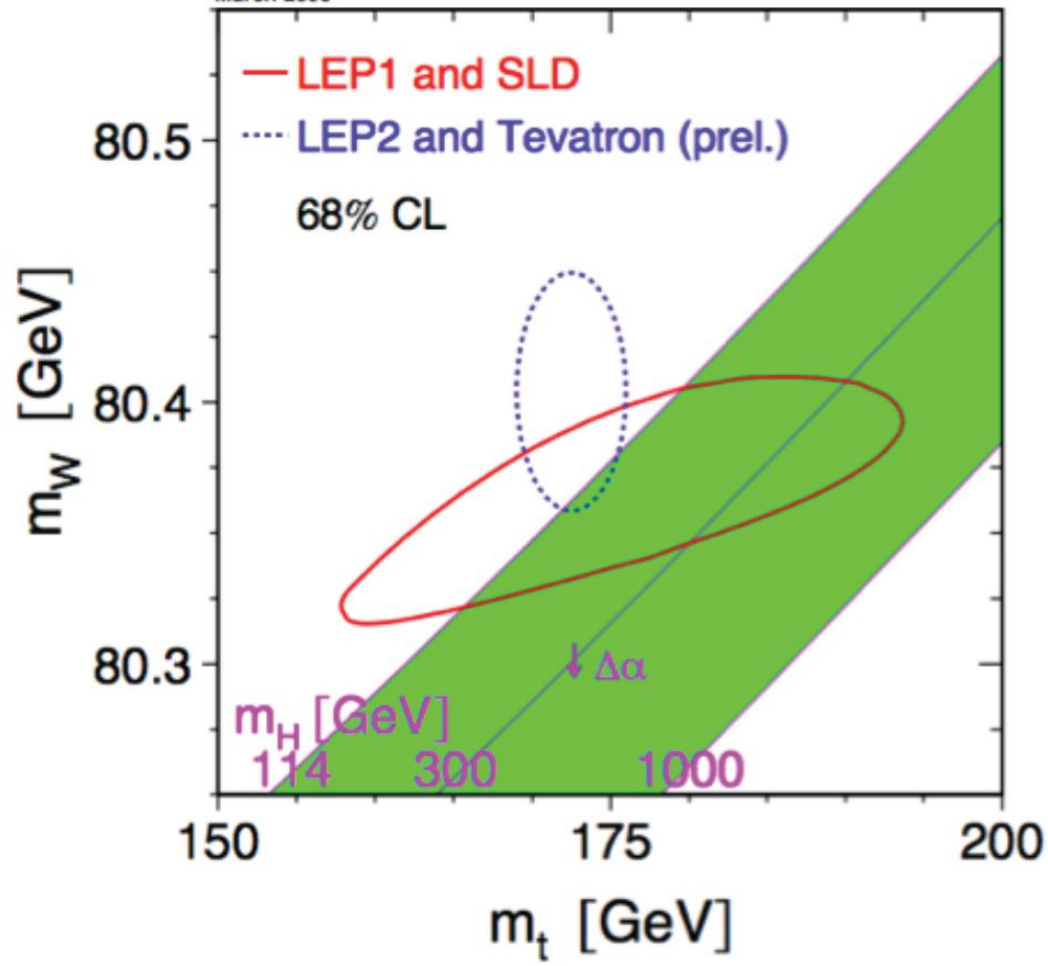
M_W is a function of M_{top} and M_H
 in SM as the quantum field theory (history)



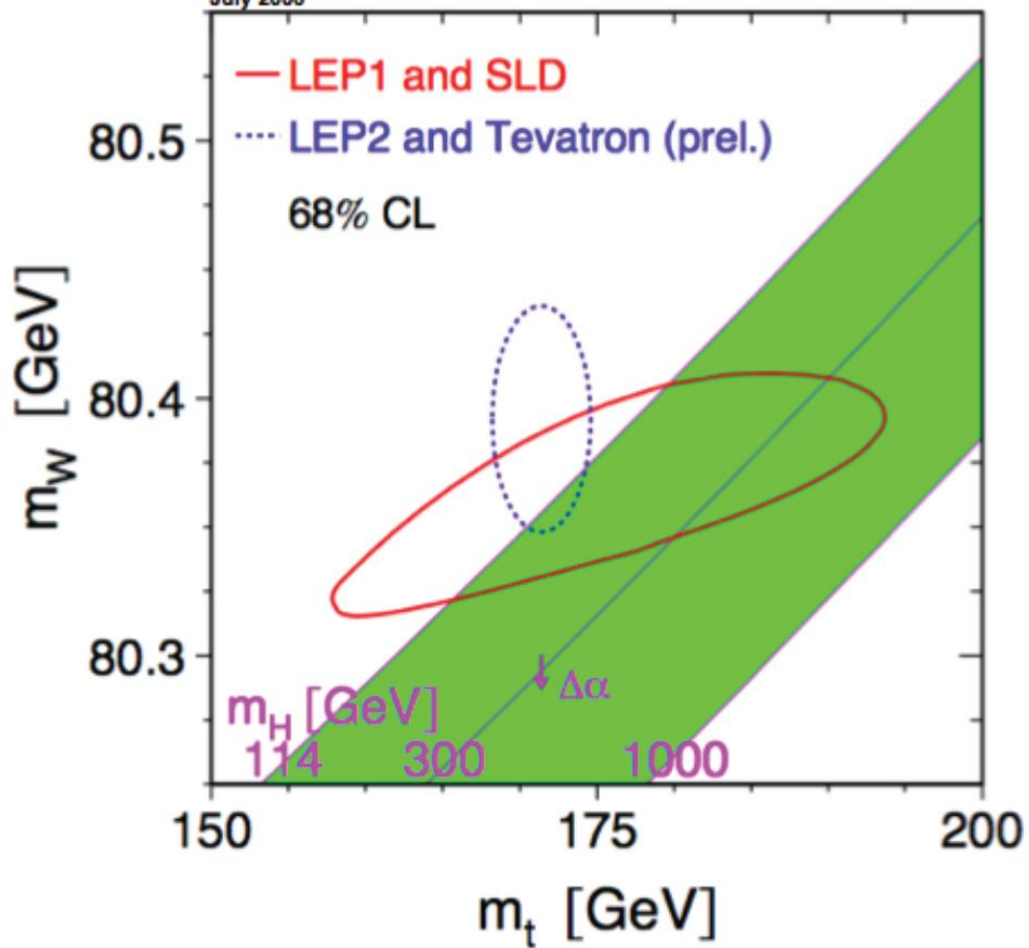
July 2005



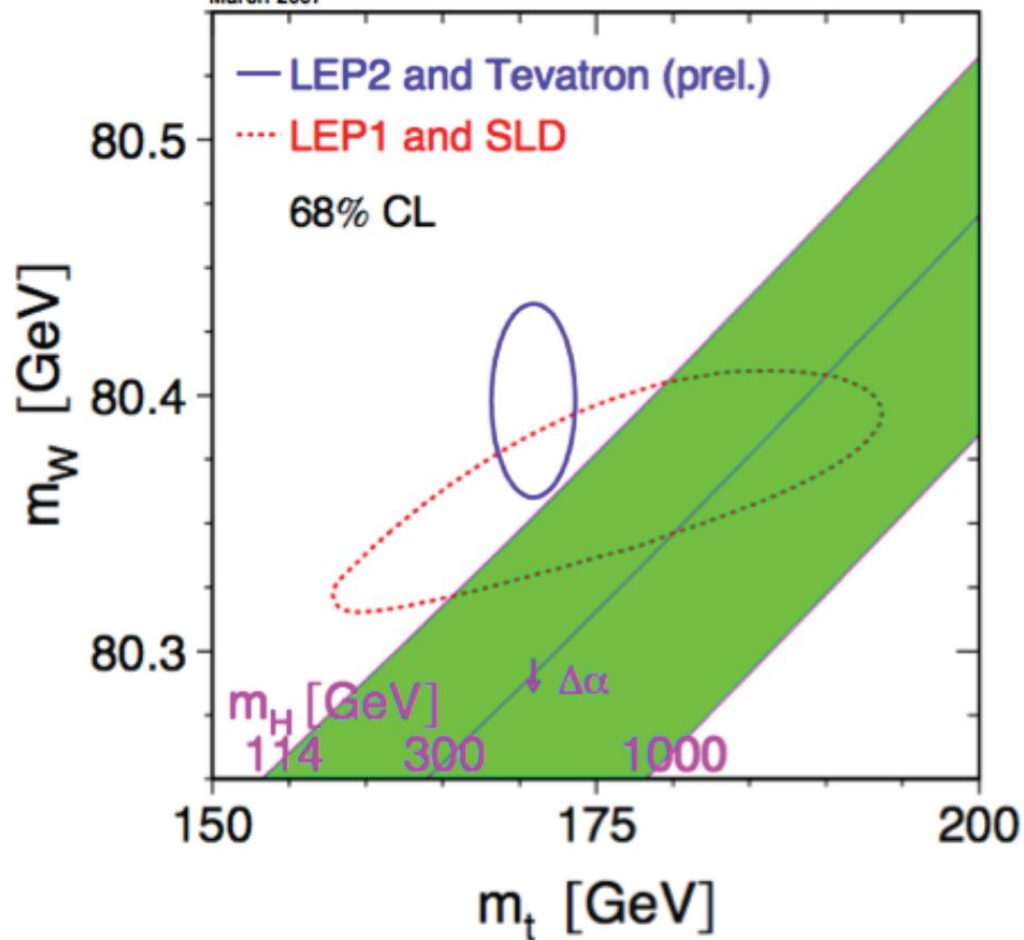
March 2006

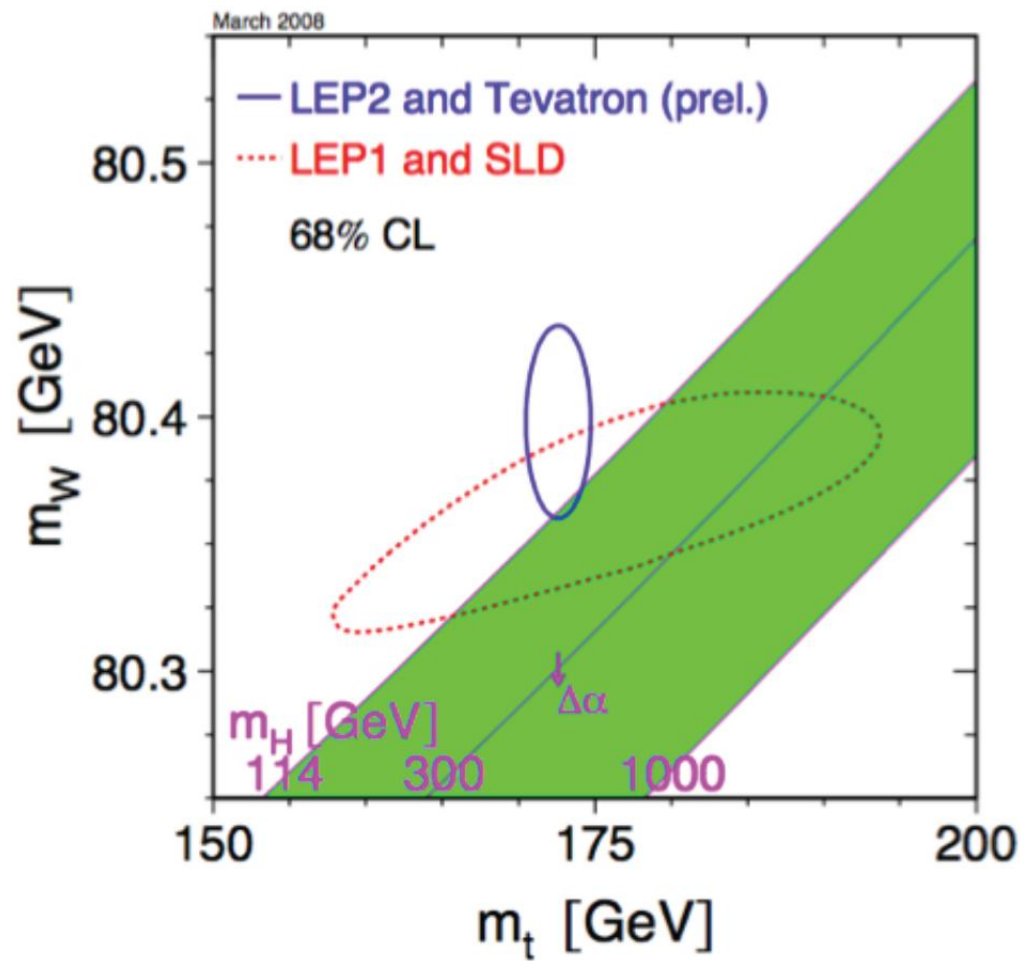


July 2006

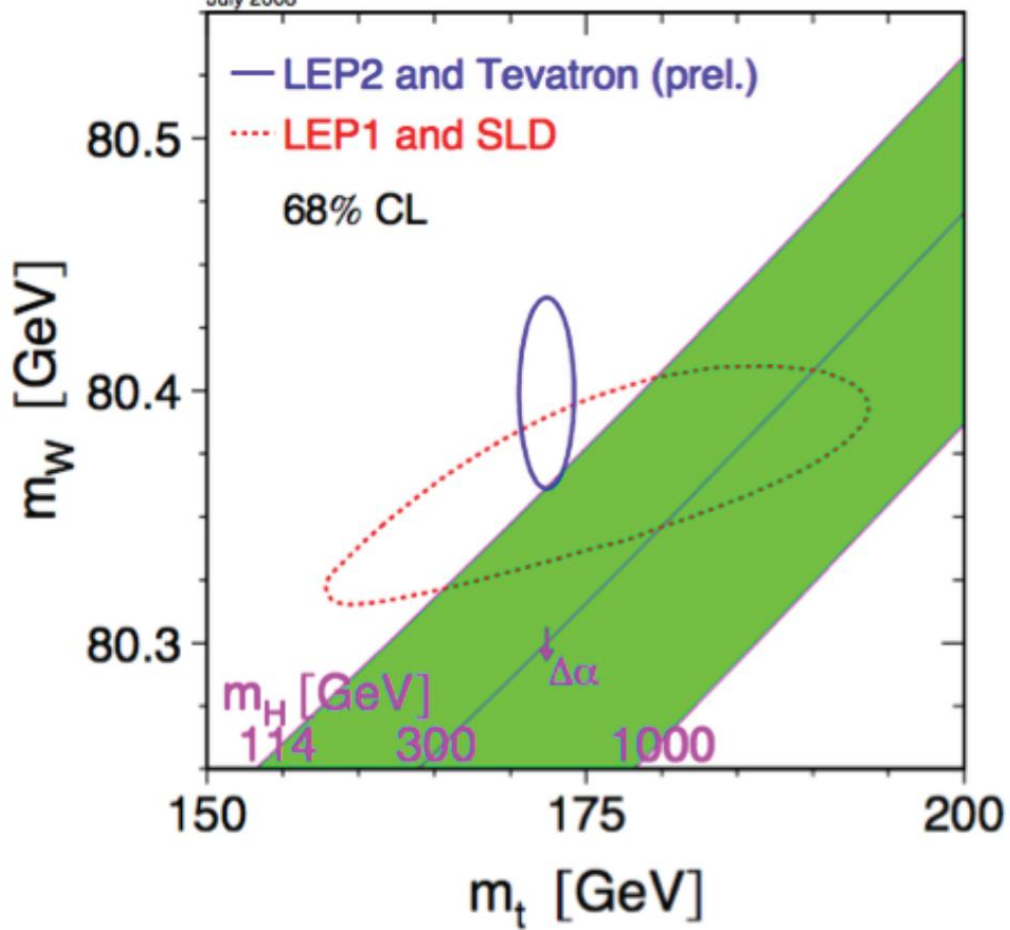


March 2007

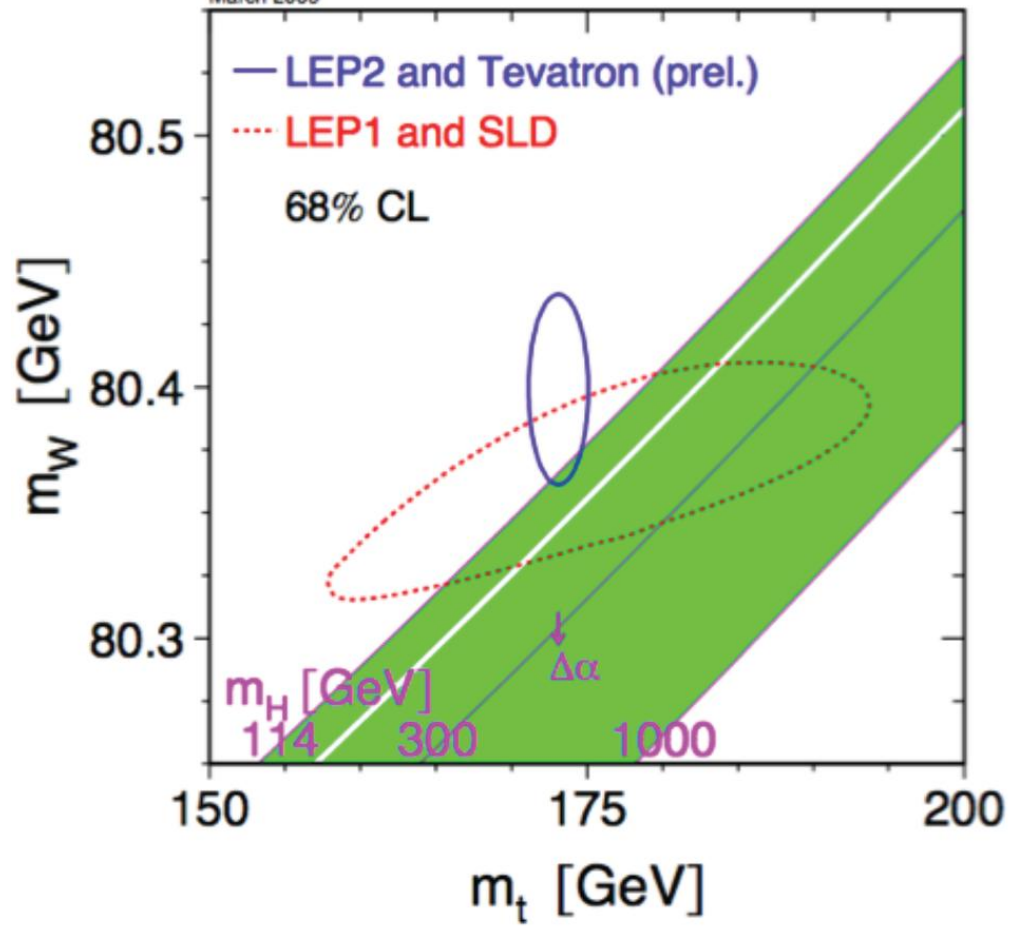


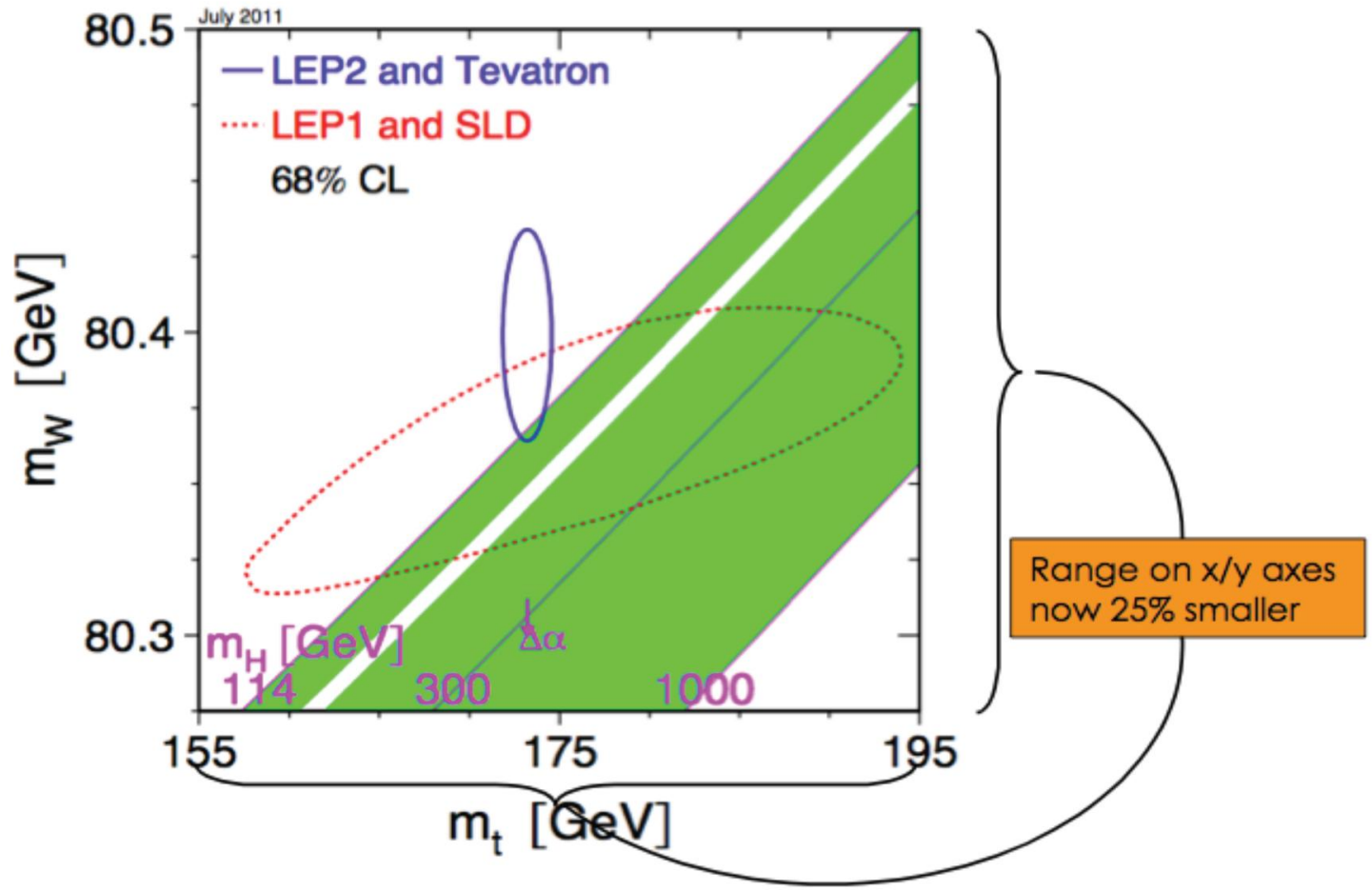


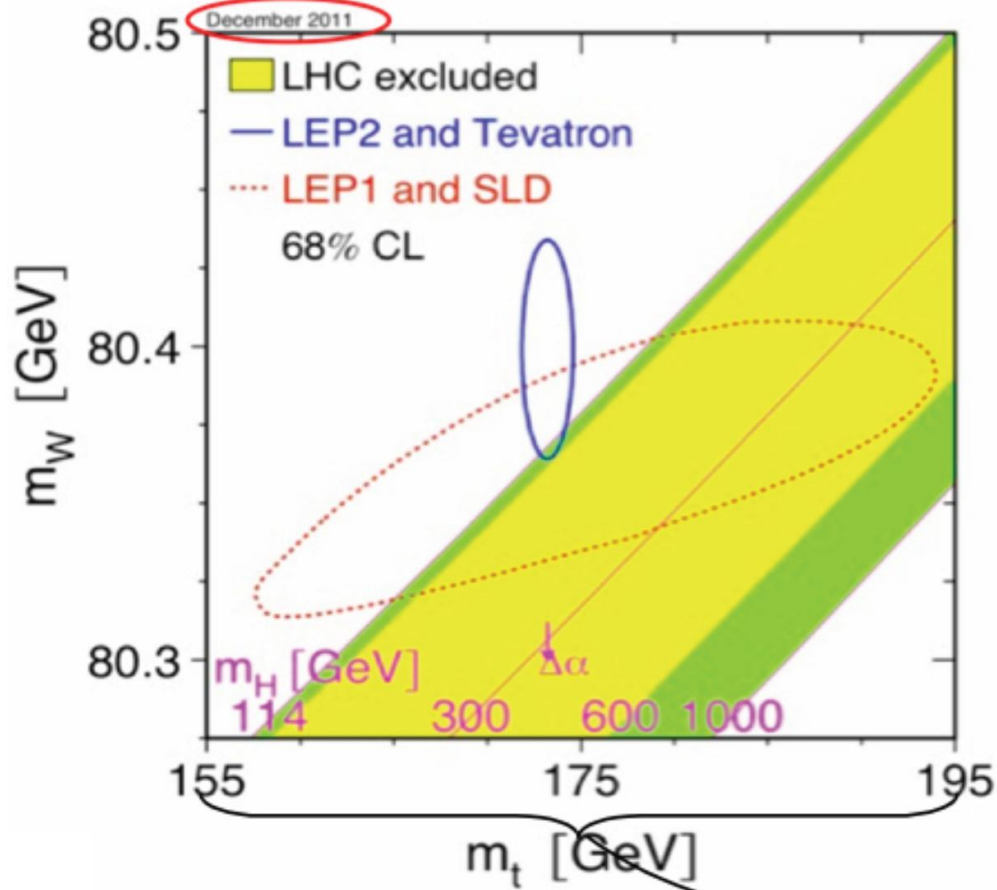
July 2008



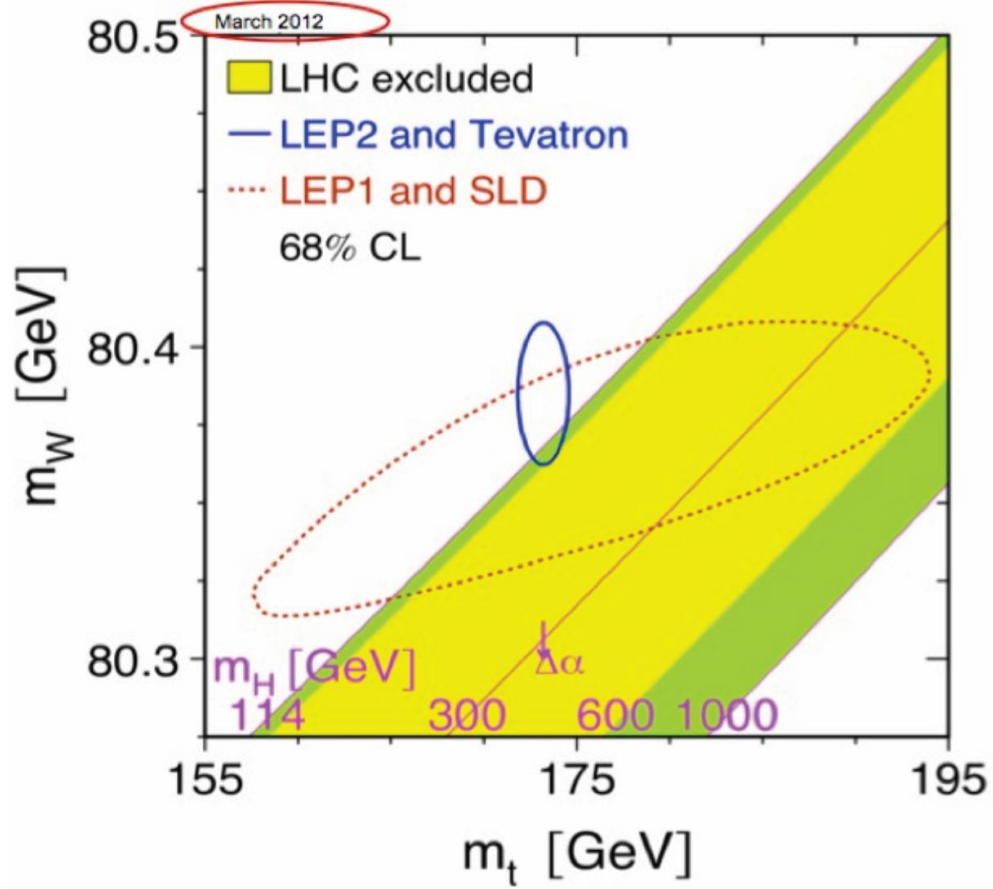
March 2009



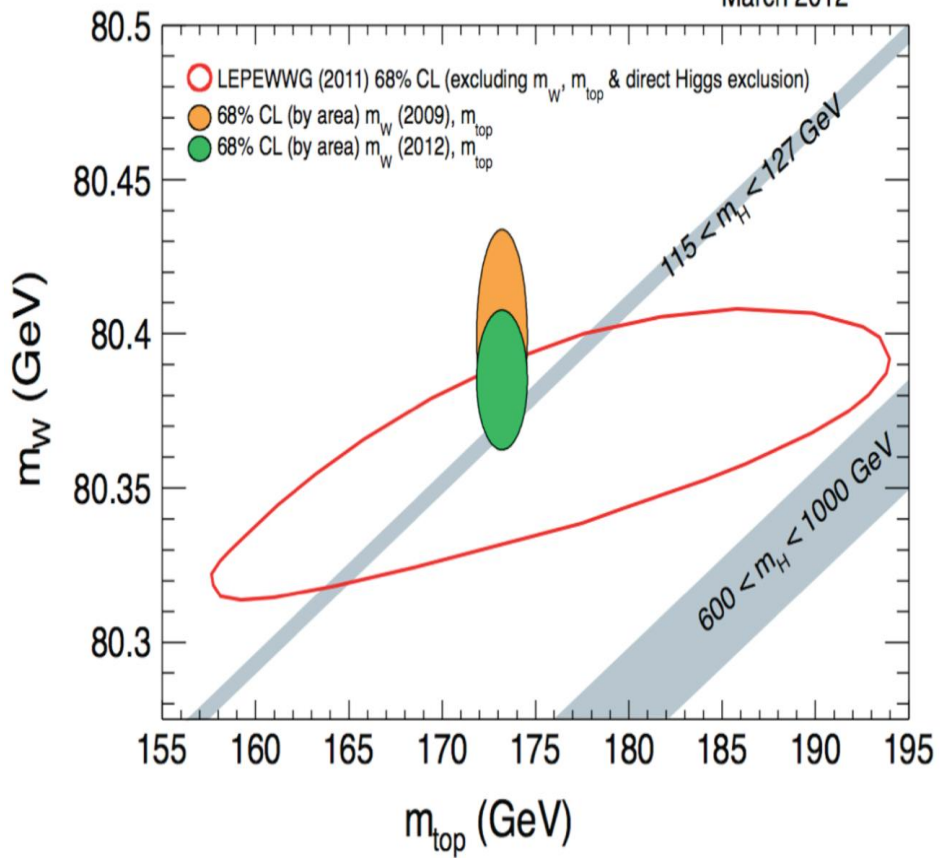




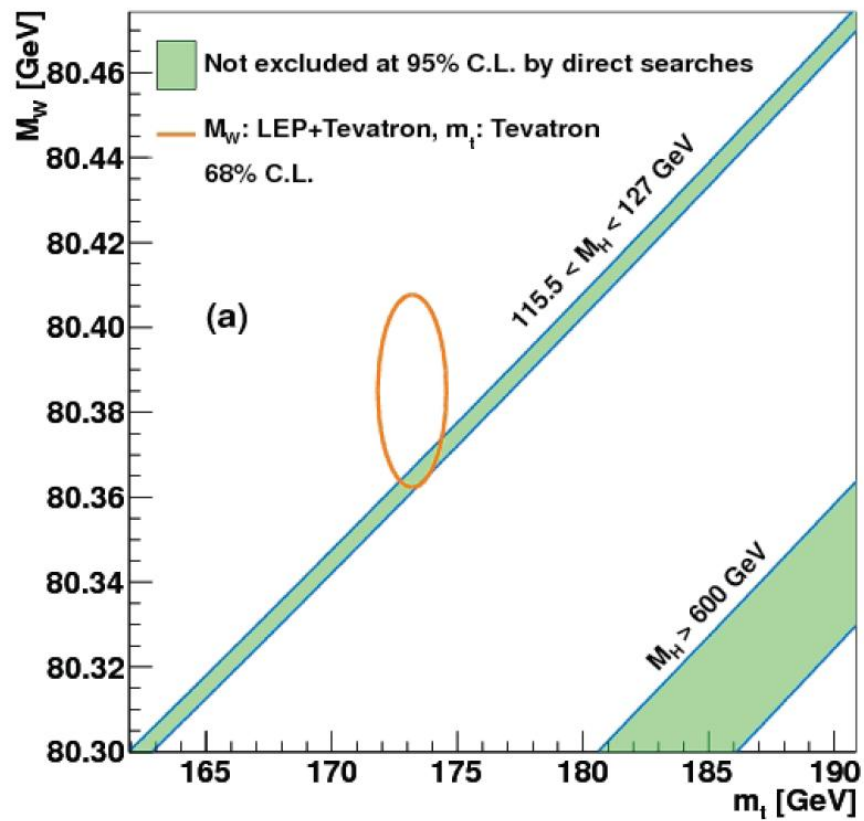
Range on x/y axes
now 25% smaller



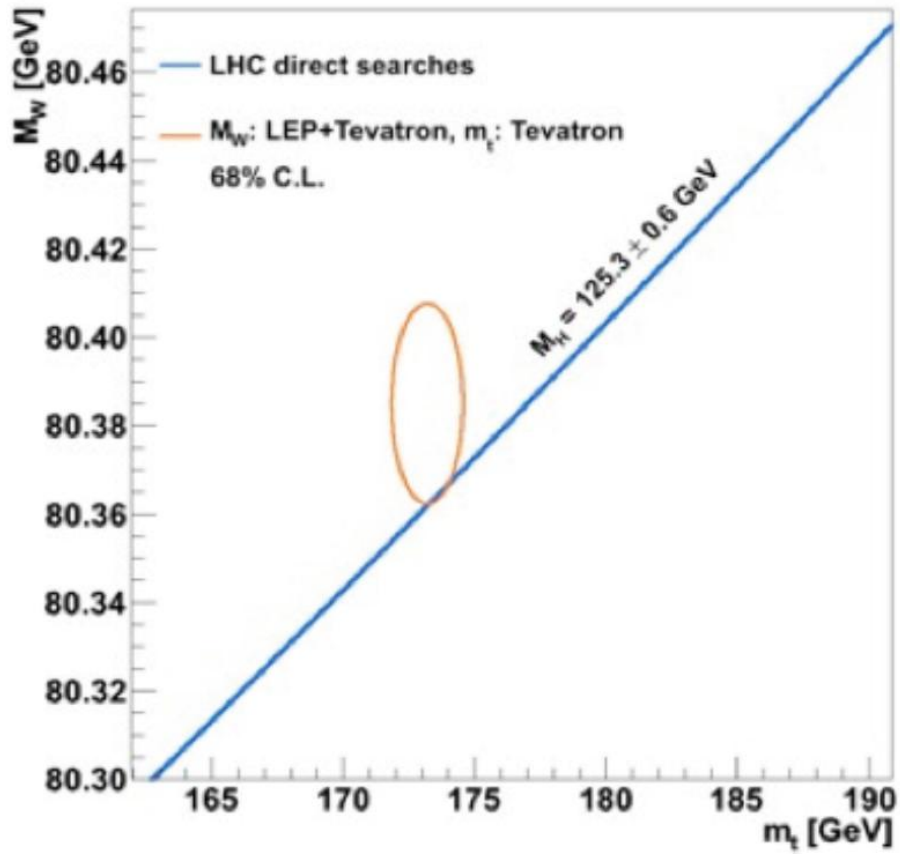
March 2012



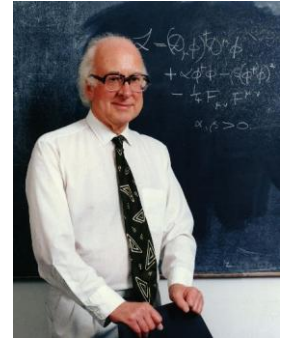
April 2012



July 2012



Higgs Boson



Masses of quarks and leptons (except neutrinos)

Masses of W and Z bosons

Unitarity and renormalizability of the SM

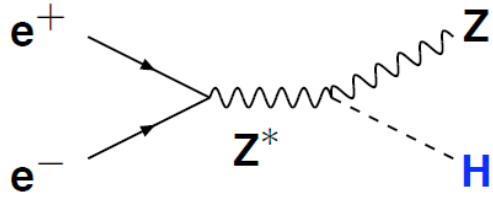


**Brout-Englert-Higgs
-Hagen-Guralnik-Kibble mechanism**

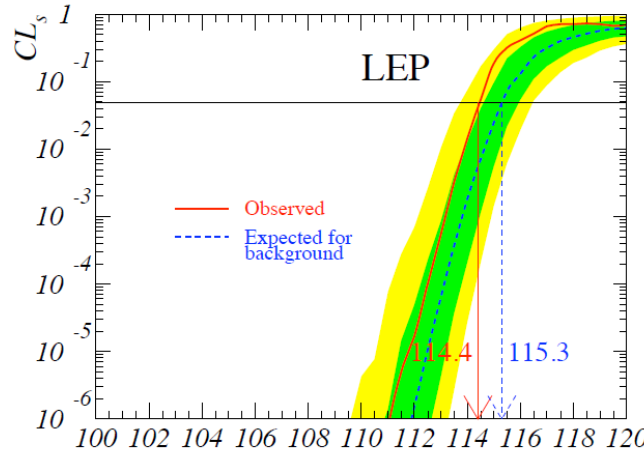
- [4] F. Englert, R. Brout, Phys. Rev. Lett. 13 (1964) 321, [doi:10.1103/PhysRevLett.13.321](https://doi.org/10.1103/PhysRevLett.13.321).
- [5] P.W. Higgs, Phys. Lett. 12 (1964) 132, [doi:10.1016/0031-9163\(64\)91136-9](https://doi.org/10.1016/0031-9163(64)91136-9).
- [6] P.W. Higgs, Phys. Rev. Lett. 13 (1964) 508, [doi:10.1103/PhysRevLett.13.508](https://doi.org/10.1103/PhysRevLett.13.508).
- [7] G. Guralnik, C. Hagen, T.W.B. Kibble, Phys. Rev. Lett. 13 (1964) 585, [doi:10.1103/PhysRevLett.13.585](https://doi.org/10.1103/PhysRevLett.13.585).
- [8] P.W. Higgs, Phys. Rev. 145 (1966) 1156, [doi:10.1103/PhysRev.145.1156](https://doi.org/10.1103/PhysRev.145.1156).
- [9] T.W.B. Kibble, Phys. Rev. 155 (1967) 1554, [doi:10.1103/PhysRev.155.1554](https://doi.org/10.1103/PhysRev.155.1554).

Что нам было известно о бозоне Хиггса до БАК?

1. Ограничения из прямых поисков:



$M_H > 114.4 \text{ GeV}$ 95% C.L.

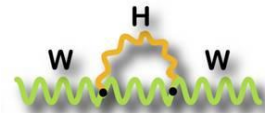
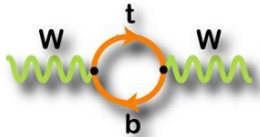


Tevatron (in gluon fusion with decay to WW):

Excluded region

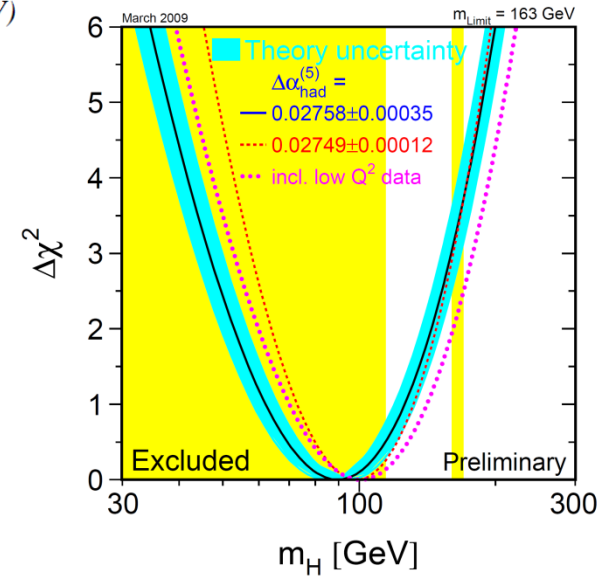
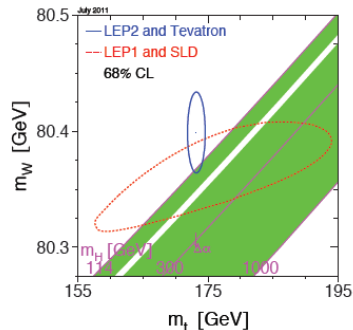
$M_H : 160-170 \text{ GeV}$

2. Ограничения из прецизионных измерений: $M_H(\text{GeV})$



$$(\Delta r)_{\text{top}} \approx -\frac{3G_F m_t^2}{8\sqrt{2}\pi^2} \frac{1}{t_W^2} \quad (\Delta r)_{\text{Higgs}} \approx \frac{11G_F M_Z^2 c_W^2}{24\sqrt{2}\pi^2} \ln \frac{m_h^2}{M_Z^2}$$

$M_H < 155 \text{ GeV}$ 95% C.L.



3. From the unitarity of $VV \rightarrow VV$ ($V: W, Z$) amplitudes:

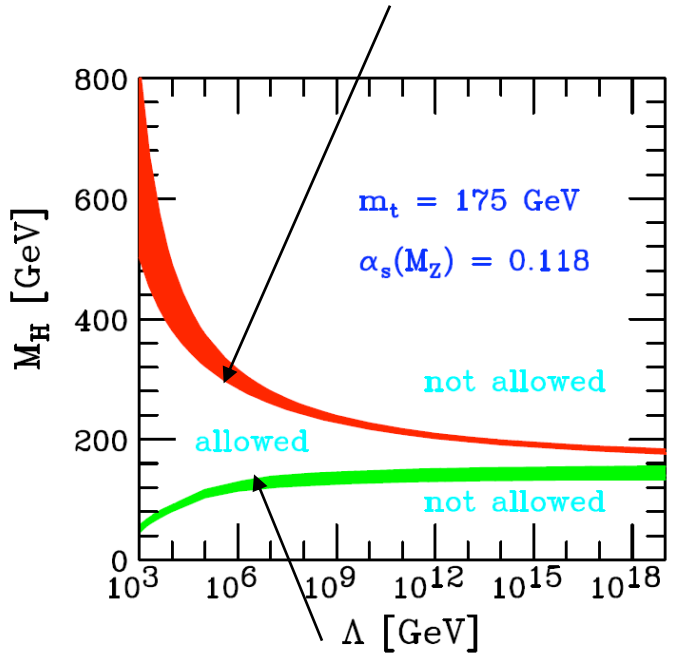
$$\text{Im}(a_1) = |a_1|^2 \quad |\text{Re}(a_1)| \leq \frac{1}{2}$$

$$M_H \lesssim 710 \text{ GeV} \quad \text{if } \sqrt{s} \gg M_H$$

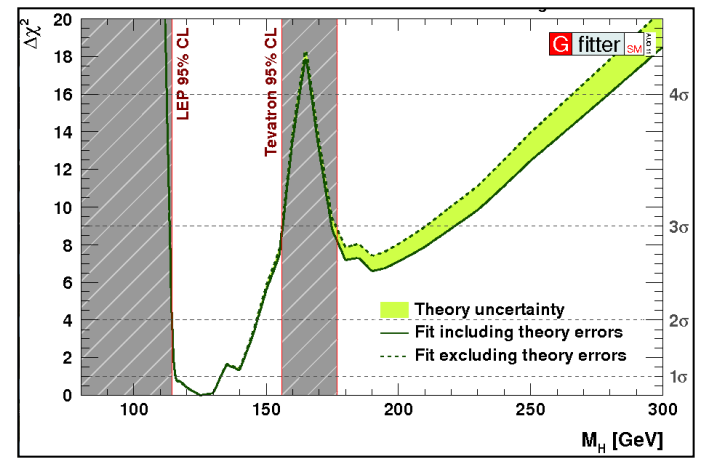
$$\sqrt{s} \lesssim 1.2 \text{ TeV} \quad \text{if } \sqrt{s} \ll M_H$$

4. From self-consistency of quantum theory:

No Landau pole (triviality)



Combining all direct and indirect constraints:



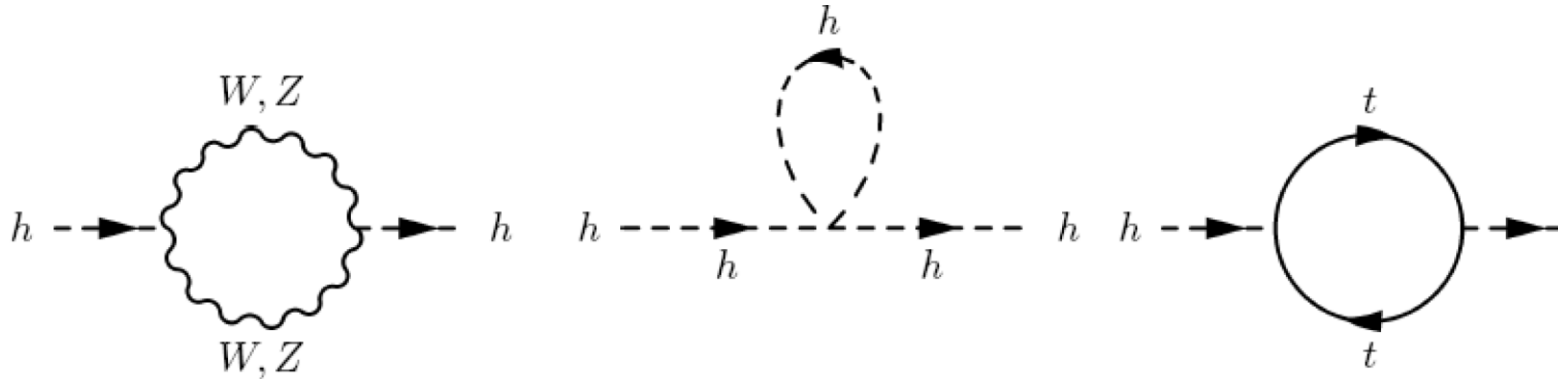
Gfitter collaboration, Aug. 2011

$$M_H = 125 \pm 10 \text{ GeV}$$

Positive self coupling $\lambda(Q^2) > 0$ (vacuum stability)

The simplest Higgs mechanism SM is not stable with respect to quantum corrections (naturalness problem)

Loop corrections to the Higgs mass



$$\delta m_H^2 = \frac{3G_F}{4\sqrt{2}\pi^2} (2m_W^2 + m_Z^2 + m_H^2 - 4m_t^2) \Lambda^2 \approx -(0.2 \Lambda)^2$$

$$\delta m_H < m_H$$

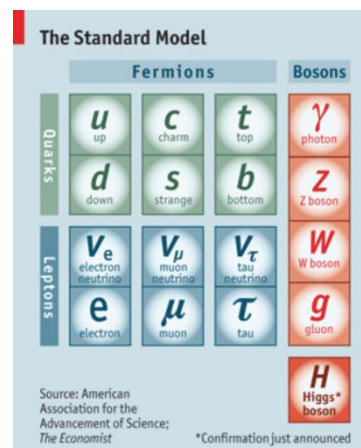
$$\Lambda < 1 \text{ TeV}$$

In SM there is no symmetry which protects a strong dependence of Higgs mass on a possible new scale

Something is needed in addition to SM...

- * Standard Model is the renormalizable anomaly free gauge quantum field theory with spontaneously broken electroweak symmetry. Remarkable agreement with many experimental measurements.
 - * The EW SM has 17 parameters (from experiments)
 - gauge-Higgs sector contains 4 parameters:
 g_1, g_2, μ^2, λ best measured α_{em}, G_F, M_Z (or α_{em}, s_W, M_W) plus M_H
-
- In addition, 6 quarks masses, 3 lepton masses, 3 mixing angles and one phase of the CKM matrix
- plus α_{QCD} **18 SM parameters**
- (+ may be masses and mixing parameters from neutrino sector)

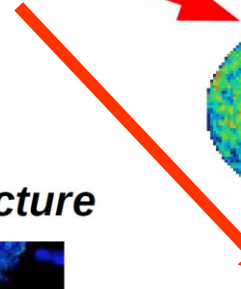
		Three generations of matter (fermions)			
		I	II	III	
mass		2.4 MeV/c ²	1.27 GeV/c ²	171.2 GeV/c ²	0
charge		$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0
spin		$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
name		u up	c charm	t top	γ photon
					H Higgs boson
Quarks	mass	4.8 MeV/c ²	104 MeV/c ²	4.2 GeV/c ²	0
	charge	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
		d down	s strange	b bottom	g gluon
Leptons	mass	<2.2 eV/c ²	<0.17 MeV/c ²	<13.5 MeV/c ²	91.2 GeV/c ²
	charge	0	0	0	0
	spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1
		ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	Z ⁰ Z boson
		e electron	μ muon	τ tau	W [±] W boson



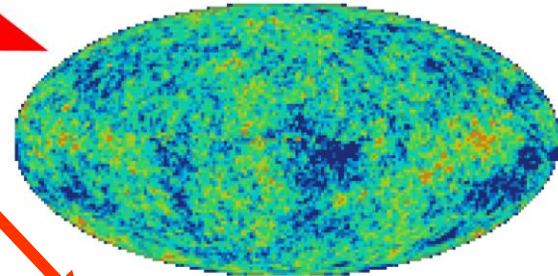
Facts which can not be explained in SM

- EW symmetry is broken - photon is massless, W and Z are massive particles
- Fermions have very much different masses
($M_{top} \approx 172 \text{ GeV}$, $m_e \approx 0.5 \text{ MeV}$, $\Delta m_\nu \approx 10^{-3} \text{ eV}$)

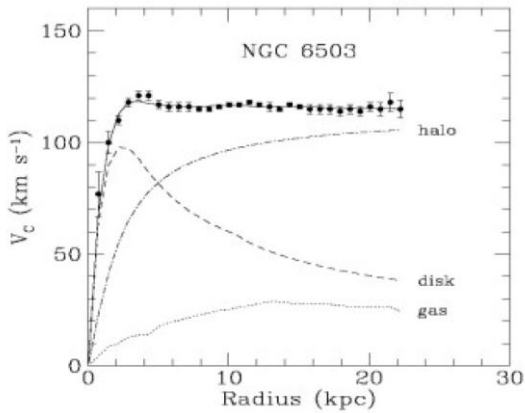
- Dark Matter exists in the Universe



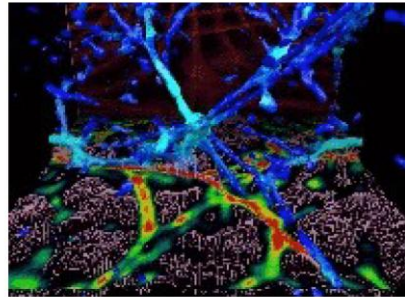
CMB



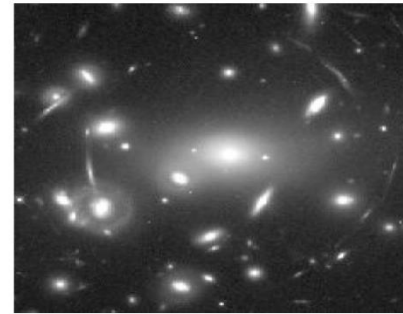
Rotation curves of galaxies



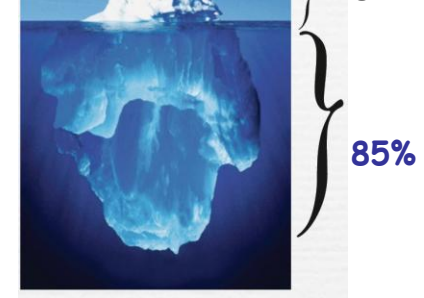
Large Scale Structure



Lensing



Barionic matter
(1% in stars, 14% in gas)



Dark unknown matter

- $(g-2)_\mu$ (about 3.5σ)
- Neutrino masses, mixing, oscillations
- Particle - antiparticle asymmetry in the Universe,

CP violation

baryon asymmetry: $\frac{n_B - n_{\bar{B}}}{n_B + n_{\bar{B}}} \sim 10^{-10}$

CKM phase - too small effect

In addition to mentioned problems (naturalness/hierarchy, dark matter content, CP violation) SM does not give answers to many questions

What is a generation? Why there are only 3 generations?

How quarks and leptons related to each other, what is a nature of quark-lepton analogy?

What is responsible for gauge symmetries, why charges are quantize?
Are there additional gauge symmetries?

What is responsible for a formation of the Higgs potential?

To which accuracy the CPT symmetry is exact?

Why gravity is so weak comparing to other interactions?

.....

LHC - Why Terascale?

Stabilization of the Higgs mechanism

→ $\Lambda \sim 1 \text{ TeV}$

Unitarization of EW vector boson and heavy quark amplitudes

→ $\Lambda \sim 1 \text{ TeV}$

If $M_h \sim 1 \text{ TeV} \rightarrow$ SM Higgs width $\sim 0.5 \text{ TeV}$, strong coupling regime

Dark Matter density: in most popular scenarios masses of DM candidates are less than 1 TeV

Main options beyond SM



1. Fundamental Higgs:

- Supersymmetric models
(MSSM, NMSSM...)

2. Composite Higgs:

- Models with new strong dynamics
(Chiral Lagrangians from holography, latest technicolor variants, Little Higgs...)

3. Mixed cases:

- Models with extra space dimensions
- Partially composite models...

4. Many more (hidden valleys, landscape)

Supersymmetry is one of the most favorite BSM ideas, relating spin $\frac{1}{2}$ fermions with spin 0,1 bosons

$$Q|\text{Boson}\rangle = |\text{Fermion}\rangle \quad Q^\dagger|\text{Boson}\rangle = |\text{Fermion}\rangle$$

Fermion degrees of freedom \leftrightarrow boson degrees of freedom

Minimal particle content

□ Gauge / Gaugino Sector

Standard Bosons	Supersymmetric Partners
$W^\pm \quad H^\pm$	Charginos $\chi_1^\pm \quad \chi_2^\pm$
$g \quad Z$ $h \quad H \quad A$	Neutralinos $\chi_1^0 \quad \chi_2^0 \quad \chi_3^0 \quad \chi_4^0$
g_i	Gluinos \tilde{g}_i

[Two Higgs doublets]

[All fermions]

And also ...

Graviton G	Gravitino \tilde{G}
--------------	-----------------------

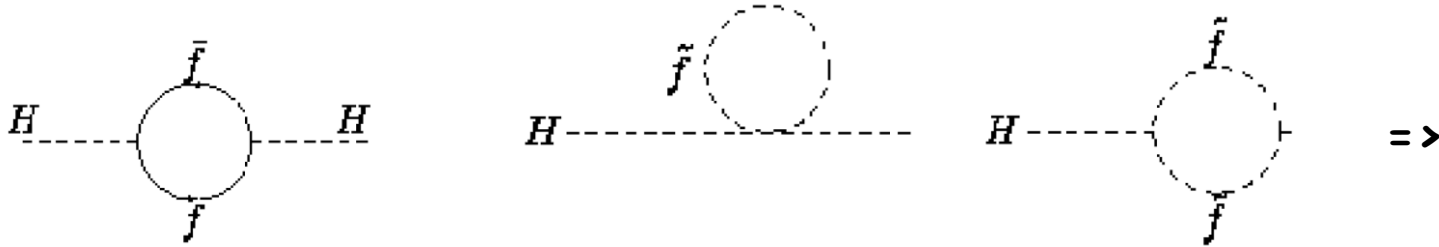
□ Particle / Sparticle Sector

Standard Particles	Supersymmetric Partners
Leptons ℓ	Sleptons $\tilde{\ell}_{R,L}$
Neutrinos ν_ℓ	Sneutrinos $\tilde{\nu}_\ell$
Quarks q	Squarks $\tilde{q}_{R,L}$

[All scalars]

SUSY

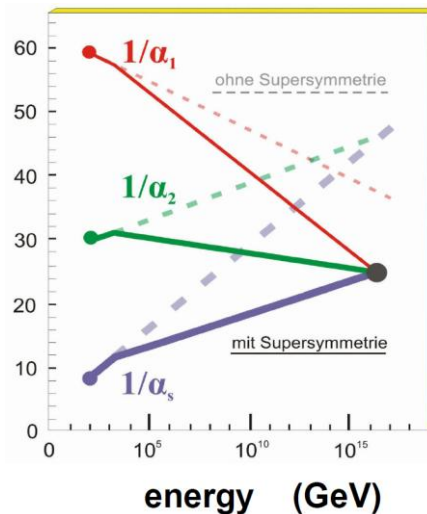
1. Cancellation of Λ^2 dependence



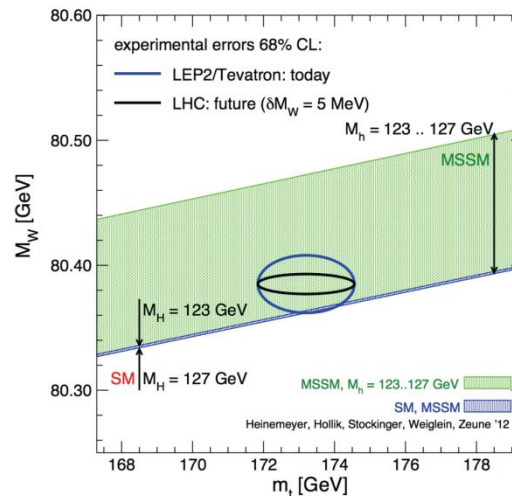
$$\Delta M_H^2|_{\text{tot}} = \frac{\lambda_f^2 N_f}{4\pi^2} \left[(m_f^2 - m_S^2) \log\left(\frac{\Lambda}{m_S}\right) + 3m_f^2 \log\left(\frac{m_S}{m_f}\right) \right] \quad M_H \text{ is protected!}$$

2. Lightest SUSY particle is stable (if R-parity) - very good Dark Matter candidate

3. Unification of couplings in contrast to SM



4. Fit of EW precision data



SUSY is one of the most attractive idea for BSM physics

SUSY, if exists, is broken, and there are many possibilities:

Gravity mediation

Gauge madiation

Gaugino mediation

Anomaly mediation

Hidden sector mediation

...

Many models:

MSSM

CMSSM

mSUGRA

mGMSB

mAMSB

Split SUSY

...

NMSSM

Natural SUSY

...

In general the unconstrained MSSM has 105 parameters
(22 with reasonable assumptions)

(many parameter space points of mSUGRA scenario are rulled out already)

Concrete predictions depend strongly on MSSM breaking scenario.
There are no theory arguments to prefer some of them.

Many nice SUSY feaches are due to additional global symmetry-
R-parity. Tiny deviations of R-parity possible leading to processes
with FCNC, lepton/barion number violation, proton decay...
But what is an origin of R-parity?...

Models with new strong dynamics

Most of composite models are based on symmetry breaking by nontrivial Top condensate

For example (assisted technicolor with top-seesaw):

$$SU(3)_1 \times SU(3)_2 \times U(1)_1 \times U(1)_2 \xrightarrow{\langle \Phi \rangle} SU(3)_{\text{QCD}} \times U(1)_Y$$

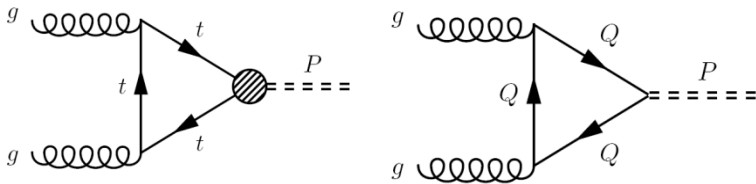
$\langle \Phi \rangle$ is the condensation of $\langle t\bar{t} \rangle = f_\pi$

3d generation quarks and 1st,2d generation quarks are charged under two different SU(3)

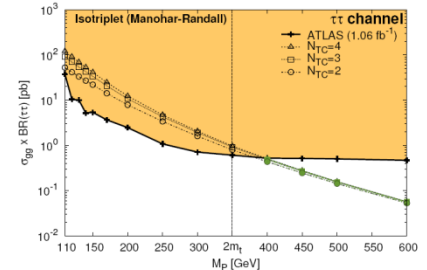
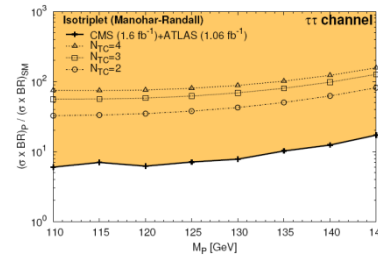
One should avoid FCNC, too large top mass, constrains from s,t,u parameters

In general, there are: techni-pions, technirhos, composite Higgs(es), vector-like top-quark partners

R.Chivikula, P.Ittisamai, E.Simmons



CMS and ATLAS searches for the Higgs in gamma-gamma and tau-tau modes exclude techni-(pseudo)scalars upto $2M_{\text{top}}$



New strong dynamics (Little Higgs, Technicolor like models ...)

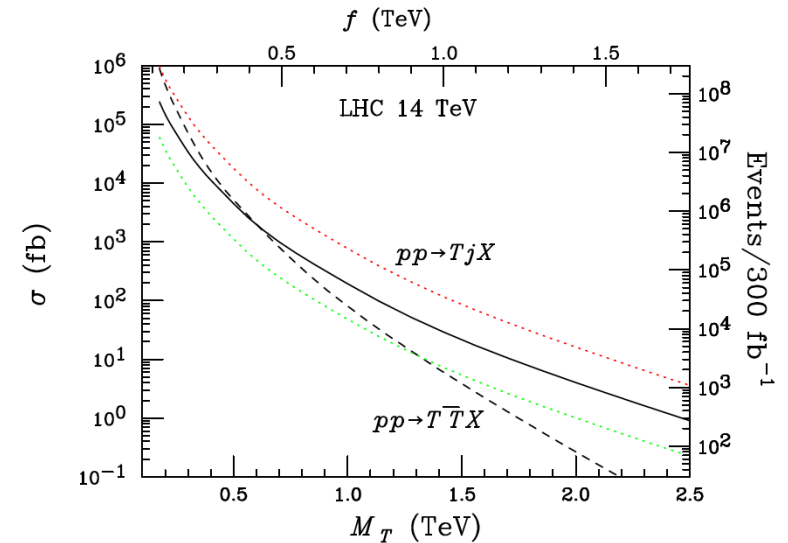
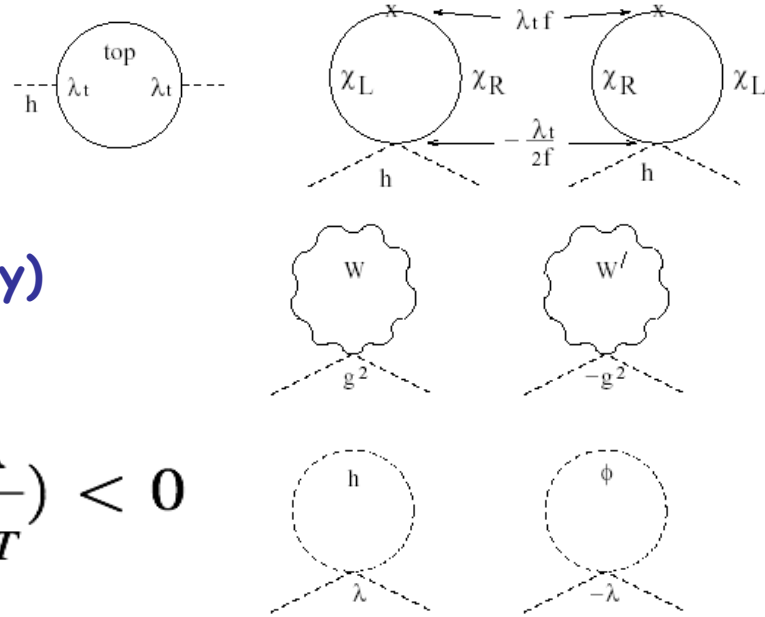
In Little Higgs models
new particle loops cancel
same spin SM particle loops
(cancellation at 1-loop level only)

$$\delta m_H^2 = -\frac{3}{8\pi^2} \lambda_t^2 m_T^2 \ln\left(\frac{\Lambda}{m_T}\right) < 0$$

(similar to SUSY)

If T-parity is assumed there is a DM
candidate

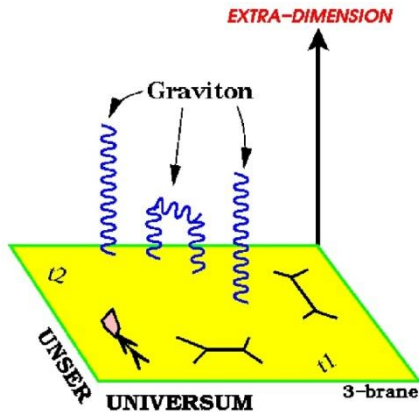
top-quark partner T can be found
at the LHC in few TeV mass range



Models with extra space dimensions

we are confined on some 4-dim. brane imbedded into higher dim. bulk

ADD type models

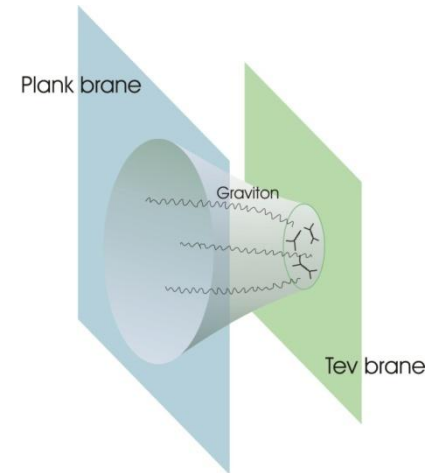


UED type scenarios

with SM fields in ADD or RS bulk

KK-parity ->
LKKP is a good DM candidate

RS type models

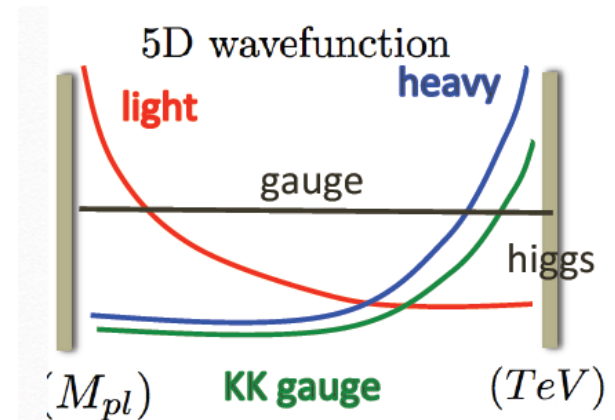
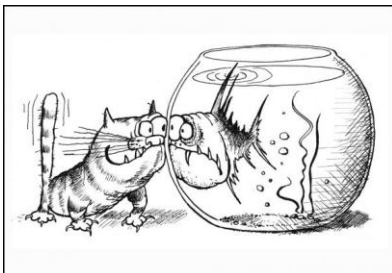


Can unify the forces

Can explain why gravity is weak (solve hierarchy problem)

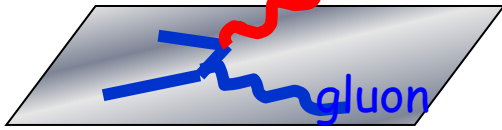
Contain Dark Matter Candidates

Can generate neutrino masses

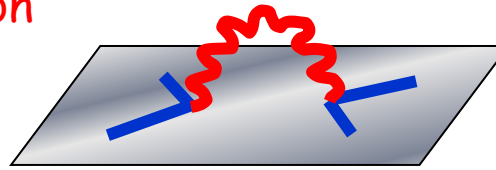


In ADD scenario typical processes:

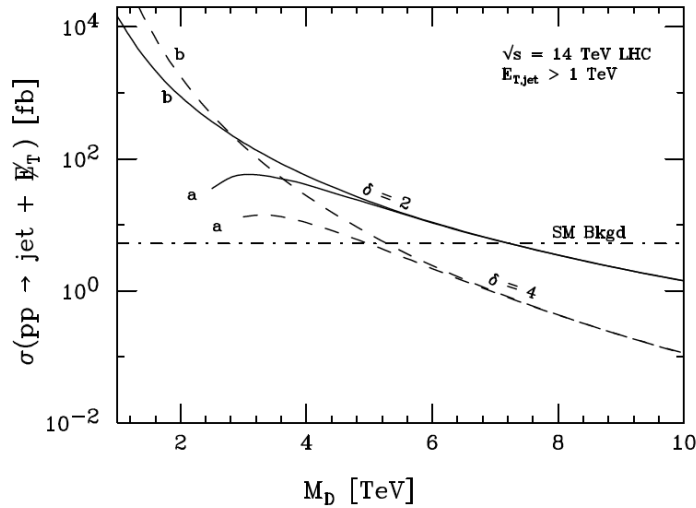
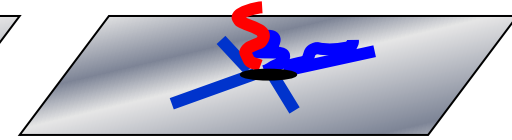
KK Gravitons emission
graviton



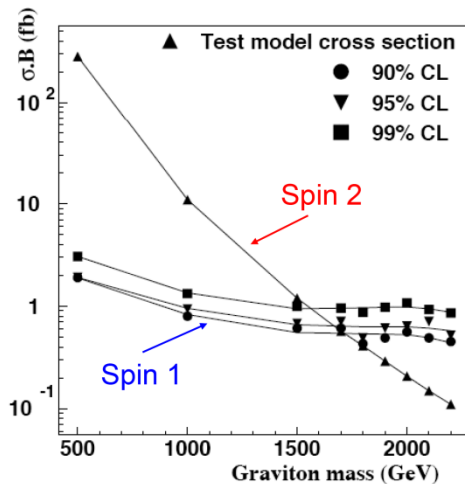
Virtual KK gravitons



TeV black holes



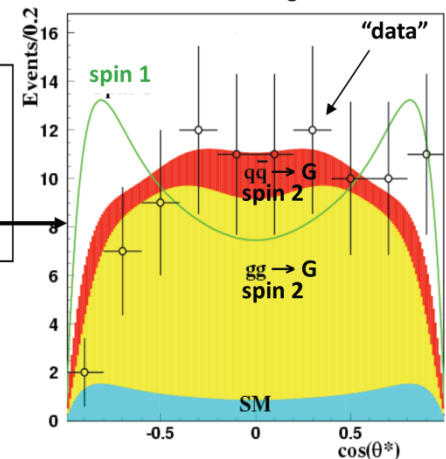
δ	Max M_D sensitivity $\mathcal{L} = 100 \text{ fb}^{-1}$	Max M_D sensitivity $\mathcal{L} = 10 \text{ fb}^{-1}$	Min M_D perturbativity
2	8.5 TeV	7.9 TeV	3.8 TeV
3	6.8	6.3	4.3
4	5.8	5.5	4.8
5	5.0	4.6	5.4



In RS scenario typical processes:
Spin 2 resonances or virtual KK gravitons below thresholds

Graviton ($s=2$) or Z' ($s=1$)?
→ look at e^\pm angular distributions

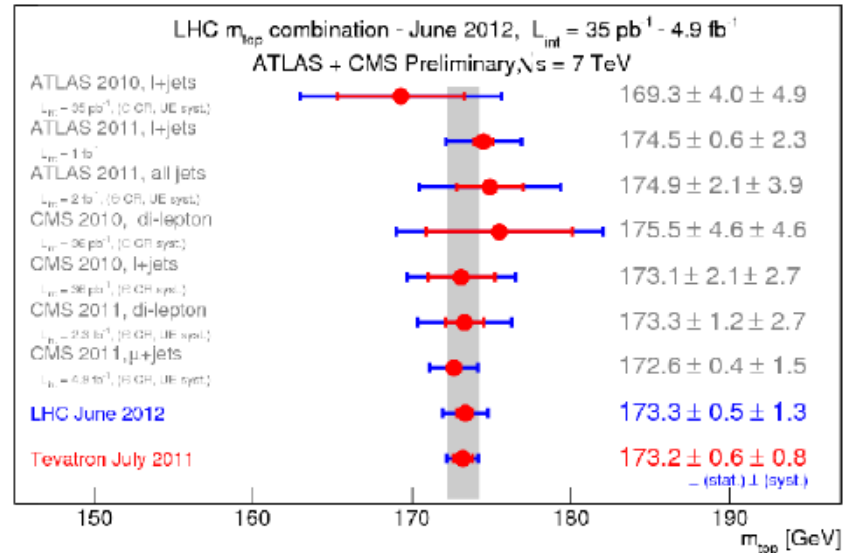
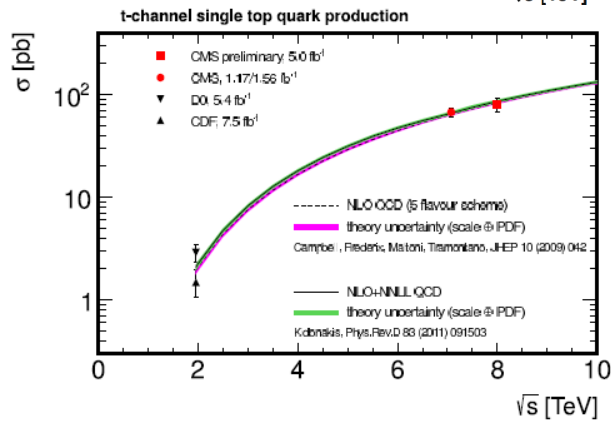
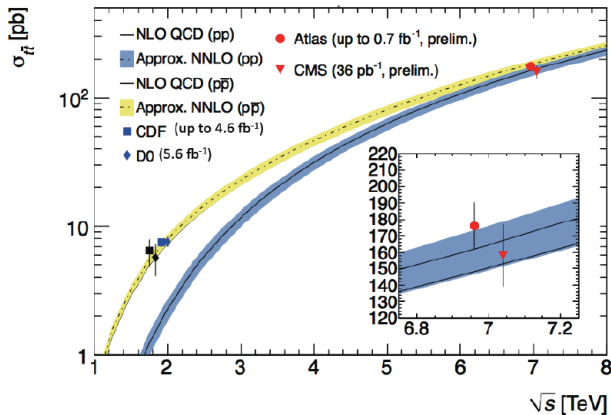
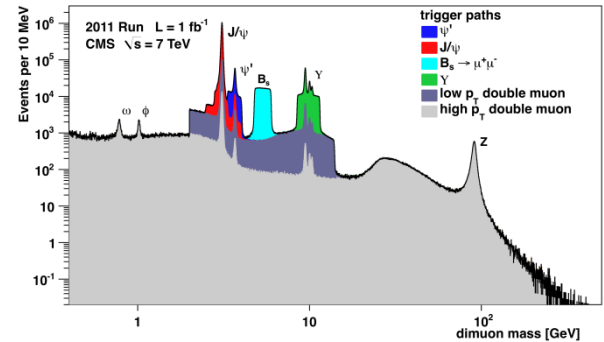
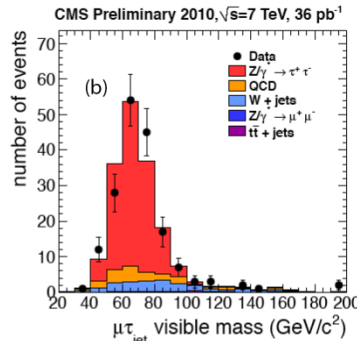
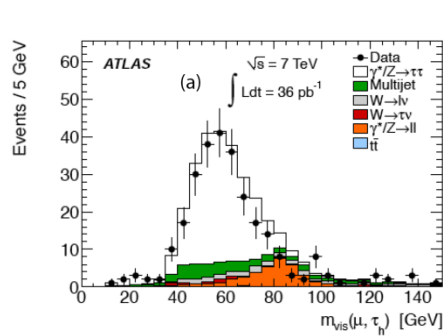
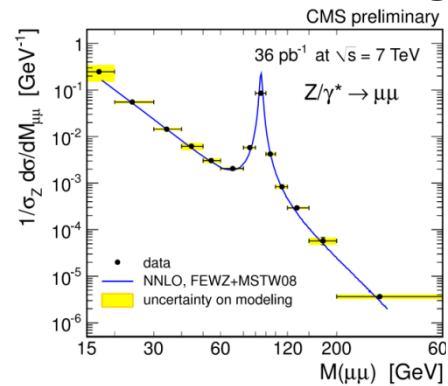
ATLAS, 100 fb^{-1} , $m_G=1.5 \text{ TeV}$



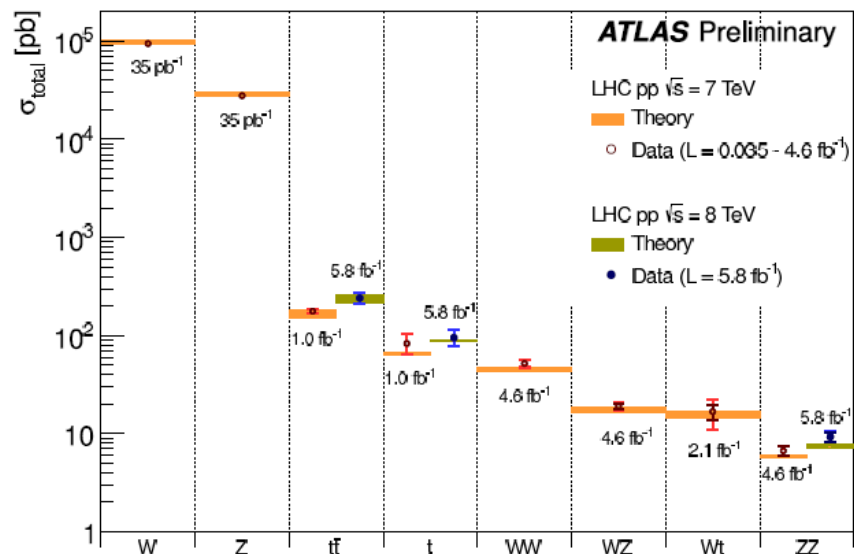
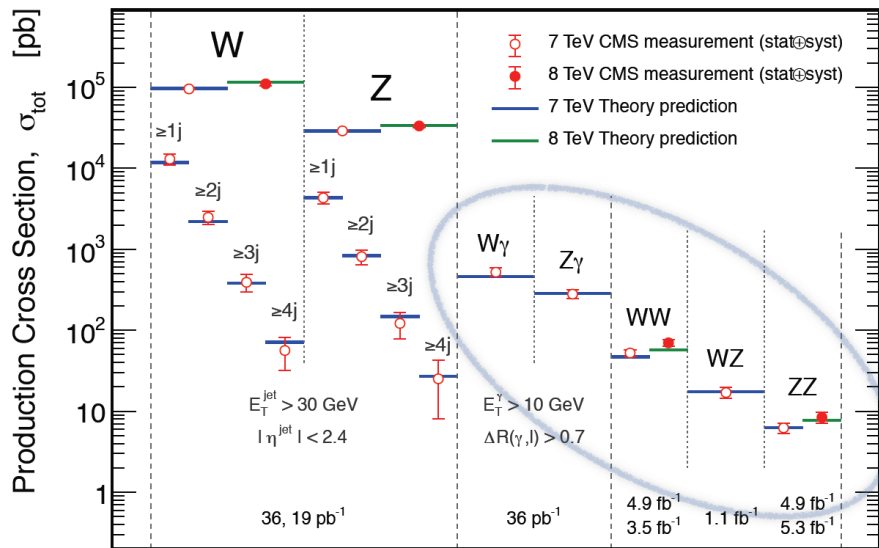
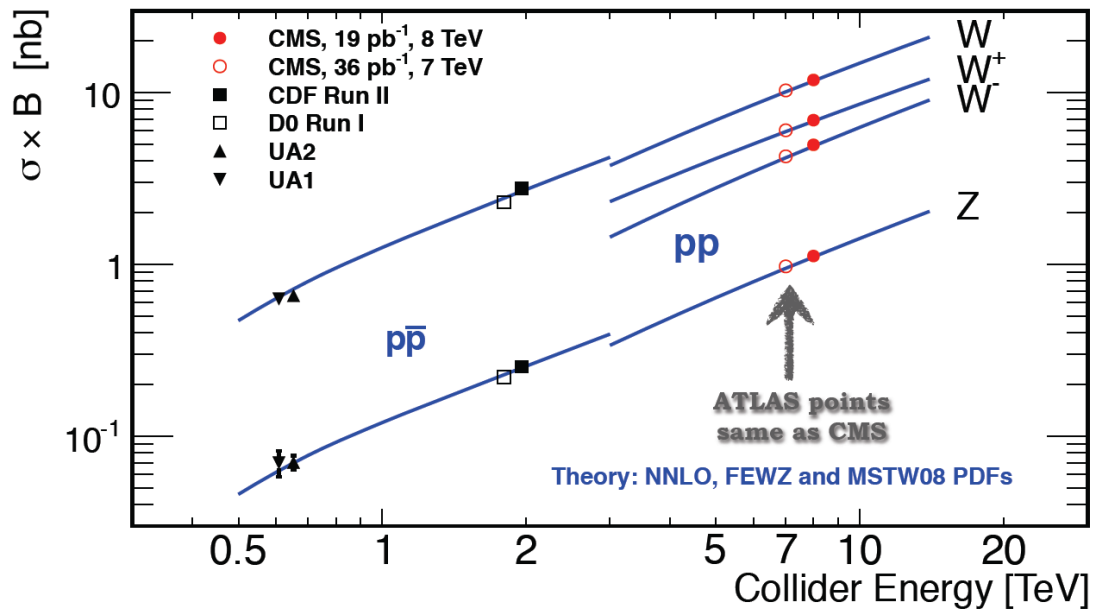
What the LHC experiments tell us?

First LHC results also confirm Standard Model

Rediscovery of the SM: W, Z, Top ... are found
 WZ, Wgamma, ZZ, Top pair, single Top ... are measured

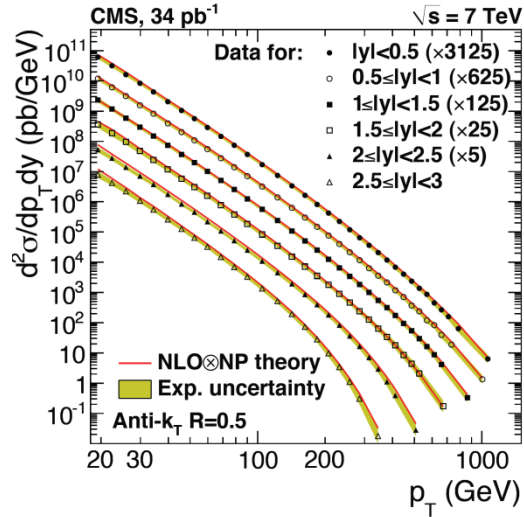


All in an agreement with the SM

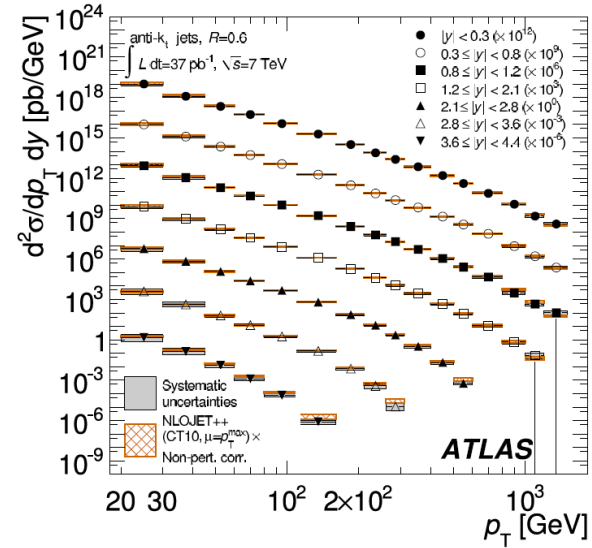
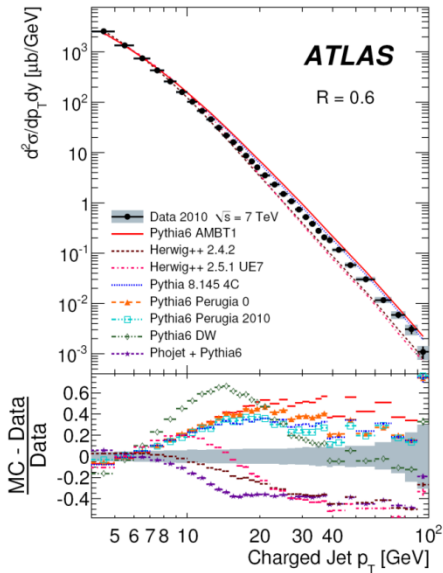
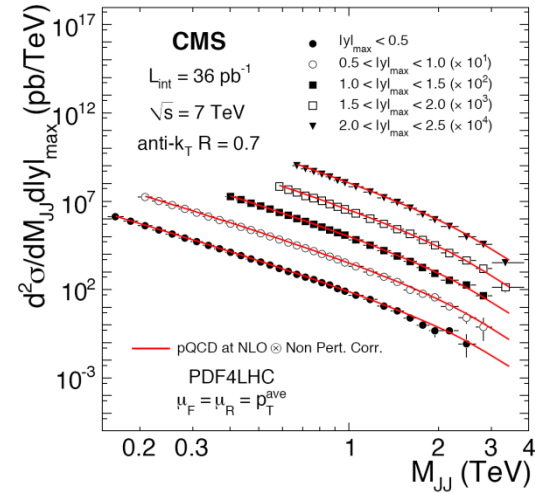


New remarkable QCD results in various kinematical regions

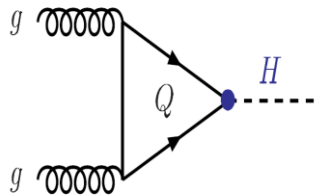
Double-differential inclusive jet production



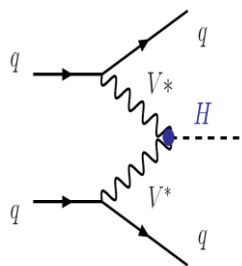
Double-differential inclusive dijet production



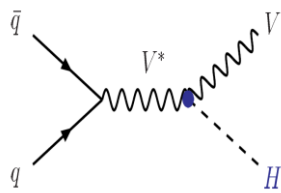
Higgs production modes, decays and signatures at LHC



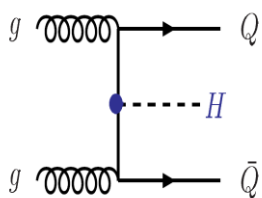
Gluon-gluon fusion



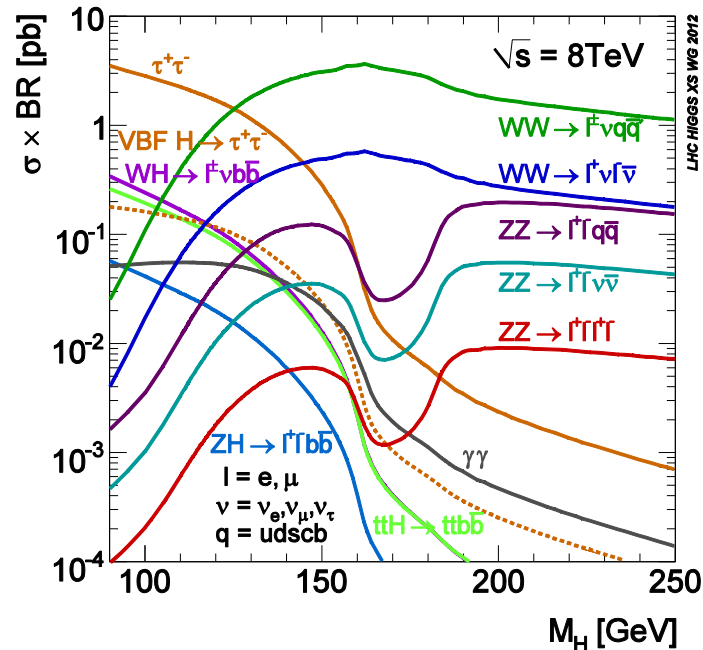
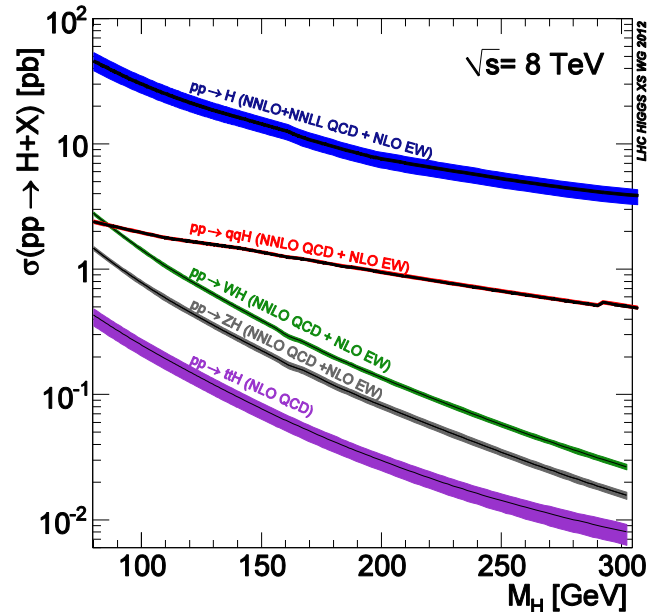
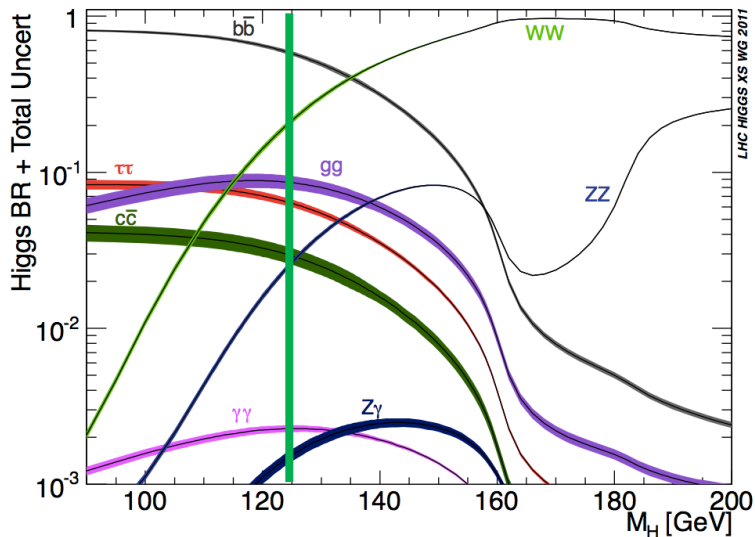
Vector boson fusion



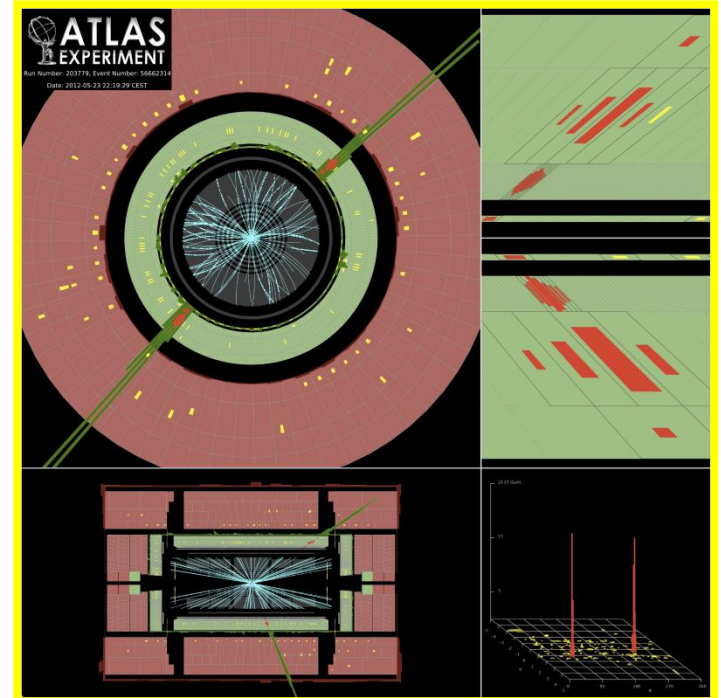
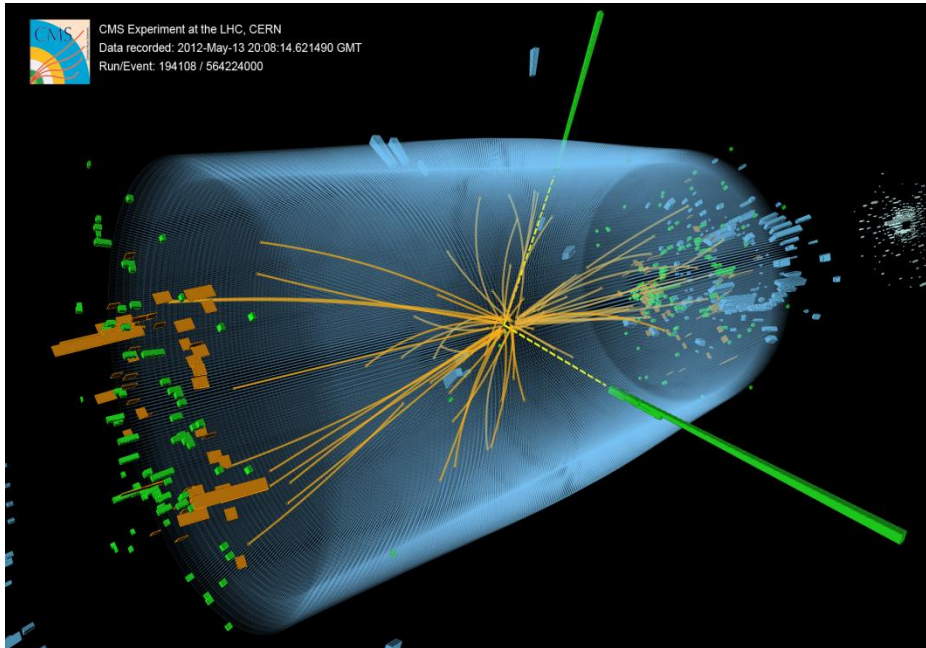
W/Z-Higgs associated



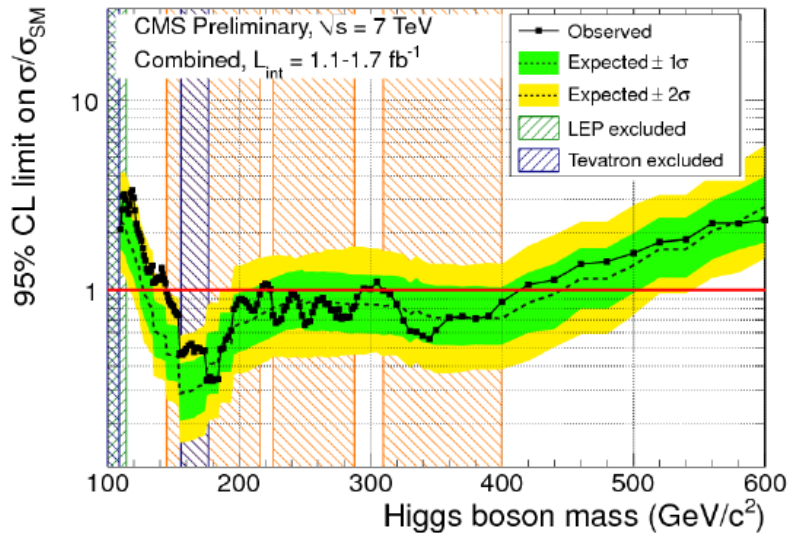
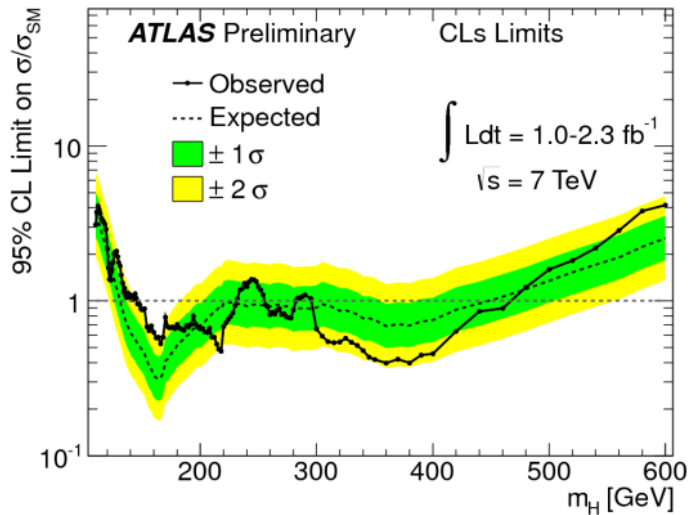
t t-bar Higgs associated



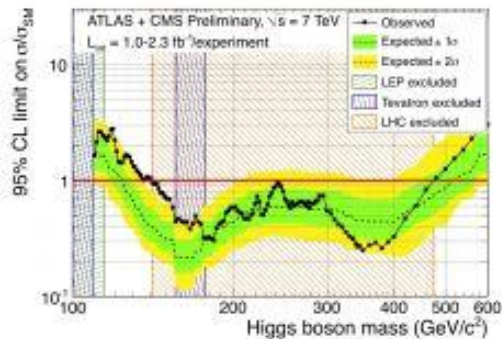
Examples of the CMS and ATLAS events with two photons (Higgs candidates)



LHC limits on Higgs mass



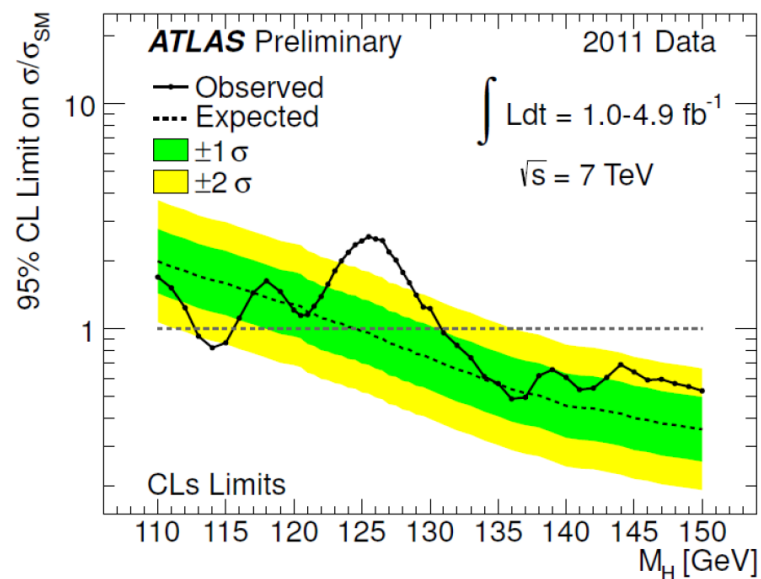
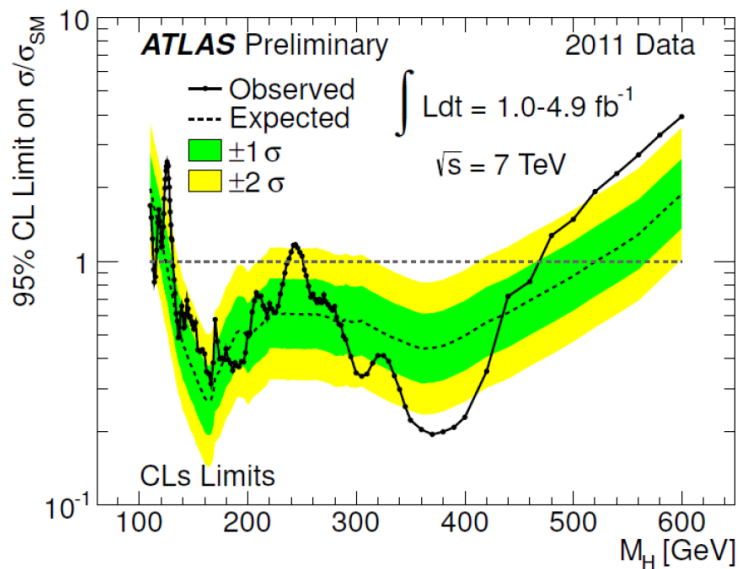
Excluded either by ATLAS or CMS 145-466 GeV (except 288-296 GeV) 95%CL



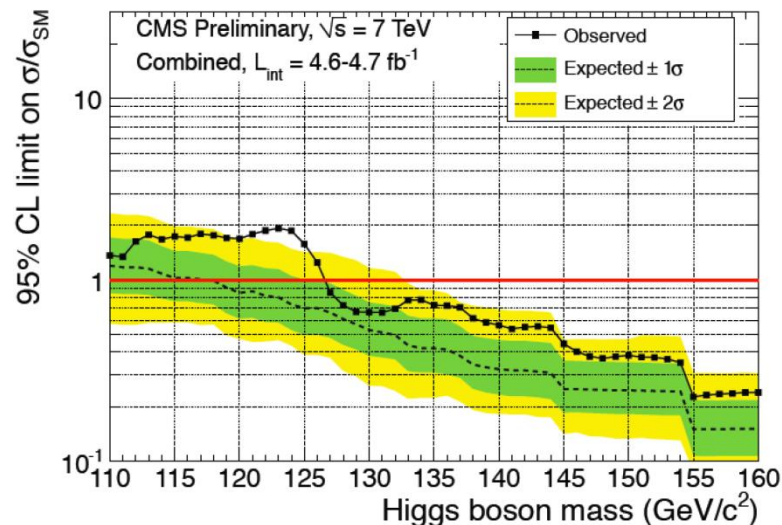
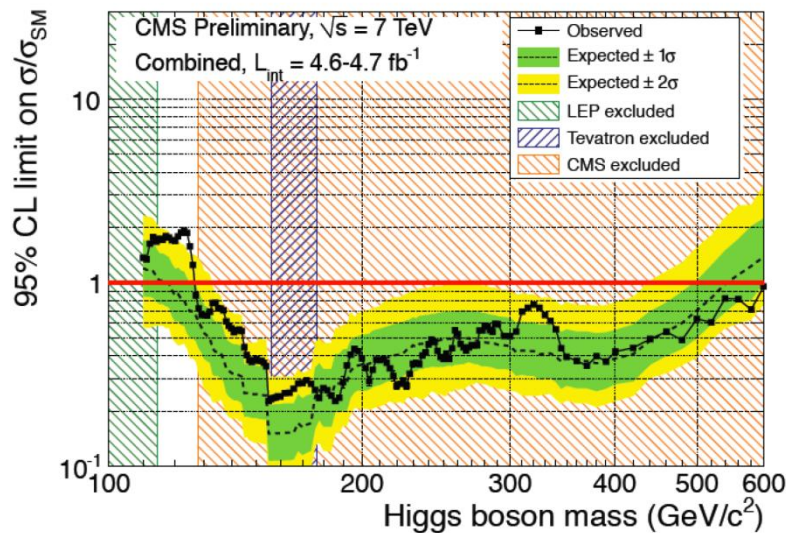
**CMS and ATLAS combined result for M_H :
141-476 GeV is excluded**

With roughly 10 1/fb per experiment at the LHC one expects to reach for SM Higgs combined 5 σ sensitivity in the interval

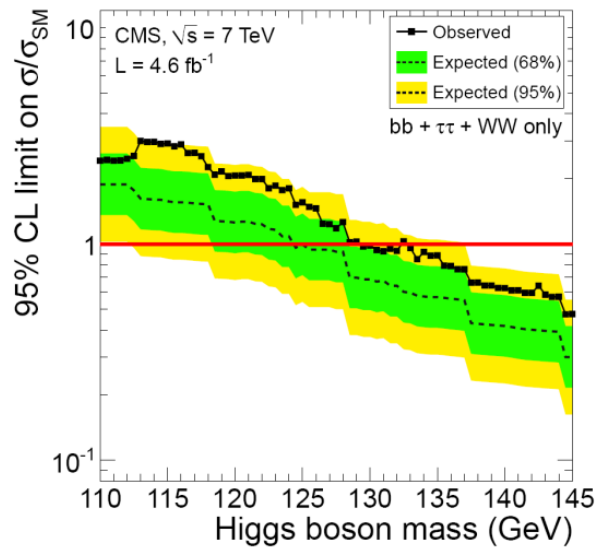
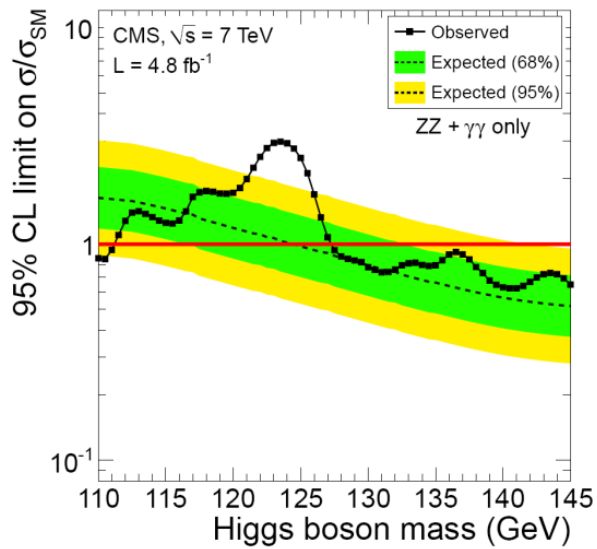
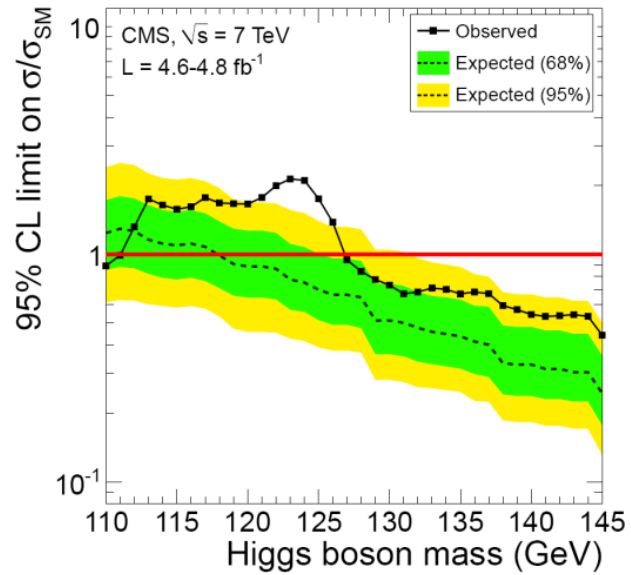
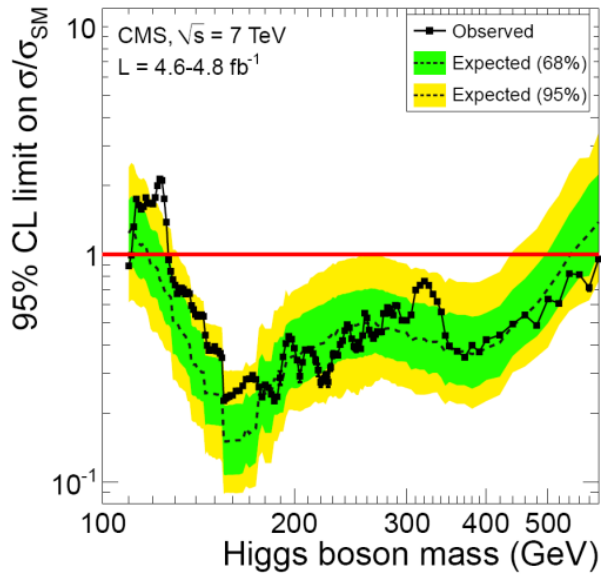
114 < M_H < 600 GeV



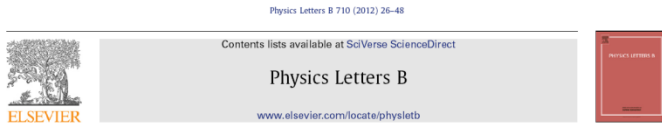
Combining only $\gamma\gamma$ and $4l$: 3.6σ



Small window from 115 GeV to 127 GeV is remaining with a small access at about 125 GeV



Observation of the Higgs-like boson in 2012 at the LHC



Combined results of searches for the standard model Higgs boson in pp collisions at $\sqrt{s} = 7$ TeV [☆]

CMS Collaboration*

CERN, Switzerland

ARTICLE INFO

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ABSTRACT

Combined results are reported from searches for the standard model Higgs boson in proton-proton collisions at $\sqrt{s} = 7$ TeV in five Higgs boson decay modes: $\gamma\gamma$, bb , $\tau\tau$, WW , and ZZ . The explored Higgs boson mass range is 110–600 GeV. The analysed data correspond to an integrated luminosity of 4.6–4.8 fb⁻¹. The expected excluded mass range in the absence of the standard model Higgs boson is 118–543 GeV at 95% CL. The observed results exclude the standard model Higgs boson in the mass range 127–600 GeV at 95% CL, and in the mass range 129–525 GeV at 99% CL. An excess of events above the expected standard model background is observed at the low end of the explored mass range making the observed limits weaker than expected in the absence of a signal. The largest excess, with a local significance of 3.1 σ , is observed for a Higgs boson mass hypothesis of 124 GeV. The global significance of observing an excess with a local significance $\geq 3.1\sigma$ anywhere in the search range 110–600 (110–145) GeV is estimated to be 1.5 σ (2.1 σ). More data are required to ascertain the origin of the observed excess.

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Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC [☆]

ATLAS Collaboration*

This paper is dedicated to the memory of our ATLAS colleagues who did not live to see the full impact and significance of their contributions to the experiment.

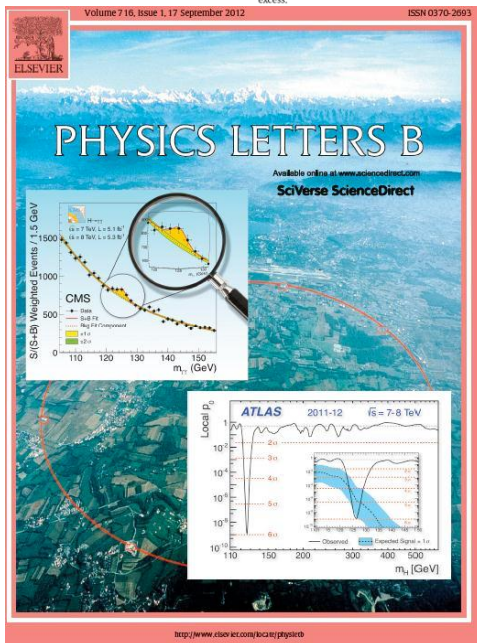
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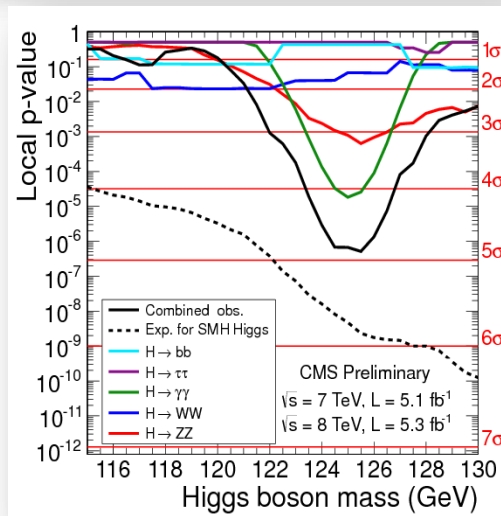
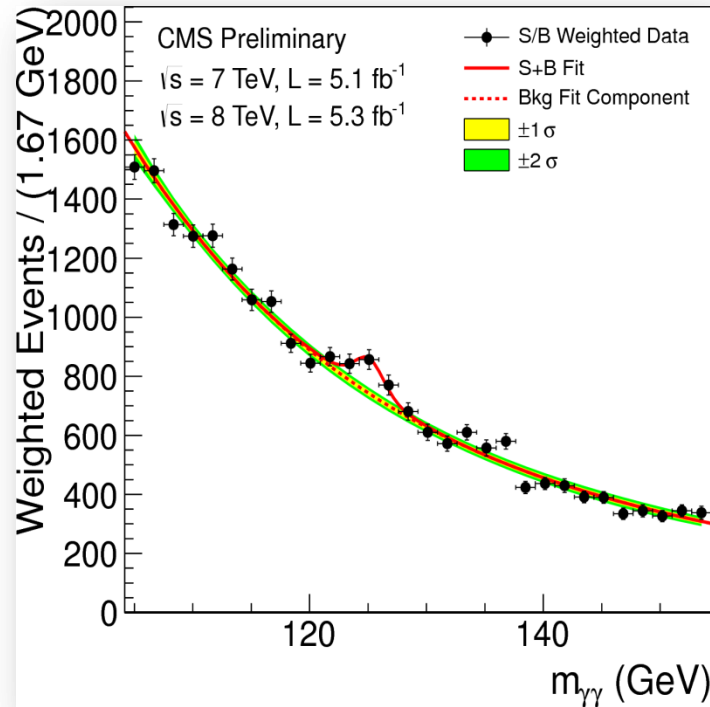
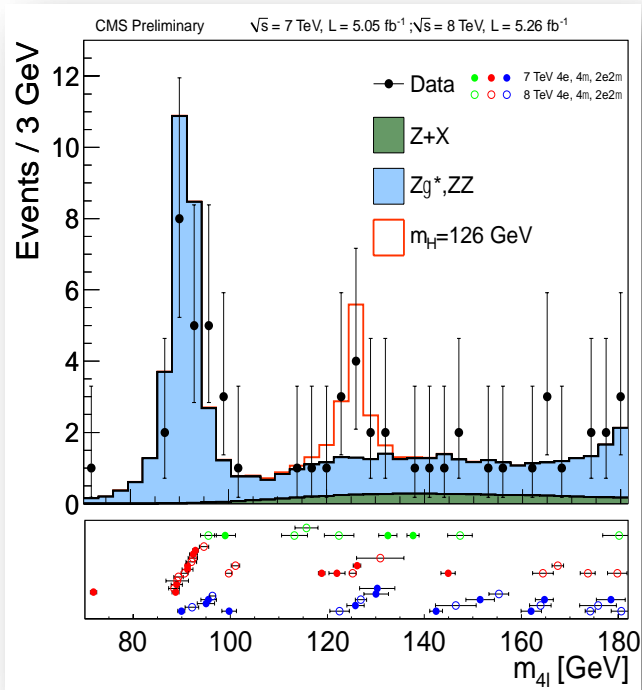
ABSTRACT

A search for the Standard Model Higgs boson in proton-proton collisions with the ATLAS detector at the LHC is presented. The datasets used correspond to integrated luminosities of approximately 4.8 fb⁻¹ collected at $\sqrt{s} = 7$ TeV in 2011 and 5.8 fb⁻¹ at $\sqrt{s} = 8$ TeV in 2012. Individual searches in the channels $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$, $H \rightarrow \gamma\gamma$ and $H \rightarrow WW^{(*)} \rightarrow e\nu\mu\nu$ in the 8 TeV data are combined with previously published results of searches for $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ in the 7 TeV data and results from improved analyses of the $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ channels in the 7 TeV data. Clear evidence for the production of a neutral boson with a measured mass of 126.0 ± 0.4 (stat) ± 0.4 (sys) GeV is presented. This observation, which has a significance of 5.9 standard deviations, corresponding to a background fluctuation probability of 1.7×10^{-5} , is compatible with the production and decay of the Standard Model Higgs boson.

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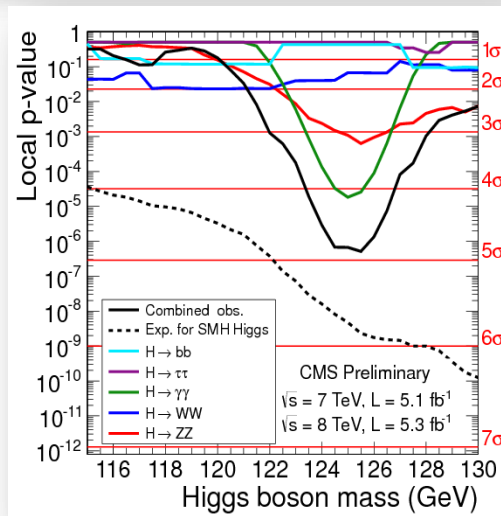
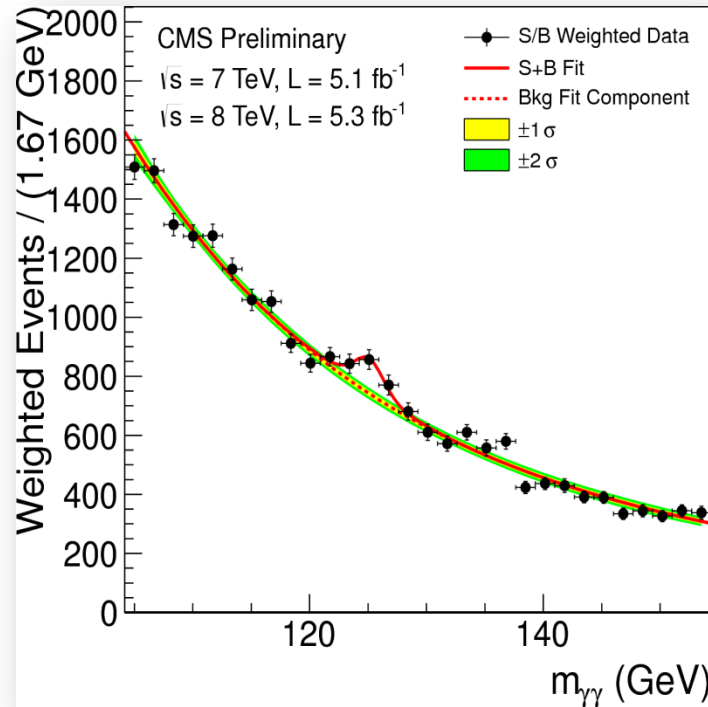
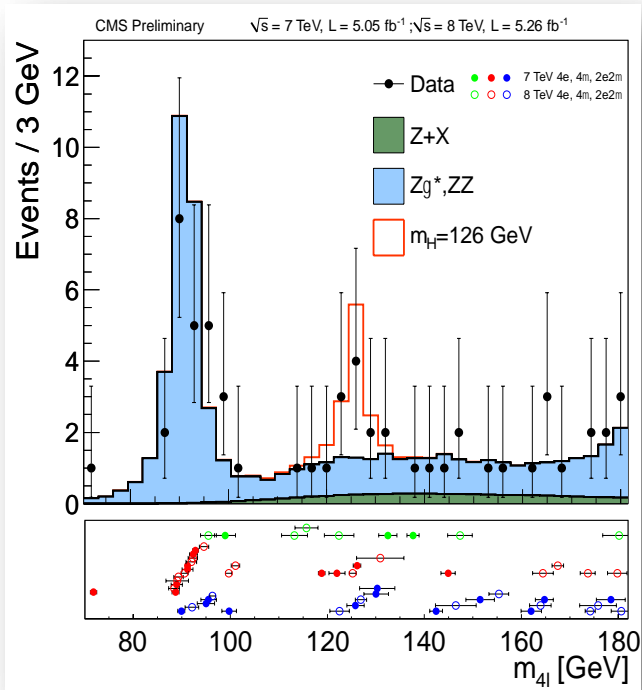


Результаты CMS на ICHEP2012

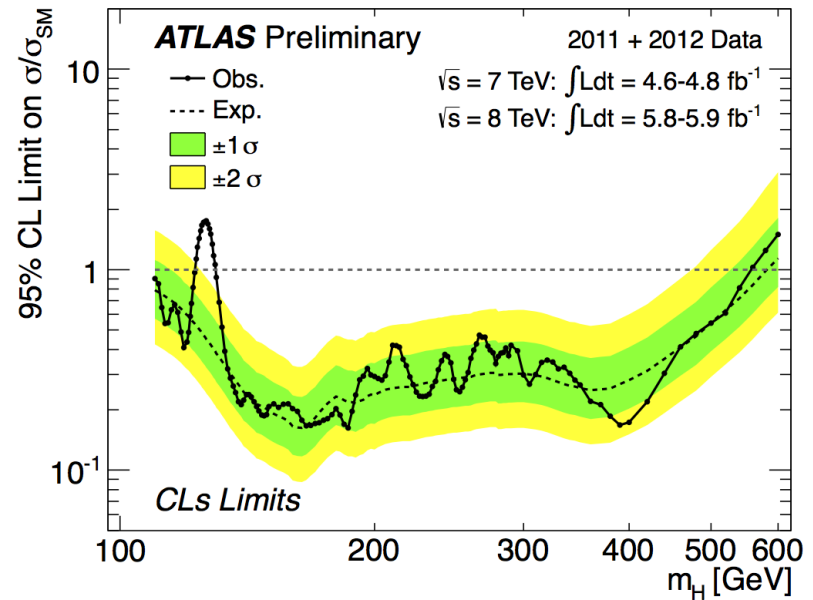
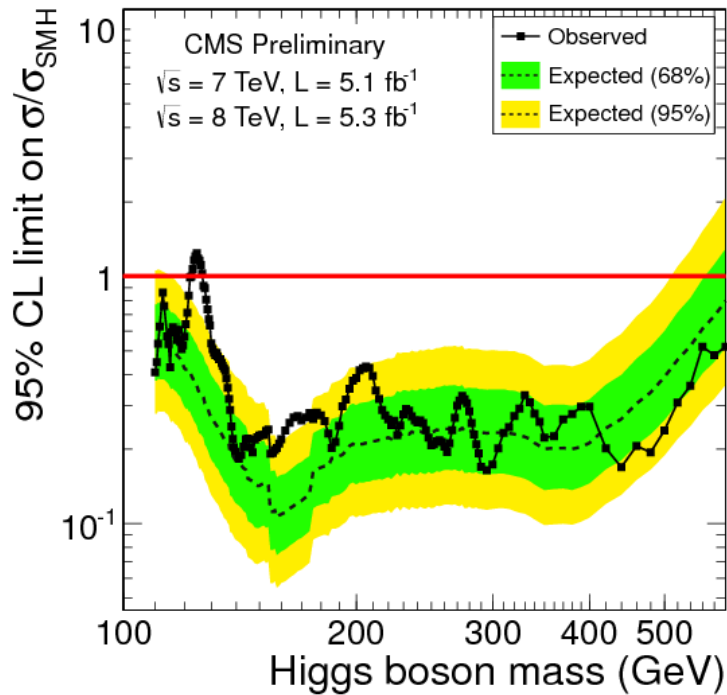


Зарегистрирован новый бозон
с массой
 $125.3 \pm 0.6 \text{ GeV}$
Достоверность результата
 4.9σ !

Результаты CMS на ICHEP2012

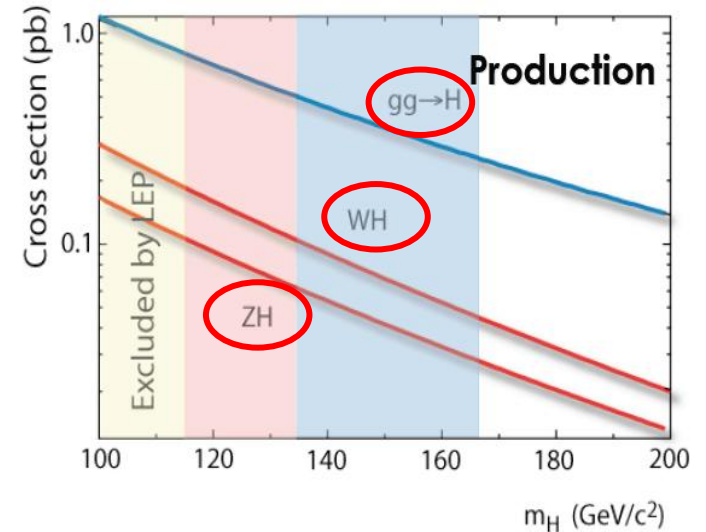
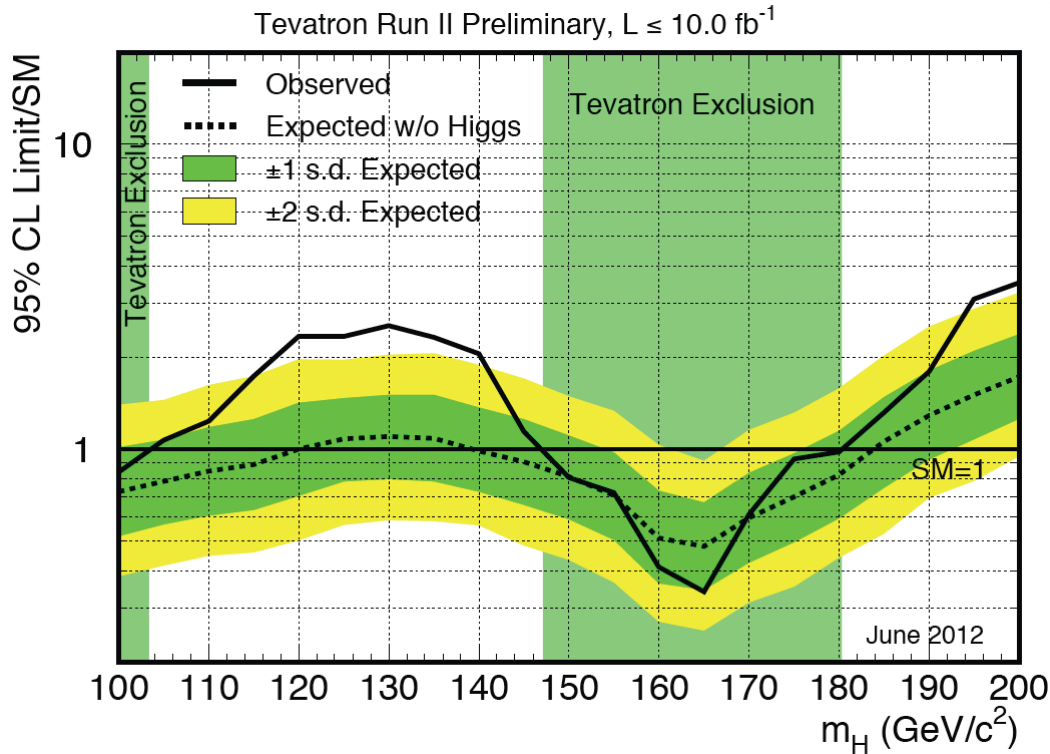


Зарегистрирован новый бозон
с массой
 $125.3 \pm 0.6 \text{ GeV}$
Достоверность результата
 4.9σ !



**Both CMS and ATLAS have excluded SM Higgs
 in the mass interval upto about 560 GeV
 except small interval where the signal was observed**

Summer 2012 Tevatron Combination



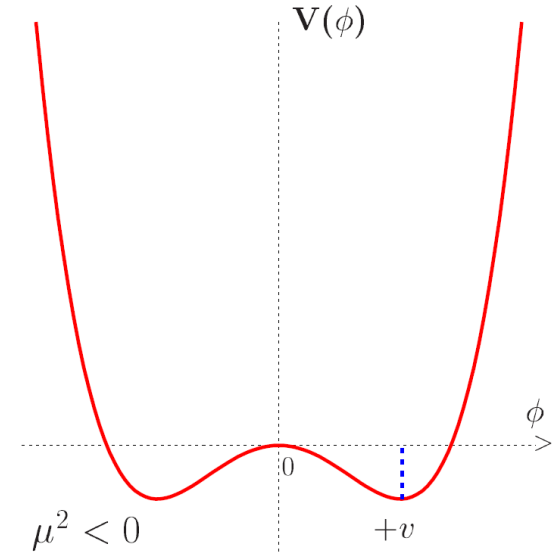
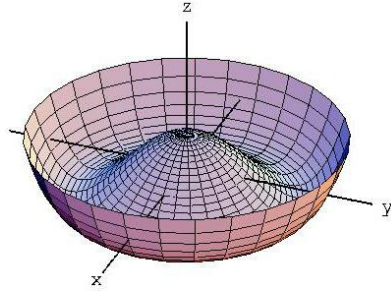
**Significant excess:
 $\sim 3\sigma$ for $115 \rightarrow 140 \text{ GeV}$**

Not only discovery and mass measurement

but also

Very good precision of the mass measurement

Exclusion of large range of higher masses



$$V(\Phi) = \mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$$

$$\Phi \rightarrow \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ H+v \end{pmatrix}$$

$$\mathcal{L}_H = \frac{1}{2} (\partial_\mu H) (\partial^\mu H) - V = \frac{1}{2} (\partial^\mu H)^2 - \lambda v^2 H^2 - \lambda v H^3 - \frac{\lambda}{4} H^4$$

$$M_H^2 = 2\lambda v^2 = -2\mu^2$$

$$\lambda \cong 0.12$$

Origin of the EWSB potential → a weakly-coupled theory

Is observed state the Standard Model Higgs?

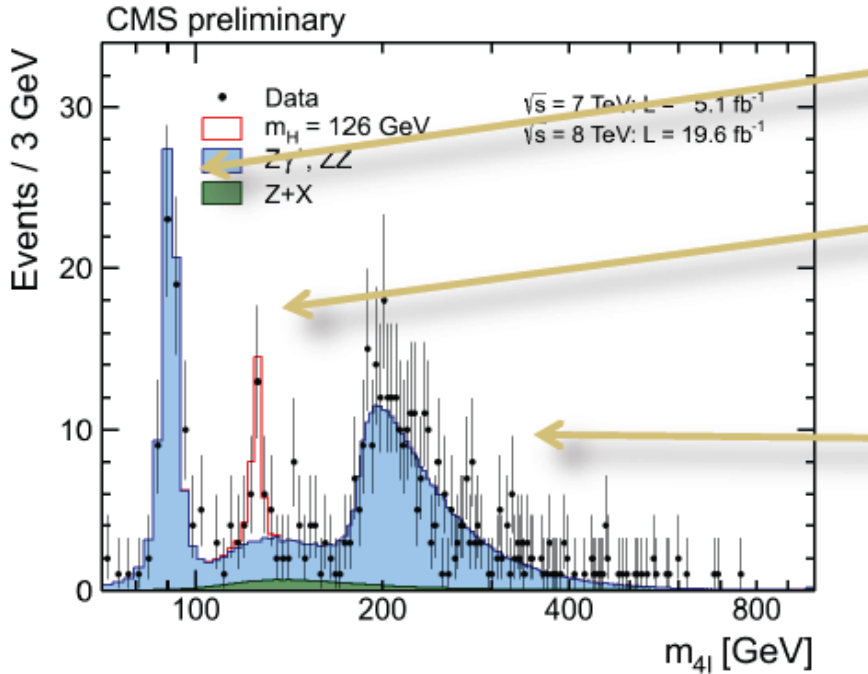
Production and decay modes?

Spin and parity?



More precise measurements are needed

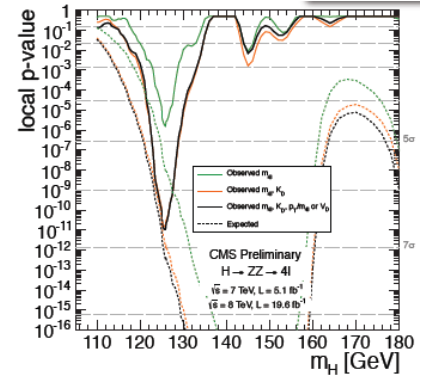
ZZ*-channel CMS



$Z\gamma \rightarrow 4l$

$H \rightarrow ZZ^* \rightarrow 4l$

$ZZ \rightarrow 4l$



p-value:
 Expected: 7.1σ
 Observed: 6.7σ

$$\mu = 0.92^{+0.30}_{-0.24}$$

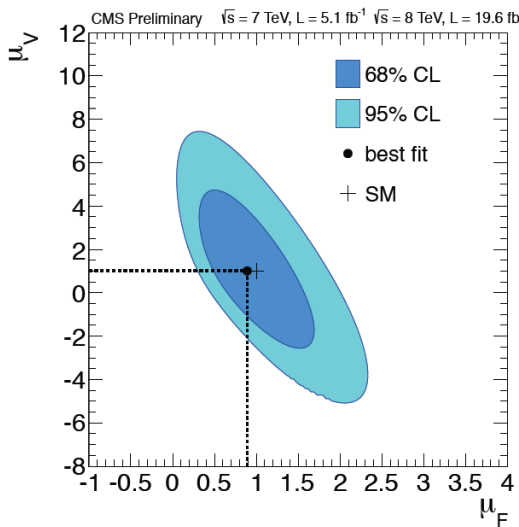
$$M_H = 125.8 \pm 0.5(\text{stat}) \pm 0.2(\text{syst}) \text{ GeV}$$

Signal strength

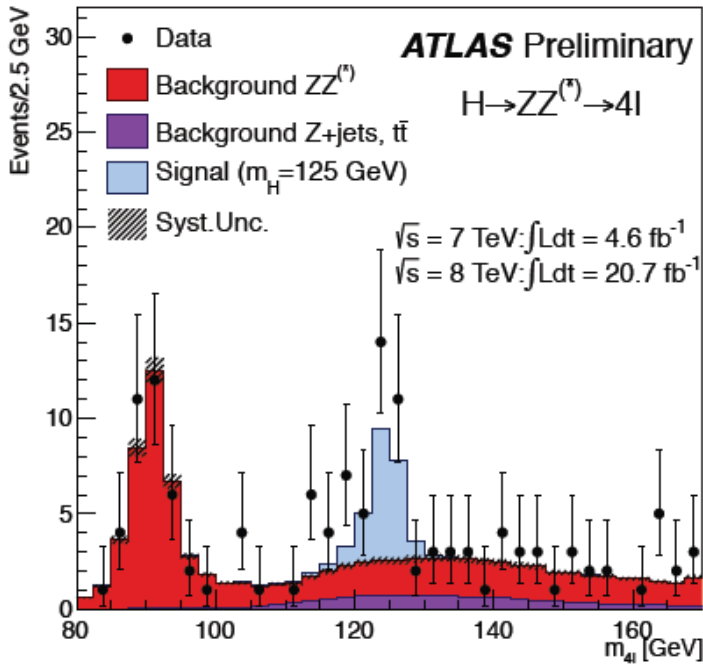
$$\mu_V (qqH, ZH, WH) = 1.0^{+2.4}_{-2.3}$$

$$\mu_F (gg \rightarrow H, t\bar{t}H) = 0.9^{+0.5}_{-0.4}$$

$$\mu = \frac{\sigma \times \text{BR}}{(\sigma \times \text{BR})_{\text{SM}}}$$



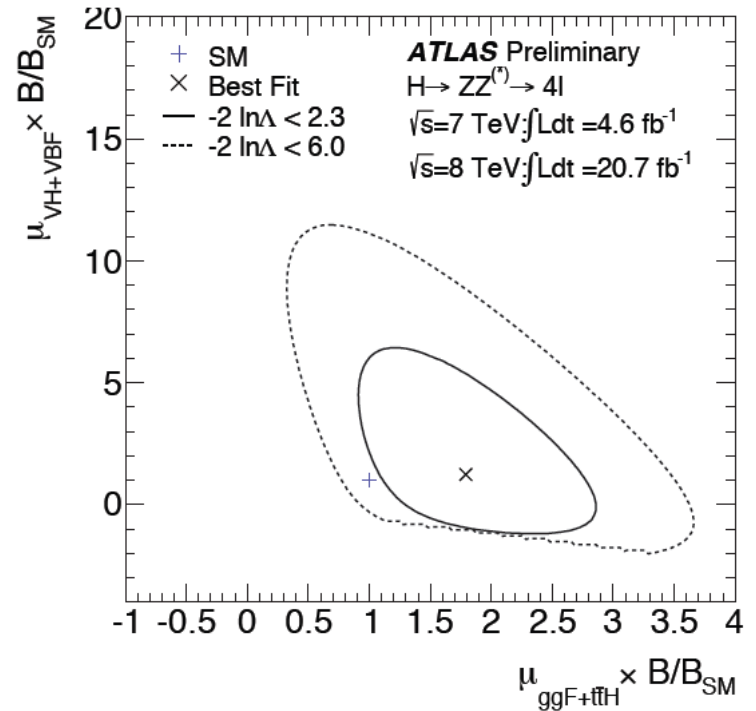
ZZ*-channel ATLAS



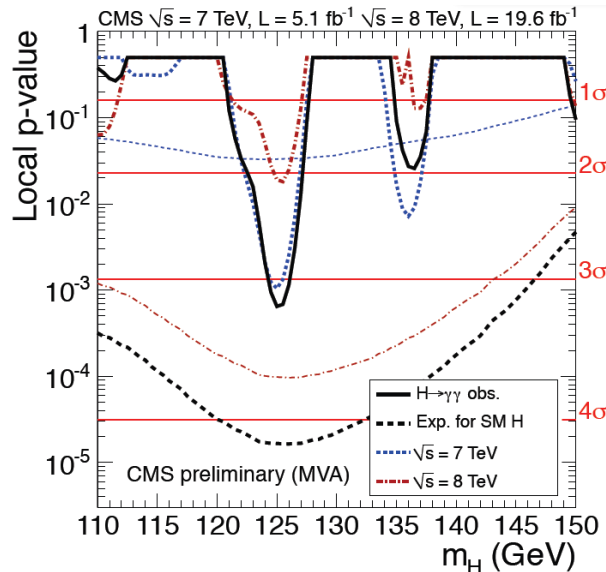
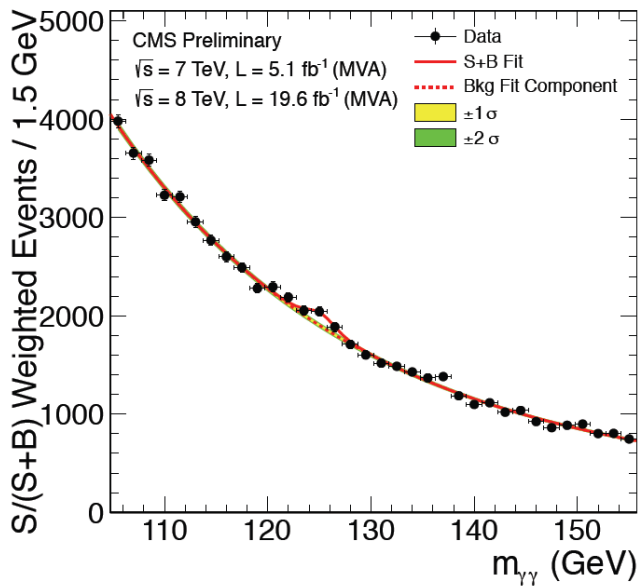
$$M_H = 124.3^{+0.6}_{-0.5}(\text{stat})^{+0.5}_{-0.3}(\text{syst}) \text{ GeV}$$

$$\text{Signal strength } \mu = 1.7^{+0.5}_{-0.4}$$

$$\mu_{\text{VBF+VH}} / \mu_{\text{ggF+ttH}} = 0.7^{+2.4}_{-1.0}$$



$\gamma\gamma$ -channel CMS



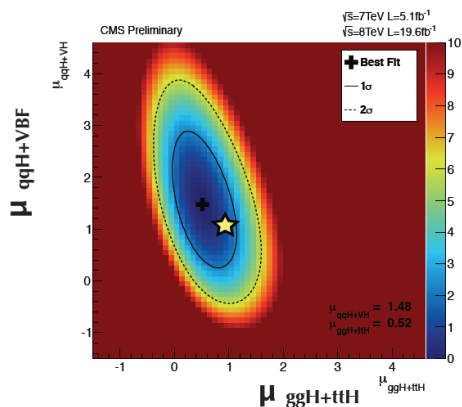
Boosted decision trees:

$$\sigma/\sigma_{\text{SM}} = 0.78^{+0.28}_{-0.26}$$

Cut based:

$$\sigma/\sigma_{\text{SM}} = 1.11^{+0.32}_{-0.30}$$

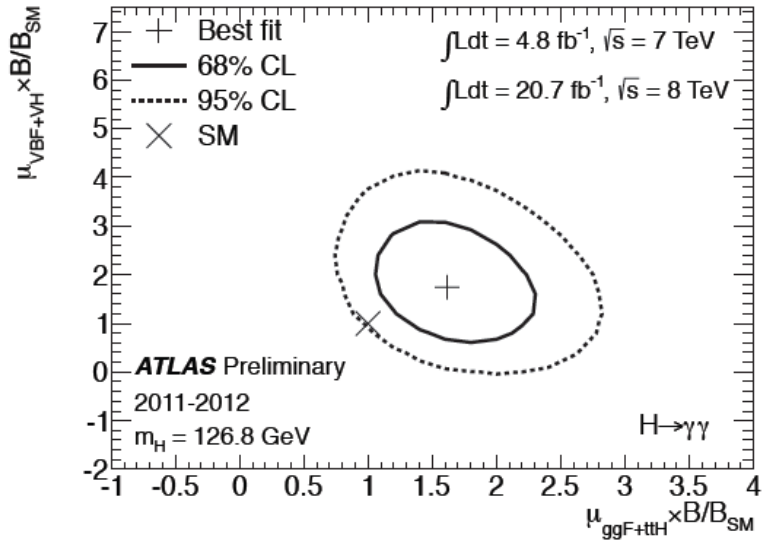
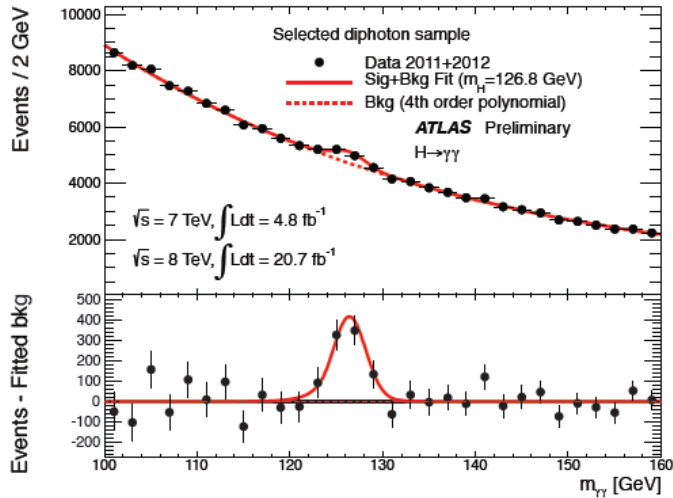
$$M_H = 125.4 \pm 0.5(\text{stat.}) \pm 0.6(\text{syst.}) \text{ GeV}$$



$$\mu_{\text{qqH+VBF}} = 1.48$$

$$\mu_{\text{ggH+ttH}} = 0.52$$

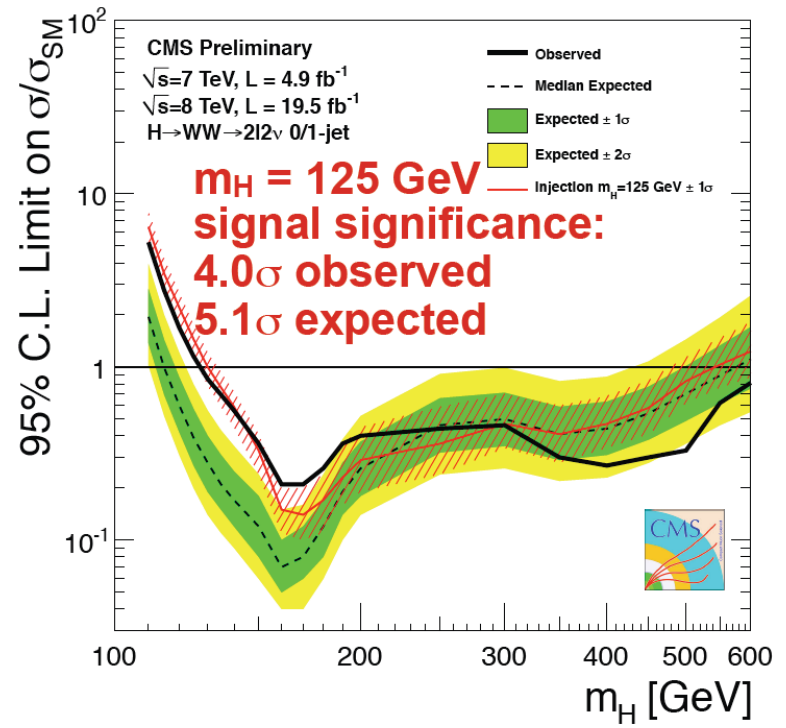
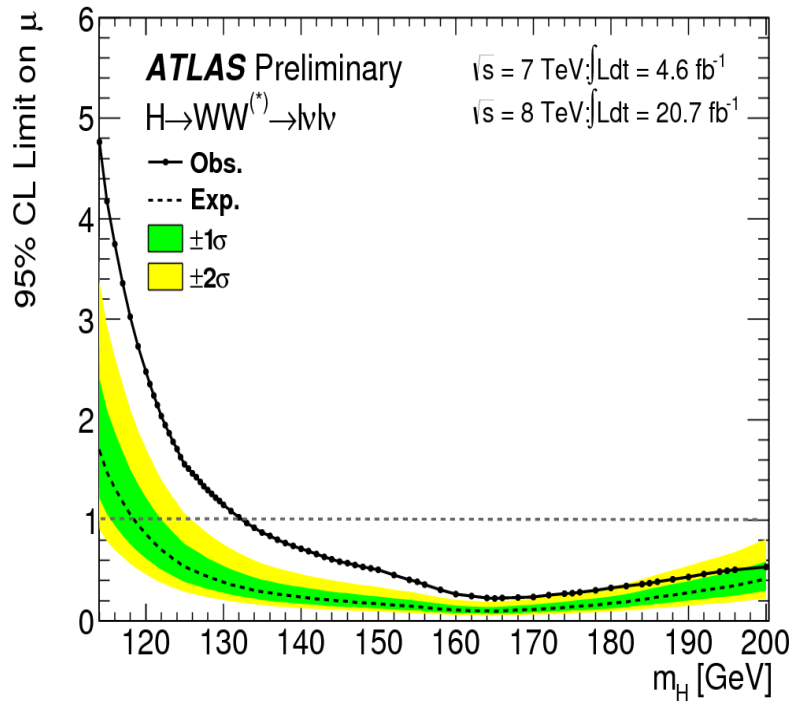
$\gamma\gamma$ -channel ATLAS



$$M_H = 126.8 \pm 0.2(\text{stat}) \pm 0.7(\text{syst}) \text{ GeV}$$

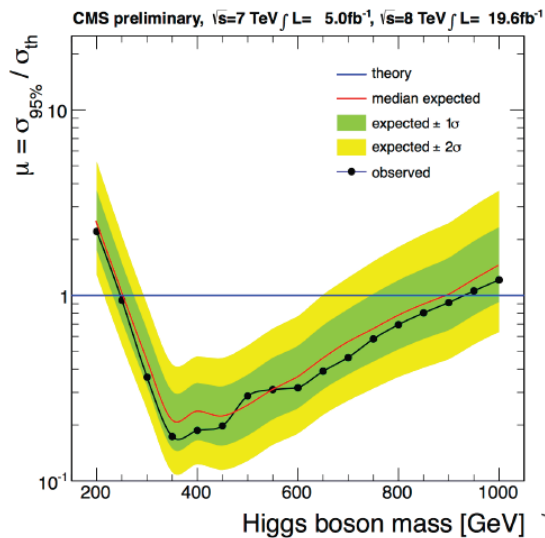
Signal strength $\mu = 1.65 \pm 0.24(\text{stat})^{+0.25}_{-0.18}(\text{syst})$
 2.3σ from the SM hypothesis

H → WW*



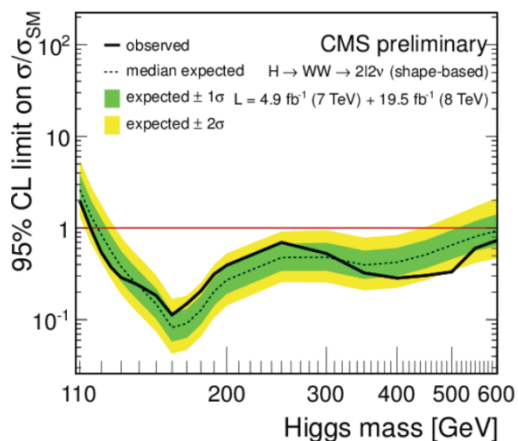
bb and $\tau\tau$ modes are not yet visible individually

First evidence 3.4 σ for combination of bb and $\tau\tau$ is given by CMS



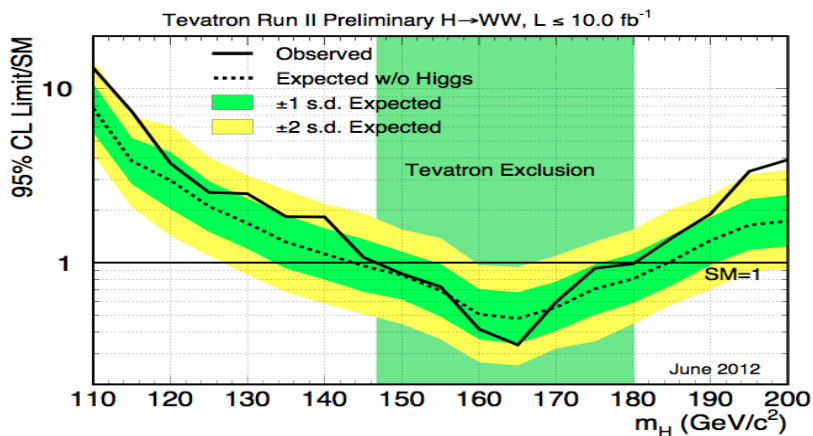
ZZ -> 2l + 2v

Excluded region 248-930 GeV



WW -> 2l + 2v

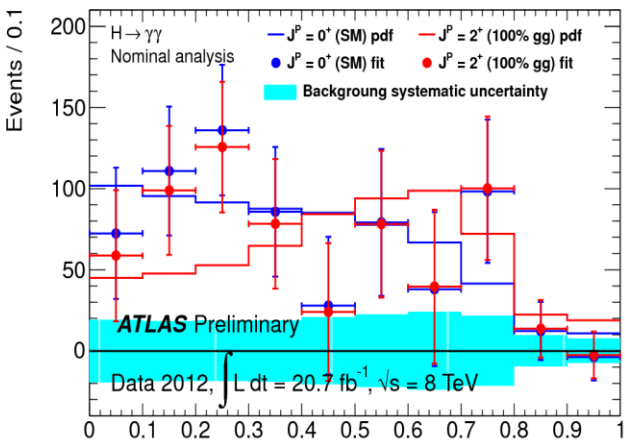
Excluded region 140-590 GeV



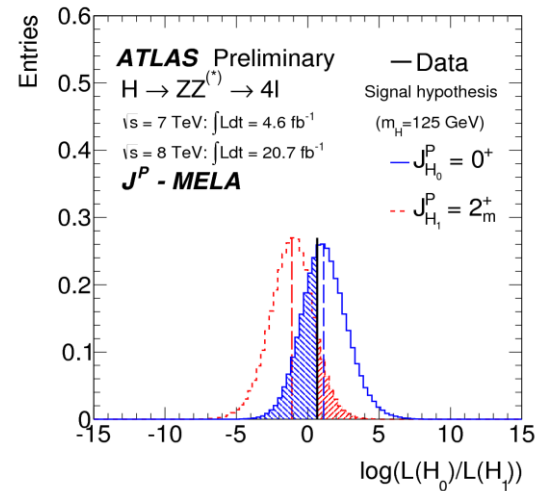
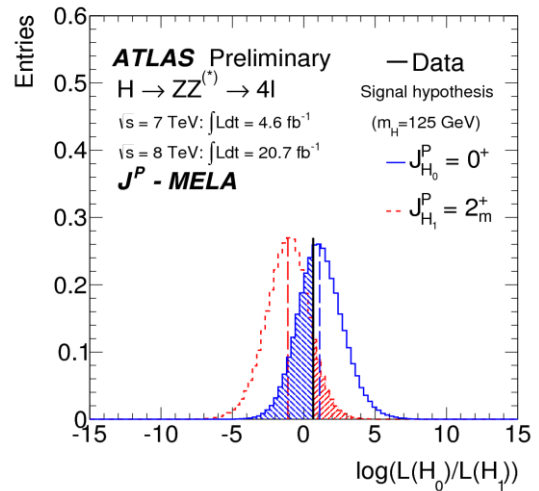
Tevatron combination: Excluded region
 $147 < M_H < 180\text{ GeV}$ at 95% CL

Spin, Parity ?

H decays to $\gamma\gamma$, can not have spin 1 - Landau, Yang theorem



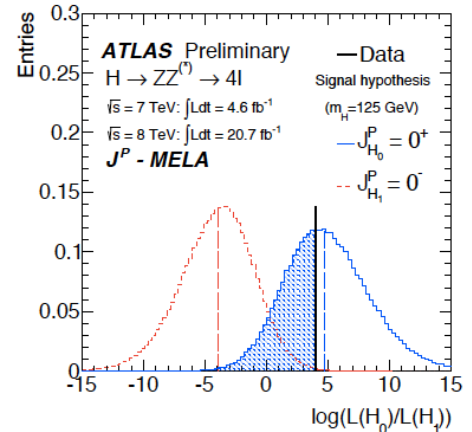
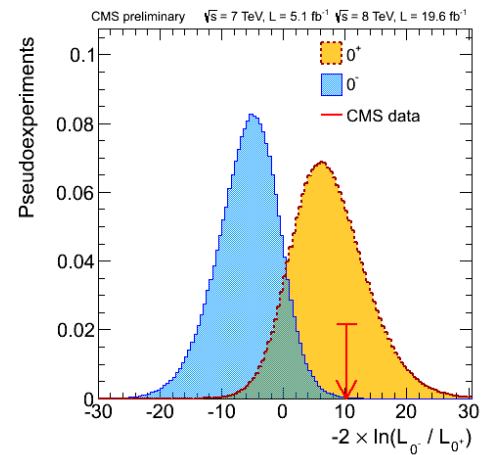
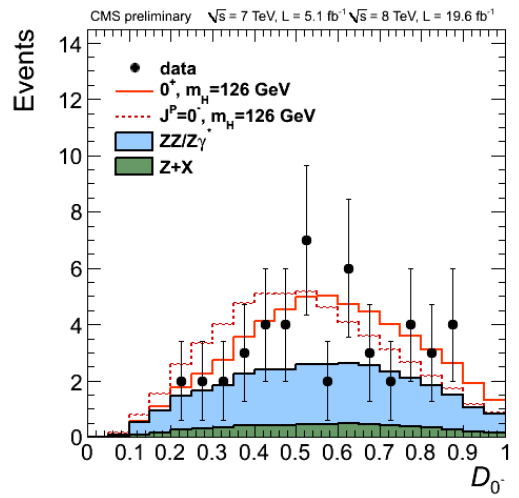
$$\cos \theta^* = \frac{\sinh(\eta_{\gamma_1} - \eta_{\gamma_2})}{\sqrt{1 + (p_T^{\gamma\gamma} / m_{\gamma\gamma})^2}} \cdot \frac{2p_T^{\gamma_1} p_T^{\gamma_2}}{m_{\gamma\gamma}^2} \quad [\cos \theta^*]$$



Graviton-like spin-2+ disfavoured at 99.9% CL

$$\mathcal{D}_{J^P} = \frac{\mathcal{P}_{SM}}{\mathcal{P}_{SM} + \mathcal{P}_{J^P}} = \left[1 + \frac{\mathcal{P}_{J^P}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell})}{\mathcal{P}_{SM}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell})} \right]^{-1}$$

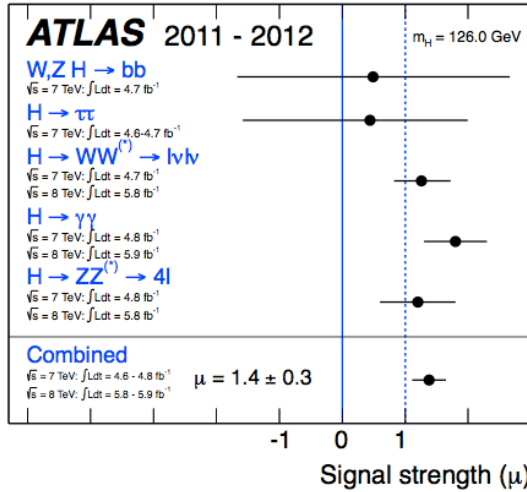
$$\text{MELA} = \left[1 + \frac{\mathcal{P}_{\text{bkg}}(m_1, m_2, \theta_1, \theta_2, \Phi, \theta^*, \Phi_1 | m_{4\ell})}{\mathcal{P}_{\text{sig}}(m_1, m_2, \theta_1, \theta_2, \Phi, \theta^*, \Phi_1 | m_{4\ell})} \right]^{-1}$$



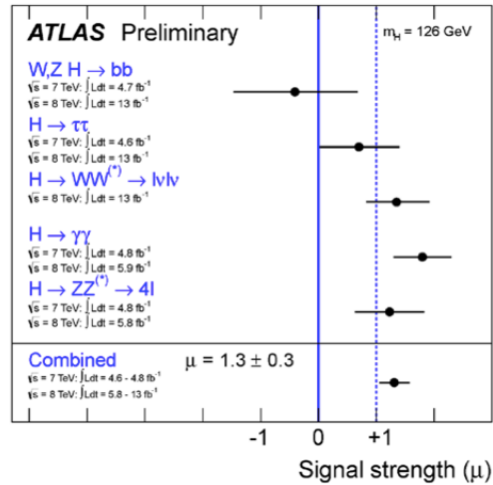
Pseudoscalar 0- disfavoured at > 99% CL

Signal strength μ

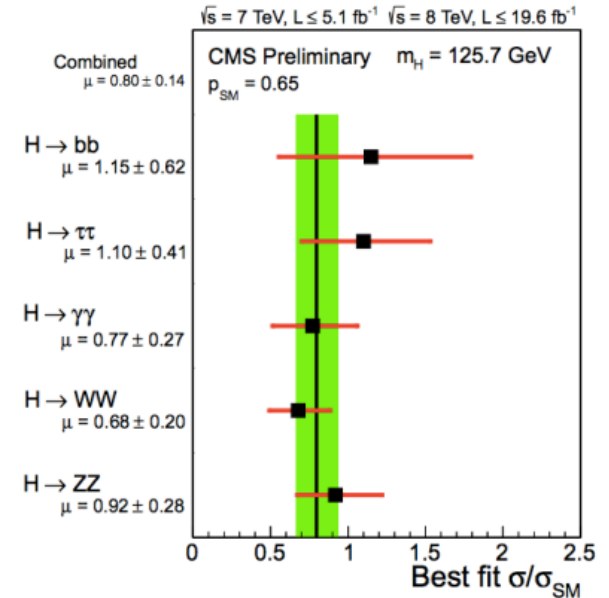
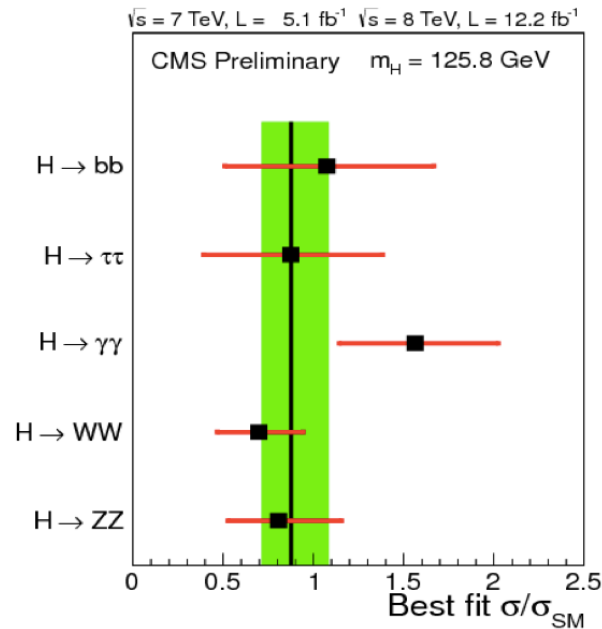
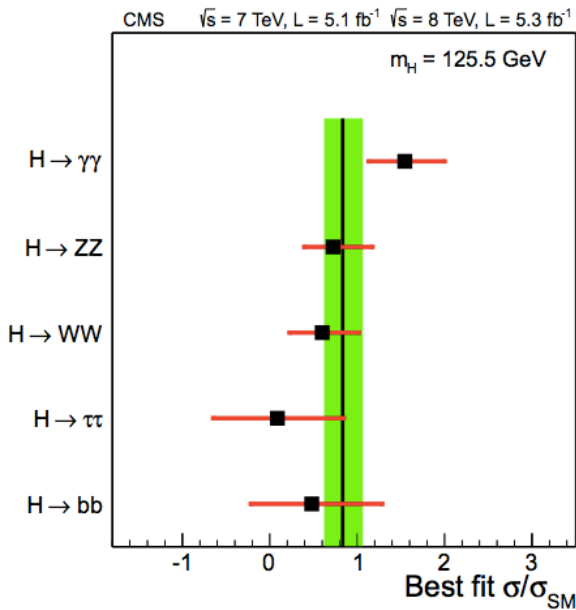
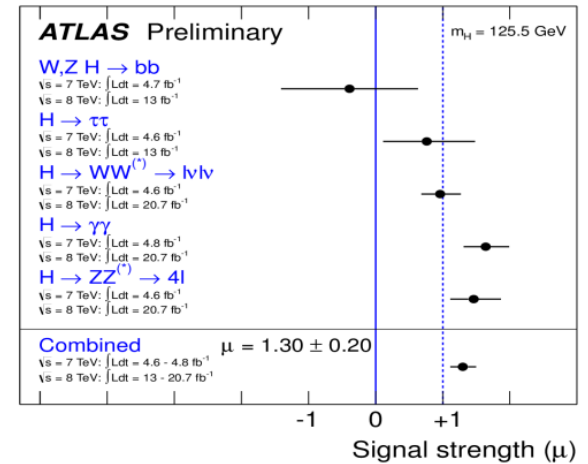
July 2012



November 2012

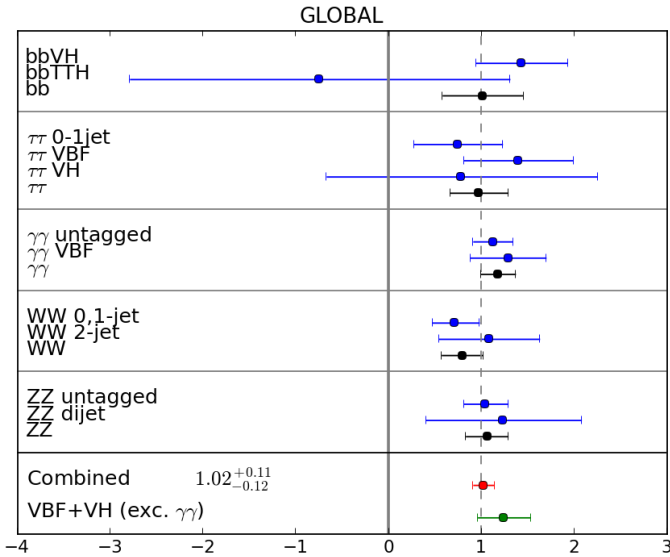


July 2013



Does the resonance couple to particle masses?

John Ellis et al 2013



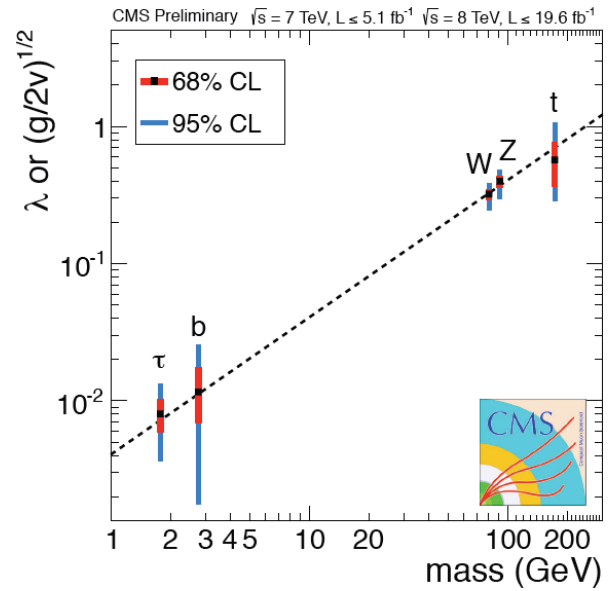
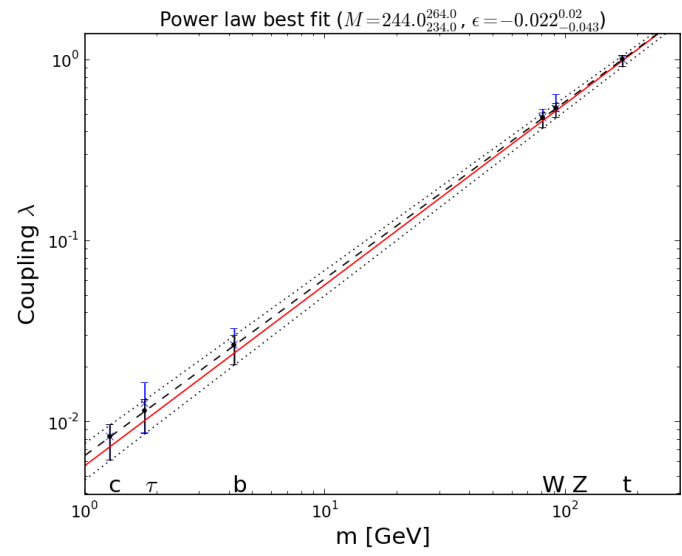
$$\lambda_f = \sqrt{2} \left(\frac{m_f}{M} \right)^{1+\epsilon}, \quad g_V = 2 \left(\frac{m_V^{2(1+\epsilon)}}{M^{1+2\epsilon}} \right)$$

In SM $\epsilon=0, M=v$

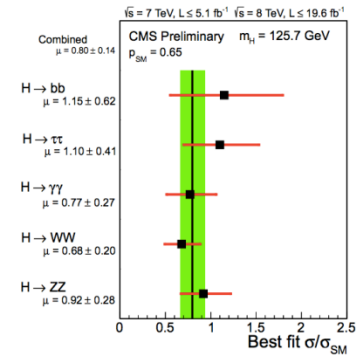
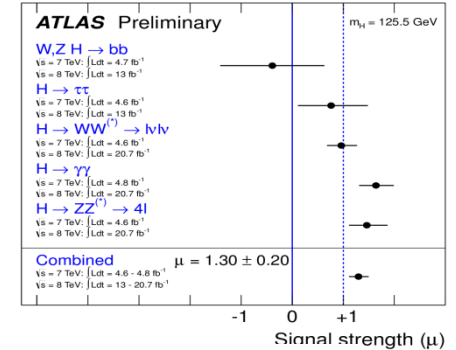
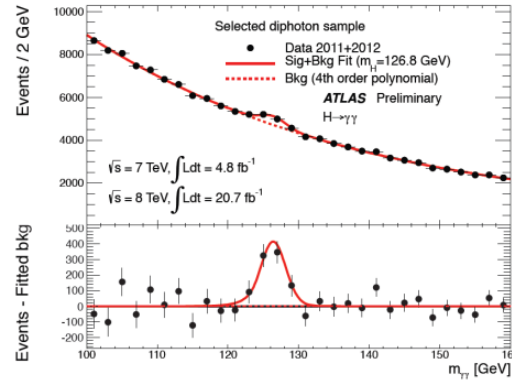
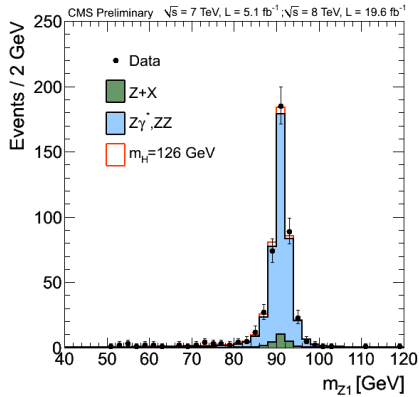
Best fit:

$$\epsilon = -0.022^{+0.042}_{-0.021}$$

$$M = 244^{+20}_{-10} \text{ GeV}$$



- Higgs boson is found. Confirmation in 2013 with more statistics



- 0^+ state is confirmed

- Signal strengths coincide with SM predictions

- But accuracy is not good enough. Not all modes observed. Lot of room for extensions

Limits on invisible mode are still very weak

ATLAS (4.7+13.0 fb⁻¹):
 $\text{Br}(H \rightarrow \chi\chi) < 65\% @ 95\% \text{ CL},$
 $M_H = 125 \text{ GeV}$

CMS (5+20 fb⁻¹):
 $\text{Br}(H \rightarrow \chi\chi) < 75\% @ 95\% \text{ CL},$
 $M_H = 125 \text{ GeV}$

$$M_h^2 \leq M_Z^2 + \Delta m^2$$

→ susy breaking term
(at one-loop)

$$(125 \text{ GeV})^2$$

$$(91 \text{ GeV})^2$$

$$(86 \text{ GeV})^2$$

$$\Delta m_h^2 = \frac{3m_t^4}{4\pi^2 v^2} \left[\ln \left(\frac{M_{\text{SUSY}}^2}{m_t^2} \right) + \frac{X_t^2}{M_{\text{SUSY}}^2} \left(1 - \frac{X_t^2}{12M_{\text{SUSY}}^2} \right) \right]$$

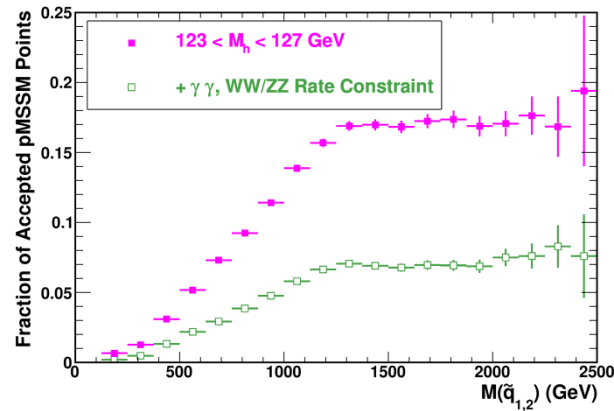
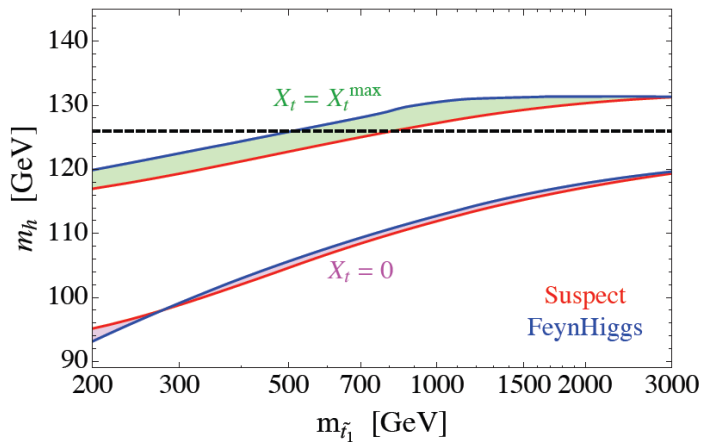
$$M_{\text{SUSY}} \equiv \sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}}$$

Heavy stop is needed

Fine tuning:

$$-\frac{m_Z^2}{2} = |\mu|^2 + m_{H_u}^2$$

MSSM Higgs Mass



Mahmoudi et al 2012

Many options for observed state in the NMSSM

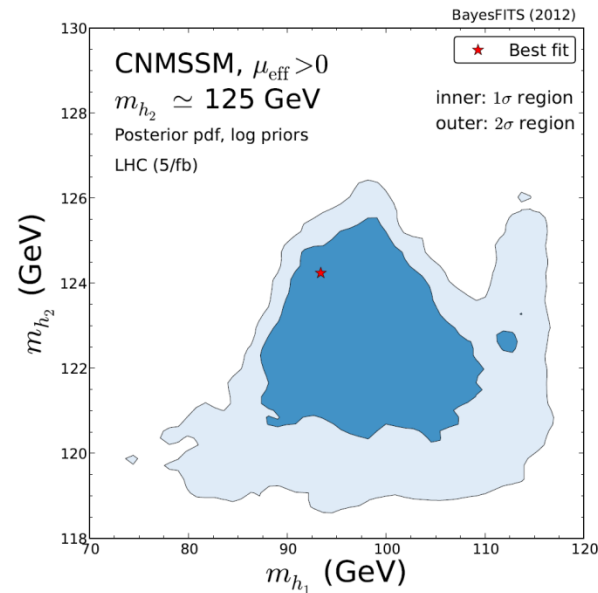
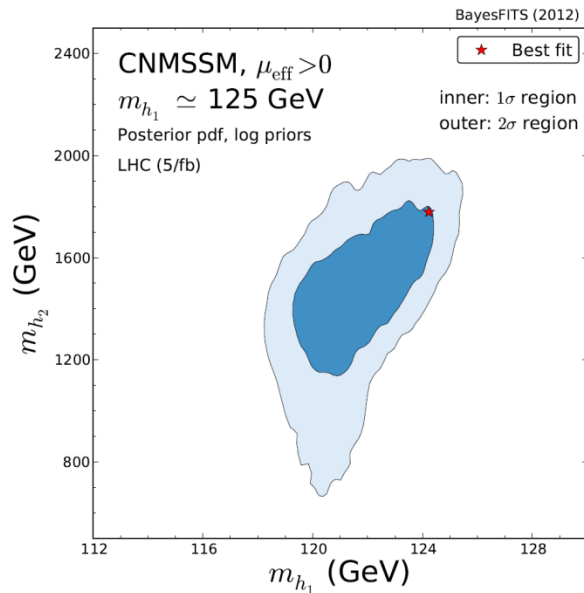
Kazakov et al

Lightest Higgs M_h 125 GeV (most of studies)

Heavy Higgs M_h 98 GeV, M_H 125 GeV (Drees, Belanger et al)

Denenerate Higgses (Gunion et al)

The aim is to make better overall χ^2 fit...



(Kowalska et al)

One of the motivation - Effective chiral Lagrangian from holographic viewpoint

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 + \frac{v^2}{4}\text{Tr}(D_\mu \Sigma^\dagger D^\mu \Sigma) \left[1 + 2a \frac{h}{v} \right]$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^{\bar{i}}) \Sigma \left[1 + c \frac{h}{v} \right] \begin{pmatrix} y_{ij}^u & u_R^j \\ y_{ij}^d & d_R^j \end{pmatrix} + h.c. + \dots$$

Crojean et al 2012

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \frac{1}{2}m_h^2 h^2 - \frac{d_3}{6} \left(\frac{3m_h^2}{v} \right) h^3 - \frac{d_4}{24} \left(\frac{3m_h^2}{v^2} \right) h^4 \dots$$

$$- \left(m_W^2 W_\mu W_\mu + \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right)$$

$$- \sum_{\psi=u,d,l} m_{\psi^{(i)}} \bar{\psi}^{(i)} \psi^{(i)} \left(1 + c_\psi \frac{h}{v} + c_{2\psi} \frac{h^2}{v^2} + \dots \right)$$

$$+ \frac{g^2}{16\pi^2} \left(c_{WW} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{ZZ} Z_{\mu\nu}^2 + c_{Z\gamma} Z_{\mu\nu} \gamma_{\mu\nu} \right) \frac{h}{v} + \dots$$

$$+ \frac{g^2}{16\pi^2} \left[\gamma_{\mu\nu}^2 \left(c_{\gamma\gamma} \frac{h}{v} + \dots \right) + G_{\mu\nu}^2 \left(c_{gg} \frac{h}{v} + c_{2gg} \frac{h^2}{v^2} \dots \right) \right]$$

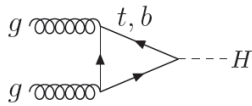
$$+ \frac{g^2}{16\pi^2} \left[\frac{c_{hhgg}}{\Lambda^2} G_{\mu\nu}^2 \frac{(\partial_\rho h)^2}{v^2} + \frac{c'_{hhgg}}{\Lambda^2} G_{\mu\rho} G_{\rho\nu} \frac{\partial_\mu h \partial_\nu h}{v^2} + \dots \right]$$

$$+ \dots$$

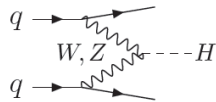
$$\mathcal{L} = \mathcal{L}_h - (M_W^2 W_\mu^+ W^{\mu-} + \frac{1}{2} M_Z^2 Z_\mu Z^\mu) [1 + 2a \frac{h}{v} + \mathcal{O}(h^2)] - m_{\psi_i} \bar{\psi}_i \psi_i [1 + c \frac{h}{v} + \mathcal{O}(h^2)] + \dots$$

Contino et al 2012,
....

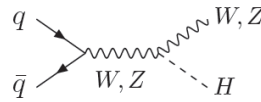
Espinosa et al 2012



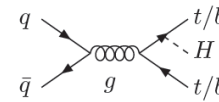
c^2



a^2



a^2



c^2

$$\frac{\sigma_{NLO}}{\sigma_{SM}^{NLO}}$$

$$\begin{aligned} \Gamma(H \rightarrow f\bar{f}) &= c^2 \Gamma^{SM}(H \rightarrow f\bar{f}), \\ \Gamma(H \rightarrow VV) &= a^2 \Gamma^{SM}(H \rightarrow VV), \\ \Gamma(H \rightarrow gg) &= c^2 \Gamma^{SM}(H \rightarrow gg), \\ \Gamma(H \rightarrow \gamma\gamma) &= \frac{(cI_\gamma + aJ_\gamma)^2}{(I_\gamma + J_\gamma)^2} \Gamma^{SM}(H \rightarrow \gamma\gamma) \end{aligned}$$

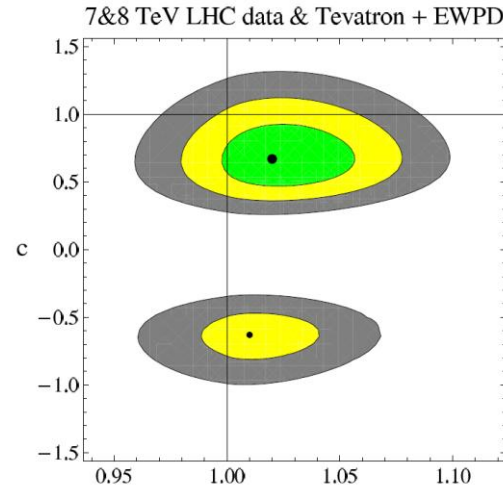
In SM $a=1$ and $c=1$

Higgs signal strength parameter:

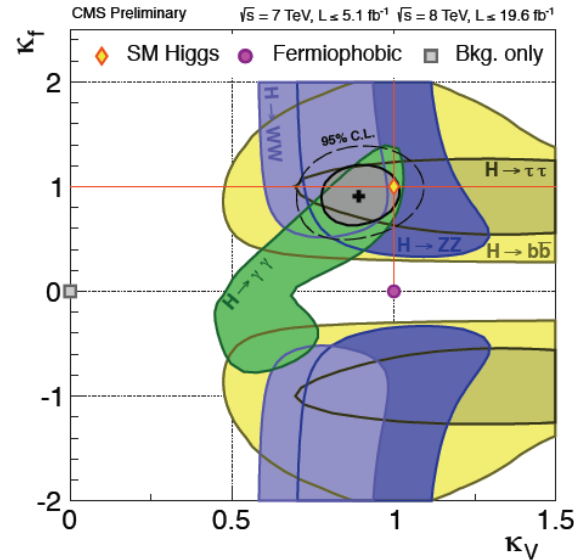
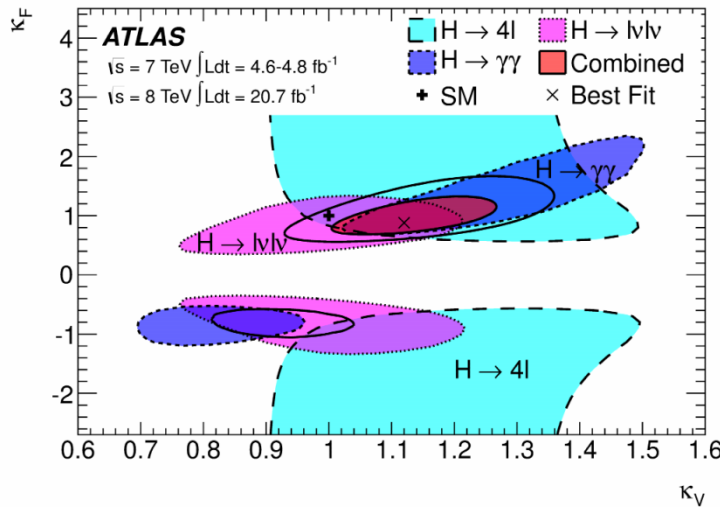
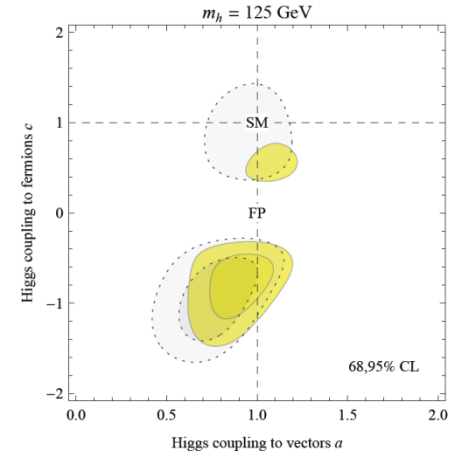
$$\mu_i = \frac{[\sigma_{j \rightarrow h} \times \text{Br}(h \rightarrow i)]_{\text{observed}}}{[\sigma_{j \rightarrow h} \times \text{Br}(h \rightarrow i)]_{\text{SM}}}$$

Global χ^2 fit
DO, CDF, CMS, ATLAS data

Espinosa et al 2012



Giardino et al 2012



All experiments compatible with SM predictions: accuracy $\sim 10\text{--}20\%$

- ATLAS: κ_V [1.05, 1.22] at 68% CL --- κ_f [0.76, 1.18] at 68% CL

- CMS: κ_V [0.74, 1.06] at 95% CL --- κ_f [0.61, 1.33] at 95% CL

Gerutti EPS 2013

$$\Gamma_H = \Gamma_{SM} + \Gamma_{BSM}$$

$$BR_{BSM} = \Gamma_{BSM} / \Gamma_H$$

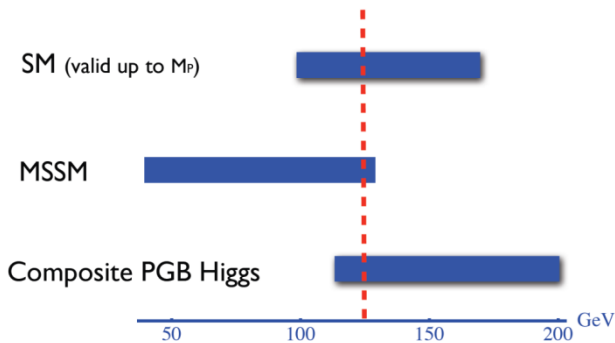
ATLAS: $BR_{BSM} < 0.60$ @ 95%CL (0.67exp.)

$\kappa_b = \kappa_W \dots = 1$ and 3 Fitted Parameters.: κ_γ κ_g BR_{BSM}

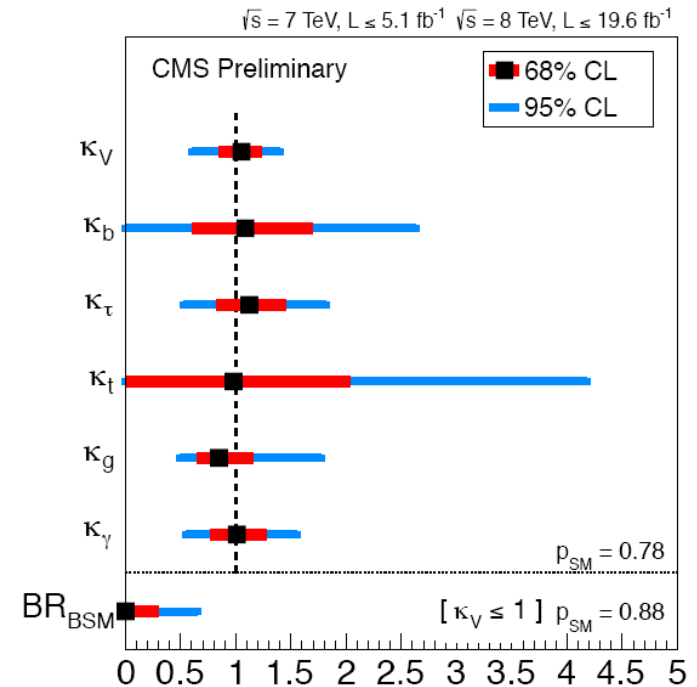
CMS: $BR_{BSM} < 0.64$ @ 95%CL (0.66exp.)

7 Fitted Parameters : κ_V κ_γ κ_g κ_t κ_τ κ_b BR_{BSM}

Different assumptions but similar limits



Pomarol 2012



BSM searches

Collision energy $>$ particle production threshold

-Searches for new particles

strongly interacting new particles with large cross sections (squarks, gluinos...)

top partners motivated by naturalness (stop, sbottom, vector like quarks, t^* ...)

new resonances predicted by many BSM extensions (Z' , W' , π_T , ρ_T , KK states, ..)

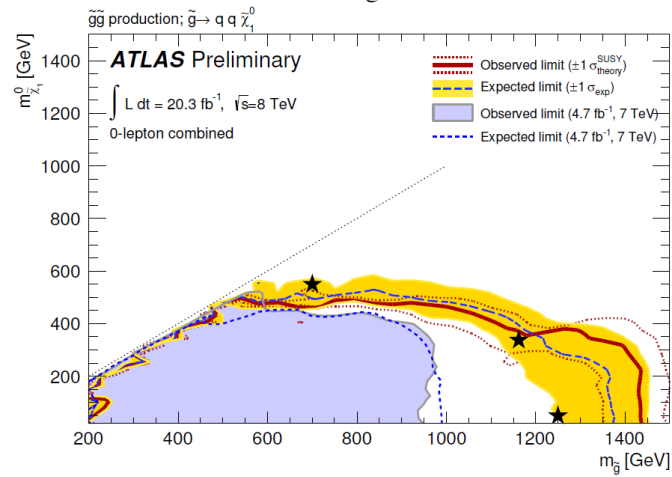
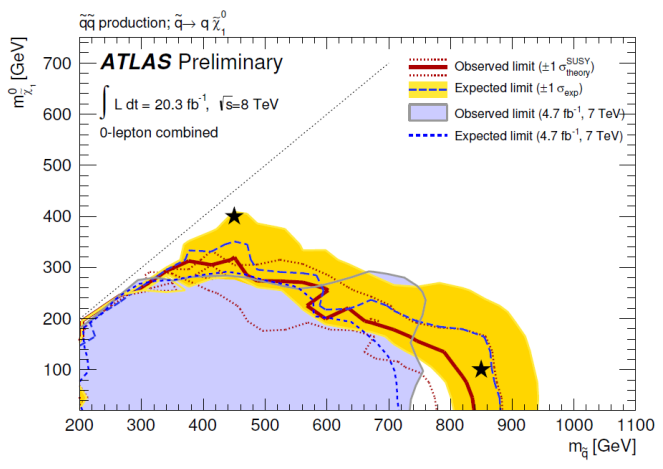
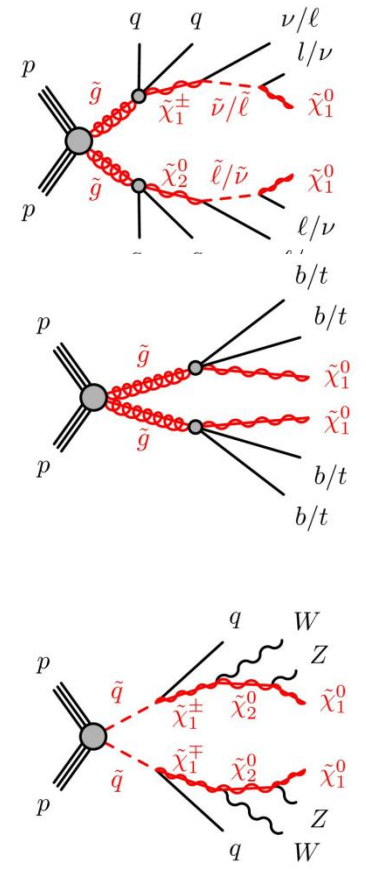
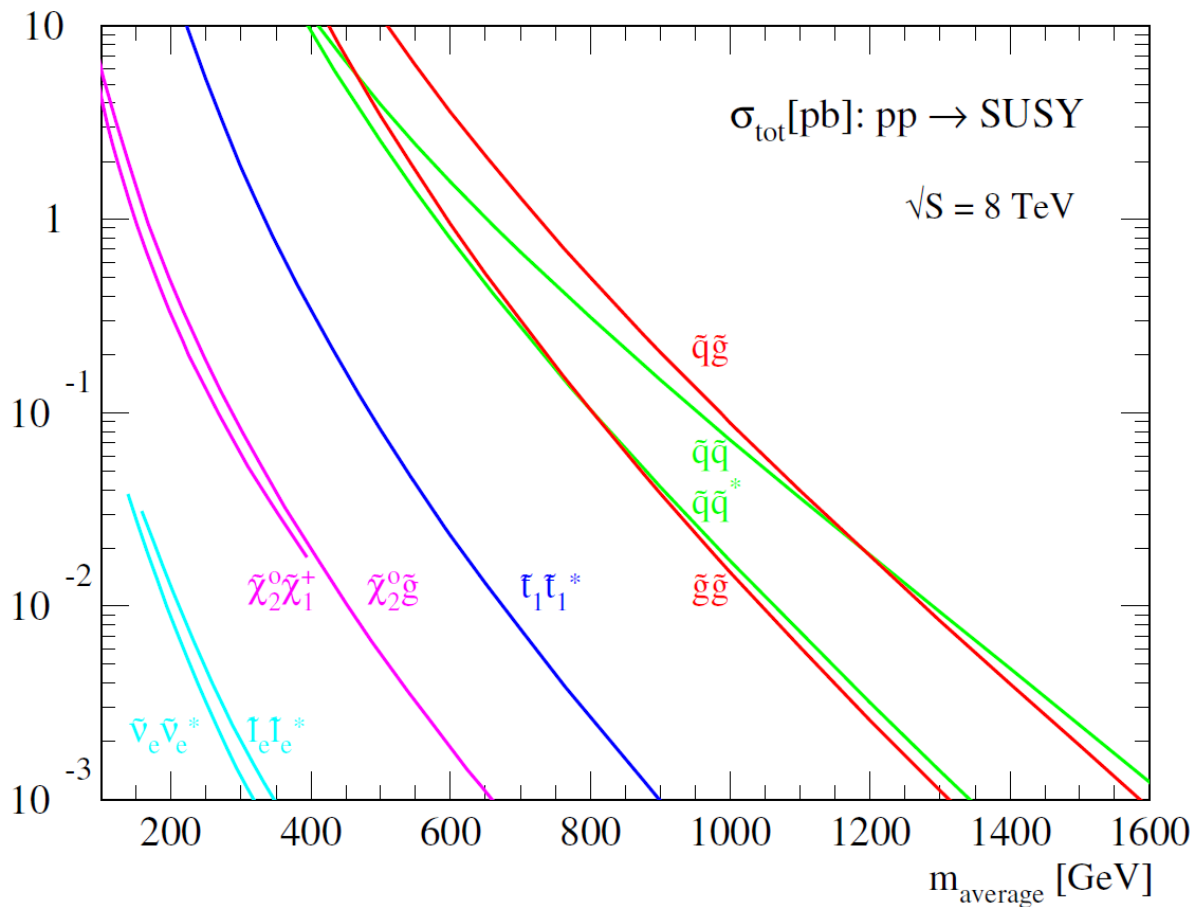
extended Higgs sector (new neutral Higgses, charged Higgs)

Collision energy $<$ particle production threshold

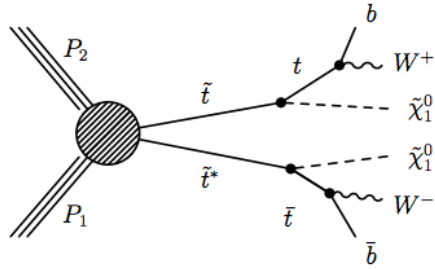
-Anomalous/new interactions of SM particles

(anom. gauge boson couplings, anom. Wtb couplings, FCNC ...)

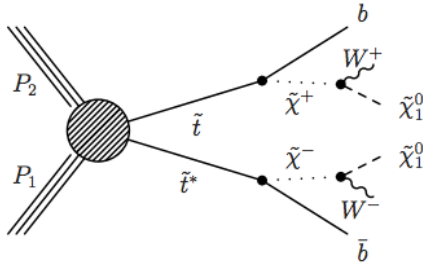
-New particle contributions via quantum loops



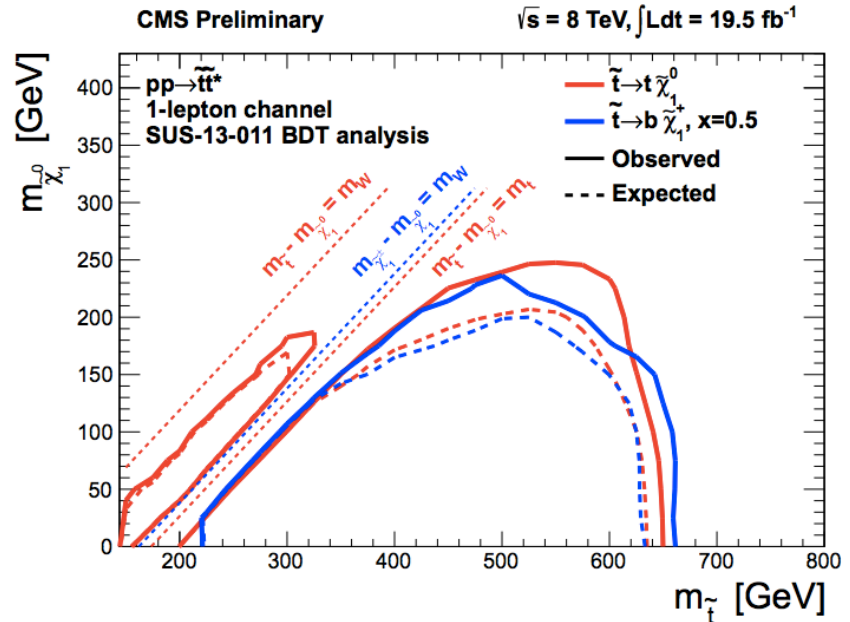
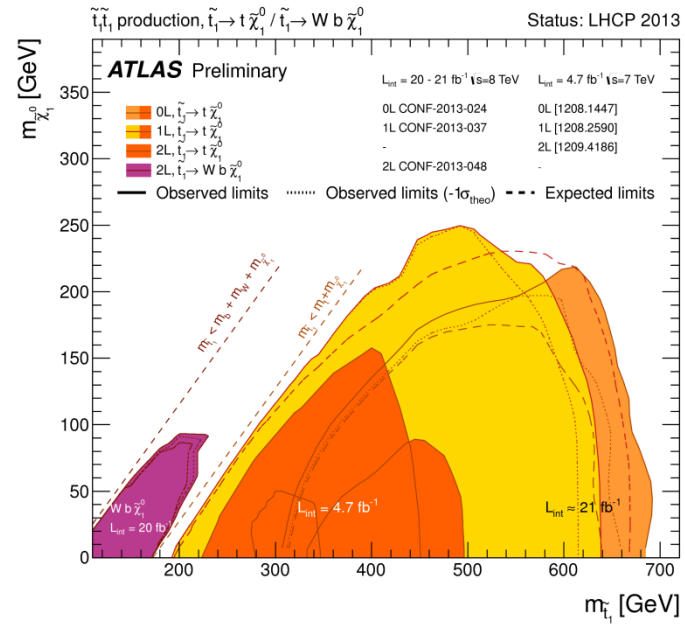
$$\tilde{t}\tilde{t} \rightarrow t\bar{t}\tilde{\chi}_1^0\tilde{\chi}_1^0$$



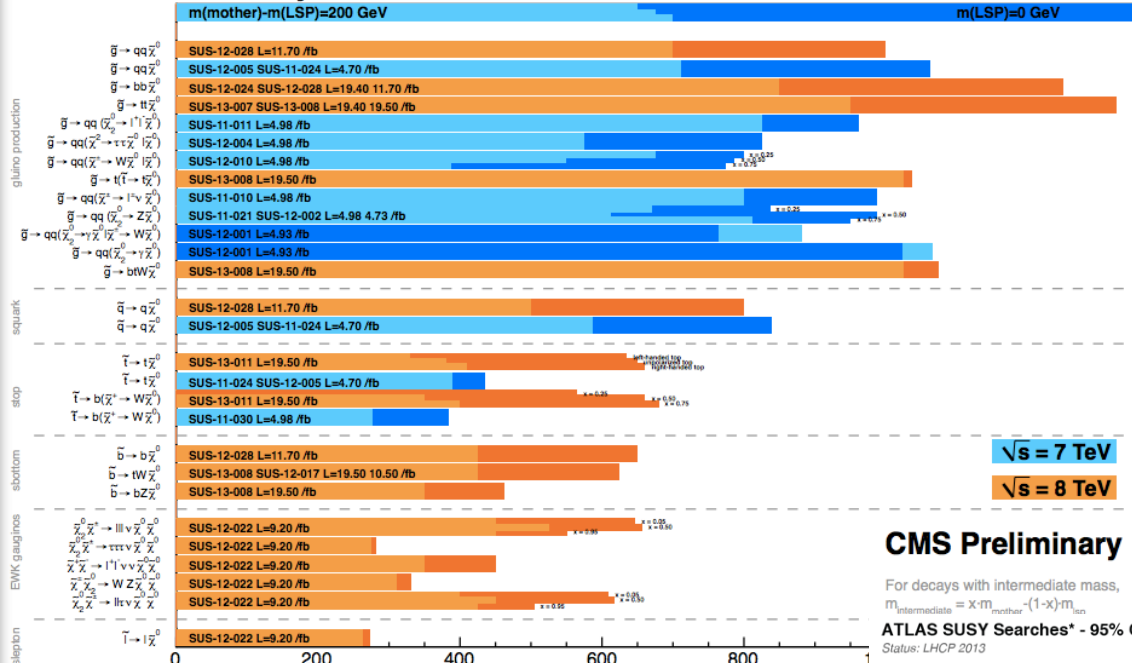
$$\tilde{t}\tilde{t} \rightarrow b\bar{b}\tilde{\chi}_1^+\tilde{\chi}_1^- \rightarrow b\bar{b}W^+W^-\tilde{\chi}_1^0\tilde{\chi}_1^0$$



Mass exclusion limits:
 $M_{stop} \sim 660 \text{ GeV}$ and
 $M_{bottom} \sim 630 \text{ GeV}$



Summary of CMS SUSY Results* in SMS framework LHC 13



CMS Preliminary

For decays with intermediate mass,
 $m_{intermediate} = x \cdot m_{mother} - (1-x) \cdot m_{LSP}$

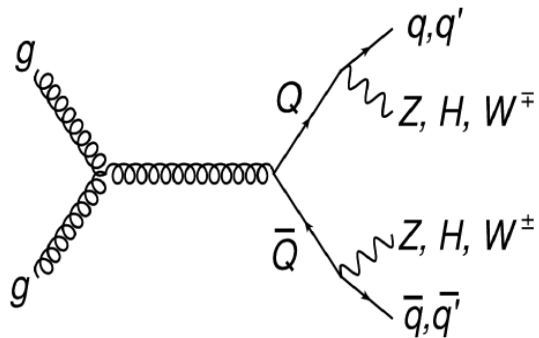
ATLAS SUSY Searches* - 95% CL Lower Limits
 Status: LHC 13

ATLAS Preliminary
 $\int L dt = (4.4 - 20.7) \text{ fb}^{-1}$ $\sqrt{s} = 7, 8 \text{ TeV}$

Model	e, μ, τ, γ	Jets	E_{miss}	[$L_{int} \text{ (fb}^{-1})$]	Mass limit	Reference	
Inclusive searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	$m_{\tilde{g}} - m_{\tilde{u}_L}$ 1.3 TeV	ATLAS-CONF-2013-047
	MSUGRA/CMSSM	1 e, μ	4 jets	Yes	5.8	$m_{\tilde{g}} - m_{\tilde{d}_L}$ 1.24 TeV	ATLAS-CONF-2012-104
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	any $m_{\tilde{g}}$ 1.1 TeV	ATLAS-CONF-2013-054
	$\tilde{g} \rightarrow q\bar{q}$	0	2-6 jets	Yes	20.3	$m_{\tilde{g}} > 0 \text{ GeV}$	ATLAS-CONF-2013-047
	Clusio med $\tilde{g} \rightarrow \tilde{g} - \tilde{g} \rightarrow \tilde{g} - \tilde{g}$	0	2-6 jets	Yes	20.3	$m_{\tilde{g}} > 0 \text{ GeV}$ 1.3 TeV	ATLAS-CONF-2013-047
	$\tilde{g} \rightarrow q\bar{q} + \tilde{g} \rightarrow \tilde{g} - \tilde{g}$	1 e, μ	2-4 jets	Yes	4.7	$m_{\tilde{g}} > 200 \text{ GeV}$ 600 GeV	1208-4688
	$\tilde{g} \rightarrow q\bar{q} + \tilde{g} \rightarrow \tilde{g} - \tilde{g}$	2 e, μ (SS)	3 jets	Yes	20.7	$m_{\tilde{g}} > 650 \text{ GeV}$ 1.1 TeV	ATLAS-CONF-2013-007
	GMSB (L NLSP)	2 e, μ	2-4 jets	Yes	4.7	$m_{\tilde{g}} > 124 \text{ TeV}$	1208-4688
	GMSB (I NLSP)	1-2 τ	0-2 jets	Yes	20.7	$m_{\tilde{g}} > 150 \text{ GeV}$ 1.4 TeV	ATLAS-CONF-2013-026
	GGM (bino NLSP)	2 τ	0	Yes	4.8	$m_{\tilde{g}} > 50 \text{ GeV}$ 1.67 TeV	1209-0753
3 σ gen. squarks direct production	$\tilde{g} \rightarrow q\bar{q}$	0	3b	Yes	12.8	$m_{\tilde{g}} > 200 \text{ GeV}$ 1.24 TeV	ATLAS-CONF-2012-145
	$\tilde{g} \rightarrow q\bar{q}$	2 e, μ (SS)	0-3b	No	20.7	$m_{\tilde{g}} > 100 \text{ GeV}$ 600 GeV	ATLAS-CONF-2013-007
	$\tilde{g} \rightarrow q\bar{q}$	0	7-10 jets	Yes	20.3	$m_{\tilde{g}} > 200 \text{ GeV}$ 1.14 TeV	ATLAS-CONF-2013-054
	$\tilde{g} \rightarrow q\bar{q}$	0	3b	Yes	12.8	$m_{\tilde{g}} > 200 \text{ GeV}$ 1.18 TeV	ATLAS-CONF-2012-145
	$b, \tilde{b}_1 \rightarrow b\bar{b}$	0	2b	Yes	20.1	$m_{\tilde{b}_1} < 100 \text{ GeV}$	ATLAS-CONF-2013-053
	$b, \tilde{b}_2 \rightarrow b\bar{b}$	2 e, μ (SS)	0-3b	Yes	20.7	$m_{\tilde{b}_2} > 2 m_{\tilde{g}} > 430 \text{ GeV}$	ATLAS-CONF-2013-049
	$\tilde{t}_1 \rightarrow t\bar{t}$	1-2 e, μ	1-2b	Yes	4.7	$m_{\tilde{t}_1} > 95 \text{ GeV}$	1208-4305, 1209-2102
	$\tilde{t}_2 \rightarrow t\bar{t}$	2 e, μ	0-2 jets	Yes	20.3	$m_{\tilde{t}_2} = m_{\tilde{t}_1} - m_{W^{\pm}}$ 50 GeV, $m_{\tilde{t}_2} \ll m_{\tilde{g}}$	ATLAS-CONF-2013-048
	$\tilde{t}_1 \rightarrow t\bar{t}$	2 e, μ	0-2 jets	Yes	20.3	$m_{\tilde{t}_1} > 0 \text{ GeV}$, $m_{\tilde{t}_2} m_{\tilde{g}} > 10 \text{ GeV}$	ATLAS-CONF-2013-048
	$\tilde{t}_1 \rightarrow t\bar{t}$	0	2b	Yes	20.1	$m_{\tilde{t}_1} > 200 \text{ GeV}$, $m_{\tilde{t}_2} m_{\tilde{g}} > 5 \text{ GeV}$	ATLAS-CONF-2013-053
EW direct	$\tilde{t}_1 \rightarrow t\bar{t}$	1 e, μ	1b	Yes	20.7	$m_{\tilde{t}_1} > 0 \text{ GeV}$ 206-810 GeV	ATLAS-CONF-2013-037
	$\tilde{t}_2 \rightarrow t\bar{t}$	0	2b	Yes	20.5	$m_{\tilde{t}_2} > 0 \text{ GeV}$ 326-566 GeV	ATLAS-CONF-2013-004
	$\tilde{t}_1 \rightarrow t\bar{t}$	2 e, μ (Z)	1b	Yes	20.7	$m_{\tilde{t}_1} > 150 \text{ GeV}$ 500 GeV	ATLAS-CONF-2013-025
	$\tilde{t}_2 \rightarrow t\bar{t}$	3 e, μ (Z)	1b	Yes	20.7	$m_{\tilde{t}_2} = m_{\tilde{t}_1} + 180 \text{ GeV}$ 320 GeV	ATLAS-CONF-2013-025
	$\tilde{t}_1 \rightarrow t\bar{t}$	2 e, μ	0	Yes	20.3	$m_{\tilde{t}_1} > 0 \text{ GeV}$ 55-315 GeV	ATLAS-CONF-2013-049
	$\tilde{t}_2 \rightarrow t\bar{t}$	2 e, μ	0	Yes	20.3	$m_{\tilde{t}_2} > 0 \text{ GeV}$, $m_{\tilde{t}_1} < 0.5 m_{\tilde{g}} + m_{\tilde{g}}$ 125-480 GeV	ATLAS-CONF-2013-049
	$\tilde{t}_1 \rightarrow t\bar{t}$	2 e, μ	0	Yes	20.7	$m_{\tilde{t}_1} > 0 \text{ GeV}$, $m_{\tilde{t}_2} < 0.5 m_{\tilde{g}} + m_{\tilde{g}}$ 180-330 GeV	ATLAS-CONF-2013-026
	$\tilde{t}_2 \rightarrow t\bar{t}$	3 e, μ	0	Yes	20.7	$m_{\tilde{t}_2} > 0 \text{ GeV}$, $m_{\tilde{t}_1} < 0.5 m_{\tilde{g}} + m_{\tilde{g}}$ 326-566 GeV	ATLAS-CONF-2013-004
	$\tilde{t}_1 \rightarrow t\bar{t}$	3 e, μ	0	Yes	20.7	$m_{\tilde{t}_1} = m_{\tilde{t}_2}$, $m_{\tilde{t}_2} < 0.5 m_{\tilde{g}} + m_{\tilde{g}}$ 315 GeV	ATLAS-CONF-2013-035
	Long-lived particles	Direct $\tilde{g} \rightarrow q\bar{q}$ prod. long-lived \tilde{Z}_1	0	1 jet	Yes	4.7	$m_{\tilde{Z}_1} > 230 \text{ GeV}$
Stable \tilde{g} , R-hadrons		0-2 e, μ	0	Yes	4.7	$m_{\tilde{g}} > 985 \text{ GeV}$	1211-1597
GMSB, stable \tilde{Z}_1 low β		2 e, μ	0	Yes	4.7	$5 < \text{tamp} < 20$	1211-1597
GMSB, $\tilde{Z}_1 \rightarrow G$ long-lived \tilde{Z}_1		2 τ	0	Yes	4.7	$0.4 < \langle \tau \rangle < 2 \text{ ns}$	1304-6310
$\tilde{Z}_1 \rightarrow q\bar{q}$ (RPV)		1 e, μ	0	Yes	4.4	$1 \text{ ms} < \langle \tau \rangle < 1 \text{ s}$, \tilde{Z}_1 decoupled	1210-7451
LFV $\tilde{g} \rightarrow q\bar{q} + X, \tilde{Z}_1 \rightarrow e\bar{e}$		2 e, μ	0	-	4.6	$\lambda_{12} > 10^{-10}, \lambda_{13} > 0.05$	1212-1272
LFV $\tilde{g} \rightarrow q\bar{q} + X, \tilde{Z}_1 \rightarrow e\bar{e} + e$		1 e, $\mu + e$	0	-	4.6	$\lambda_{12} > 10^{-10}, \lambda_{13} > 0.05$	1212-1272
Blinear RPV CMSSM		1 e, μ	7 jets	Yes	4.7	$m_{\tilde{g}} = m_{\tilde{b}_1}, c_{1\mu} < 1 \text{ mm}$	ATLAS-CONF-2012-140
$\tilde{Z}_1 \rightarrow t\bar{t} - W\tilde{Z}_1, \tilde{Z}_1 \rightarrow e\bar{e} + \tilde{g} + \nu_{\tau}$		4 e, μ	0	Yes	20.7	$m_{\tilde{Z}_1} > 300 \text{ GeV}, \lambda_{1\mu} > 0$	ATLAS-CONF-2013-036
$\tilde{Z}_1 \rightarrow t\bar{t} - W\tilde{Z}_1, \tilde{Z}_1 \rightarrow e\bar{e} + \nu_{\tau} + \nu_{\tau}$		3 e, $\mu + \tau$	0	Yes	20.7	$m_{\tilde{Z}_1} > 80 \text{ GeV}, \lambda_{1\mu} > 0$	ATLAS-CONF-2013-036
RPV	$\tilde{g} \rightarrow q\bar{q}$	0	6 jets	-	4.6	$m_{\tilde{g}} > 694 \text{ GeV}$	1210-4813
	$\tilde{g} \rightarrow q\bar{q}$	2 e, μ (SS)	0-3b	Yes	20.7	$m_{\tilde{g}} > 880 \text{ GeV}$	ATLAS-CONF-2013-007
	Scalar gluon	0	4 jets	-	4.6	sgluon 100-287 GeV	1210-4826
Other	WIMP interaction (D5, Dirac χ)	0	mono-jet	Yes	10.5	M scale 704 GeV	ATLAS-CONF-2012-147
						incl. limit from 1110.2690 $m_{\tilde{g}} > 80 \text{ GeV}$, limit of $< 687 \text{ GeV}$ for D8	

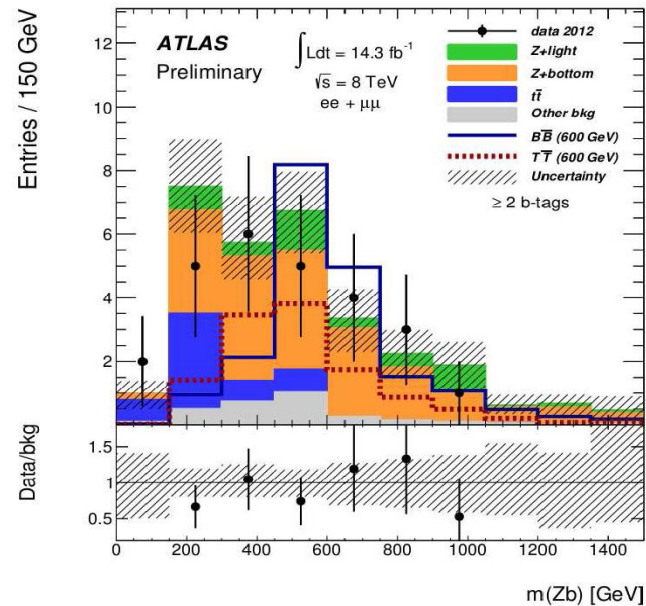
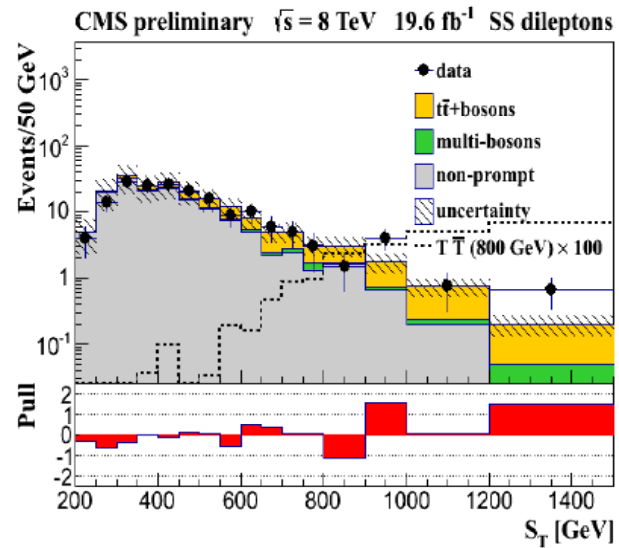
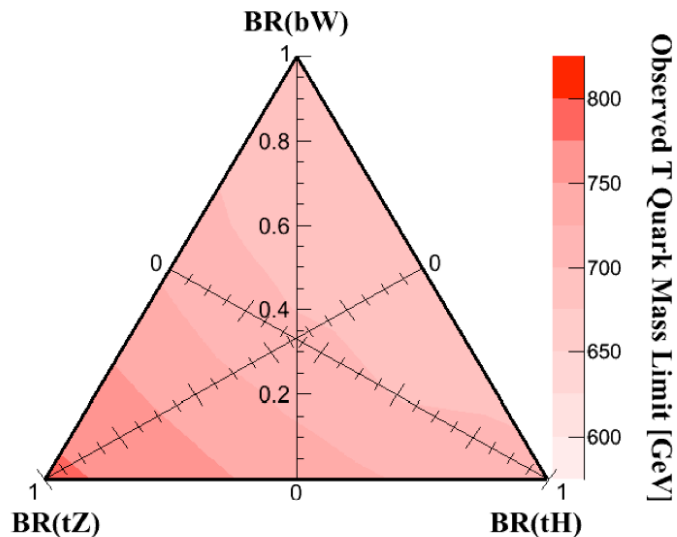
*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Vector like top partners

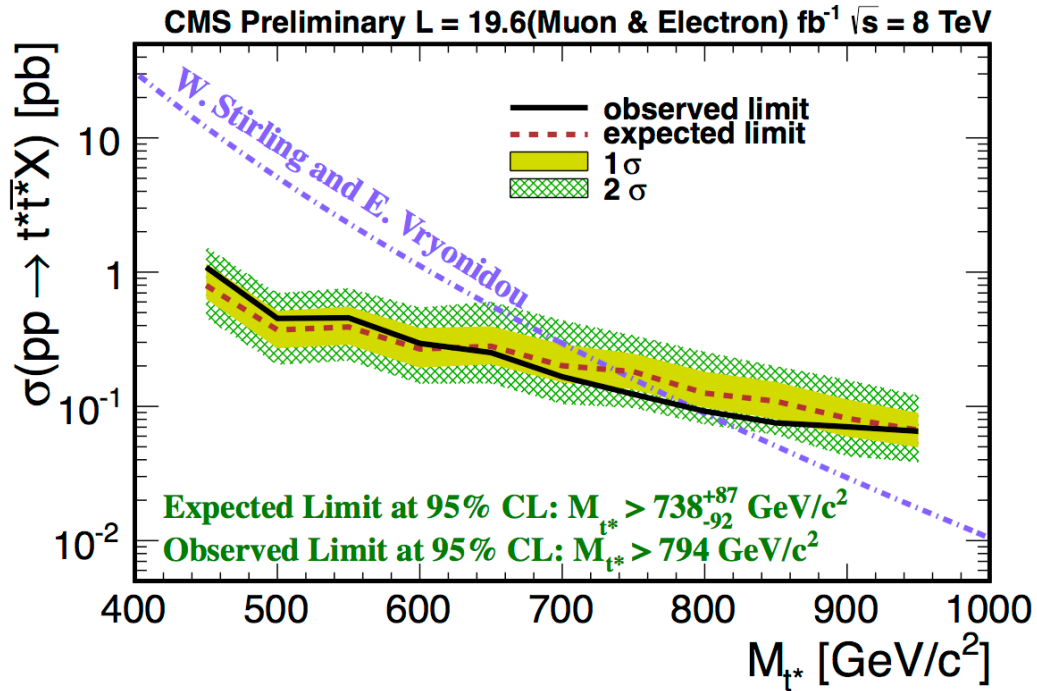


Mass limits ~ 700 GeV
 ($T \rightarrow Zt$, $T \rightarrow Wb$, $T \rightarrow Ht$)

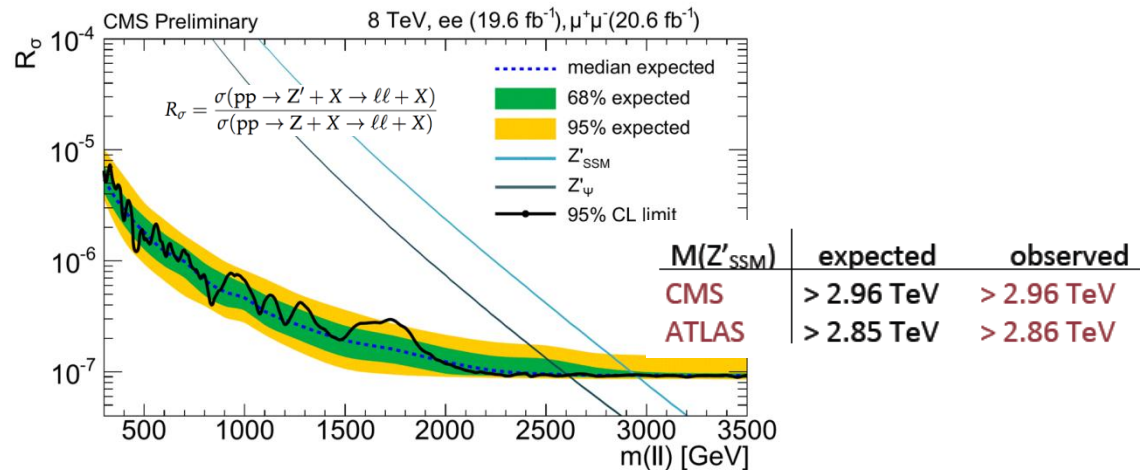
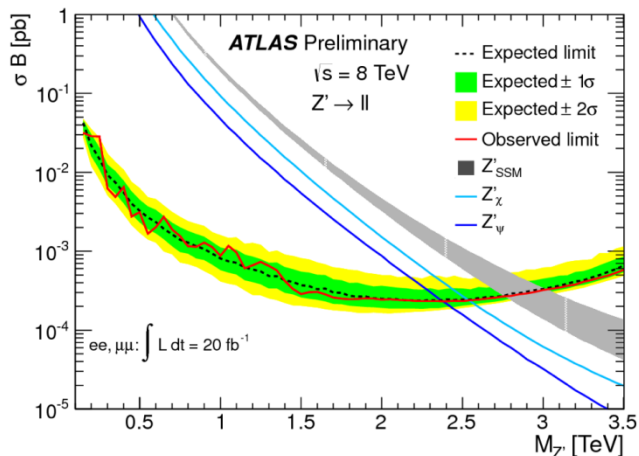
CMS preliminary $\sqrt{s} = 8$ TeV 19.6 fb⁻¹



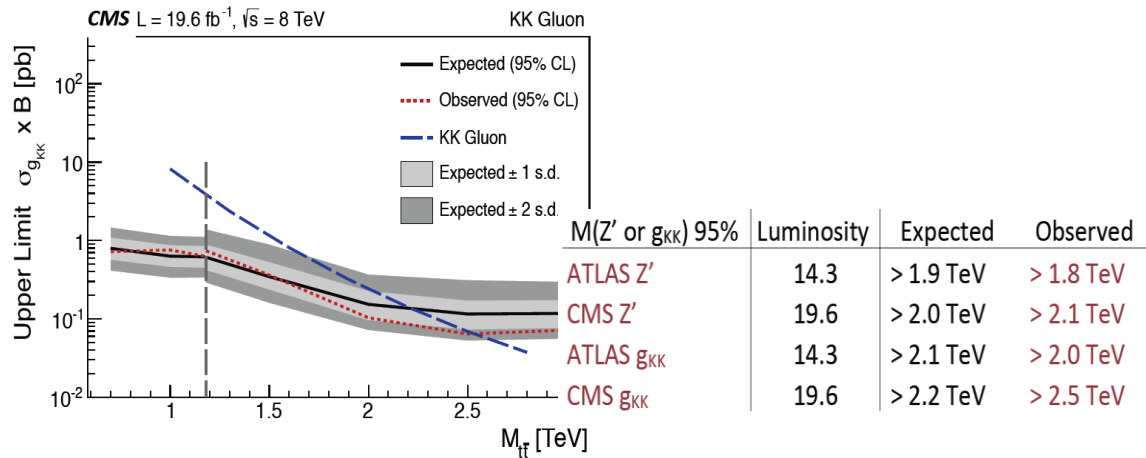
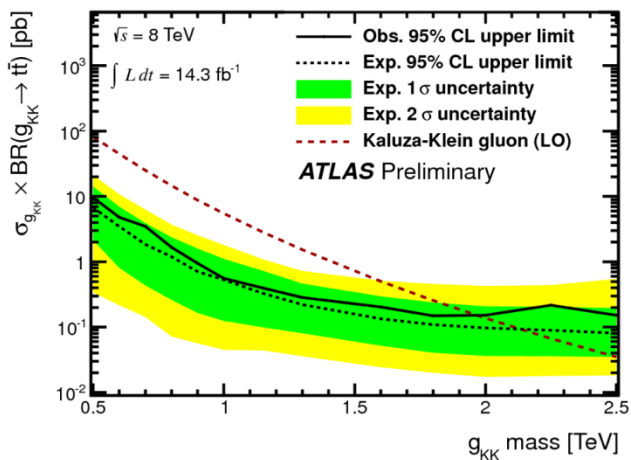
Searches for excited Top in top+jet



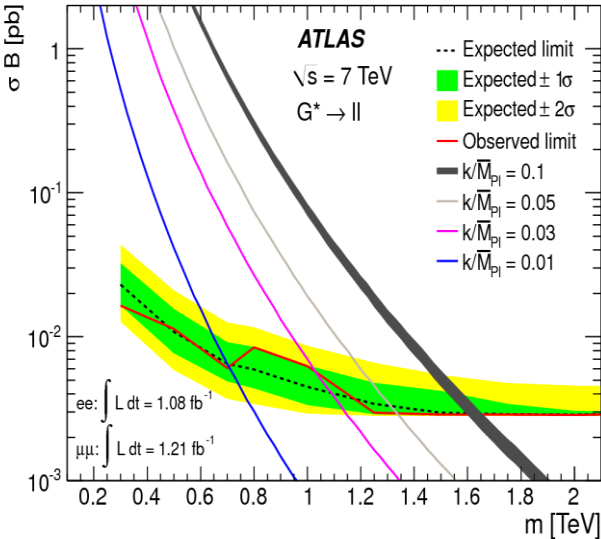
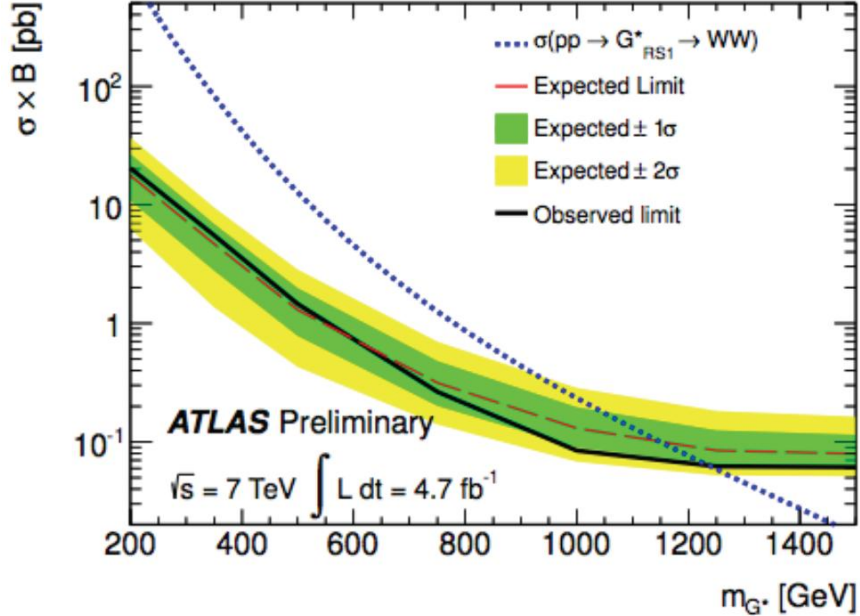
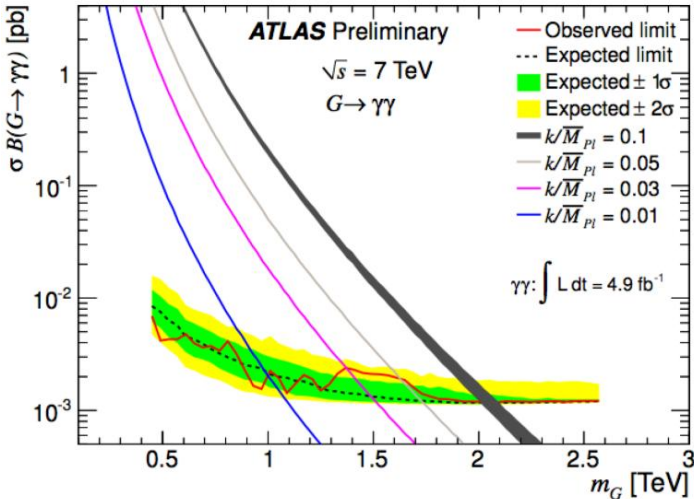
Searches for Z' in dileptons



Searches for Z' and KK resonances in top pair

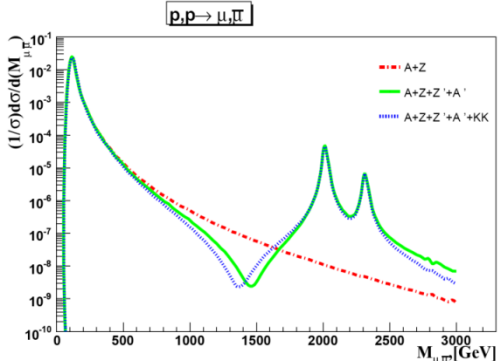


Searches for RS gravitons

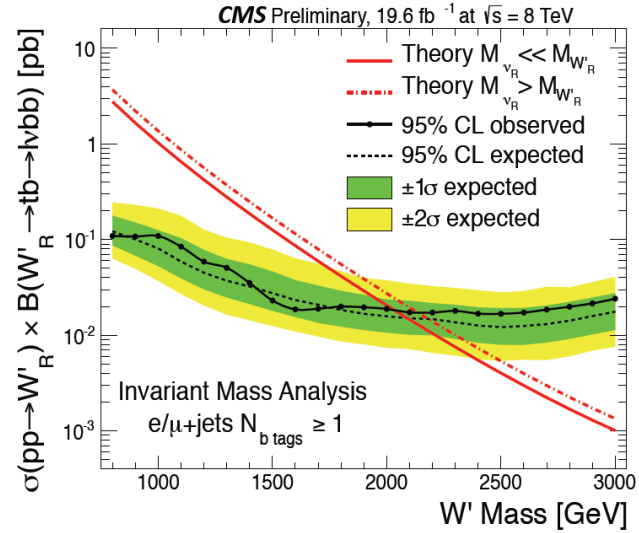
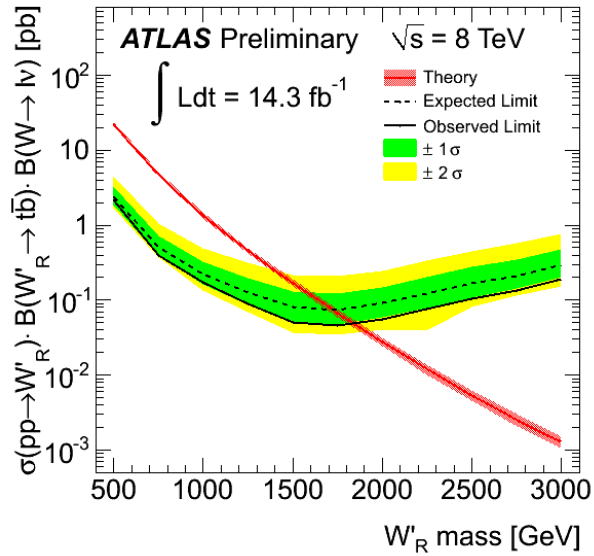
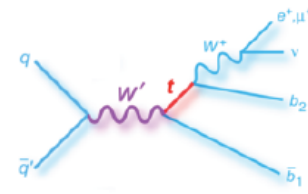


No interferences yet
 The interferences
 should be included

Boos, Bunichev, Smolyakov, Volobuev



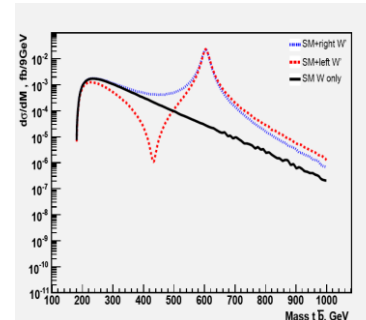
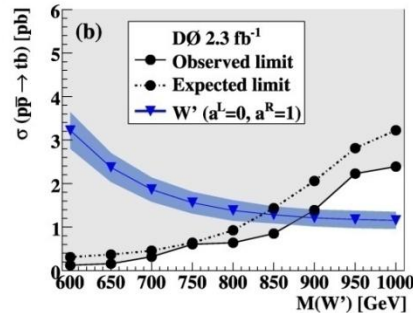
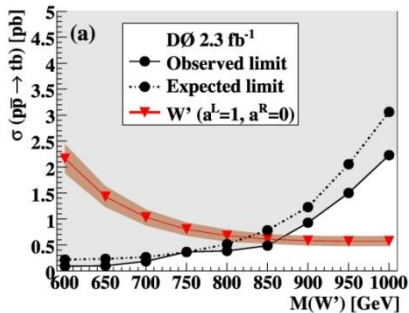
Searches for W' in top+b



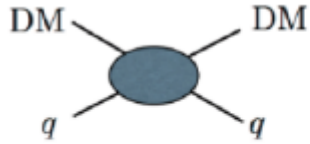
$M_{W'_R}(tb)$ 95% CL	Luminosity	Expected	Observed
ATLAS W'_R	14.3	> 1.72 TeV	> 1.84 TeV
CMS W'_R	19.6	> 2.09 TeV	> 2.03 TeV

D0 limits: $M_{W'} > 830$ (860) GeV L(R)

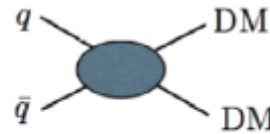
Boos, Bunichev, Dudko, Perfilov



Negative interference

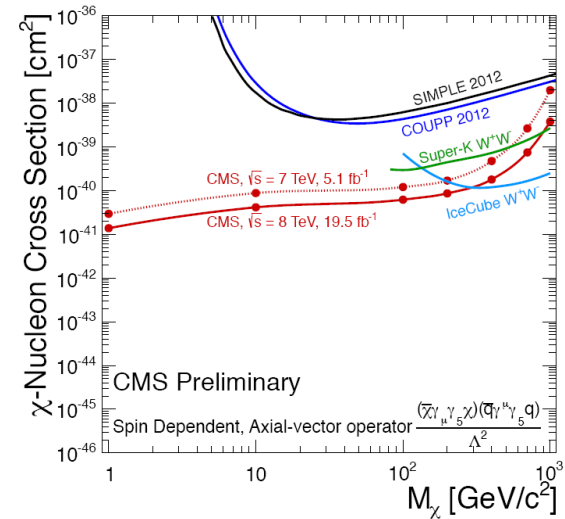
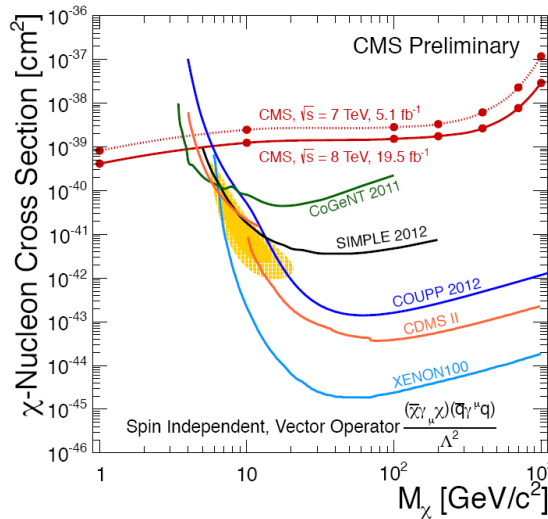
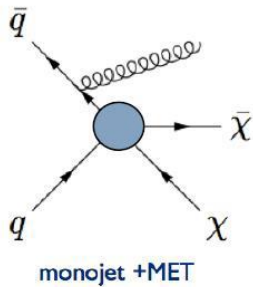


Direct Detection (t-channel)



Collider Searches (s-channel)

Dark Matter searches in monojets

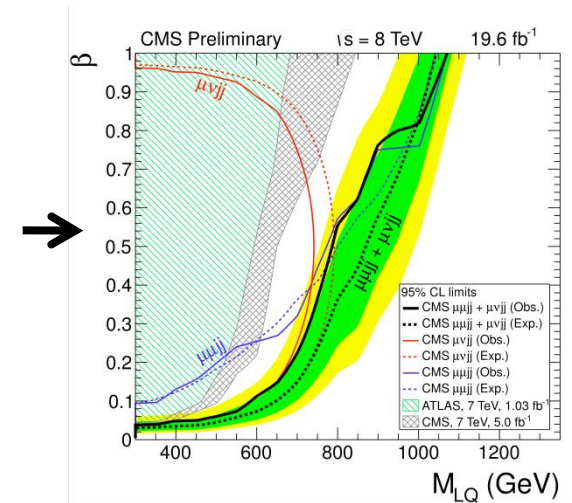
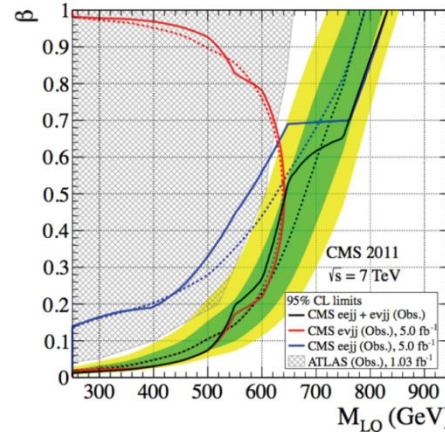
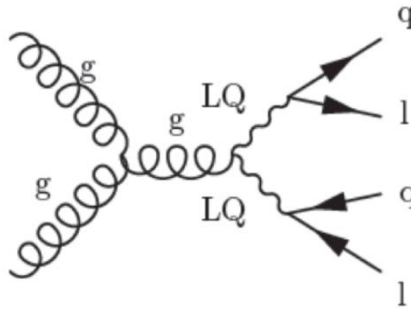


Leptoquark searches

LQs are predicted by composite models, GUT ...

$$\beta = \text{BR}(\text{LQ} \rightarrow lq)$$

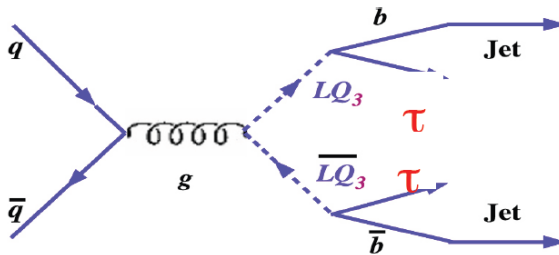
1st and 2d generation LQs



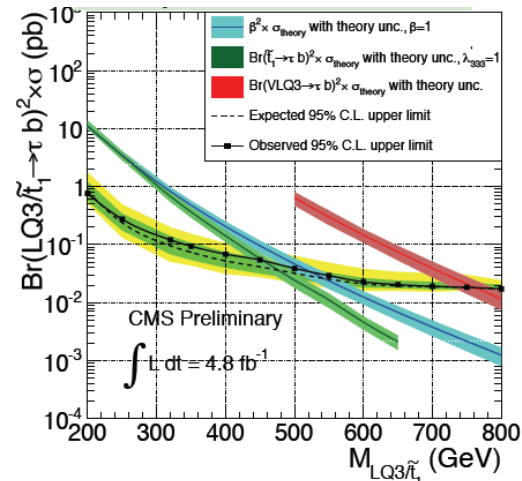
$M_{LQ1} > 830$ (640) GeV for $\beta=1$ (0.5)
 $M_{LQ2} > 840$ (650) GeV for $\beta=1$ (0.5)

→ 1070 (785) GeV excluded for $\beta=1$ (0.5)

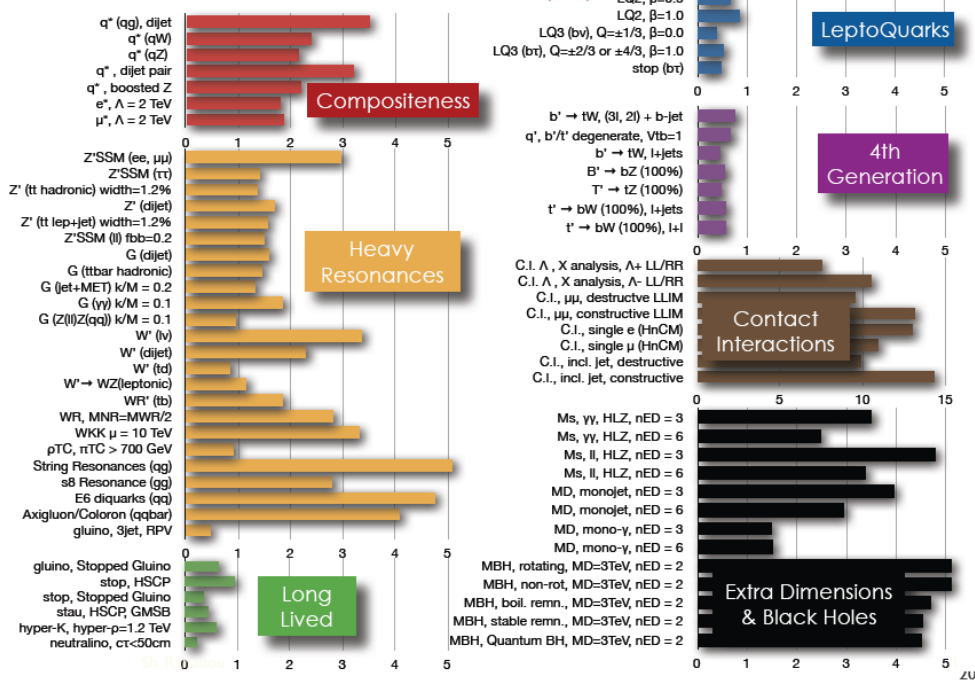
3d generation LQs



$M_{SLQ3} > 525$ GeV for $\beta=1$
 $M_{VLQ3} > 763$ GeV for $\beta=1$



CMS EXOTICA 95% CL EXCLUSION LIMITS (TeV)

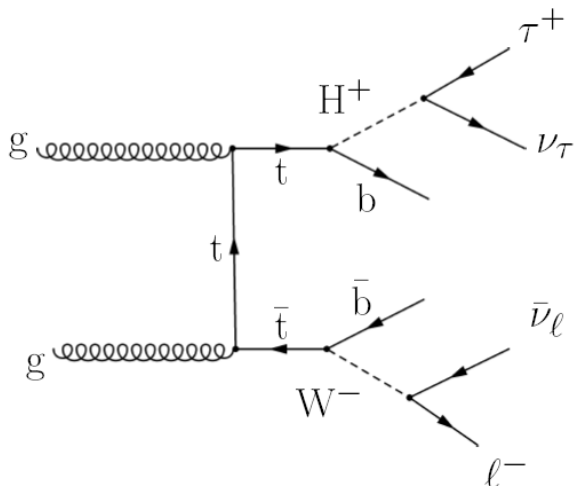


"There are more things in heaven and earth, Horatio, than are dreamt of in your philosophy." -- Hamlet

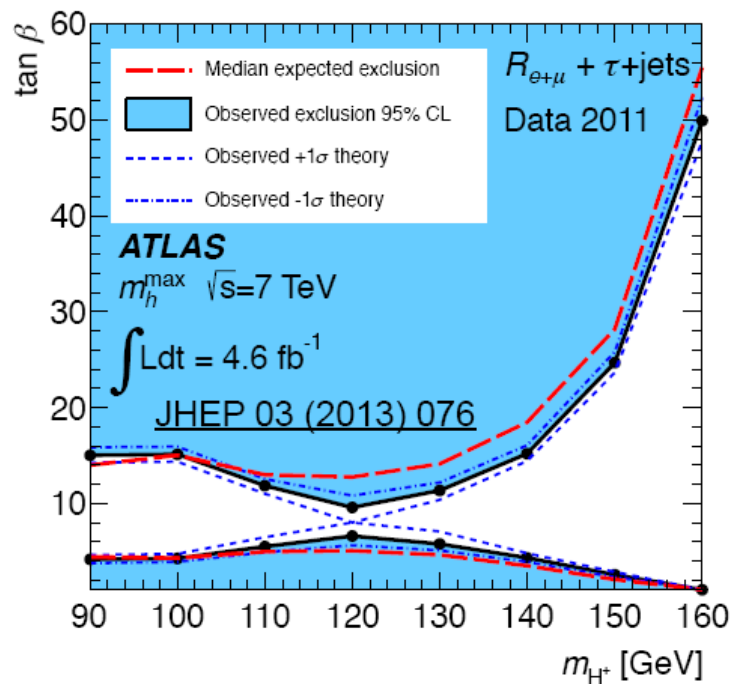
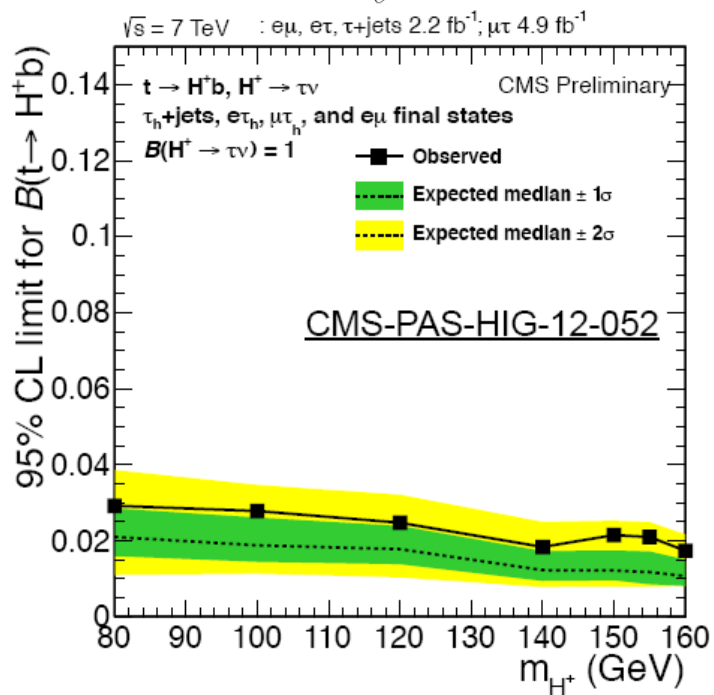
ATLAS Exotics Searches* - 95% CL Lower Limits (Status: May 2013)			
Search	Reference	Lower Limit (TeV)	Notes
Large ED (ADD) : monojet + E_T^{miss}	ATLAS-CONF-2012-049	4.37 TeV	M_{*2} (n=2)
Large ED (ADD) : monophoton + E_T^{miss}	ATLAS-CONF-2012-075	1.92 TeV	M_{*2} (n=2)
Large ED (ADD) : diphoton + dilepton	ATLAS-CONF-2012-116	4.18 TeV	M_{*2} (HLZ S=3, NLO)
UED : diphoton + E_T^{miss}	ATLAS-CONF-2012-075	1.68 TeV	Compact, scale R^{-1}
S/Z, ED : dilepton	ATLAS-CONF-2012-253	4.71 TeV	$M_{*2} - R^{-1}$
RS1 : dilepton	ATLAS-CONF-2013-017	2.47 TeV	Graviton mass ($k/M_{*2} = 0.1$)
RS1 : WW resonance	ATLAS-CONF-2012-049	1.28 TeV	Graviton mass ($k/M_{*2} = 0.1$)
Bulk RS ZZ resonance	ATLAS-CONF-2012-130	830 GeV	Graviton mass ($k/M_{*2} = 1.0$)
UED : diphoton + jets	ATLAS-CONF-2012-075	1.68 TeV	g_{*2} mass
ADD BH ($M_{*2}/M_{*3}=3$) : SS dimuon	ATLAS-CONF-2012-049	1.25 TeV	M_{*2} (S=6)
ADD BH ($M_{*2}/M_{*3}=3$) : leptons + jets	ATLAS-CONF-2012-049	1.57 TeV	M_{*2} (S=6)
Quantum black hole : dijet	ATLAS-CONF-2012-171	4.11 TeV	M_{*2} (S=6)
qqqq contact interaction : ll	ATLAS-CONF-2012-171	7.8 TeV	Λ
qqll CI : ee & $\mu\mu$	ATLAS-CONF-2012-116	13.9 TeV	Λ (constructive int.)
uuft CI : SS dilepton + jets + E_T^{miss}	ATLAS-CONF-2012-130	3.11 TeV	Λ (C=1)
Z' (SSM) : $m_{Z'}$	ATLAS-CONF-2012-121	2.60 TeV	Z' mass
Z' (SSM) : $m_{Z'}$	ATLAS-CONF-2012-121	1.4 TeV	Z' mass
Z' (leptophobic topcolor) : $ll \rightarrow l$ -jets	ATLAS-CONF-2012-130	1.4 TeV	Z' mass
W' (SSM) : $m_{W'}$	ATLAS-CONF-2012-049	2.55 TeV	W' mass
W' ($\rightarrow tq, g=1$) : $m_{W'}$	ATLAS-CONF-2012-049	439 GeV	W' mass
W' ($\rightarrow tb, LRSM$) : $m_{W'}$	ATLAS-CONF-2012-049	1.84 TeV	W' mass
Scalar LQ pair ($\beta=1$) : kin. vars. in eejj, evjj	ATLAS-CONF-2012-049	680 GeV	1 st gen. LQ mass
Scalar LQ pair ($\beta=1$) : kin. vars. in $\mu\mu jj, \nu\nu jj$	ATLAS-CONF-2012-049	680 GeV	2 nd gen. LQ mass
Scalar LQ pair ($\beta=1$) : kin. vars. in $\tau\tau jj, \nu\tau jj$	ATLAS-CONF-2012-049	554 GeV	3 rd gen. LQ mass
4 th generation : $ll \rightarrow l$ -jets	ATLAS-CONF-2012-049	650 GeV	t^* mass
4 th generation : $ll \rightarrow l$ -jets	ATLAS-CONF-2012-049	720 GeV	b^* mass
Vector-like quark : $TT \rightarrow H+X$	ATLAS-CONF-2012-019	790 GeV	T mass (isospin doublet)
Vector-like quark : CC, m_{*2}	ATLAS-CONF-2012-019	1.12 TeV	VLQ mass (charge = -1/3, coupling $\kappa_{q2} = \sqrt{1/3}$)
Excited quarks : γ -jet resonance	ATLAS-CONF-2012-049	2.46 TeV	q^* mass
Excited quarks : dijet resonance	ATLAS-CONF-2012-148	3.84 TeV	q^* mass
Excited b quark : W-t resonance	ATLAS-CONF-2012-148	870 GeV	b^* mass (left-handed coupling)
Excited leptons : γ -jet resonance	ATLAS-CONF-2012-148	2.2 TeV	l^* mass
Techni-hadrons (LSTC) : dilepton	ATLAS-CONF-2012-049	890 GeV	ρ_{*2} mass ($m_{\rho_{*2}}/m_{*2} = m(\rho_{*2}) = M_{*2}$)
Techni-hadrons (LSTC) : WZ resonance	ATLAS-CONF-2012-049	920 GeV	ρ_{*2} mass ($m_{\rho_{*2}}/m_{*2} = m(\rho_{*2}) + m_{*2} = 1.1m(\rho_{*2})$)
Major neutr. (LRSM, no mixing) : 2-lag + jets	ATLAS-CONF-2012-049	1.6 TeV	N mass ($ V_{*2} = 2$ TeV)
Heavy lepton N^c (type III seesaw) : Z-l resonance	ATLAS-CONF-2012-049	1.05 TeV	N^c mass ($ V_{*2} = 0.055, V_{*2} = 0.063, V_{*2} = 0$)
H_c^{\pm} (DY prod., BR($H^{\pm} \rightarrow ll$)) : SS ee ($\mu\mu$)	ATLAS-CONF-2012-049	400 GeV	H_c^{\pm} mass (limit at 300 GeV for $\mu\mu$)
Color octet scalar : dijet resonance	ATLAS-CONF-2012-049	1.80 TeV	Scalar resonance mass
Multi-charged particles (DY prod.) : highly ionizing tracks	ATLAS-CONF-2012-049	460 GeV	mass ($ \eta = 4e$)
Magnetic monopoles (DY prod.) : highly ionizing tracks	ATLAS-CONF-2012-049	187 GeV	mass

ATLAS Preliminary
 $\Lambda_{*2} = (1-20) \text{ fb}^{-1}$
 $\sqrt{s} = 7, 8 \text{ TeV}$

Charged Higgs searches



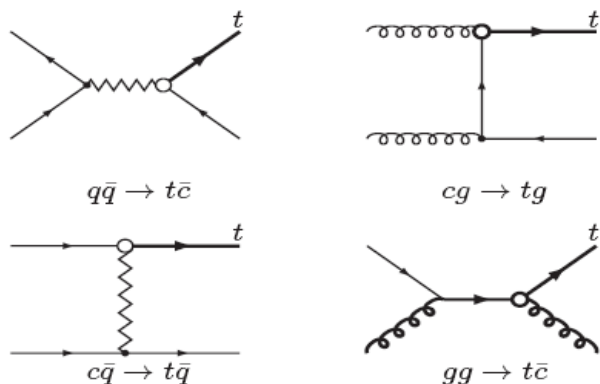
For $m(H^+) < m(t) - m(b)$
 the decay $t \rightarrow H^+ b$ is allowed,
 $H^+ \rightarrow \tau \nu$ dominates



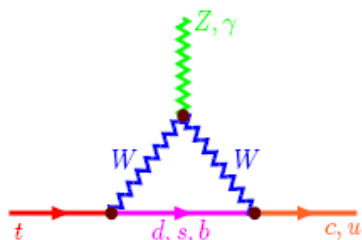
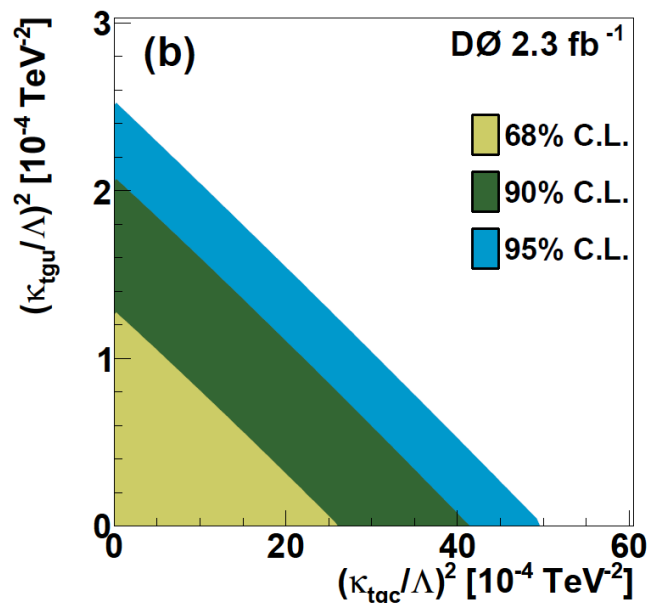
FCNC anomalous top couplings

- Couplings: tqg , $tq\gamma$, tqZ , where $q = u, c$

$$\Delta\mathcal{L}^{eff} = \frac{1}{\Lambda} [\kappa_{tq}^{\gamma,Z} e\bar{t}\sigma_{\mu\nu}qF_{\gamma,Z}^{\mu\nu} + \kappa_{tq}^g g_s\bar{t}\sigma_{\mu\nu}\frac{\lambda^i}{2}qG^{i\mu\nu}] + h.c.$$



W' boson and FCNC MC event samples from SingleTop (CompHEP) generator



	SM	two-Higgs	SUSY
$BR(t \rightarrow cg)$	$5 \cdot 10^{-11}$	10^{-6}	10^{-3}
$BR(t \rightarrow c\gamma)$	$5 \cdot 10^{-13}$	10^{-6}	10^{-5}
$BR(t \rightarrow cZ)$	$\sim 10^{-13}$	10^{-9}	10^{-4}

FCNC decays are highly suppressed in SM $t \rightarrow qg, t \rightarrow q\gamma, t \rightarrow qZ$

To compare FCNC limits from top decays and top production one can express limits on FCNC couplings in term of Br fractions

$$\Gamma(t \rightarrow qg) = \left(\frac{\kappa_{tq}^g}{\Lambda}\right)^2 \frac{8}{3} \alpha_s m_t^3, \quad \Gamma(t \rightarrow q\gamma) = \left(\frac{\kappa_{tq}^\gamma}{\Lambda}\right)^2 2\alpha m_t^3,$$

$$\Gamma(t \rightarrow qZ)_\gamma = (|v_{tq}^Z|^2 + |a_{tq}^Z|^2) \alpha m_t^3 \frac{1}{4M_Z^2 \sin^2 2\theta_W} \left(1 - \frac{M_Z^2}{m_t^2}\right)^2 \left(1 + 2\frac{M_Z^2}{m_t^2}\right),$$

$$\Gamma(t \rightarrow qZ)_\sigma = \left(\frac{\kappa_{tq}^Z}{\Lambda}\right)^2 \alpha m_t^3 \frac{1}{\sin^2 2\theta_W} \left(1 - \frac{M_Z^2}{m_t^2}\right)^2 \left(2 + \frac{M_Z^2}{m_t^2}\right)$$

Expectations:

$t \rightarrow$	Tevatron	LHC		ILC
	Run II	decay	production	
gq	0.06%	1.6×10^{-3}	1×10^{-5}	–
γq	0.28%	2.5×10^{-5}	3×10^{-6}	4×10^{-6}
Zq	1.3%	1.6×10^{-4}	1×10^{-4}	2×10^{-4}

CDF:

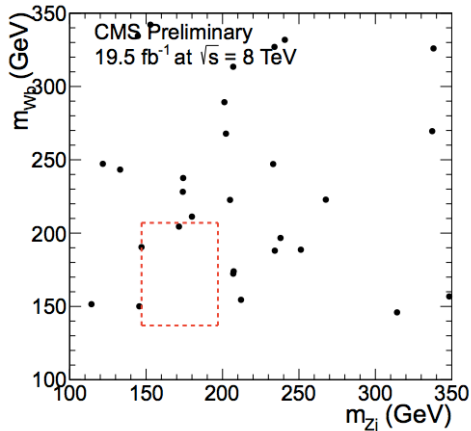
$$\mathcal{B}(t \rightarrow u + g) < 3.9 \cdot 10^{-4}$$

$$\mathcal{B}(t \rightarrow c + g) < 5.7 \cdot 10^{-3}$$

DO:

LHC limits are about to come!

	tgu	tgc
Cross section	0.20 pb	0.27 pb
κ_{tgf}/Λ	0.013 TeV ⁻¹	0.057 TeV ⁻¹
$\mathcal{B}(t \rightarrow qg)$	2.0×10^{-4}	3.9×10^{-3}



CMS limit: $\mathcal{B}(t \rightarrow Zq) < 0.07\%$ @ 95% C.L.

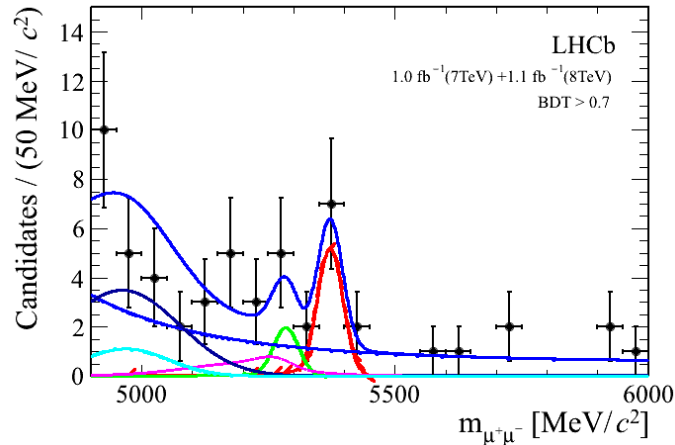
Many interesting new results

Indirect search for BSM physics - the main goal of the LHCb experiment.

Rare $B_s \rightarrow \mu\mu$ decay

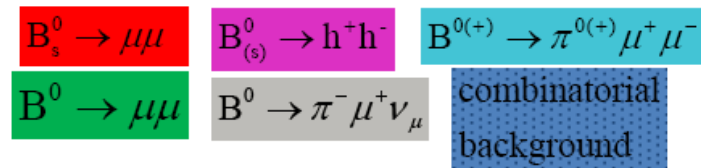
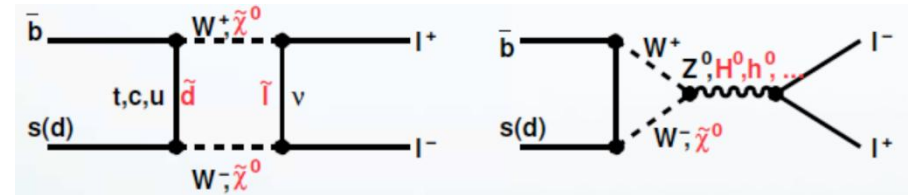
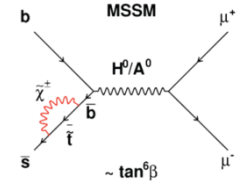
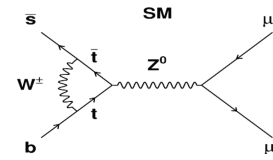
LHCb:

$$Br(B_s^0 \rightarrow \mu\mu) = (3.2_{-1.2}^{+1.4} (\text{stat})_{-0.3}^{+0.5} (\text{syst})) \times 10^{-9}$$



CMS:

$$\mathcal{B}(B_s^0 \rightarrow \mu^+\mu^-) = (3.0_{-0.9}^{+1.0}) \times 10^{-9}$$



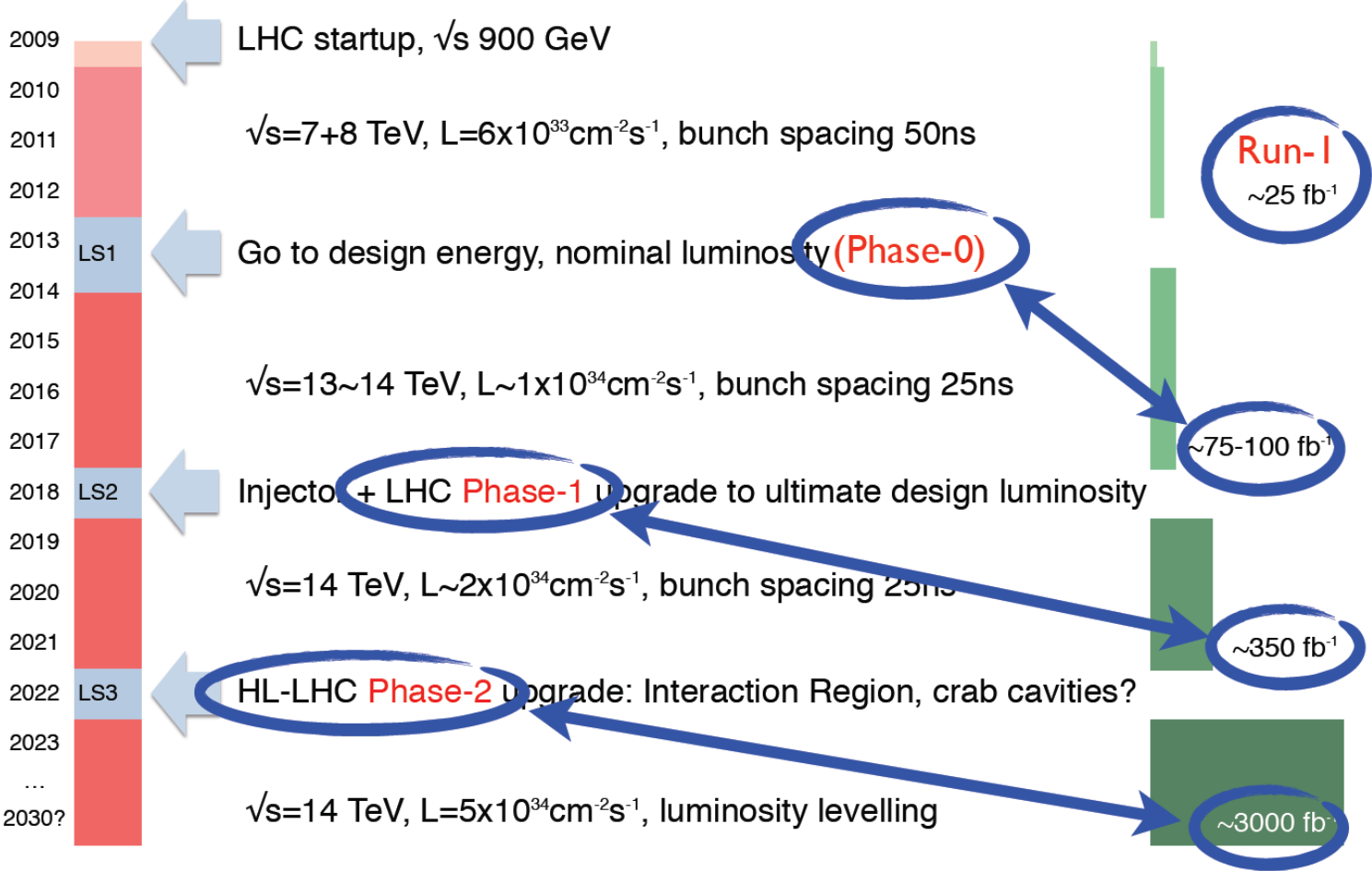
Observed decay rate is compatible with the SM expectation:

$$Br(B_s \rightarrow \mu\mu) = (3.53 \pm 0.38) \times 10^{-9}$$



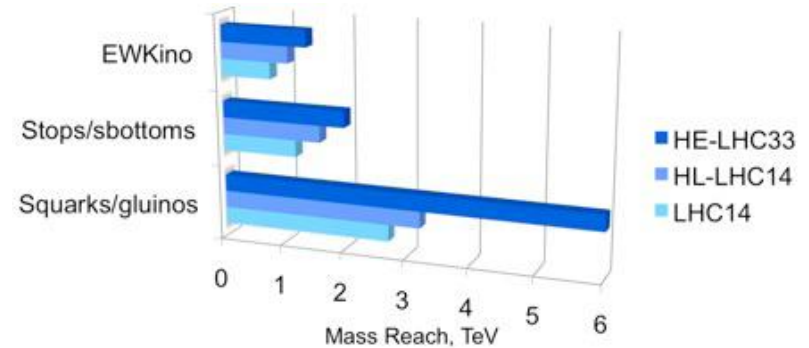
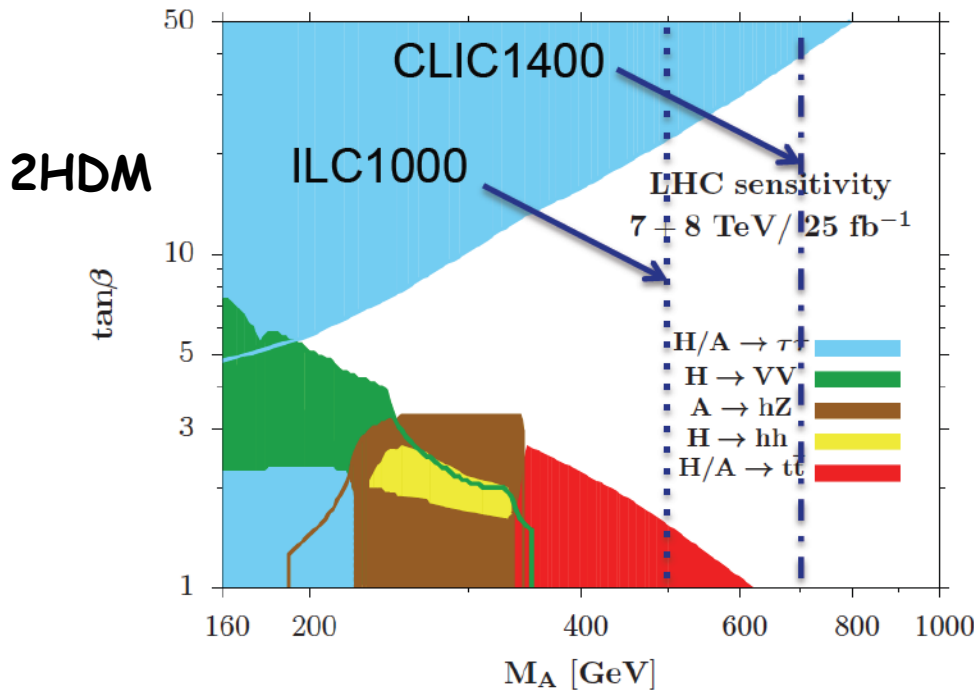
“It doesn't matter how beautiful your theory is, it doesn't matter how smart you are. If it doesn't agree with experiment, it's wrong”.

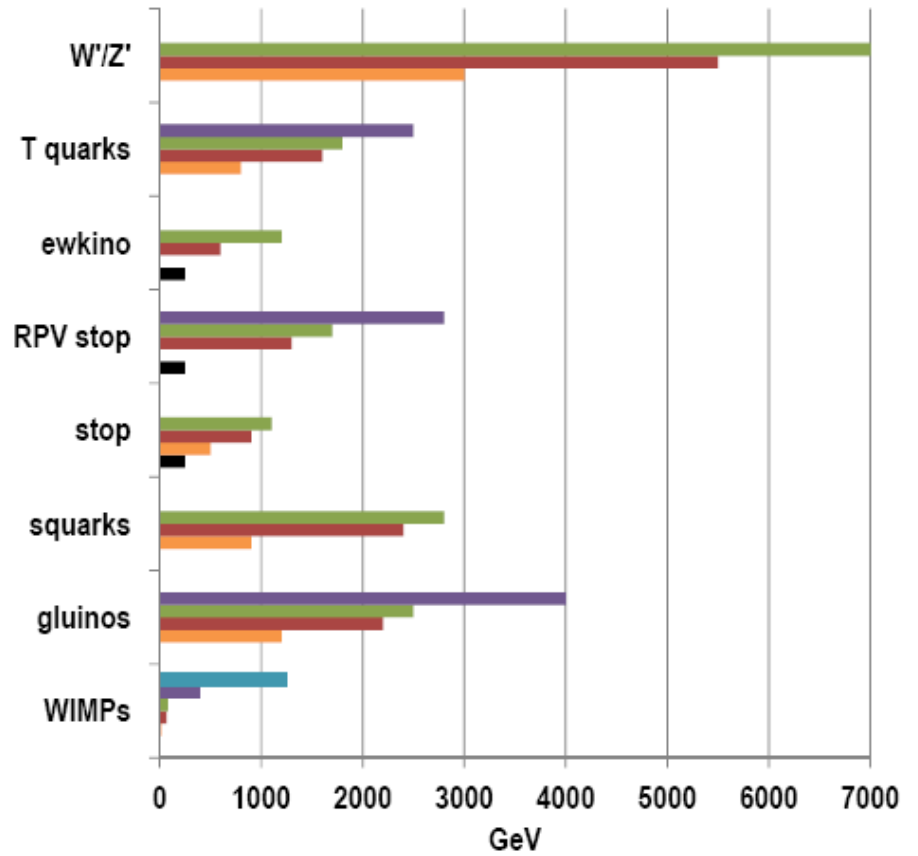
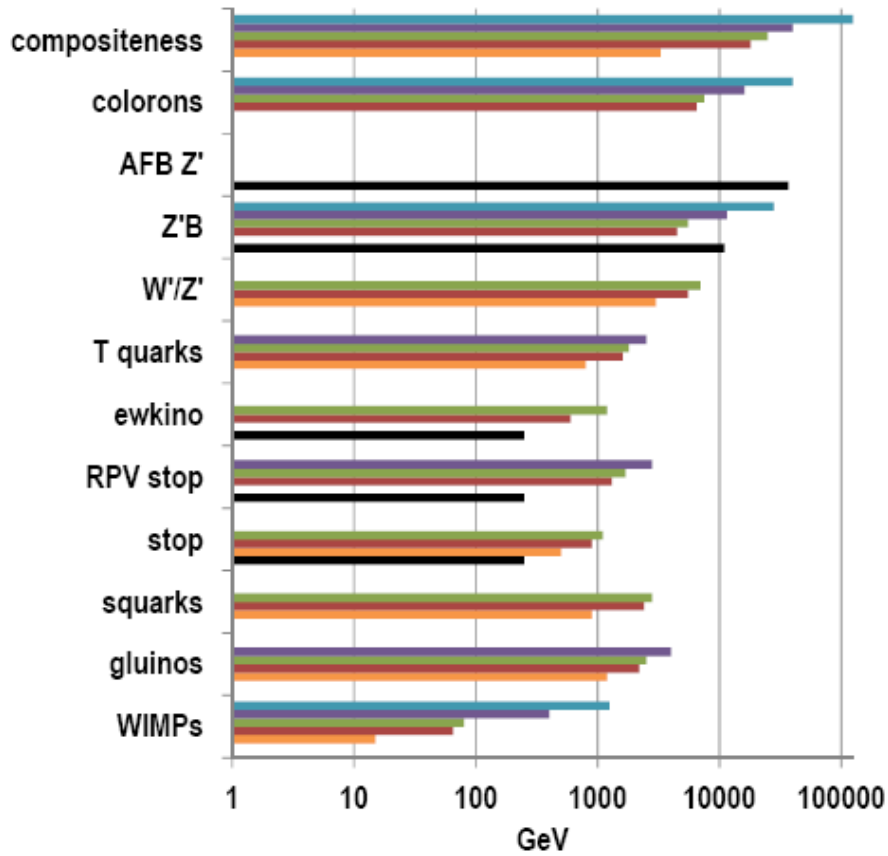
Richard P. Feynman



Expected precisions for Higgs couplings

Facility	LHC	HL-LHC	ILC500	ILC500-up	ILC1000	ILC1000-up	CLIC	TLEP (4 IPs)
\sqrt{s} (GeV)	14,000	14,000	250/500	250/500	250/500/1000	250/500/1000	350/1400/3000	240/350
$\int \mathcal{L} dt$ (fb ⁻¹)	300/expt	3000/expt	250+500	1150+1600	250+500+1000	1150+1600+2500	500+1500+2000	10,000+2600
κ_γ	5 – 7%	2 – 5%	8.3%	4.4%	3.8%	2.3%	–/5.5/<5.5%	1.45%
κ_g	6 – 8%	3 – 5%	2.0%	1.1%	1.1%	0.67%	3.6/0.79/0.56%	0.79%
κ_W	4 – 6%	2 – 5%	0.39%	0.21%	0.21%	0.13%	1.5/0.15/0.11%	0.10%
κ_Z	4 – 6%	2 – 4%	0.49%	0.24%	0.44%	0.22%	0.49/0.33/0.24%	0.05%
κ_ℓ	6 – 8%	2 – 5%	1.9%	0.98%	1.3%	0.72%	3.5/1.4/<1.3%	0.51%
κ_d	10 – 13%	4 – 7%	0.93%	0.51%	0.51%	0.31%	1.7/0.32/0.19%	0.39%
κ_u	14 – 15%	7 – 10%	2.5%	1.3%	1.3%	0.76%	3.1/1.0/0.7%	0.69%





- pp, 100 TeV, 3000/fb
- pp, 33 TeV, 3000/fb
- pp, 14 TeV, 3000/fb
- pp, 14 TeV, 300/fb
- pp, 8 TeV, 20/fb
- ee, 500 GeV, 500/fb

Very simplified plot does not show holes in searches

Discovering of new Higgs-like state

“Theorist”

“Experimentalist”



Not discovering anything else



???????



New phases of the LHC and new colliders !!!!!

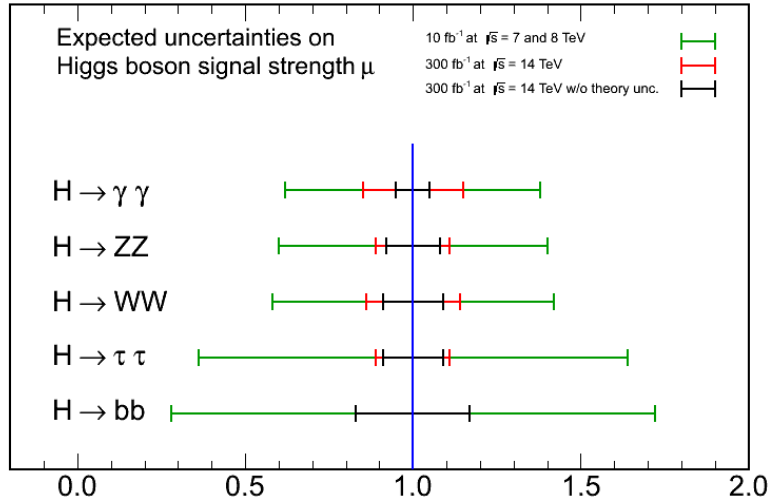
LHC (phase 0,1), HL-LHC (phase 2), HE-LHC, VLHC, ILC, CLIC

We are in a very beginning of exploration of the Terascale !

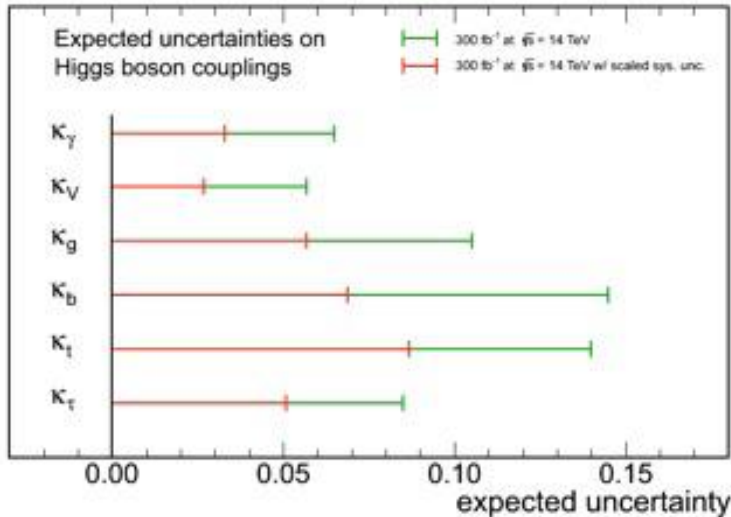
BACKUP SLIDES

Expected uncertainties in CMS

CMS Projection

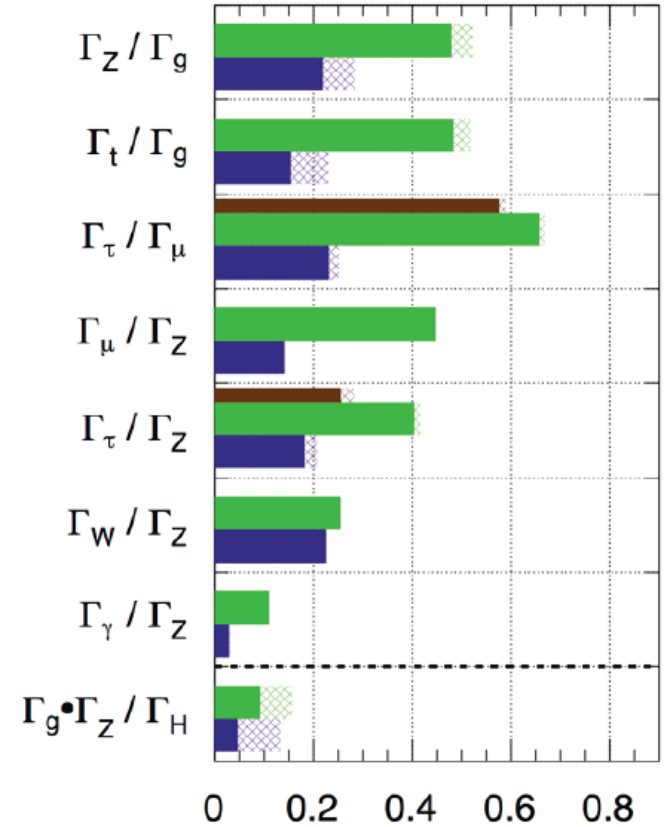


CMS Projection



ATLAS

$\sqrt{s} = 14$ TeV: $\int L dt = 300 \text{ fb}^{-1}$; $\int L dt = 3000 \text{ fb}^{-1}$
 $\int L dt = 300 \text{ fb}^{-1}$ extrapolated from 7+8 TeV



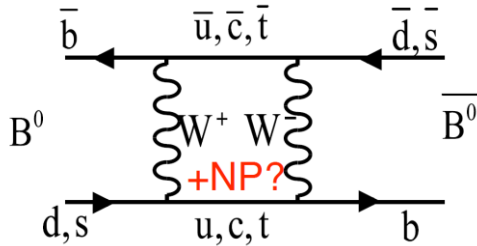
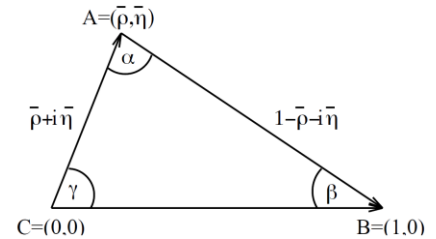
Indirect search for BSM physics - the main goal of the LHCb experiment.

Flavor & CP
in CKM matrix

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \sim \begin{pmatrix} 1 & +\lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 & +A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

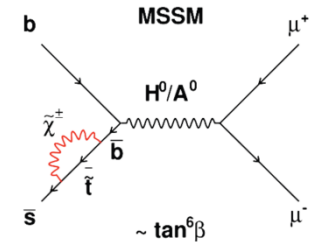
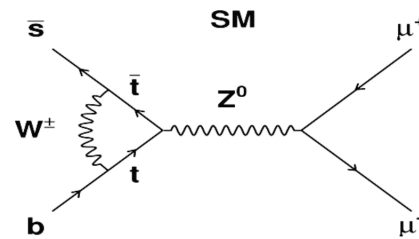
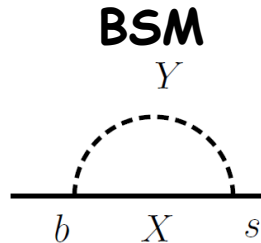
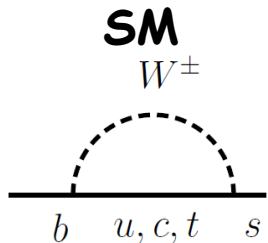
One of unitarity triangles $V_{ub}V_{ud}^* + V_{cb}V_{cd}^* + V_{tb}V_{td}^* = 0$

Two body B-meson hadronic decays ($B_d^- \rightarrow J/\psi K_S^0, \dots$) and phase of B_0 oscillations for CP violation studies



Deviations in Br fractions in rare b-decays

A-penguin, where A is added in all places radiative boson, $A = \gamma, Z, g, h^0$



Key measurements: rare B-meson decays ($B_S^- \rightarrow \mu\mu, B_S^- \rightarrow K^* \mu\mu, B_d^- \rightarrow K^* ee, B_S^- \rightarrow \phi\gamma$)