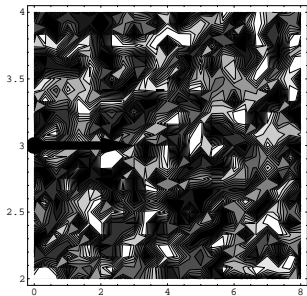


Quantum particle in a random media

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Random media



Localization length, or free path ℓ ?

- Functional representation
- Unordered media
- Random potential
- Random force
- Quantum field as random media

■ Random media

$$H = \frac{p^2}{2m} + W(x)$$

Potential $W(x)$ \implies a random media

■ Localization length $\ell(k) = \frac{1}{\Gamma(k)}$

$$\clubsuit \quad \mathbf{G}_k(\mathbf{x}, \mathbf{x}' | \mathbf{W}) = \frac{1}{\frac{p^2}{2m} + W(x) - \frac{k^2}{2m} + i0} \delta(\mathbf{x} - \mathbf{x}')$$

$$\clubsuit \quad \mathbf{G}_k(\mathbf{r}) = \langle \mathbf{G}_k(\mathbf{x}, \mathbf{x}' | \mathbf{W}) \rangle_{\mathbf{W}}, \quad \mathbf{r} = |\mathbf{x} - \mathbf{x}'|$$

$$\clubsuit \quad \frac{1}{\ell(k)} = \Gamma(k) = - \lim_{r \rightarrow \infty} \frac{\ln |\mathbf{G}_k(\mathbf{r})|}{r}$$

Functional representation

$$\begin{aligned}
 \clubsuit \quad \mathbf{G}_k(\mathbf{r}|\mathbf{W}) &= i \int_0^\infty dt e^{-it} \left[-\frac{1}{2m} \frac{d^2}{dx^2} + \mathbf{W}(x) - \frac{k^2}{2m} - i0 \right] \delta(\mathbf{x}) \\
 &= \mathbf{B} \int_0^\infty \frac{dt}{t^{\frac{3}{2}}} e^{i \left(\frac{tk^2}{m} + \frac{x^2 m}{t} \right)} \int \frac{D\eta}{\mathbf{C}} e^{i \int_0^t d\nu \left[\frac{m}{2} \dot{\eta}^2(\nu) - \mathbf{W}\left(x \frac{\nu}{t} + \eta(\nu)\right) \right]} \\
 &= \mathbf{B} \int_0^\infty \frac{ds}{s^{\frac{3}{2}}} e^{i \frac{rk}{2} \left(s + \frac{1}{s} \right)} \int \frac{D\xi}{\mathbf{C}} e^{i \frac{1}{2} \int_0^r d\tau \dot{\xi}^2(\tau) - i m \frac{s}{k} \int_0^r d\tau \mathbf{W}\left(n\tau - \sqrt{\frac{s}{k}} \xi(\tau) \right)}
 \end{aligned}$$

$$\xi(0) = \xi(r) = 0$$

$$\clubsuit \left\langle e^{-i\frac{ms}{k} \int_0^r d\tau W(n\tau + x' - \sqrt{\frac{s}{k}} \xi(\tau))} \right\rangle_w = e^{Z_r[\frac{s}{k}, \xi]}$$

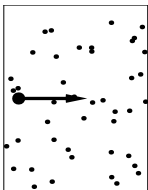
$$\clubsuit I_r(\mathbf{k}, s) = \int \frac{D\xi}{C} e^{\frac{i}{2} \int_0^r d\tau \dot{\xi}^2(\tau) + Z_r[\frac{s}{k}, \xi]} \rightarrow e^r A\left(\frac{s}{k}\right)$$

$$\clubsuit G_{\mathbf{k}}(r) = B \int_0^\infty \frac{ds}{s^{\frac{3}{2}}} e^{r\left[\frac{ik}{2}\left(s + \frac{1}{s}\right) + A\left(\frac{s}{k}\right)\right]} \sim e^r \Phi(\mathbf{k})$$

$$\clubsuit \Phi(\mathbf{k}) = \min_s \left[\frac{ik}{2} \left(s + \frac{1}{s} \right) + A\left(\frac{s}{k}\right) \right] = \min_{\mathbf{v}} \left[\frac{ik^2}{2} \mathbf{v} + \frac{i}{2\mathbf{v}} + A(\mathbf{v}) \right]$$

$$\frac{1}{\ell(\mathbf{k})} = \Gamma(\mathbf{k}) = -\text{Re } \Phi(\mathbf{k})$$

Unordered media



N - atoms, Λ - volume

$n = \frac{N}{\Lambda} = \text{const}$ - density

$P = \frac{1}{\Lambda^N}$ - probability

$$\mathbf{W}(\mathbf{x}) = \mathbf{W}(\mathbf{x}|\{\mathbf{y}_j\}) = \sum_j^N \mathbf{V}(\mathbf{x} - \mathbf{y}_j)$$

$$\mathbf{V}(\mathbf{x}) = \mathbf{g} \mathbf{U}\left(\frac{\mathbf{r}}{\mathbf{R}}\right), \quad \int_0^\infty d\mathbf{s} \mathbf{U}(\mathbf{s}) = \mathbf{1}, \quad nR^3 \ll 1$$

$$e^{Z_r\left[\frac{\mathbf{s}}{\mathbf{k}}, \xi\right]} = \left\langle e^{-im\frac{\mathbf{s}}{\mathbf{k}} \int_0^r d\tau \mathbf{W}(n\tau + \mathbf{x}' - \sqrt{\frac{\mathbf{s}}{\mathbf{k}}} \xi(\tau) | \mathbf{y}_j)} \right\rangle_{\{\mathbf{y}_j\}}$$

$$= \exp \left\{ nR^3 \int \frac{d\mathbf{y}}{R^3} \left[e^{-im\frac{\mathbf{s}}{\mathbf{k}} \int_0^r d\tau \mathbf{V}(\mathbf{y} + n\tau - \sqrt{\frac{\mathbf{s}}{\mathbf{k}}} \xi(\tau))} - \mathbf{1} \right] \right\}$$

$$e^{r A(\frac{s}{k})} = \int \frac{D\xi}{C} e^{i \frac{1}{2} \int_0^r d\tau \xi^2(\tau) + Z_r[\frac{s}{k}, \xi]}$$

$$A\left(\frac{s}{k}\right) \approx n \int \frac{D\xi}{C} e^{i \frac{1}{2} \int_0^r d\tau \xi^2(\tau)} \int dy \left[e^{-i \frac{ms}{k} \int_0^r d\tau V(\mathbf{y} + n\tau - \sqrt{\frac{s}{k}} \xi(\tau))} - 1 \right]$$

$$\mathbf{A}\left(\frac{s}{k}\right) = i \frac{n}{2} \cdot 4\pi \frac{s}{k} \mathbf{F}\left(\frac{k}{s}, \mathbf{0}\right) = \frac{n}{2} \left[i\phi\left(\frac{k}{s}\right) - \sigma\left(\frac{k}{s}\right) \right]$$

$\mathbf{F}\left(\frac{k}{s}, \mathbf{0}\right)$ - elastic scattering amplitude for momentum $\frac{k}{s}$ and scattering angle $\theta = 0$.

$$\Phi(\mathbf{k}) = \min_s \left[\frac{i\mathbf{k}}{2} \left(s + \frac{1}{s} \right) + i n \cdot 2\pi \frac{s}{k} \mathbf{F}\left(\frac{k}{s}, \mathbf{0}\right) \right] \Rightarrow s = 1.$$

$$\frac{1}{\ell(\mathbf{k})} = \Gamma(\mathbf{k}) = -\text{Re } \Phi(\mathbf{k}) = n\sigma(\mathbf{k}) \Rightarrow I(\mathbf{r}) \sim |\mathbf{G}_E(\mathbf{r})|^2 \sim e^{-r n\sigma(\mathbf{k})}$$

$\sigma(k)$ - total cross section, $\ell = \frac{1}{n\sigma(k)}$ - free path.

Small momenta $k \rightarrow 0$

$$4\pi \frac{s}{k} F\left(\frac{k}{s}, \mathbf{0}\right) = 4\pi \frac{s}{k} \operatorname{Re} F\left(\frac{k}{s}, \mathbf{0}\right) + i\sigma\left(\frac{k}{s}\right) \Rightarrow 4\pi \frac{s}{k} a + i\sigma\left(\frac{k}{s}\right)$$

a - scattering length, $\sigma\left(\frac{k}{s}\right) \sim a^2 \ll \frac{a}{k}$

$$\Phi(\mathbf{k}) = \min_s \left[i \left[\frac{k}{2} \left(s + \frac{1}{s} \right) + \frac{2\pi n a}{k} s \right] - n \sigma\left(\frac{k}{s}\right) \right] \Rightarrow n \sigma\left(\sqrt{k^2 + 4\pi n a}\right)$$

$\frac{1}{k^2} a = \lambda^2 a = V_{\text{int}}$ - "interaction" volume,

$\frac{1}{n} \sim L^3 = V_{\text{atom}}$ - "atom" volume

multiple scattering $k^2 \ll 4\pi n a \Rightarrow \lambda^2 a \gg L^3, \quad V_{\text{int}} \gg V_{\text{atom}}$

$$I(\mathbf{r}) \sim e^{-n \sigma(\sqrt{k^2 + 4\pi n a}) \cdot r} = \begin{cases} e^{-n \sigma(k) \cdot r}, & k \rightarrow \infty, \quad V_{\text{int}} \ll V_{\text{atom}} \\ e^{-n \sigma(\sqrt{4\pi n a}) \cdot r}, & k \rightarrow 0, \quad V_{\text{int}} \gg V_{\text{atom}} \end{cases}$$

Random potential. Linear path approximation

Potential $W(x)$ is a random function with correlator

$$\langle W(x)W(y) \rangle_W = K(x-y) = g^2 U\left(\frac{|x-y|}{R}\right)$$

R - correlation length, g^2 - coupling constant

$$e^{rA\left(\frac{k}{s}\right)} = \int \frac{D\xi}{C} e^{\frac{i}{2} \int_0^r d\tau \dot{\xi}^2(\tau) - \frac{m^2 s^2}{2k^2} \iint_0^r d\tau d\tau' K(n(\tau-\tau') - \sqrt{\frac{s}{k}}(\xi(\tau) - \xi(\tau')))}$$

Large $r \Rightarrow \xi(\tau) \sim \tau \Rightarrow \xi(\tau) = b\tau + \eta(\tau)$

$$e^{rA\left(\frac{k}{s}\right)} \approx \int \frac{db}{C_b} e^{\frac{i}{2}rb^2 - \frac{m^2 s^2}{2k^2} \iint_0^r d\tau d\tau' K\left((n+b\sqrt{\frac{s}{k}})(\tau-\tau')\right)} = \int \frac{db}{C_b} e^{r\left[\frac{i}{2}b^2 - \frac{s^2}{k^2} \frac{G}{|n+\sqrt{\frac{s}{k}}b|}\right]}$$

$$A\left(\frac{k}{s}\right) \approx \min_b \left[\frac{i}{2}b^2 - \frac{s^2}{k^2} \frac{G}{|n+\sqrt{\frac{s}{k}}b|} \right], \quad G = \int_0^\infty dv K(v)$$

$$\begin{aligned}
\Phi(\mathbf{k}) &= \min_{s, \mathbf{b}} \left\{ i \frac{\mathbf{k}}{2} \left(s + \frac{1}{s} (1 + (\mathbf{b} + 1)^2) \right) - \left(\frac{s}{\mathbf{k}} \right)^2 \frac{\mathbf{G}}{|\mathbf{b}|} \right\} \\
&= i \frac{\mathbf{k}}{2} \min_{x > 0} \left\{ \sqrt{2 \left(1 + x^2 + i \frac{2\mathbf{G}}{x\mathbf{k}^3} \right)} - x \right\} \\
\frac{1}{\ell(\mathbf{k})} = \Gamma(\mathbf{k}) = -\text{Re } \Phi(\mathbf{k}) &= \begin{cases} O\left(\frac{\mathbf{G}}{k^2}\right), & k \rightarrow \infty, \\ O\left(\mathbf{G}^{\frac{1}{3}}\right), & k \rightarrow 0. \end{cases}
\end{aligned}$$

Random force $\mathcal{F}(\mathbf{x})$. Perturbation calculation.

$$H = \frac{\mathbf{p}^2}{2m} + (\mathcal{F}(\mathbf{x})\mathbf{x})$$

$$\langle \mathcal{F}_i(\mathbf{x})\mathcal{F}_j(\mathbf{y}) \rangle = \delta_{ij}\mathbf{K}(\mathbf{x} - \mathbf{y}) = \delta_{ij}g^2\mathbf{U}\left(\frac{|\mathbf{x} - \mathbf{y}|}{R}\right), \quad \int_0^\infty ds \mathbf{U}(s) < \infty.$$

$$\mathbf{G}_k(\mathbf{r}) = \int_0^\infty \frac{ds}{s^{\frac{3}{2}}} e^{\frac{irk}{2}(s + \frac{1}{s})} \int \frac{D\xi}{C} e^{\frac{i}{2} \int_0^r d\tau \dot{\xi}^2(\tau) - \frac{1}{2} \left(\frac{ms}{k}\right)^2 \iint_0^r d\tau_1 d\tau_2 \mathbf{K}(\mathbf{b}(\tau_1) - \mathbf{b}(\tau_2)) (\mathbf{b}(\tau_1)\mathbf{b}(\tau_2))}$$

$$\mathbf{b}(\tau) = \mathbf{n}\tau + \sqrt{\frac{s}{k}}\xi(\tau)$$

$$\ln \mathbf{G}_k(\mathbf{r}) = \min_{\mathbf{v}} \left[\frac{irk^2}{2}\mathbf{v} + \frac{i\mathbf{r}}{2\mathbf{v}} - \frac{1}{2}\mathbf{v}^2 \mathbf{r}^3 \frac{m^2}{3} \int_0^{\left(\frac{2}{\mathbf{v}}\right)^2} \frac{dz}{4\pi} \tilde{\mathbf{K}}(z) \right] \rightarrow -\frac{3}{8}(2H)^{\frac{1}{3}} \cdot \mathbf{r}^{\frac{5}{3}}$$

Quantum field as a random media

1. Quantum field $\Phi(\mathbf{x})$ - external gaussian random field

$$H = \frac{\mathbf{p}^2}{2m} + h\Phi(\mathbf{x})$$

$$G_k(\mathbf{x} - \mathbf{y}) = \left\langle \frac{1}{\frac{\mathbf{p}^2}{2m} + h\Phi(\mathbf{x}) - \frac{\mathbf{k}^2}{2m}} \delta(\mathbf{x} - \mathbf{y}) \right\rangle_{\Phi}$$

2. Quantum field and quantum particle - parts of unique system

$$H = H_0[\Phi] + \frac{\mathbf{p}^2}{2m} + h\Phi(\mathbf{x})$$

$$G_k(\mathbf{x} - \mathbf{y}) = \left\langle \frac{1}{H_0[\Phi] + \frac{\mathbf{p}^2}{2m} + h\Phi(\mathbf{x}) - \frac{\mathbf{k}^2}{2m}} \delta(\mathbf{x} - \mathbf{y}) \right\rangle_{\Phi}$$

Field $\Phi(\mathbf{x})$

$$H_0[\Phi] = \int d\mathbf{p} \rho a_{\mathbf{p}}^+ a_{\mathbf{p}}$$

$$\Phi(\mathbf{x}, t) = \int \frac{d\mathbf{p}}{(2\pi)^{\frac{3}{2}} \sqrt{2\rho}} \left[a_{\mathbf{p}} e^{-ipt + i\mathbf{p}\mathbf{x}} + a_{\mathbf{p}}^+ e^{ipt' - i\mathbf{p}\mathbf{x}} \right], \quad \rho = |\mathbf{p}|$$

$$\begin{aligned} K(\mathbf{x} - \mathbf{x}') &= \langle 0 | T \{ \Phi(\mathbf{x}, t) \Phi(\mathbf{x}', t') \} | 0 \rangle = D_c(\mathbf{x} - \mathbf{x}', t - t') \\ &= \int \frac{d\mathbf{p}}{(2\pi)^3 2\rho} e^{-i\rho|t-t'| + i\mathbf{p}(\mathbf{x} - \mathbf{x}')} \end{aligned}$$

1. Quantum particle and quantum field as random media

$$G_k(\mathbf{x} - \mathbf{y}) = \left\langle \frac{1}{\frac{p^2}{2m} + h\Phi(\mathbf{x}) - \frac{k^2}{2m}} \delta(\mathbf{x} - \mathbf{y}) \right\rangle_{\Phi} \sim e^{-r \frac{g^2}{8\pi k}}$$

2. Quantum field and quantum particle - parts of unique system

$$G_k(\mathbf{x} - \mathbf{y}) = \left\langle \frac{1}{H_0[\Phi] + \frac{p^2}{2m} + h\Phi(\mathbf{x}) - \frac{k^2}{2m}} \delta(\mathbf{x} - \mathbf{y}) \right\rangle_{\Phi} \sim \begin{cases} 1 & k < m \\ e^{-r \frac{g^2 m^2}{k^2} \frac{k-m}{8\pi^2}} & k > m \end{cases}$$

Relativistic particle

$$[\square_x + m^2 + \mathbf{g}\Phi(\mathbf{x})] \mathbf{G}(\mathbf{x} - \mathbf{x}') = \delta(\mathbf{x} - \mathbf{x}'), \quad \square = \partial_t^2 - \nabla^2, \quad \mathbf{x}^2 = \mathbf{x}_0^2 - \mathbf{x}^2$$

$$\mathbf{G}_E(\mathbf{x} - \mathbf{y}) = \langle \mathbf{G}_E(\mathbf{x}, \mathbf{y} | \Phi) \rangle_\Phi = \left\langle \frac{1}{\square_x + m^2 + \mathbf{g}\Phi(\mathbf{x})} \delta(\mathbf{x} - \mathbf{y}) \right\rangle_\Phi$$

$$\ln |\mathbf{G}(\mathbf{x} - \mathbf{y})| \rightarrow 0 \quad \Rightarrow \quad \frac{1}{\ell} = \Gamma = 0$$

Conclusion

- The path integral representation - the best method to investigate differential equations with random coefficients.
- Green function - representation in the form of the Gaussian path integral.
- Large momenta k - perturbation calculations.
- Unordered media

$$I(r) \sim e^{-n \sigma(\sqrt{k^2+4\pi na}) \cdot r} = \begin{cases} e^{-n \sigma(k) \cdot r}, & k \rightarrow \infty, \quad \text{single scatt.;} \\ e^{-n \sigma(\sqrt{4\pi na}) \cdot r}, & k \rightarrow 0, \quad \text{multiple scatt.} \end{cases}$$

- Random potential - Linear path approximation for small k .
- Random force - stable bound states.
- Quantum field as a random media - localization is absent in the relativistic case only.