

$D = 4$ Black Holes From Geodesics

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[A. Ploegh](#), [A. Sorin](#), [T. Van Riet](#), [J. Rosseel](#)

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Seed solution: Simplest solution with all duality invariant properties of the most general one

Static, Asymptotically Flat Black Holes in D=4 SUGRAS

Bosonic field content

- n_S scalar fields ϕ^r ($r = 1, \dots, n_S$)
- n_V vector fields A_μ^Λ ($\Lambda = 0, \dots, n_V - 1$)
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The ansatz

$$ds^2 = -e^{2U} dt^2 + e^{-2U} \left[\frac{c^4}{\sinh^4(c\tau)} d\tau^2 + \frac{c^2}{\sinh^2(c\tau)} (d\theta^2 + \sin(\theta) d\varphi^2) \right]$$

- $\phi^r = \phi^r(\tau)$, $U = U(\tau)$, $\frac{d\tau}{dr} = \frac{\sinh^2(c\tau)}{c^2} = \frac{1}{(r-r_0)^2 - c^2}$;
- c extremality parameter, two horizons: $r_\pm = r_0 \pm c$
- electric and magnetic charges e_Λ , m^Λ : $\Gamma^M \equiv (m^\Lambda, e_\Lambda)$
- Extreme solutions $c = 0$: $\lim_{\tau \rightarrow -\infty} e^{-2U} = \frac{A_H}{4\pi} \tau^2$

Seed solution in maximal SUGRA

- 70 scalar fields $\phi^r \in \mathcal{M}_{scal} = \frac{G_4}{H_4} = \frac{E_{7(7)}}{SU(8)}$
- 28 vector fields A_{μ}^{Λ}
- Duality group is $E_{7(7)}$; $\Gamma^M = (m^{\Lambda}, e_{\Lambda}) \in \mathbf{56}$ symplectic representation
- Parameters of a black-hole encoded in central charges computed at infinity:
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$$Z_{AB} \xrightarrow{SU(8)} \begin{pmatrix} Z_1 \epsilon & 0 & 0 & 0 \\ 0 & Z_2 \epsilon & 0 & 0 \\ 0 & 0 & Z_3 \epsilon & 0 \\ 0 & 0 & 0 & Z_4 \epsilon \end{pmatrix}$$

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- (Z_k) can be identified with the charges $(Z, \bar{Z}_s, \bar{Z}_t, \bar{Z}_u)$ of a $\mathcal{N} = 2$ STU truncation
- Five $SU(8)$ invariants: $\rho_k = |Z_k|$, $\theta = \text{Arg}(Z_1 Z_2 Z_3 Z_4)$
- Seed solution, also solution to the STU truncation, has 5 parameters

$D = 3$ description of $D = 4$ stationary solutions

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- $\mathcal{M}_{scal}^{(3)}$ **pseudo-Riemannian**, *negative signature* directions along $\mathcal{Z}^M = (\zeta^\Lambda, \tilde{\zeta}_\Lambda)$. $H \subset G$ is **non-compact** semisimple
- Spherical symmetry: $\phi^I = \phi^I(\tau)$, **solution is a geodesic** on $\mathcal{M}_{scal}^{(3)}$
- Invariant measure along the geodesic coincides with the extremality parameter: $G_{IJ}(\phi) \dot{\phi}^I \dot{\phi}^J = 2c^2$

Mathematical description of the geodesic

Definitions...

- Let \mathfrak{g} , \mathfrak{h} be the Lie algebras of G and H . Involution $\sigma(\mathfrak{h}) = -\eta\mathfrak{h}^T\eta = \mathfrak{h}$ induces the (pseudo-) Cartan decomposition:

$$\mathfrak{g} = \mathfrak{h} \oplus \mathfrak{k}$$

with $\sigma(\mathfrak{k}) = -\mathfrak{k}$

- Given coset representative $\mathcal{V}(\phi^I) \in e^{Solv}$ and geodesic $\phi^I(\tau)$, define $\mathcal{V}(\tau) \equiv \mathcal{V}(\phi^I(\tau))$:

$$\mathcal{V}^{-1}\dot{\mathcal{V}}(\tau) = V(\tau) + W(\tau) , \quad V(\tau) \in \mathfrak{k} , \quad W(\tau) \in \mathfrak{h}$$

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- $\dot{V} + [W, V] = 0$, Lax Pair equation [Liouville i. system (1007.3209)]

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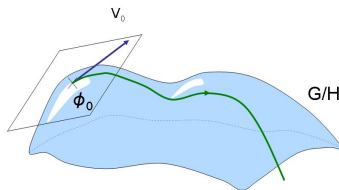
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- $\dot{V} + [W, V] = 0$, Lax Pair equation [Liouville i. system (1007.3209)]
- Q is the Noether charge matrix: $Q = 2\mathcal{V}^{-T}V^T\mathcal{V}^T$

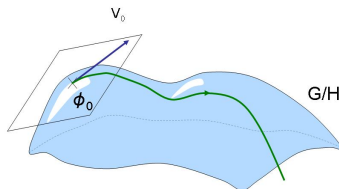
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The global symmetry in $D = 3$

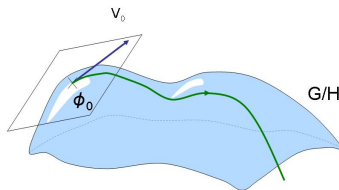
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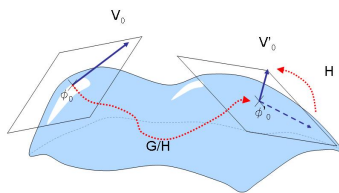
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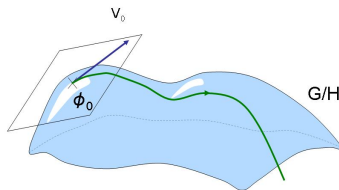


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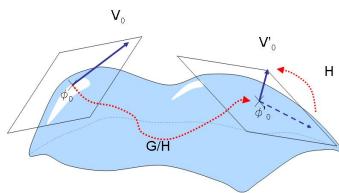


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Some algebra....

- G is larger than G_4 : $G_4 \times SL(2, \mathbb{R})_E \subset G$. Its algebra decomposes as follows:

$$\mathfrak{g} = \mathfrak{sl}(2, \mathbb{R})_E \oplus \mathfrak{g}_4 \oplus (\mathbf{2}, \mathbf{R})$$

where \mathbf{R} is the (symplectic) representation of the electric and magnetic charges under G_4

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- $\mathcal{N} = 8$ example:

$$G_4 = E_{7(7)}, \quad H_4 = \mathrm{SU}(8), \quad G = E_{8(8)}, \quad H = \mathrm{SO}^*(16)$$

$$\mathbf{R} = \mathbf{56} \text{ of } E_{7(7)} \text{ and } \hat{\mathbf{R}} = \mathbf{28}_{-1} + \overline{\mathbf{28}}_{+1} \text{ of } \mathrm{U}(1)_E \times H_4 = \mathrm{U}(8)$$

Seed Geodesic in Universal Submanifold

- Any element of $\mathfrak{K}^{(\hat{R})}$ or $\mathfrak{J}^{(\hat{R})}$ can be rotated by $U(1)_E \times H_4$ into minimal *abelian* subalgebras $\mathfrak{K}^{(N)} = \text{Span}(k_k)$ and $\mathfrak{J}^{(N)} = \text{Span}(J_k)$, where $k = 0, \dots, p-1$ and

$$p = \text{rank} \left(\frac{H}{U(1)_E \times H_4} \right)$$

Together with $H_k = [J_k, k_k]$ they generate $SL(2, \mathbb{R})^p \subset G$

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$$V_0 \cap \mathfrak{K}^{(\hat{R})} = Z_{AB} k^{AB} - \bar{Z}^{AB} k_{AB} \xrightarrow{U(8)} \sum_{k=1}^4 \rho_k k_k \quad (\text{normal form of } \mathbf{28}_{+1})$$

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- Seed geodesic is a p -charge solution within (arXiv:0806.2310)

$$\mathcal{M}_N = \left(\frac{SL(2, \mathbb{R})}{SO(1, 1)} \right)^p \times SO(1, 1)^{r-p} \subset \frac{G}{H}$$

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$$p = \text{rank} \left(\frac{H}{U(1)_E \times H_4} \right)$$

Together with $H_k = [J_k, k_k]$ they generate $SL(2, \mathbb{R})^p \subset G$

- $\mathcal{N} = 8$ example: $p = \text{rank} \left(\frac{SO^*(16)}{U(8)} \right) = 4$

$$V_0 \cap \mathfrak{K}^{(\hat{R})} = Z_{AB} k^{AB} - \bar{Z}^{AB} k_{AB} \xrightarrow{U(8)} \sum_{k=1}^4 \rho_k k_k \quad (\text{normal form of } \mathbf{28}_{+1})$$

Notice θ rotated away! **Seed geodesic is a 4-parameter solution**

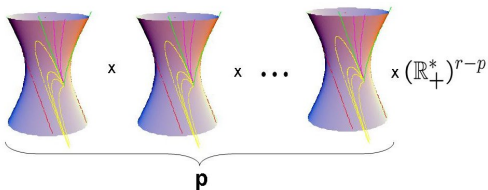
- Seed geodesic is a p -charge solution within (arXiv:0806.2310)

$$\mathcal{M}_N = \left(\frac{SL(2, \mathbb{R})}{SO(1, 1)} \right)^p \times SO(1, 1)^{r-p} \subset \frac{G}{H}$$

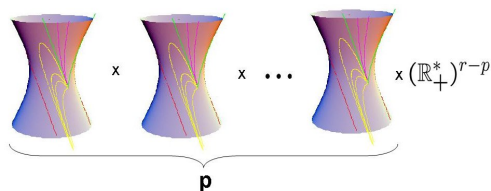
- True for all V_0 -diagonalizable cases. V_0 non-diagonalizable (e.g. extremal solutions) only geodesics originating from regular solutions ($A_H > 0$).

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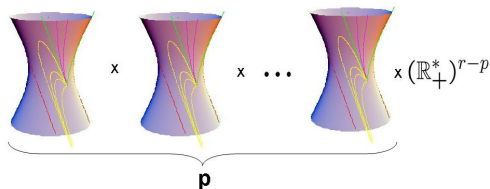


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- Construction of the seed geodesic within a universal submanifold common to a broad class of models. E.g. $p = 4$: $N = 8$, $N = 2$ with rank-3 symmetric SK manifolds (STU) etc.

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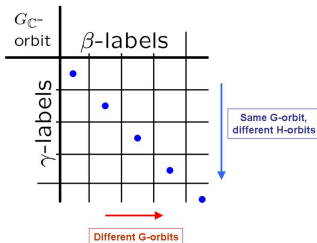
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- $D = 4, N = 2$ SUGRA coupled to 1 vector multiplet
- Complex scalar t in $\frac{SL(2, \mathbb{R})}{SO(2)}$ [$\mathcal{F}(t) = t^3$] coupled to 2 vectors; 4 charges
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- Explicit matrix representation: $J = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, k = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \mathcal{H} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

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- Solution in terms of U, t, \mathcal{Z}^M :

$$e^{-2U} = \sqrt{\mathbf{H}_0(\mathbf{H}_1)^3}; \quad t = -i \sqrt{\frac{\mathbf{H}_0}{\mathbf{H}_1}}, \quad \mathcal{Z}^0 = \frac{\varepsilon_0 a_0 \tau}{\mathbf{H}_0}; \quad \mathcal{Z}^1 = \sqrt{3} \frac{\varepsilon_1 a_1 \tau}{\mathbf{H}_1}$$

$\mathbf{H}_k = 1 - \sqrt{2} a_k \tau$. Charges are: $e_0 = \varepsilon_0 a_0$, $m^1 = -\varepsilon_1 a_1$

- Regular solution: $a_k > 0 \quad \Rightarrow \quad \beta\text{-label} = \gamma\text{-label}$.

- At the horizon $\tau \rightarrow -\infty$:

$$\varphi_i \rightarrow \varphi_i^{\text{fix}} \text{ (stable attractor), } e^{-2U} \rightarrow \frac{A_H}{4\pi} \tau^2$$

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$$\frac{A_H}{4\pi} = \sqrt{4 a_0 (a_1)^3} = \sqrt{4 \varepsilon e_0 (m^1)^3} = \sqrt{\varepsilon l_4(e, m)}$$

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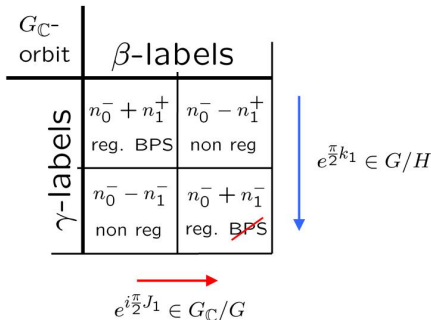
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- $|a_k| = 1$ modulo action of $\text{SO}(1, 1)^2 \subset H$



Other Orbits and Tensor Classifiers

- Other two orbits with $l_4 = 0$, X step-2 nilpotent:
 - $\mathcal{O}_1: m^1 \rightarrow 0$ (doubly-critical) small-bh
 - $\mathcal{O}_2: e_0 \rightarrow 0$ (lightlike) small-bh
- Last orbit: X step-7 nilpotent. No regular single-center, regular multicenter solution [Bossard, Ruef, 1106.5806]
- Describe orbits through H -invariants from Lax matrix. Define symmetric H -covariant tensors whose signature is H -invariant [Fre', Sorin, M.T., 1103.0848]

Tensor Classifiers

$$V_0 = \Delta^{\alpha A} K_{\alpha A}, \quad \Delta^{\alpha A} \in (\mathbf{2}, \mathbf{4}) \text{ of } H = \mathrm{SL}(2, \mathbb{R})_1 \times \mathrm{SL}(2, \mathbb{R})_2$$

$$\mathcal{T}^{xy} = \epsilon_{\alpha\beta} \Delta^{\alpha A} \Delta^{\beta B} \Pi_{AB}^{xy} \in (\mathbf{1}, \mathbf{3}) \times_s (\mathbf{1}, \mathbf{3})$$

$$\mathfrak{T}^{xy} = (s^a)_{\alpha\beta} (s_a)_{\gamma\delta} (t^x)_{AB} (t^y)_{CD} \Delta^{\alpha A} \Delta^{\beta B} \Delta^{\gamma C} \Delta^{\delta D} \in (\mathbf{1}, \mathbf{3}) \times_s (\mathbf{1}, \mathbf{3})$$

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where $\mathfrak{sl}(2, \mathbb{R})_1 = \mathrm{Span}\{s_a\}$, $\mathfrak{sl}(2, \mathbb{R})_2 = \mathrm{Span}\{t_x\}$. BPS solution $\Leftrightarrow \mathcal{T}^{xy} \equiv 0$

Conclusions

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