

Hadronic molecules with heavy quarks

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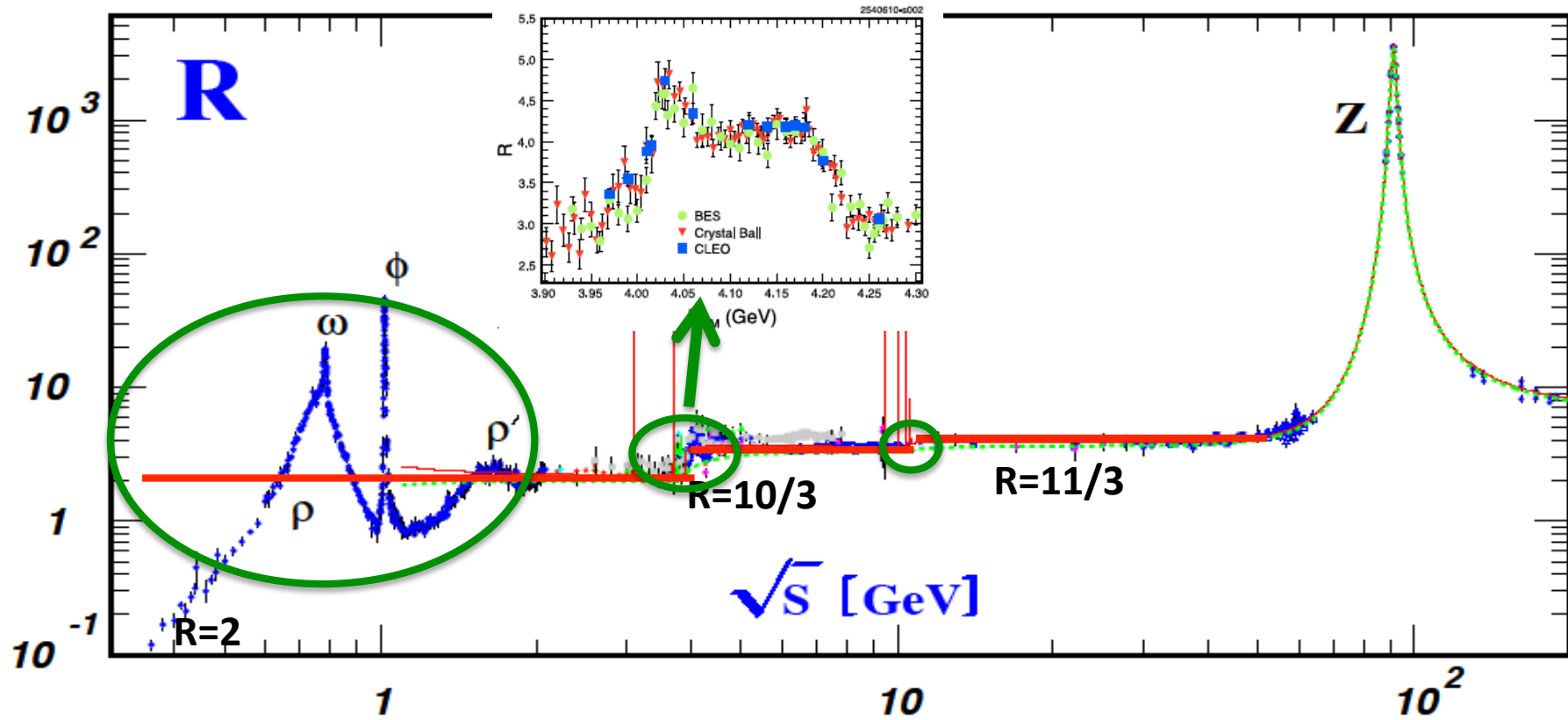
1. Introduction

— *Hadron Physics* —

- Evidence of Higgs (like) has been observed
Low energy QCD for Hadron is perhaps least understood
- Lagrangian is **simple** but **not easy to solve**
- Long history: Experiments, Models (Empirical rules),
Computer simulations (**Lattice QCD, Kei SC@Kobe**)
- Recent (last decade) observations has opened
Unexpectedly **rich spectrum** near **thresholds**
at KEK, Spring-8, J-PARCBES, RHIC, LHC, ...

Hadrons are around thresholds

$$R(s) = \sigma(e^+e^- \rightarrow \text{hadrons}, s) / \sigma(e^+e^- \rightarrow \mu^+\mu^-, s).$$



1 GeV →

1 GeV →

GeV in Log scale

u d s

Mass
generation

c

b

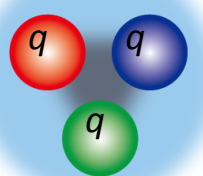
Hadrons are composite

Many resonant states

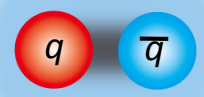
Particle data book (PDG)

				LIGHT UNFLAVORED (S = C = B = 0)				STRANGE (S = ±1, C = B = 0)		CHARMED, STRANGE (C = S = ±1)		cτ F(J ^{PC})			
				F(J ^{PC})		F(J ^{PC})		I(J ^{PC})		I(J ^{PC})					
p	P ₁₁	****	Δ(1232)	P ₃₃	****	Σ ⁺	P ₁₁	****	Ξ ⁰	P ₁₁	****	Λ _c ⁺	****	• η _c (1S)	0 ⁺ (0 ⁻ +)
n	P ₁₁	****	Δ(1600)	P ₃₃	***	Σ ⁰	P ₁₁	****	Ξ ⁻	P ₁₁	****	Λ _c (2595) ⁺	***	• J/ψ(1S)	0 ⁻ (1 ⁻ -)
N(1440)	P ₁₁	****	Δ(1620)	S ₃₁	****	Σ ⁻	P ₁₁	****	Ξ(1530)	P ₁₃	****	Λ _c (2625) ⁺	***	• X _{c0} (1P)	0 ⁺ (0 ⁺ +)
N(1520)	D ₁₃	****	Δ(1700)	D ₃₃	****	Σ(1385)	P ₁₃	****	Ξ(1620)	*		Λ _c (2765) ⁺	*	• X _{c1} (1P)	0 ⁺ (1 ⁺ +)
N(1535)	S ₁₁	****	Δ(1750)	P ₃₁	*	Σ(1480)	*		Ξ(1690)	***		Λ _c (2880) ⁺	***	• X _{c1} (1P)	0 ⁺ (1 ⁺ +)
N(1650)	S ₁₁	****	Δ(1900)	S ₃₁	**	Σ(1560)	**		Ξ(1820)	D ₁₃	***	Λ _c (2940) ⁺	***	• h _c (1P)	? ⁺ (1 ⁺ +)
N(1675)	D ₁₅	****	Δ(1905)	F ₃₅	****	Σ(1580)	D ₁₃	*	Ξ(1950)	***		Λ _c (2940) ⁺	***	• X _{c2} (1P)	0 ⁺ (2 ⁺ +)
N(1680)	F ₁₅	****	Δ(1910)	P ₃₁	****	Σ(1620)	S ₁₁	**	Ξ(2030)	***		Σ _c (2520)	***	• η _c (2S)	0 ⁺ (0 ⁻ +)
N(1700)	D ₁₃	***	Δ(1920)	P ₃₃	***	Σ(1660)	P ₁₁	***	Ξ(2120)	*		Σ _c (2800)	***	• ψ(2S)	0 ⁻ (1 ⁻ -)
N(1710)	P ₁₁	***	Δ(1930)	D ₃₅	***	Σ(1670)	D ₁₃	****	Ξ(2250)	**		Ξ _c ⁺	***	• X(3770)	0 ⁻ (1 ⁻ -)
N(1720)	P ₁₃	****	Δ(1940)	D ₃₃	*	Σ(1690)	**		Ξ(2370)	**		Ξ _c ⁰	***	• X(3872)	0 ⁺ (? ⁺ +)
N(1900)	P ₁₃	**	Δ(1950)	F ₃₇	****	Σ(1750)	S ₁₁	***	Ξ(2500)	*		Ξ _c ⁻	***	• X _{c2} (2P)	0 ⁺ (2 ⁺ +)
N(1990)	F ₁₇	**	Δ(2000)	F ₃₅	**	Σ(1770)	P ₁₁	*	Ω ⁻	****		Ξ _c ⁺	***	• X(3940)	? ⁺ (? ⁺ +)
N(2000)	F ₁₅	**	Δ(2100)	S ₃₁	*	Σ(1775)	D ₁₅	****	Ω(2250) ⁻	***		Ξ _c ⁰	***	• X(3945)	? ⁺ (? ⁺ +)
N(2080)	D ₁₃	**	Δ(2200)	G ₃₇	*	Σ(1840)	P ₁₃	*	Ω(2380) ⁻	**		Ξ _c ⁻	***	• ψ(4040)	0 ⁻ (1 ⁻ -)
N(2090)	S ₁₁	*	Δ(2300)	H ₃₉	**	Σ(1880)	P ₁₁	**	Ω(2470) ⁻	**		Ξ _c ⁻	***	• B [±] /B ⁰ ADMIXTURE	0 ⁻ (1 ⁻ -)
N(2100)	P ₁₁	*	Δ(2350)	D ₃₅	**	Σ(1915)	F ₁₅	****				Ξ _c ⁻	***	• B [±] /B ⁰ /B _s ⁰ /b-baryon ADMIXTURE	0 ⁻ (1 ⁻ -)
N(2190)	G ₁₇	****	Δ(2390)	F ₃₇	*	Σ(1940)	D ₁₃	***				Ξ _c ⁻	***	V _{ub} and V _{ub} CKM Matrix Elements	0 ⁻ (1 ⁻ -)
N(2200)	D ₁₅	**	Δ(2400)	G ₃₉	**	Σ(2000)	S ₁₁	*				Ξ _c ⁻	***	• B*	0 ⁻ (1 ⁻ -)
N(2220)	H ₁₉	****	Δ(2420)	H _{3,11}	****	Σ(2030)	F ₁₇	****				Ξ _c ⁻	***	B _s ⁺ (5732)	? ⁺ (? ⁺ +)
N(2250)	G ₁₉	****	Δ(2750)	I _{3,11}	**	Σ(2070)	F ₁₅	*				Ξ _c ⁻	***	• B _s ⁺ (5747) ⁰	1/2(1 ⁺ +)
N(2600)	I _{3,11}	***	Δ(2950)	K _{3,15}	**	Σ(2080)	P ₁₃	*				Ξ _c ⁻	***	• B _s ⁺ (5747) ⁰	1/2(2 ⁺ +)
N(2700)	K _{1,13}	**				Σ(2100)						Ξ _c ⁻	***		
			Λ	P ₀₁	****	Σ(2100)						Ξ _c ⁻	***		
			Λ(1405)	S ₀₁	***	Σ(2100)						Ξ _c ⁻	***		
			Λ(1520)			Σ(2100)						Ξ _c ⁻	***		
			Λ(1600)			Σ(2100)						Ξ _c ⁻	***		
			Λ(1670)			Σ(2100)						Ξ _c ⁻	***		
			Λ(1690)			Σ(2100)						Ξ _c ⁻	***		
			Λ(1800)	S ₀₁	***	Σ(2100)						Ξ _c ⁻	***		
			Λ(1810)	P ₀₁	***	Σ(2100)						Ξ _c ⁻	***		
			Λ(1820)	F ₀₅	***	Σ(2100)						Ξ _c ⁻	***		
			Λ(1830)	D ₀₅	***	Σ(2100)						Ξ _c ⁻	***		
			Λ(1890)	P ₀₃	****	Σ(2100)						Ξ _c ⁻	***		
			Λ(2000)	*		Σ(2100)						Ξ _c ⁻	***		
			Λ(2020)	F ₀₇	*	Σ(2100)						Ξ _c ⁻	***		
			Λ(2100)	G ₀₇	****	Σ(2100)						Ξ _c ⁻	***		
			Λ(2110)	F ₀₅	***	Σ(2100)						Ξ _c ⁻	***		
			Λ(2325)	D ₁₃	*	Σ(2100)						Ξ _c ⁻	***		
			Λ(2350)	H ₀₉	***	Σ(2100)						Ξ _c ⁻	***		
			Λ(2585)	**		Σ(2100)						Ξ _c ⁻	***		

Baryons ~ qqq
No exception



Mesons ~ qq̄
Except two



But all of them seem excitations of minimum number (2 or 3) of valence quarks

But why not **exotics** such as

- **Glueon** excitations (glueballs, hybrids, ...)
- Hadrons of more quarks,
tetraquarks, pentaquarks, ...
- Multi-hadron hadrons (**hadronic molecules**)
other than atomic nuclei

But all of them seem excitations of minimum number (2 or 3) of valence quarks

But why not **exotics** such as

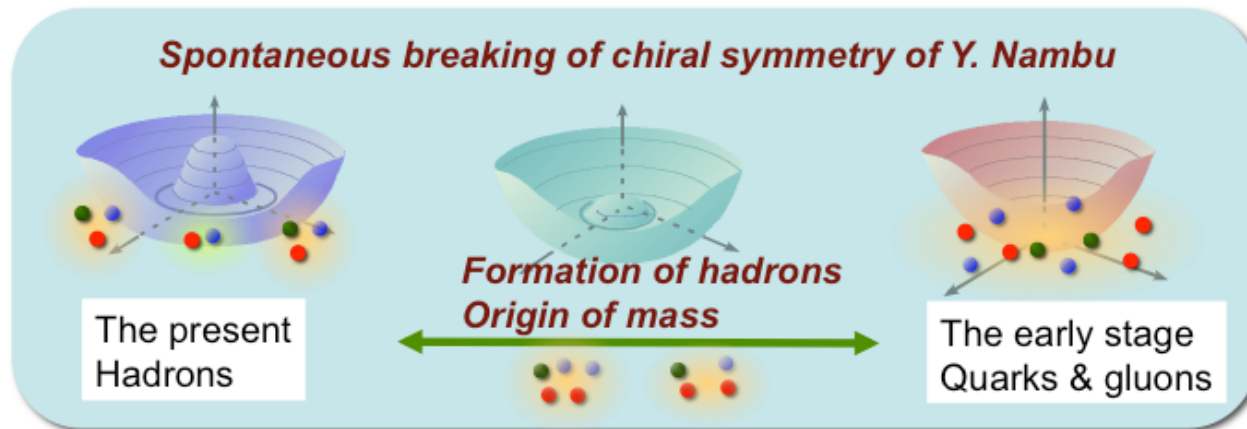
- **Glueon** excitations (glueballs, hybrids, ...)
- Hadrons of more quarks,
tetraquarks, pentaquarks, ...
- Multi-hadron hadrons (**hadronic molecules**)
other than atomic nuclei

Variety of flavors (**Heavy quarks**) may be a good probe

Nambu's SSB χ & Yukawa's pion



- (1) Light bare quarks
=> **Massive constituent quarks**
- (2) Appearance of the **massless pion**



OPEP

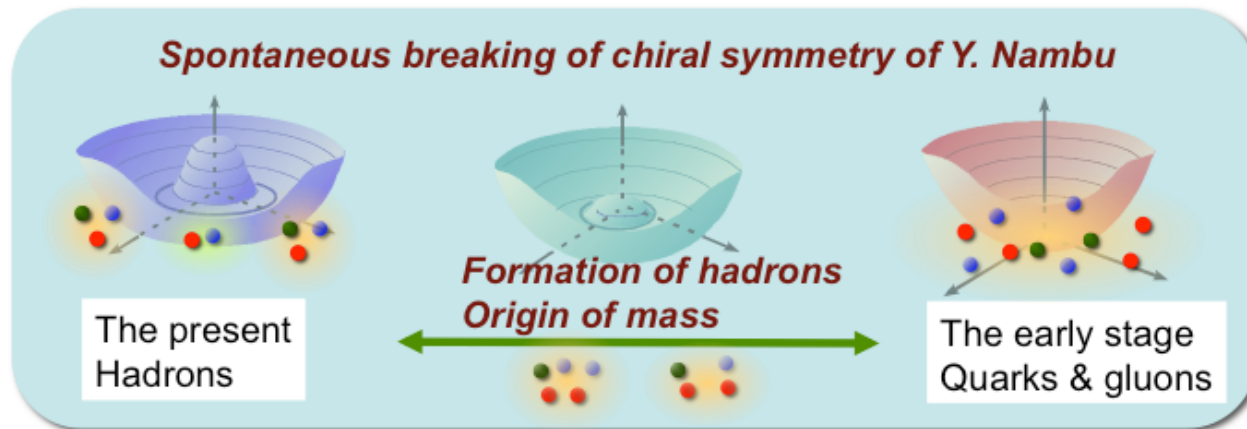
Long range **tensor force**
spin and orientation dependent

$$\begin{array}{ccc}
 \vec{\sigma} \cdot \vec{q} & \xrightarrow{\pi} & \vec{\sigma} \cdot \vec{q} \\
 | & & | \\
 N & & N
 \end{array}$$

Nambu's SSB χ & Yukawa's pion

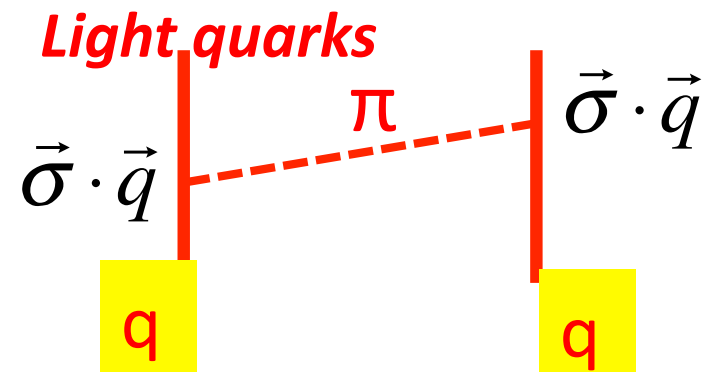


- (1) Light bare quarks
=> **Massive constituent quarks**
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OPEP

Long range **tensor force**
spin and orientation dependent



2. Exotic hadrons with many quarks

Near flavor thresholds



Mesons

X(3872)

State	Mass (MeV)	Width (MeV)	Decay	Production
Y _s (2175)	2175±8	58±26	ff ₀	ISR
X(3872)	3871.84±0.33	<0.95	J/ψpp, J/ψg	B decay
X(3872)	3872.8 +0.7/-0.6	3.9 +2.8/-1.8	D* ⁰ D ⁰	B decay
Z(3940)	3929±5	29±10	DD	gg
X(3940)	3942±9	37±17	DD*	Double-charm
Y(3940)	3942±17	87±34	J/ψw	B decay
Y(4008)	4008 +82/-49	226 +97/-80	J/ψpp	ISR
Z(4051) ⁺	4051 +24/-43	82 +51/-28	ρc ₁	B decay
X(4160)	4156±29	139 +113/-65	D*D*	Double-charm
Z(4248) ⁺	4248 +185/-45	177 +320/-72	ρc ₁	B decay
Y(4260)	4264±12	83±22	J/ψpp	ISR
Y(4350)	4361±13	74±18	γ'pp	ISR
Z(4430) ⁺	4433±5	45 +35/-18	γ'p	B decay
Y(4660)	4664±12	48±15	γ'pp	ISR
Y _b (10890)	10889.6±2.3	54.7 +8.9/-7.6	ppΥ(nS)	e ⁺ e ⁻ annihilation
Y(3915)	3915±4	17±10	J/ψw	gg
X(4350)	4350 +4.7/-5.1	13 +18/-14	J/ψf	gg
h _b (1P)	9898.3±1.5		MM(pp)	Υ(5S) /Y _b decay
h _b (2P)	10259.3 +1.6/-1.2		MM(pp)	Υ(5S) /Y _b decay
Z _b (10610)	10608.4±2.0	15.6±2.5	(Υ(nS) or h _b)p	Υ(5S) /Y _b decay
Z _b (10650)	10653.2±1.5	14.4±3.2	(Υ(nS) or h _b)p	Υ(5S) /Y _b decay

Z_b(10610)

Z_b(10650)

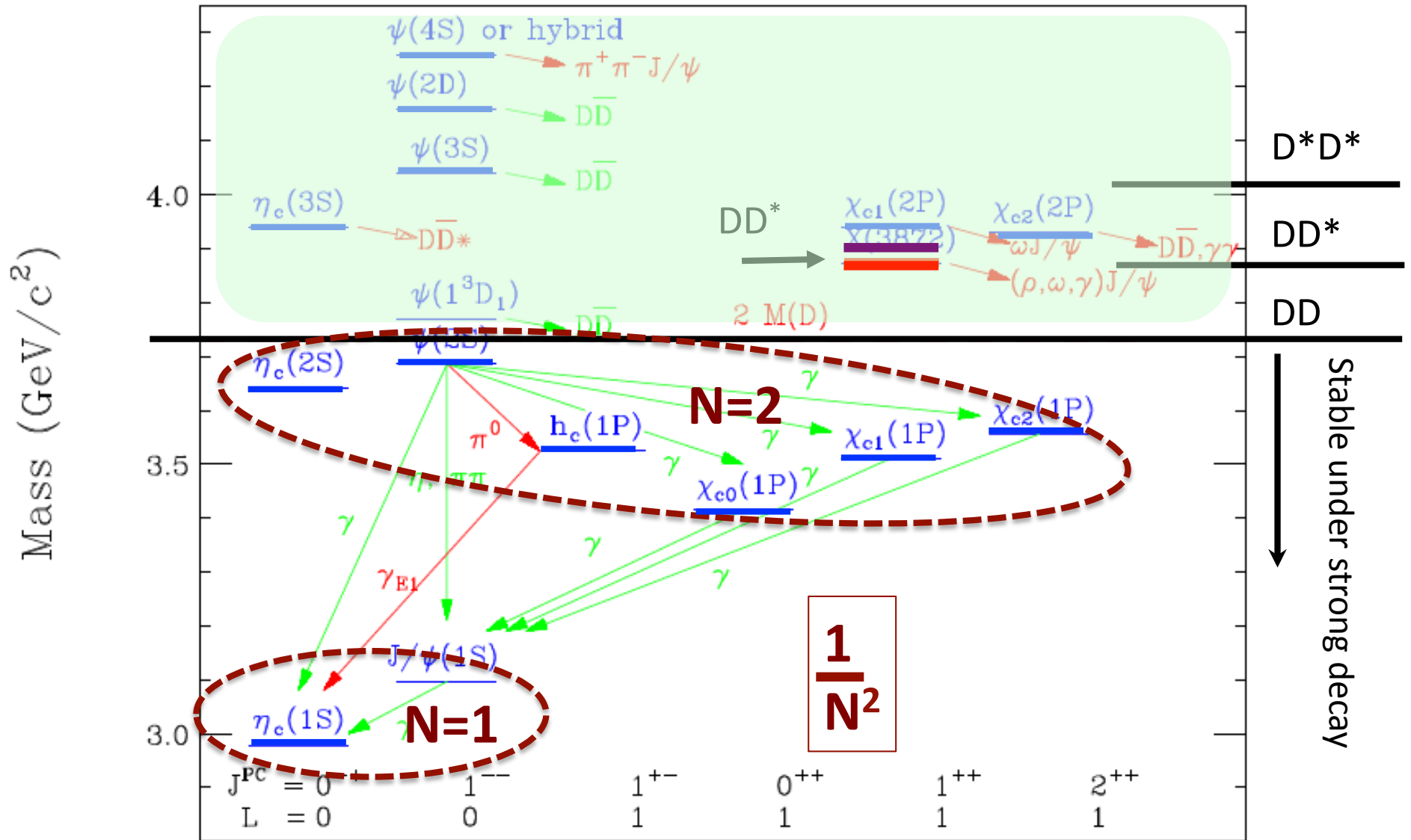
Z_c(3900), Phys. Rev. Lett. 110, 252001 (2013)

July 2013

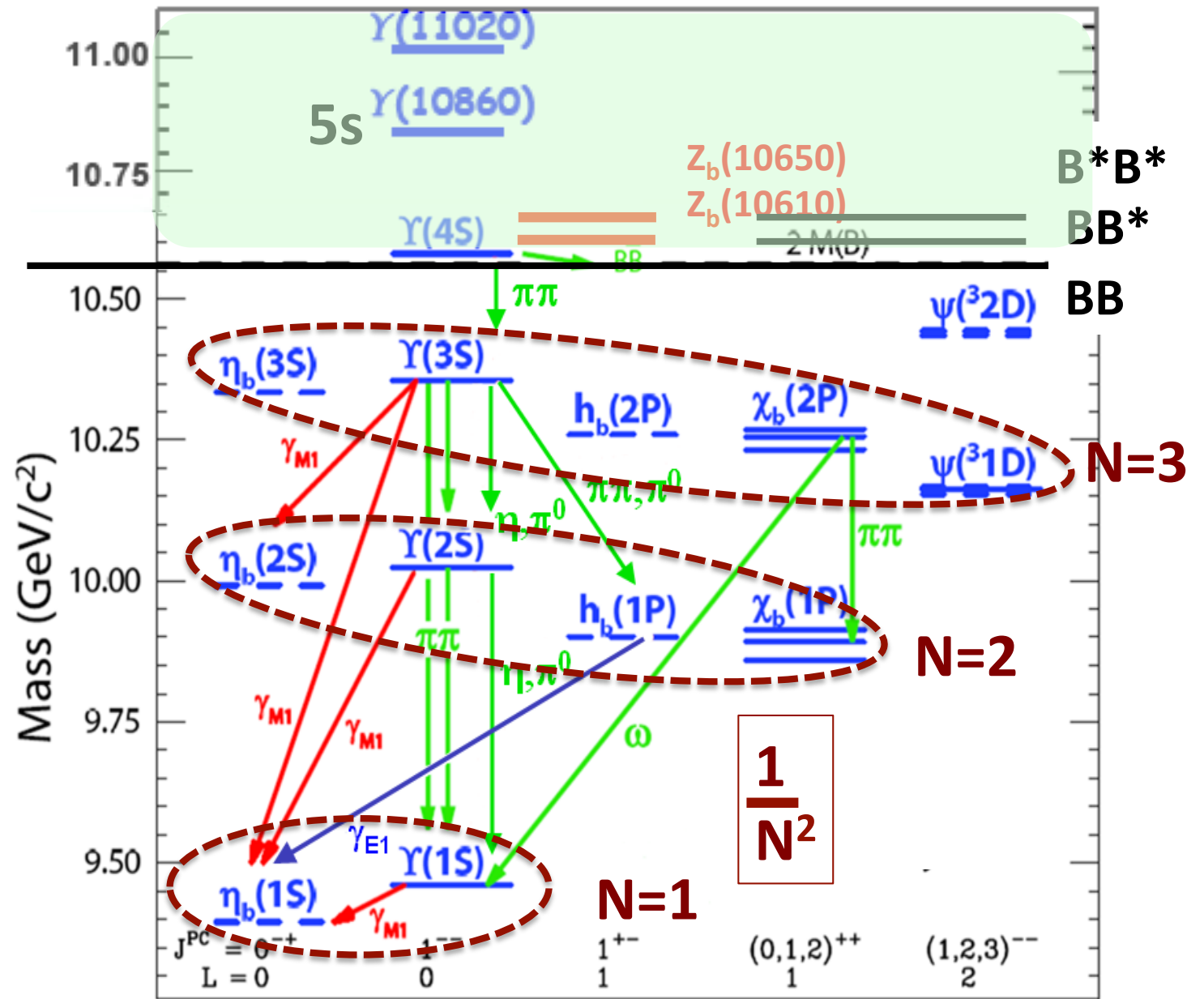
Baikal WS

C quark spectroscopy

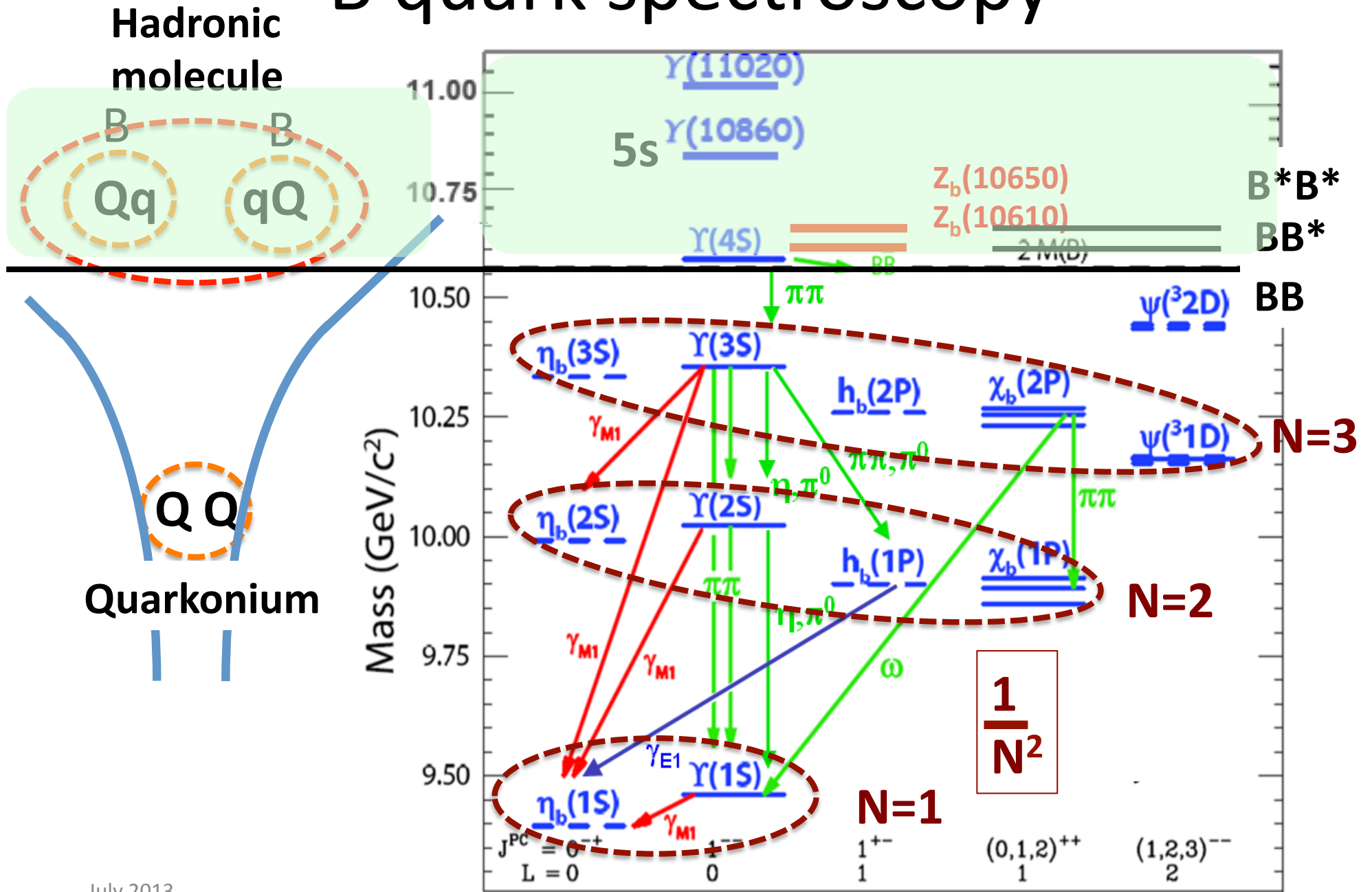
— Z(4430)
— Z_c(3900)



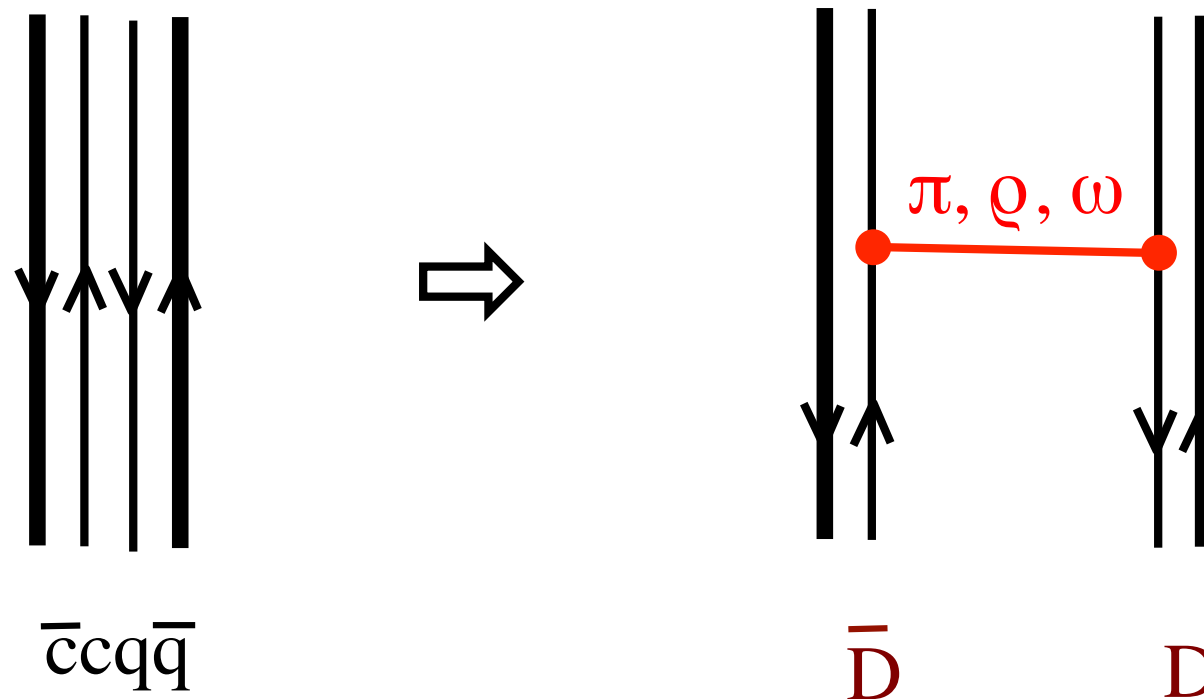
B quark spectroscopy



B quark spectroscopy



Clusterized multiquarks and **interactions**



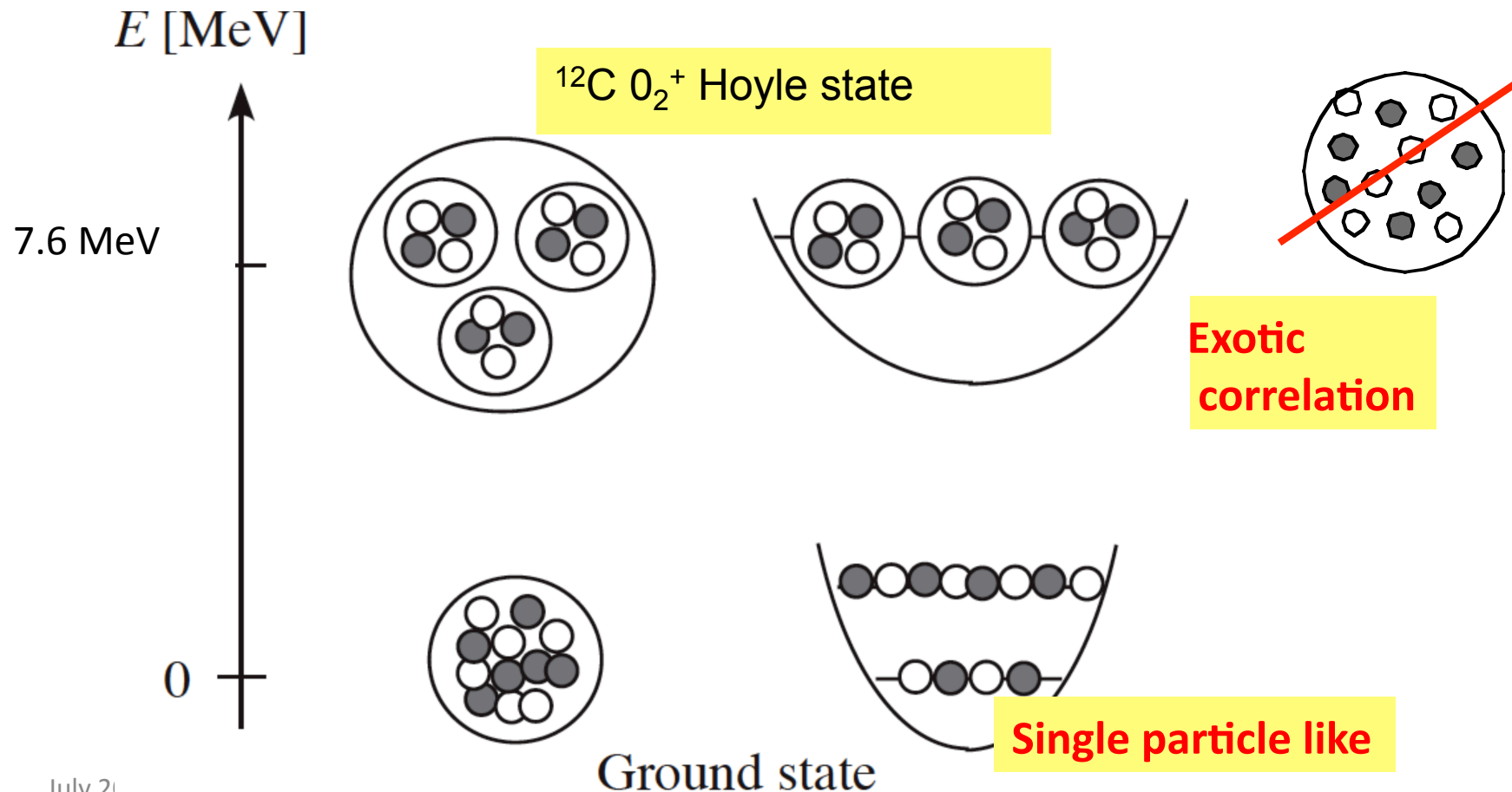
Multiquarks (with light flavor) rearrange into hadrons
Chiral dynamics among light quarks

Two issues

- (1) $q\bar{q}$ creation near the threshold \rightarrow Multiquark configuration
- (2) **Rearrangements of the multiquark configuration \rightarrow Clustering**

M. Itoh et al, PRC84,054308(2011) "Physics Viewpoint"
(RCNP experiment)

Talk by S. Ershov



3. Hadronic molecules with heavy quarks

(1) $D^{\text{bar}}N$ and BN

(2) Z_b and related states

(1) $D^{\text{bar}}N$ and BN

Another form of *exotic* baryons
Hadronic molecules

Yamaguchi, Yamaguchi, Yasui and Hosaka
Phys.Rev.D84:014032,2011,
D85,054003,2012

Ohkoda, Yamaguchi, Sudoh, Yasui and Hosaka
arXiv:1210.3170, in PRD
Phys.Rev. D86: 034019, 2012
D86: 014004, 2012

OPEP induced HQ molecules

Loosely bound or resonance of $\bar{D}N$ ($\bar{c}qqqq$)

- **Heavy Q symmetry:** $\bar{D} \sim \bar{D}^* \rightarrow$ Coupled channel of $\bar{D}N$ and \bar{D}^*N

Yasui-Sudoh, PRD80, 034008, 2009

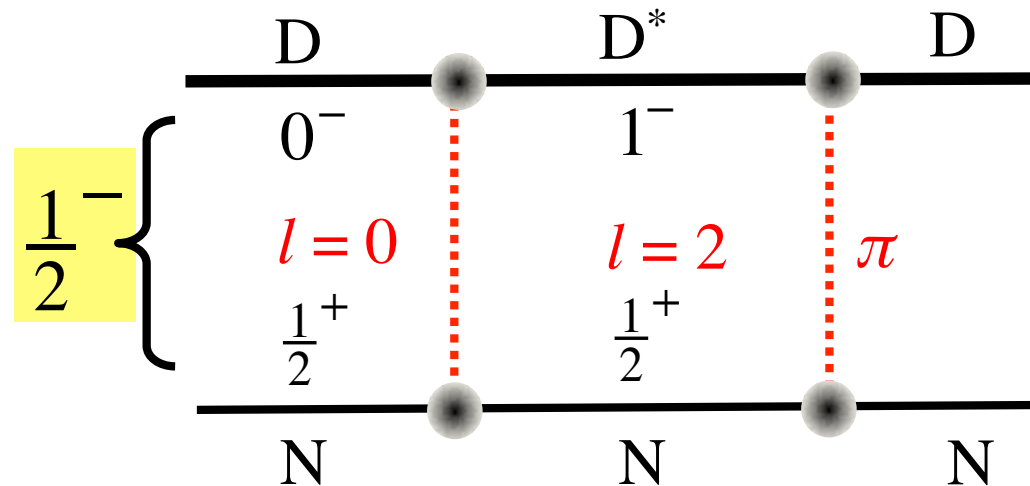
Yamaguchi-Ohkoda-Yasui and Hosaka, PRD84:014032,2011

- **Tensor of OPEP**

$$m_{K^*} - m_K \sim 400 \text{ MeV}$$

$$m_{D^*} - m_D \sim 140 \text{ MeV}$$

$$m_{B^*} - m_B \sim 35 \text{ MeV}$$

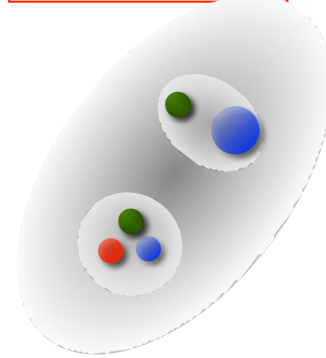


- **Short range**

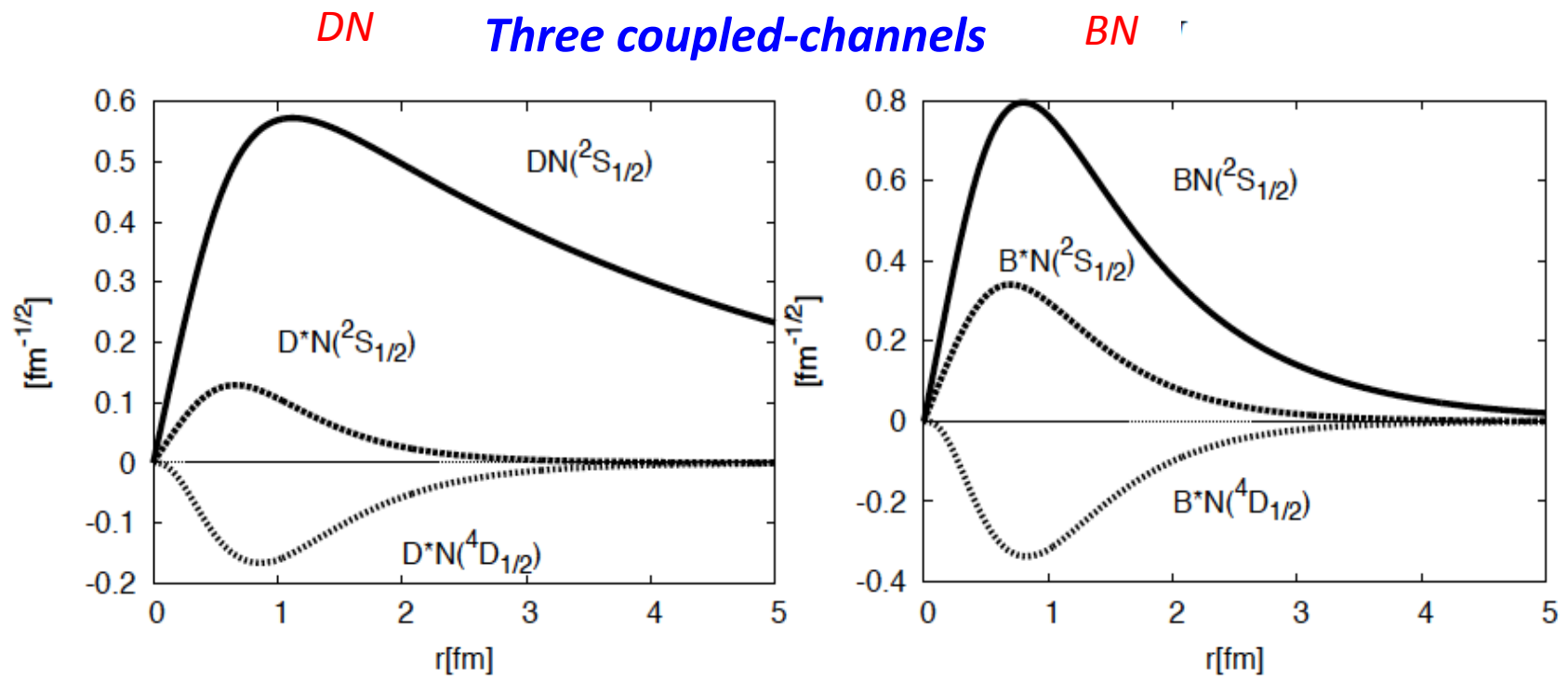
		qq	qq ^{bar} (annihilation)
π	+	ρ (attraction)	attraction
	+	ω (repulsion)	attraction

Loosely bound states: $I, J^P = 0, 1/2^-$

$\bar{Q}q-qqq$



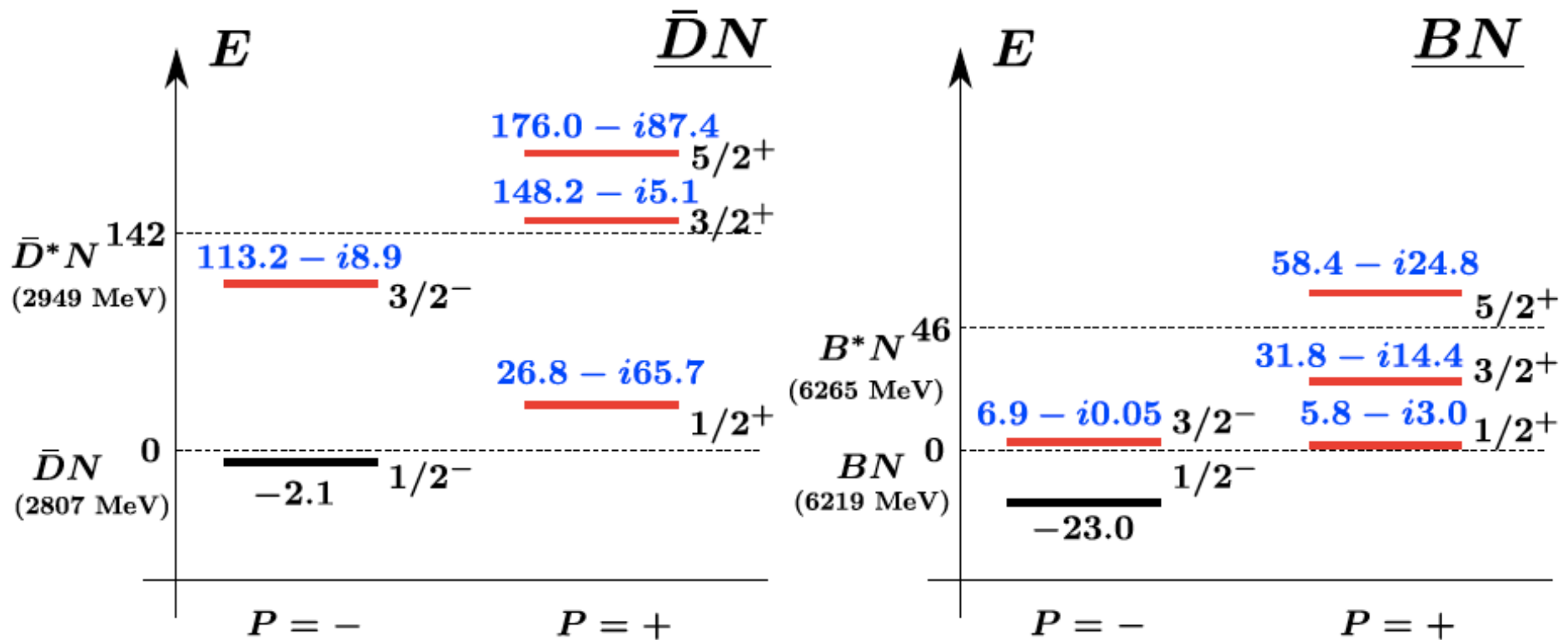
	$\bar{D}N(\pi)$	$\bar{D}N(\pi\rho\omega)$	$BN(\pi)$	$BN(\pi\rho\omega)$
E_B [MeV]	1.60	2.14	19.50	23.04
$\langle r^2 \rangle^{1/2}$ [fm]	3.5	3.2	1.3	1.2



More low-lying states

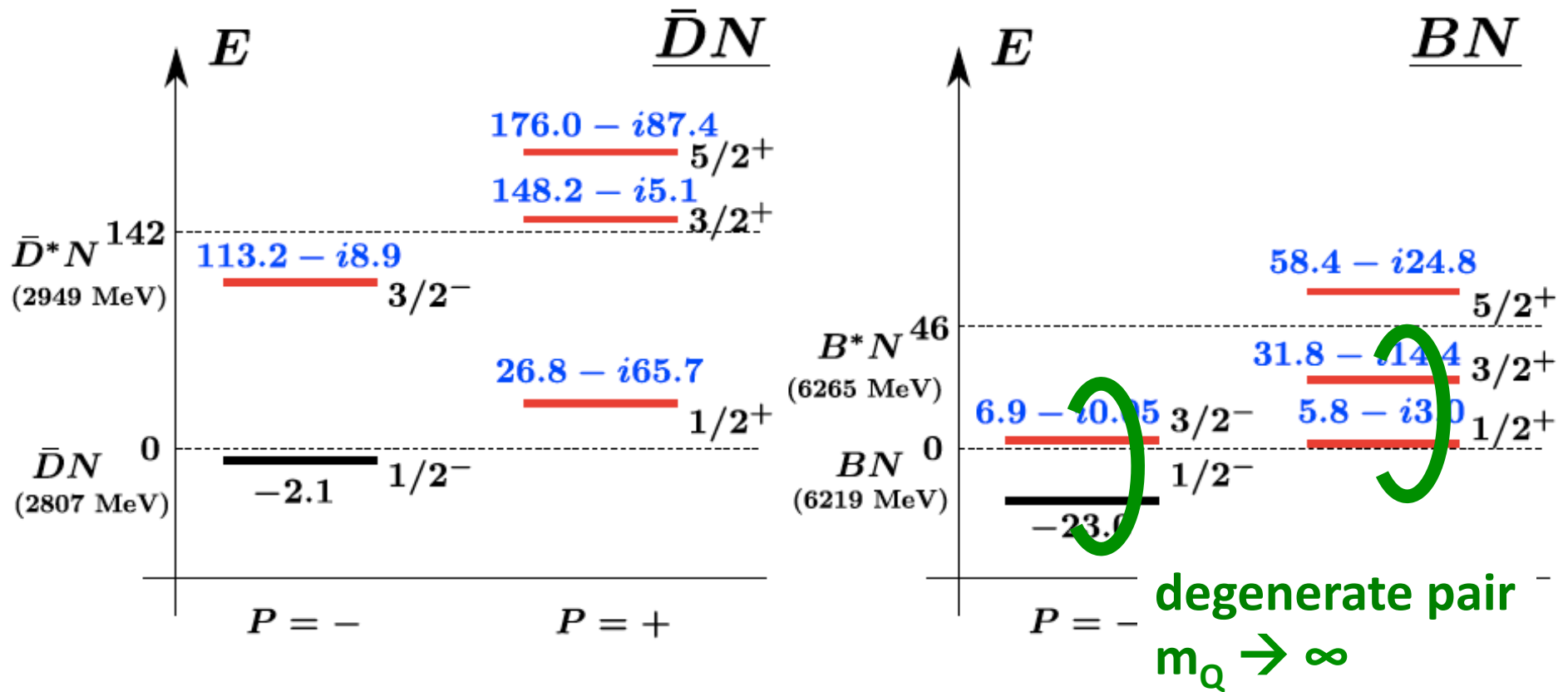
Yamaguchi et al, D85,054003,2012

Various states near the thresholds



More low-lying states

Yamaguchi et al, D85,054003,2012
Various states near the thresholds



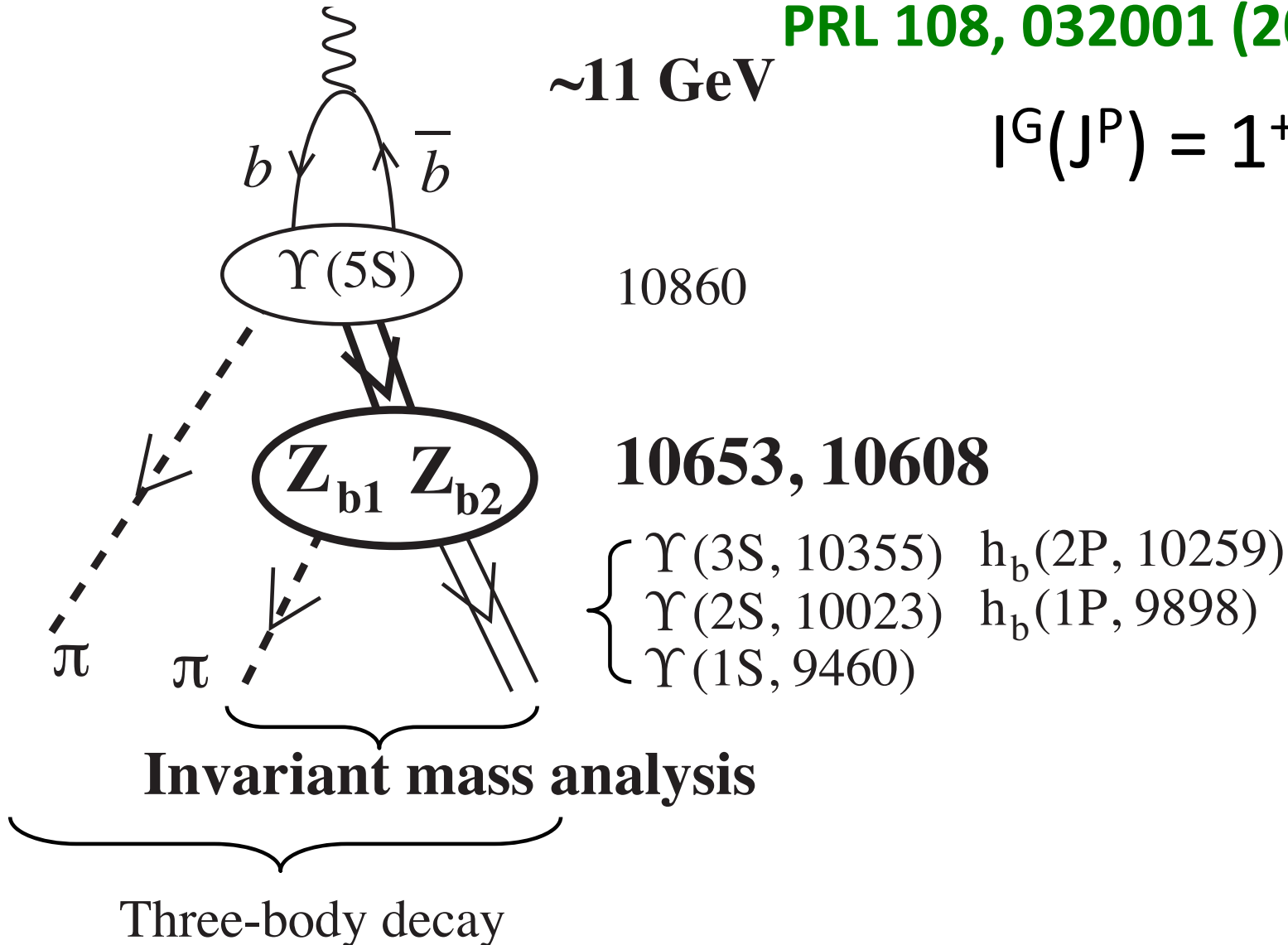
(2) $Z_b(10610, 10650)$ and related states

arXiv:1105.4583v1 [hep-ex];

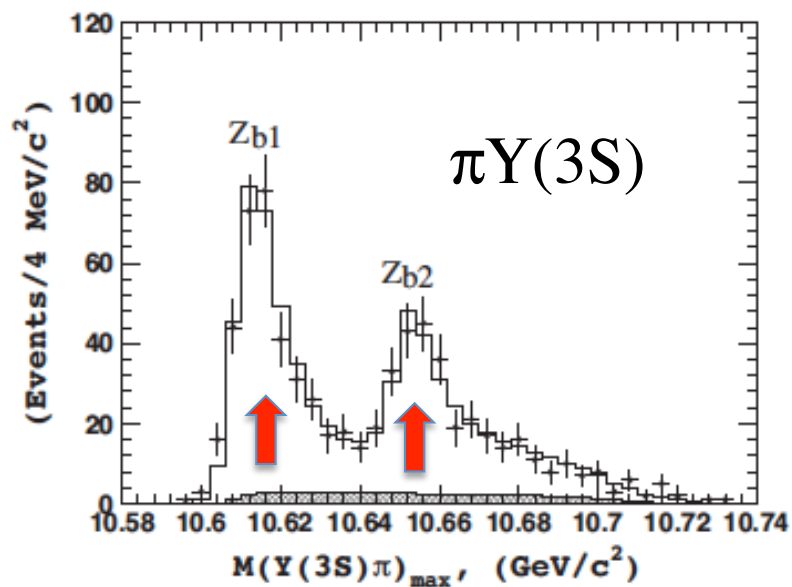
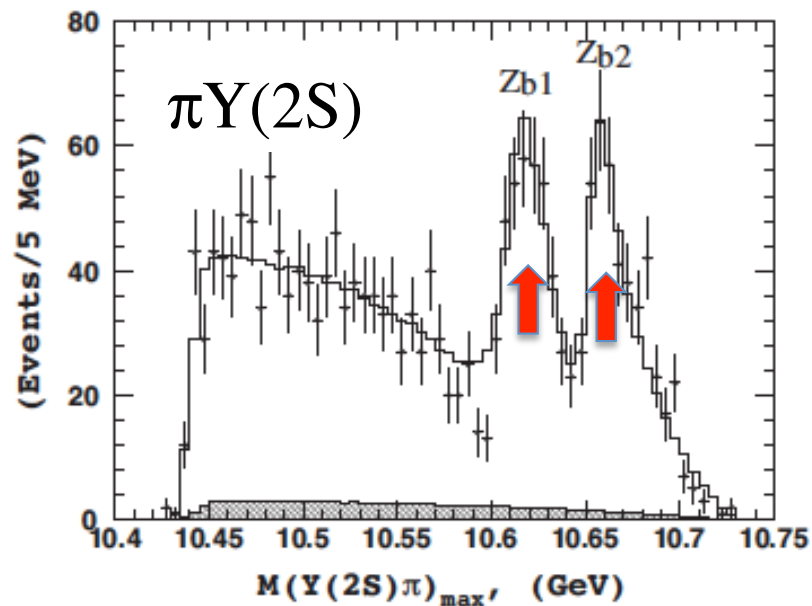
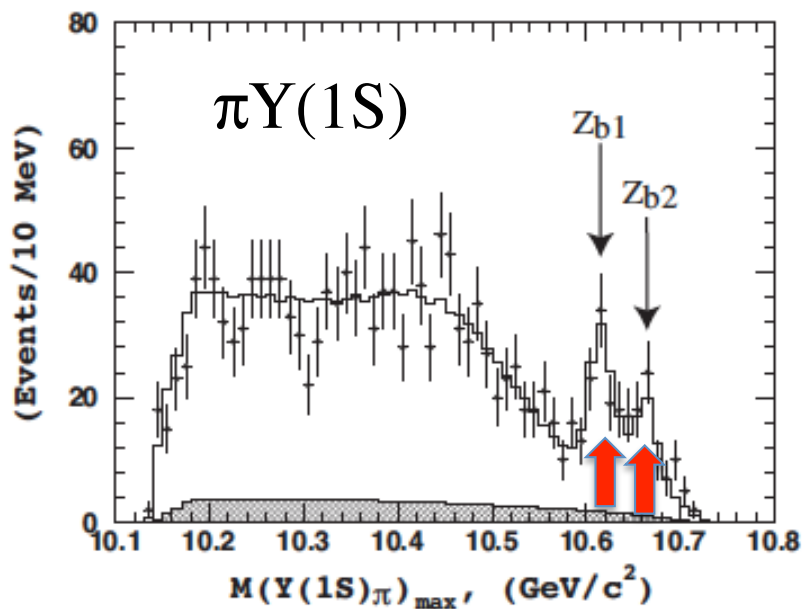
PRL 108, 032001 (2012)

~ 11 GeV

$$I^G(J^P) = 1^+(1^+)$$



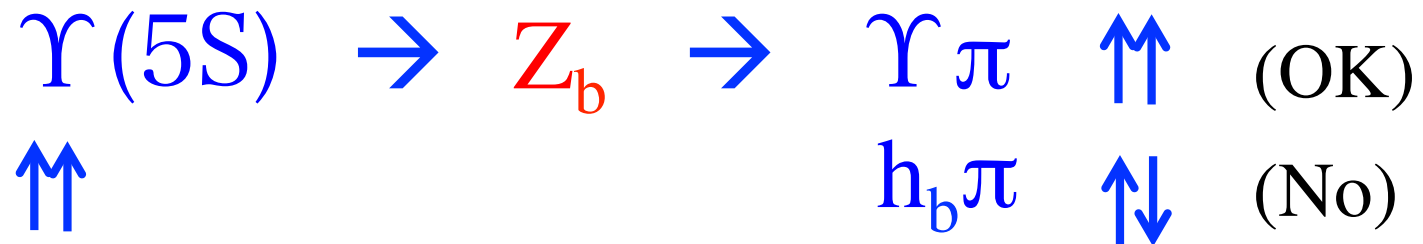
Invariant mass of $\pi Y(nS)$



In all cases, ***twin*** peaks are observed!

Characters

- States appear near the thresholds
- Masses of $Z_b(10610)$, $Z_b(10650)$ are similar
- Heavy spin changing processes occur

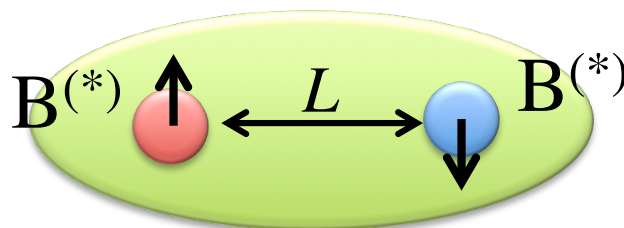


Explained by Z_b molecules

$B^{(*)}\bar{B}^{(*)}$ molecules

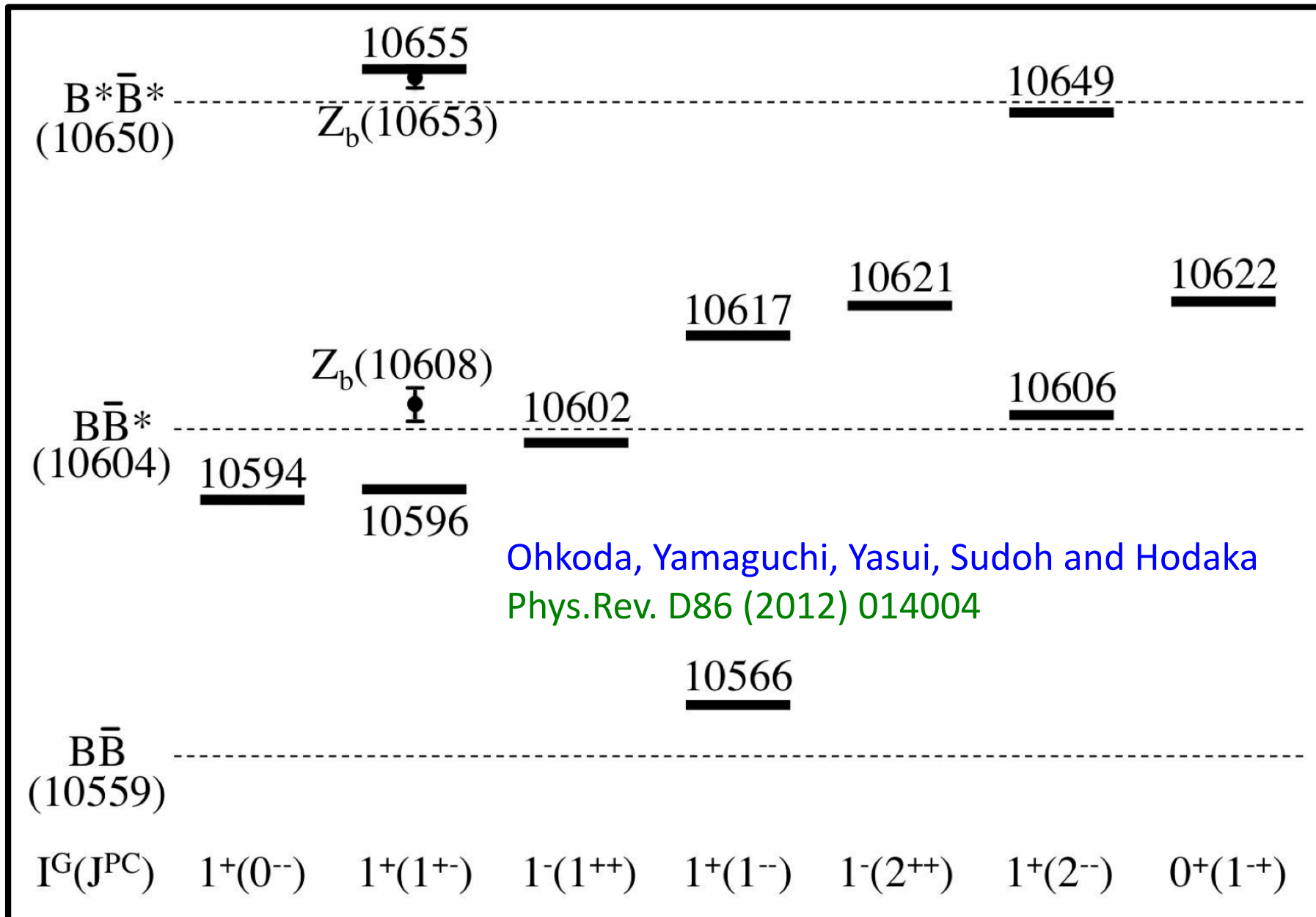
Ohkoda, Yamaguchi, Yasui, Sudoh and Hodaka
 arXiv:1111.2921 [hep-ph], Phys.Rev. D86 (2012) 014004

Rich structure due to $B\bar{B}$ (or $Q\bar{Q}q\bar{q}$)

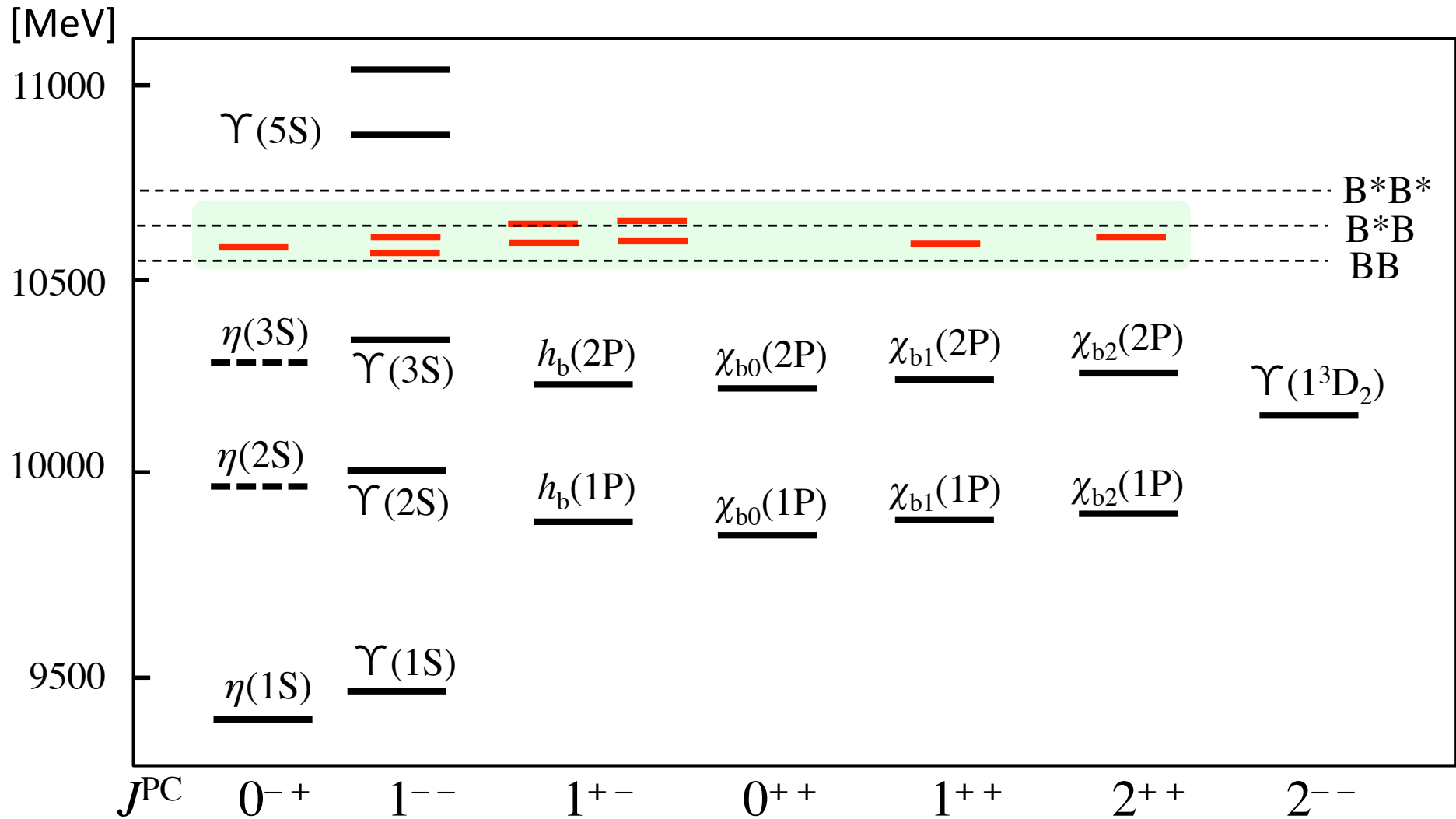


$[S, L]^J$		spin		
		$B\bar{B}$ $S = 0$	$B\bar{B}^*$ $S = 1$	$B^*\bar{B}^*$ $S = 0, 1, 2$
Orbital L				
$L = 0$		0^-	1^-	$0^-, 1^-, 2^-$
$L = 1$		1^+	$0^+, 1^+, 2^+$	$0^+, 1^+, 2^+, 3^+$
.....			

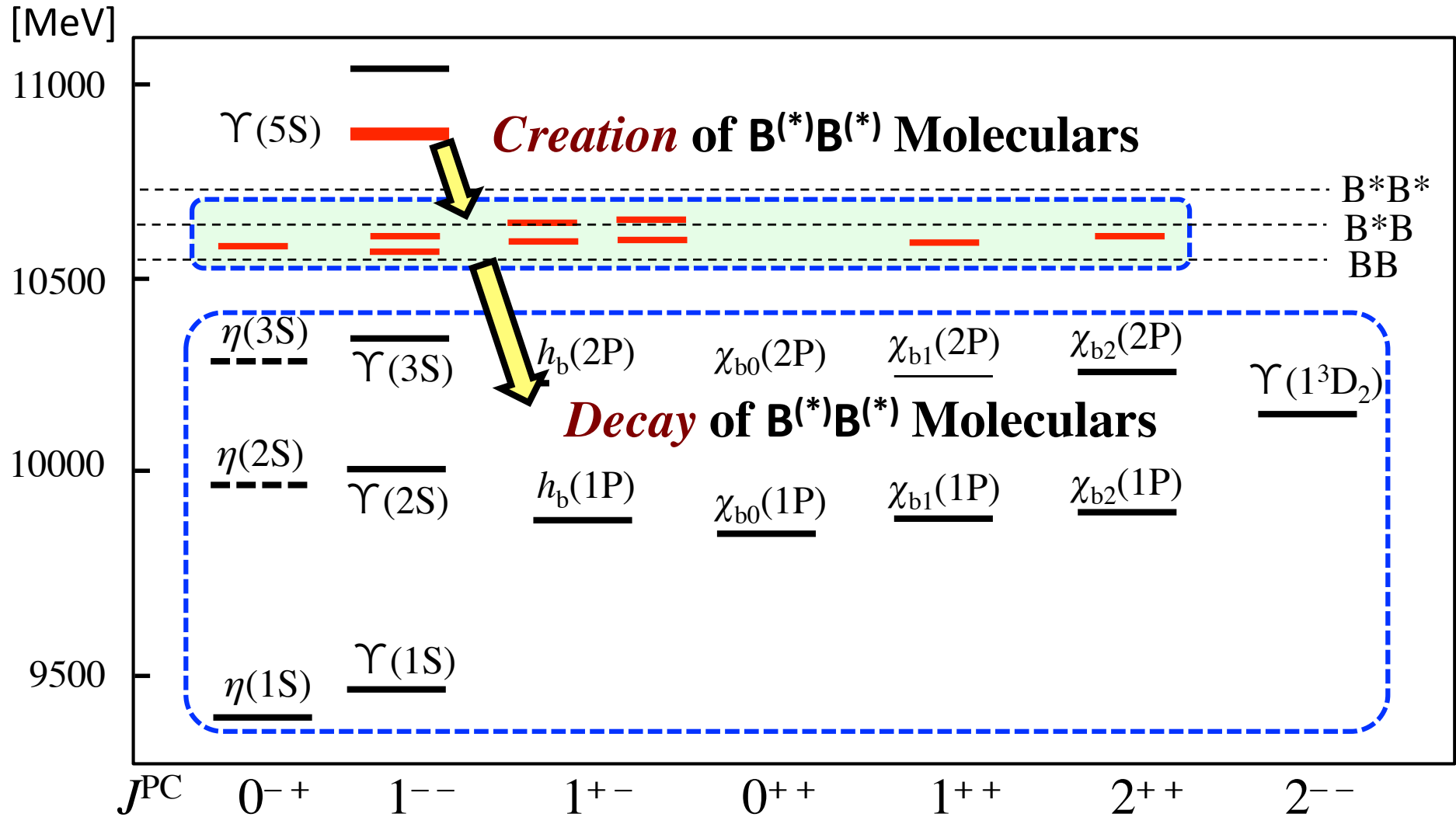
Predictions: Low-lying states, $L \leq 2$



Bottomonia and molecules near the thresholds



Transitions and selection rules



HQ selection rules

Ohkoda-Yamaguchi-Yasui-Hosak,
arXiv:1210.3170, To appear in PRD

$$J_{tot} = J_H + j_l$$

Heavy-light Recoupling: $[[\underline{J_{H1} j_{l1}}][\underline{J_{H2} j_{l2}}]]^{J_{tot}} \rightarrow [[\underline{J_{H1} J_{H2}}][\underline{j_{l1} j_{l2}}]]^{J_{tot}}$

$$\begin{aligned} Z_b(10650) \quad B^* \bar{B}^* (^3S_1) &: \underline{[[b\bar{q}]^1, [\bar{b}q]^1]^1} \\ &= \sum_{H,l} \hat{1}\hat{1}\hat{H}\hat{l} \begin{Bmatrix} 1/2 & 1/2 & 1 \\ 1/2 & 1/2 & 1 \\ H & l & 1 \end{Bmatrix} [[b\bar{b}]^H, [\bar{q}q]^l]^1 \\ &= \frac{1}{\sqrt{2}} (\underline{0_H^-} \otimes \underline{1_l^-}) + \frac{1}{\sqrt{2}} (\underline{1_H^-} \otimes \underline{0_l^-}) \end{aligned}$$

$$Z_b(10610) \quad \frac{1}{\sqrt{2}} (B\bar{B}^* - B^*\bar{B}) (^3S_1) : \frac{1}{\sqrt{2}} (\underline{0_H^-} \otimes \underline{1_l^-}) - \frac{1}{\sqrt{2}} (\underline{1_H^-} \otimes \underline{0_l^-})$$

Example: $Z_b^0 \rightarrow \chi_{bJ} \gamma$

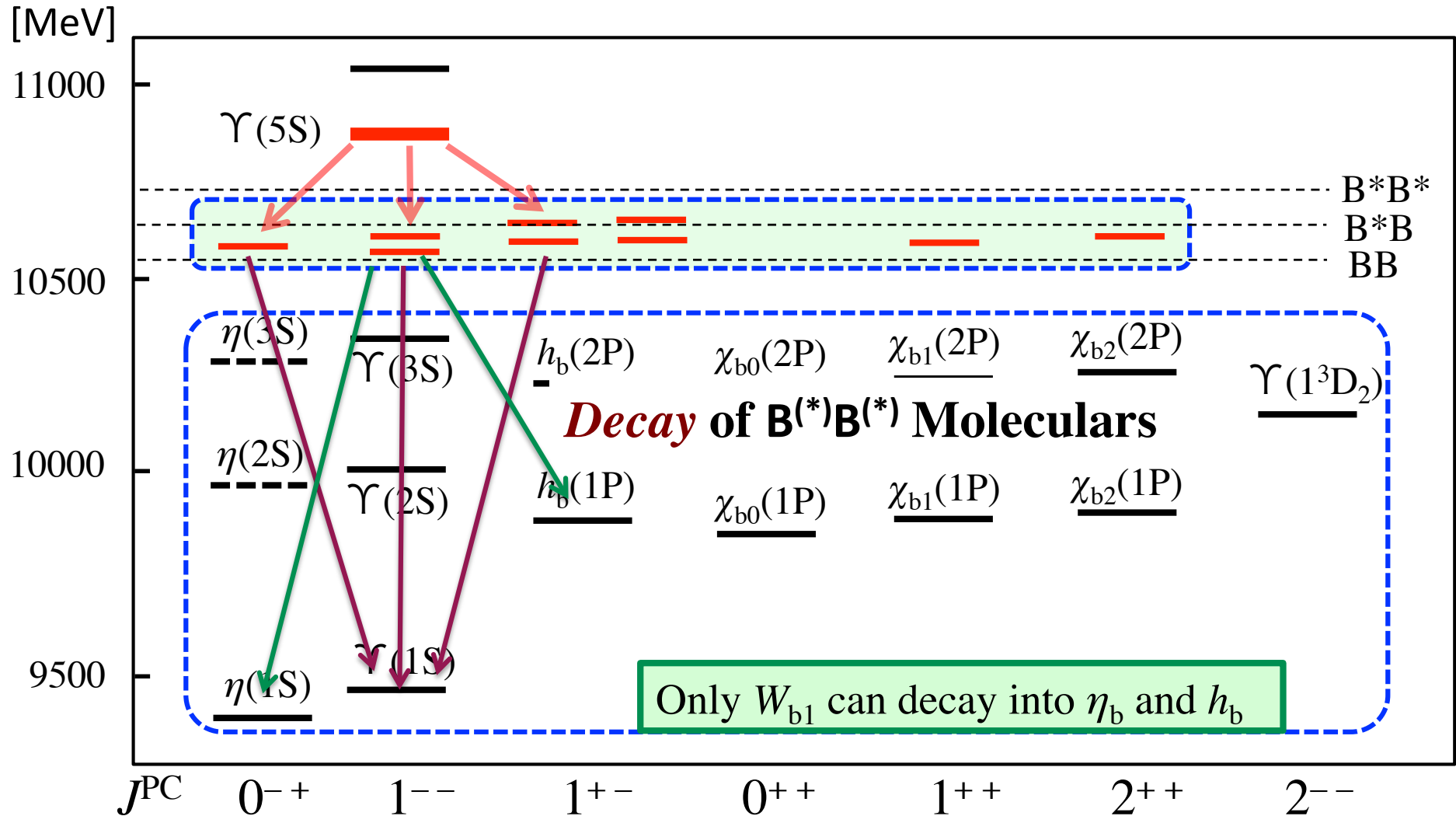
Heavy-light recoupling

$$\begin{aligned}
 \text{M1 } \chi_{b0} \gamma (1^+) &= (1_H \otimes 1_l)|_{J=0} \otimes (0_H \otimes 1_l) \\
 &= \frac{1}{3} \underline{(1_H \otimes 0_l)} - \frac{1}{\sqrt{3}} (1_H \otimes 1_l)|_{J=1} + \frac{\sqrt{5}}{3} (1_H \otimes 2_l)|_{J=1} \\
 \dots & \\
 \text{M1 } \chi_{b1} \gamma (1^+) &\sim -\frac{1}{\sqrt{3}} \underline{(1_H \otimes 0_l)} + \frac{1}{2} (1_H \otimes 1_l)|_{J=1} + \frac{15}{6} (1_H \otimes 2_l)|_{J=1} \\
 \text{E2 } \chi_{b1} \gamma (2^+) &\sim -\frac{1}{2} (1_H \otimes 1_l)|_{J=1} + \frac{\sqrt{3}}{2} (1_H \otimes 2_l)|_{J=1} \\
 \dots & \\
 \text{M1 } \chi_{b2} \gamma (1^+) &\sim \frac{\sqrt{5}}{3} \underline{(1_H \otimes 0_l)} + \frac{\sqrt{15}}{6} (1_H \otimes 1_l)|_{J=1} + \frac{1}{6} (1_H \otimes 2_l)|_{J=1} \\
 \text{E2 } \chi_{b2} \gamma (2^+) &\sim \frac{\sqrt{3}}{2} (1_H \otimes 1_l)|_{J=1} + \frac{1}{2} (1_H \otimes 2_l)|_{J=2}
 \end{aligned}$$

$\Gamma(Z_b^0 \rightarrow \chi_{b0} \gamma)$:	$\Gamma(Z_b^0 \rightarrow \chi_{b1} \gamma)$:	$\Gamma(Z_b^0 \rightarrow \chi_{b2} \gamma)$
1	:	3	:	5

$$f(W_{b0}^{--}\pi) : f(W_{b1}'^{--}\pi) : f(W_{b1}^{--}\pi) : f(W_{b2}'^{--}\pi) : f(W_{b2}^{--}\pi)$$

$$2 : 9 : 4.5 : 9 : 12$$



$$\Gamma(W_{b0}^{--} \rightarrow \Upsilon\pi) : \Gamma(W_{b1}'^{--} \rightarrow \Upsilon\pi) : \Gamma(W_{b1}^{--} \rightarrow \Upsilon\pi) : \Gamma(W_{b2}'^{--} \rightarrow \Upsilon\pi) : \Gamma(W_{b2}^{--} \rightarrow \Upsilon\pi)$$

$$4 : 1 : 1 : 3 : 1$$

Summary

- In the heavy quark region we have seen many unusual states
- Some (many) of them are candidates of **hadronic molecules**
Chiral dynamics through the pion and **heavy quark symmetry**
- Where are other excitations (gluons and multiquarks other than HM)
- Heavy quarks may probe diquarks
Manifest in hadron spectroscopy, production-decay properties
Could explain (non) existence of multiquark hadrons