# Walls of a massive Kähler sigma model on SO(2N)/U(N)

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#### Plan

- SO(2N)/U(N) as a gauge theory [Higashijima, Nitta 00]
- Moduli matrices for Grassmannian manifold [Isozumi, Nitta, Ohashi, Sakai 04]
- Moduli matrices for SO(2N)/U(N)
- Walls in SO(6)/U(3)

### SO(2N)/U(N)

- $\mathcal{N} = 1 \text{ SO}(2N)/U(N)$  in 4D as a gauge theory [Higashijima, Nitta 00]
- Bosonic Lagrangian for a massive SO(2N)/U(N) in 3D

$$\mathcal{L}_{\text{bos }3D} = -|D_m \phi_a^{\ i}|^2 - |i\phi_a^{\ j}M_j^{\ i} - i\Sigma_a^{\ b}\phi_b^{\ i}|^2 + |F_a^{\ i}|^2 + \frac{1}{2}(D_a^{\ b}\phi_b^{\ i}\bar{\phi}_i^{\ a} - D_a^{\ a}) + \left((F_0)^{ab}\phi_b^{\ i}J_{ij}\phi_a^{\text{T}j} + (\phi_0)^{ab}F_b^{\ i}J_{ij}\phi_a^{\text{T}j} + (\phi_0)^{ab}\phi_b^{\ i}J_{ij}F_a^{\text{T}j} + \text{c.c.}\right)$$

$$J = \mathbf{1} \otimes \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \qquad M_j^{\ i} = \operatorname{diag}(m_1, m_2, \cdots, m_N) \otimes \sigma_3$$
$$\Sigma = \operatorname{diag}(\Sigma_1, \Sigma_2, \cdots, \Sigma_N)$$

• Vacuum condition

$$\phi_a^{\ j} M_j^{\ i} - i \Sigma_a^{\ b} \phi_b^{\ i} = 0, \quad (\phi_0)^{ab} \phi_b^{\ i} = 0$$

• Constraints

$$\phi_{a}^{\ i}\bar{\phi}_{i}^{\ b} - \delta_{a}^{\ b} = 0, \quad \phi_{a}^{\ i}J_{ij}\phi_{\ b}^{\mathrm{T}j} = 0$$

# Moduli matrices $H_0$ [Isozumi, Nitta, Ohashi, Sakai 04]

• BPS equation

$$D\phi_a^{\ i} - (\phi_a^i M_j^{\ i} - \Sigma_a^{\ b} \phi_b^{\ i}) = 0$$

Assumption

- Fields : static, depends only on the  $x_1 \equiv x$
- Poincaré invariance on the two-dimensional world volume of walls to set  $v_0 = v_2 = 0$
- Introduce  $S_a^{\ b}(x)$  and  $f_a^{\ i}(x)$  defined by

$$\Sigma_{a}^{\ b} - iv_{a}^{\ b} = (S^{-1}\partial S)_{a}^{\ b}, \quad \phi_{a}^{\ i} = (S^{-1})_{a}^{\ b}f_{b}^{\ i}$$

then the BPS equation and the solutions are

$$\partial f_a^{\ i} = f_a^{\ j} M_j^{\ i}, \quad f_a^{\ i} = H_{0a}^{\ j} (e^{Mx})_j^{\ i}$$

• BPS solution

$$\phi_a^{\ i} = (S^{-1})_a^{\ b} H_{0b}^{\ j} (e^{Mx})_j^{\ i}$$

• Everything is invariant under

$$S'_{a}{}^{b} = V_{a}{}^{c}S_{c}{}^{b}, \quad H'_{0a}{}^{i} = V_{a}{}^{c}H_{0c}{}^{i}, \quad V \in GL(N, \mathbf{C})$$

#### Moduli matrices $H_0$

Using  $\phi_a^{\ i} = (S^{-1})_a^{\ b} H_{0b}^{\ j} (e^{Mx})_j^{\ i}$ 

• Vacuum condition :  $\phi_a^{\ j} M_j^{\ i} - i \Sigma_a^{\ b} \phi_b^{\ i} = 0$ 

 $(\Sigma_1, \Sigma_2, \Sigma_3, \cdots, \Sigma_N) = (\pm m_1, \pm m_2, \pm m_3, \cdots, \pm m_N) \rightarrow 2^{N-1} H_0's$ Parity has been removed.

• World volume symmetry

$$S_{a}^{\prime b} = V_{a}^{\ c}S_{c}^{\ b}, \quad H_{0a}^{\prime \ i} = V_{a}^{\ c}H_{0c}^{\ i}, \quad V \in GL(N, \mathbf{C})$$

The moduli matrix space is an equivalent class of the sets of  $(S, H_0)$  defined by V. • D-term constraint

$$\phi_a^{\ i}\bar{\phi}_i^{\ b} - \delta_a^{\ b} = 0 \to H_{0a}^{\ i}(e^{2Mx})_i^{\ j}H_{0j}^{\dagger \ b} = (S\bar{S})_a^{\ b} \equiv \Omega_a^{\ b}$$

• F-term constraint

$$\phi_a^{\ i}J_{ij}\phi_{\ b}^{\mathrm{T}j} = 0 \rightarrow H_{0a}^{\ i}J_{ij}H_{\ b}^{\mathrm{T}j} = 0$$

#### Moduli matrices for walls in the Grassmannian manifold

[Isozumi, Nitta, Ohashi, Sakai 04]

$$S_{a}^{\prime b} = V_{a}^{\ c}S_{c}^{\ b}, \quad H_{0a}^{\prime \ i} = V_{a}^{\ c}H_{0c}^{\ i}, \quad V \in GL(N, \mathbf{C})$$
$$H_{0a}^{\ i}(e^{2Mx})_{i}^{\ j}H_{0j}^{\dagger \ b} = (S\bar{S})_{a}^{\ b} \equiv \Omega_{a}^{\ b}$$

Walls are constructed algebraically from elementary walls.

- Elementary walls
  - :  $H_{0\langle A \leftarrow B \rangle} = H_{0\langle A \rangle} e^{a_i(r)}$ ,  $a_i(r) \equiv e^r a_i(r \in \mathbb{C})$  $\langle A \rangle$  and  $\langle B \rangle$  are the vacua in the flavor i and i + 1 respectively in the same color.

$$[cM, a_i] = c(m_i - m_{i+1})a_i = T_{\langle i \leftarrow i+1 \rangle}a_i$$

The  $a_i$  has a nonzero component only in the (i, i + 1)-th element.

This cannot be defined in the SO(2N)/U(N) manifold.

We introduce an additional element with an opposite sign into the  $a_i$ .

#### What it means that $\cdots$



# Walls in SO(6)/U(3)

• Vacua

$$H_{0\langle 1\rangle} = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \end{pmatrix}$$
$$H_{0\langle 3\rangle} = \begin{pmatrix} 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 \\ 0 & 1 & 0 & 0 & 0 & 0 \end{pmatrix}$$

• Operators generating elementary walls

# Walls in SO(6)/U(3)

- Single walls
  - Elementary walls  $(a_i(r) \equiv e^r a_i)$

$$H_{0\langle 1\leftarrow 2\rangle} = H_{0\langle 1\rangle}e^{a_1(r_1)}, \quad H_{0\langle 2\leftarrow 3\rangle} = H_{0\langle 2\rangle}e^{a_2(r_1)}, \quad H_{0\langle 3\leftarrow 4\rangle} = H_{0\langle 3\rangle}e^{a_3(r_1)}$$

Compressed walls\* level one

$$E_1 = [a_1, a_2] \neq 0, \quad E_2 = [a_2, a_3] \neq 0$$

$$H_{0\langle 1\leftarrow 3\rangle} = H_{0\langle 1\rangle} e^{E_1(r_1)}, \quad H_{0\langle 2\leftarrow 4\rangle} = H_{0\langle 2\rangle} e^{E_2(r_1)}$$

\* level two

$$E_3 = [a_1, E_2] = [E_1, a_3] \neq 0$$

$$H_{0\langle 1\leftarrow 7\rangle} = H_{0\langle 1\rangle} e^{E_3(r_1)}$$

# Walls in SO(6)/U(3)

- Multiwalls
  - Double walls

$$\begin{split} H_{0\langle 1\leftarrow 2\leftarrow 3\rangle} &= H_{0\langle 1\leftarrow 2\rangle} e^{a_2(r_2)}, \quad H_{0\langle 2\leftarrow 3\leftarrow 4\rangle} = H_{0\langle 2\leftarrow 3\rangle} e^{a_3(r_2)} \\ H_{0\langle 1\leftarrow 2\leftarrow 4\rangle} &= H_{0\langle 1\leftarrow 2\rangle} e^{E_2(r_2)}, \quad H_{0\langle 1\leftarrow 3\leftarrow 4\rangle} = H_{0\langle 1\leftarrow 3\rangle} e^{a_3(r_2)} \end{split}$$

- Triple walls

$$H_{0\langle 1 \leftarrow 2 \leftarrow 3 \leftarrow 4 \rangle} = H_{0\langle 1 \leftarrow 2 \leftarrow 3 \rangle} e^{a_3(r_3)}$$

#### Summay

- Walls of a massive K\u00e4hler sigma model on SO(2N)/U(N) are studied in the moduli matrix approach.
- $2^{N-1}$  discrete vacua are observed.
- Elementary walls in the SO(2N)/U(N) manifold can be defined by generating operators with an extra unit component.
- Compressed walls, multiwalls are obtained.

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- Compressed walls, multiwalls are obtained.

Thank you for your attention.