

Higher order vertices in HS theory

S. Didenko

(based on work with O.A Gelfond, A.V. Korybut and M. A. Vasiliev)

August 29, 2019 Yerevan

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 - $\Upsilon(\omega, \omega, C, C)$ vertex
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AdS/CFT in HS

- Weak-weak duality which does not require supersymmetry (Sundborg, Klebanov, Polyakov, Leigh, Petkou, Sezgin, Sundell)

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- Giombi and Yin tests from equations of motion: substantial piece of evidence that many of 3pt functions match.
- Generic structure of 3pt-correlators (Maldacena, Zhiboedov)

$$\langle JJJ \rangle = \cos^2 \phi \langle JJJ \rangle_b + \sin^2 \phi \langle JJJ \rangle_f + \frac{1}{2} \sin(2\phi) \langle JJJ \rangle_o$$

Locality

$$S = \int \phi \square \phi + \overbrace{D..D\phi D..D\phi D..D\phi}^{\text{local}} + \overbrace{D..D\phi D..D\phi D..D\phi D..D\phi}^{(\text{non-})\text{local?}}$$

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- Quartic order indicates possible non-localities from
holographic reconstruction Bekaert, Erdmenger, Ponomarev,
Sleight; Sleight, Taronna

(Non)locality

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$$d\omega(Y|x) + \omega * \omega = \Upsilon(\omega, \omega, C) + \Upsilon(\omega, \omega, C, C) + \dots,$$

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Perturbative vertices on AdS background Ω

$$\Upsilon(\Omega, \Omega, C) \text{ - free; } \Upsilon(\Omega, \omega, C), \Upsilon(\Omega, C, C), \Upsilon(\Omega, \Omega, C, C) \text{ - cubic}$$

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General solution is given by Vasiliev equations

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$$f(y) * g(y) = f(y) e^{i\epsilon^{\alpha\beta} \overleftarrow{\partial}_\alpha \overrightarrow{\partial}_\beta} g(y),$$

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- How come cubic vertex is local?

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$$D_\Omega \omega = (e \wedge e)^{\alpha\beta} \partial_\alpha \partial_\beta C(0).$$

Goals and summary

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- Calculate explicitly some lower and higher order interaction vertices.

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- Vertices $\Upsilon(\omega, \omega, C)$ and $\Upsilon(\omega, \omega, C, C)$ are found to have ultra-local form.
- The meaning of parameter β is revealed.

Vasiliev equations

Vasiliev equations in $d = 4$

$$d_x W + W * W = 0,$$

$$d_x S + [W, S]_* = 0,$$

$$d_x B + [W, B]_* = 0,$$

$$S * S = -i\theta_\alpha \theta^\alpha + i\eta B * \gamma + c.c.,$$

$$[S, B]_* = 0, \quad \gamma = e^{iz_\alpha y^\alpha} \theta^\beta \theta_\beta$$

$$(f * g)(z, y) = \frac{1}{(2\pi)^2} \int dudv f(z + u; y + u) g(z - v; y + v) e^{iu_\alpha v^\alpha}$$

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$$W(Z, Y|x) = \omega(Y|x) + \dots, \quad B(Z, Y|x) = C(Y|x) + \dots$$

Perturbation theory

Vacuum

$$S^0 = z_\alpha \theta^\alpha, \quad B^0 = 0, \quad W^0 = \omega(Y|x)$$
$$d_x \omega + \omega * \omega = 0.$$

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Higher orders

$$[S^0, f] \sim d_z f$$

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Common equation to solve:

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Conventional homotopy:

$$\Delta_0 J = z^\alpha \frac{\partial}{\partial \theta^\alpha} \int_0^1 \frac{dt}{t} J(tz, y; t\theta)$$

Shifted homotopy

Alternative way of writing solution

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$$q_\alpha = v_1 \partial_{1\alpha} + \dots + v_n \partial_{n\alpha} + \beta \frac{\partial}{\partial y^\alpha}$$

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$$\Delta_{q,\beta} J = \int e^{iu_\alpha v^\alpha} \int_0^1 \frac{d\tau}{\tau} (z + q - v)^\alpha \frac{\partial}{\partial \theta^\alpha} J(\tau z - (1-\tau)(q - v); y + \beta u; \tau \theta)$$

Homotopy properties

- Resolution of identity

$$\{d_z, \Delta_{q,\beta}\} = 1 - h_{q,\beta}, \quad h_{q,\beta} J(z, y; \theta) = \int e^{iu_\alpha v^\alpha} J(-q+v; y+\beta u; \mathbf{0})$$

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- anti-commutativity

$$\Delta_{p,\beta_1} \Delta_{q,\beta_2} = -\Delta_{q,\beta_2} \Delta_{p,\beta_1}, \quad h_{p,\beta_1} \Delta_{q,\beta_2} = -h_{q,\beta_2} \Delta_{p,\beta_1}$$

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- triangle identity

$$h_{d,\beta_1} \Delta_{c,\beta_2} \Delta_{b,\beta_3} + h_{d,\beta_1} \Delta_{b,\beta_3} \Delta_{a,\beta_4} = h_{c,\beta_2} \Delta_{b,\beta_3} \Delta_{a,\beta_4} + h_{d,\beta_1} \Delta_{c,\beta_2} \Delta_{a,\beta_4}$$

Homotopy properties

- Scaling property

$$h_{(1-\beta)a,0}\Delta_{(1-\beta)b,0}\Delta_{(1-\beta)c,0}f(z,y) = (1-\beta)^2 h_{a,0}\Delta_{b,0}\Delta_{c,0}f((1-\beta)z,y)$$

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- Klein property

$$\Delta_{q,\beta\gamma} = \Delta_{\frac{q}{1-\beta},0}\gamma.$$

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- Star exchange

$$\Delta_{\hat{q},\beta}(f(y) * J(z; y; \theta)) = f(y) * \Delta_{\hat{q}+\hat{p},\beta} J(z; y; \theta), \quad \hat{q} = (1-\beta)q, \quad \hat{p} = (1-\beta)p$$

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$$p_\alpha f(y) \equiv f(y) p_\alpha := -i \frac{\partial}{\partial y^\alpha} f(y)$$

Locality theorem

- **Structure theorem (O.A. Gelfond)**. There is the *even* one-form sector of HS equations (**W** and **S**) that is respected by $\Delta_{s_e, \beta}$ – homotopies and the *odd* zero-form one (**B**) respected by $\Delta_{s_o, \beta}$

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$$\det P_{ij} = 0, \quad \exp P_{ij} \partial_\alpha^i \partial^{j\alpha} \Rightarrow \text{measure of non-locality}$$

Strategy

- Same homotopy operators for all fields. $\beta = 1$ (De Filippi, lazeolla, Sundell)

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- Keeping fixed arbitrary β . Locality is reached in the limit $\beta \rightarrow -\infty$

$\Upsilon(\omega, \omega, C)$ - vertex

$$d_z S_1 = -\frac{\eta}{2} C * \gamma \quad \Rightarrow \quad S_1 = -\frac{\eta}{2} \Delta_{0,\beta}(C * \gamma) = -\frac{\eta}{2} C * \Delta_p \gamma,$$

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$$d_x C + [\omega, C]_* = 0,$$

$$W_1 = -\frac{\eta}{4i} (C * \omega * \Delta_{p+t} \Delta_{p+2t} \gamma - \omega * C * \Delta_{p+t} \Delta_p \gamma),$$

$$p_\alpha C(Y) := i \frac{\partial}{\partial y^\alpha} C(Y) \quad t_\alpha \omega(Y) := i \frac{\partial}{\partial y^\alpha} \omega(Y)$$

$\Upsilon(\omega, \omega, C)$ - vertex

$$d_x \omega + \omega * \omega = \frac{\eta}{4i} (\omega * \omega * C * X_{\omega\omega} C + C * \omega * \omega * X_{C\omega\omega} + \omega * C * \omega * X_{\omega C\omega}),$$

$\Upsilon(\omega, \omega, C)$ - vertex

$$d_X \omega + \omega * \omega = \frac{\eta}{4i} (\omega * \omega * C * X_{\omega\omega C} + C * \omega * \omega * X_{C\omega\omega} + \omega * C * \omega * X_{\omega C\omega}),$$

where

$$X_{\omega\omega C} = h_{p+t_1+t_2} \Delta_p \Delta_{p+t_2} \gamma,$$

$$X_{C\omega\omega} = h_{p+t_1+t_2} \Delta_{p+t_1+2t_2} \Delta_{p+2t_1+2t_2} \gamma,$$

$$X_{\omega C\omega} = -h_{p+t_1+t_2} \Delta_{p+t_1+2t_2} \Delta_{p+t_2} \gamma - h_{p+t_1+2t_2} \Delta_{p+2t_2} \Delta_{p+t_2} \gamma.$$

$\Upsilon(\omega, \omega, C)$ - vertex

$$d_X \omega + \omega * \omega = \frac{\eta}{4i} (\omega * \omega * C * X_{\omega\omega C} + C * \omega * \omega * X_{C\omega\omega} + \omega * C * \omega * X_{\omega C\omega}),$$

where

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explicitly

$$\Upsilon_{\omega\omega C} = \frac{\eta}{2i} \int_{[0,1]^3} d^3\tau \delta(1 - \tau_1 - \tau_2 - \tau_3) e^{i(1-\tau_3)\partial_1^\alpha \partial_2^\alpha}$$

$$\partial^\alpha \omega((1 - \tau_1)y) \partial_\alpha \omega(\tau_2 y) C(-i\tau_1 \partial_1 - i(1 - \tau_2)\partial_2),$$

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$$\partial^\alpha \omega((1 - \tau_1)y) \partial_\alpha \omega(\tau_2 y) C(-i\tau_1 \partial_1 - i(1 - \tau_2) \partial_2),$$

No y - dependence in $C(y)$ \Rightarrow ultra local form

$\Upsilon(\omega, C, C)$ - vertex

Solving for $B(z, y)$ to the second order

$$B_2 := B_2^q = \frac{\eta}{4i} \Delta_{(1-\beta)q} (C * C * (\Delta_{p_2} - \Delta_{p_1+2p_2}) \gamma), \quad q = v_1 p_1 + v_2 p_2$$

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$$v_2 - v_1 = 1$$

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$$X_{\omega CC} = h_{p_2} \Delta_{p_1+2p_2} \Delta_{p_1+2p_2+t} \gamma,$$

$$X_{CC\omega} = h_{p_2+2t} \Delta_{p_2+t} \Delta_{p_1+2p_2+2t} \gamma,$$

$$X_{C\omega C} = (h_{p_1+2p_2+2t} - h_{p_2}) \Delta_{p_2+t} \Delta_{p_1+2p_2+t} \gamma$$

$\Upsilon(\omega, C, C)$ – vertex

Explicitly,

$$\Upsilon_{\omega CC} = \frac{\eta}{2i} \int_{[0,1]^3} d^3\tau \delta(1 - \tau_1 - \tau_2 - \tau_3) (\partial_1^\alpha + \partial_2^\alpha) \partial_\alpha^\omega \omega((1 - \tau_3)y) C(\tau_1 y - i(1 - \tau_2)\partial^\omega) C(-(1 - \tau_1)y + i\tau_2\partial^\omega),$$

The result is perfectly local in accordance with PLT

Generalities

Extracting $\Upsilon(\omega, \omega, C, C)$ from equations of motion

$$d_x \omega + \omega * \omega = -\{\omega, W_1\}_* - (d_x W_1 + d_x W_2 + W_1 * W_1 + \{\omega, W_2\}_*) + O(C^3)$$

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$$\Upsilon_{C\omega\omega C}^{\eta\eta}(\beta) = -h_{0,\beta}(W_1 * W_1)|_{C\omega\omega C}$$

calculation

$$h_{0,\beta}(W_1 * W_1) \sim \int_0^1 d\tau_{1,2} \frac{\beta^2(1-\tau_1)(1-\tau_2)}{(1-\beta(\tau_1 \circ \tau_2))^4} e^{\frac{i}{1-\beta(\tau_1 \circ \tau_2)} A^\alpha (y + \beta B)_\alpha + iC}$$

$$\tau_\circ \equiv \tau_1 \circ \tau_2 = \tau_1(1-\tau_2) + \tau_2(1-\tau_1)$$

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$\lim_{\beta \rightarrow -\infty} h_{0,\beta}(W_1 * W_1)$ exists and is ultra-local!

$$I(\beta) = \int_{[0,1]^2} d\tau_{1,2} \frac{\beta^2(1-\tau_1)(1-\tau_2)}{(1-\beta\tau_\circ)^4} \exp\left(-\frac{i\beta}{\xi} \tau_1(1-\tau_2)A - \frac{i\beta}{\xi} \tau_2(1-\tau_1)B + \frac{i}{\xi} C\right)$$

calculation

$$h_{0,\beta}(W_1 * W_1) \sim \int_0^1 d\tau_{1,2} \frac{\beta^2(1-\tau_1)(1-\tau_2)}{(1-\beta(\tau_1 \circ \tau_2))^4} e^{\frac{i}{1-\beta(\tau_1 \circ \tau_2)} A^\alpha (y+\beta B)_\alpha + iC}$$

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$$I(\beta) = \int_{[0,1]^2} d\tau_{1,2} \frac{\beta^2(1-\tau_1)(1-\tau_2)}{(1-\beta\tau_\circ)^4} \exp\left(-\frac{i\beta}{\xi} \tau_1(1-\tau_2)A - \frac{i\beta}{\xi} \tau_2(1-\tau_1)B + \frac{i}{\xi} C\right)$$

$$I(-\infty) = \int d\Delta^3 \tau \tau_3 \exp(i\tau_1 A + i\tau_2 B + i\tau_3 C)$$

Result

$$\Upsilon_{C\omega\omega C}^{\eta\eta} = \frac{\eta^2}{4} \int_{[0,1]^2} d\sigma d\sigma' \sigma\sigma' \int d^3_{\Delta}\tau (t_{1\alpha}t_2^\alpha)^2$$

$$\tau_3 \exp \left[-i(\tau_2\sigma + \tau_1\sigma' + \tau_3\sigma\sigma')t_{1\alpha}t_2^\alpha \right] C(-\tau_1\sigma't_2 - (1 - \sigma(1 - \tau_1))t_1)$$

$$\omega(-\tau_3\sigma y)\omega(-\tau_3\sigma'y)C(\tau_2\sigma t_1 + (1 - \sigma'(1 - \tau_2))t_2),$$

$$t_{1,2} = -i\partial_{\omega_{1,2}}$$

Properties

- Ultra-locality of (anti)holomorphic vertices
 $\Upsilon(\omega, \omega, C(0), C(0))$
- Vanish on the gravitational background

$$\Upsilon = 0 \quad \text{for} \quad \omega = \omega_{s \leq 2}$$

What is β ?

- Introduce the following reordering

$$O_{\beta} f(z, y) = \int \frac{dudv}{(2\pi)^2} f(z + v, y + \beta u) \exp(iu_{\alpha} v^{\alpha})$$

$$\Delta_{q,\beta} = O_{\beta}^{-1} \Delta_{q,0} O_{\beta}$$

(For $\beta = 1$ - De Filippi, Iazeolla, Sundell)

- β - star-product

$$f \star_{\beta} g = \int f(z + u', y + u) g(z - (1 - \beta)v - v', y + v + (1 - \beta)v') e^{iu_{\alpha} v^{\alpha} + iu'_{\alpha} v'^{\alpha}}$$

Conclusion

- A β – class of homotopies is introduced for solving Vasiliev equations.
- Part of the quartic and quintic vertices $\Upsilon(\omega, \omega, C)$ and $\Upsilon(\omega, \omega, C, C)$ were explicitly calculated and shown to have ultra local form in the limit $\beta \rightarrow -\infty$. In addition it was shown that lower order vertices $\Upsilon(\omega, \omega, C)$ and $\Upsilon(\omega, C, C)$ are β -independent
- Holomorphic part of $\Upsilon(\omega, \omega, C, C)$ vanishes on any gravitational background.
- Interpretation of β as the reordering parameter is proposed.