

# Recurrent relations for branching coefficients of affine Lie algebra modules

Anton Nazarov

Partially based on arXiv:1007.0318, 1102.1702,  
in collaboration with V.D. Lyakhovsky

Department of theoretical physics,  
Saint-Petersburg State University,  
198904, Saint-Petersburg, Russia  
e-mail: anton.nazarov@hep.phys.spbu.ru

Supersymmetry and Quantum Symmetries, 18-23 July 2011

# Simple and affine Lie algebras

Consider semisimple Lie algebra  $\mathfrak{g}$  with commutator  $[\cdot, \cdot] : \mathfrak{g} \otimes \mathfrak{g} \rightarrow \mathfrak{g}$

## Definition

*Affine Lie algebra*  $\hat{\mathfrak{g}}$  is central extension of loop algebra, corresponding to semisimple Lie algebra  $\mathfrak{g}$  with commutation relations

$$[Xt^n + \alpha c, Yt^m + \beta c] = t^{n+m}[X, Y] + (X, Y)n\delta_{n+m,0}c$$

$$Vir \subset U(\hat{\mathfrak{g}}) \quad (\text{Sugawara construction})$$

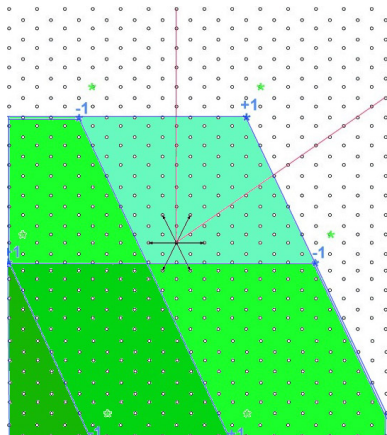
Current algebra of WZNW models in CFT is

$$\hat{\mathfrak{g}} \ltimes Vir$$

Primary fields are indexed by highest weights of irreducible  $\hat{\mathfrak{g}}$ -modules

## Weyl character formula

$$ch(L^\mu) = \frac{\sum_{w \in W} \epsilon(w) e^{w \circ (\mu + \rho)}}{\prod_{\alpha \in \Delta^+} (1 - e^{-\alpha})^{\text{mult}(\alpha)}} = \frac{\Psi(\mu)}{\Psi(0)}; \quad \frac{1}{R} = \frac{1}{\prod_{\alpha \in \Delta^+} (1 - e^{-\alpha})^{\text{mult}(\alpha)}}$$



# Recurrent relation for weight multiplicities

We can use Weyl character formula to obtain recurrent relation on weight multiplicities which can be used for calculations:

$$m_{\xi} = - \sum_{w \in W \setminus e} \epsilon(w) m_{\xi - (w(\rho) - \rho)} + \sum_{w \in W} \epsilon(w) \delta_{(w(\mu + \rho) - \rho), \xi}. \quad (1)$$

The set of weights to sum over is obtained through expansion of denominator:

$$\sum_{w \in W} \epsilon(w) e^{w\rho - \rho} = \prod_{\alpha \in \Delta^+} (1 - e^{-\alpha})^{\text{mult}(\alpha)}$$

# Branching

Consider reductive subalgebra  $\hat{\mathfrak{g}} \supset \hat{\mathfrak{a}}$ .

Each  $\hat{\mathfrak{g}}$ -module is also an  $\hat{\mathfrak{a}}$ -module, although  $L_{\hat{\mathfrak{g}}}^{\mu}$  is not irreducible as  $\hat{\mathfrak{a}}$ -module. It can be decomposed into the sum of irreducible  $\hat{\mathfrak{a}}$ -modules:

$$L_{\hat{\mathfrak{g}} \downarrow \hat{\mathfrak{a}}}^{\mu} = \bigoplus_{\nu \in P_{\hat{\mathfrak{a}}}^+} b_{\nu}^{\mu} L_{\hat{\mathfrak{a}}}^{\nu}.$$

For characters we have

$$\pi_{\hat{\mathfrak{a}}} ch(L^{\mu}) = \sum_{\nu \in P_{\hat{\mathfrak{a}}}^+} b_{\nu}^{\mu} ch(L_{\hat{\mathfrak{a}}}^{\nu}).$$

We want to write recurrent relations.

Denote by  $k_{\xi}^{(\mu)}$  signed branching coefficients.

$$\begin{aligned} k_{\xi}^{(\mu)} &= b_{\xi}^{(\mu)} & \text{if } \xi \in \bar{C}_{\hat{\mathfrak{a}}} \\ k_{\xi}^{(\mu)} &= \epsilon(w) b_{w(\xi + \rho_{\hat{\mathfrak{a}}}) - \rho_{\hat{\mathfrak{a}}}}^{(\mu)} & \text{where } w \in W_{\hat{\mathfrak{a}}} : w(\xi + \rho_{\hat{\mathfrak{a}}}) - \rho_{\hat{\mathfrak{a}}} \in \bar{C}_{\hat{\mathfrak{a}}} \end{aligned}$$

The Weyl-Kac character formula leads to the relation

$$\pi_{\hat{\mathfrak{a}}} \left( \frac{\sum_{\omega \in W} \epsilon(\omega) e^{\omega(\mu+\rho)-\rho}}{\prod_{\alpha \in \Delta^+} (1 - e^{-\alpha})^{\text{mult}(\alpha)}} \right) = \sum_{\nu \in P_{\hat{\mathfrak{a}}}^+} b_{\nu}^{(\mu)} \frac{\sum_{\omega \in W_{\hat{\mathfrak{a}}}} \epsilon(\omega) e^{\omega(\nu+\rho_{\hat{\mathfrak{a}}})-\rho_{\hat{\mathfrak{a}}}}}{\prod_{\beta \in \Delta_{\hat{\mathfrak{a}}}^+} (1 - e^{-\beta})^{\text{mult}_{\hat{\mathfrak{a}}}(\beta)}}.$$

Now we want to multiply by  $R$  and rewrite this as recurrent relation.

Consider the roots orthogonal to  $\Delta_{\hat{\mathfrak{a}}}$ .

Let  $\Delta_{\mathfrak{b}}^+ = \{\alpha \in \Delta_{\mathfrak{g}}^+ : \forall \beta \in \Delta_{\hat{\mathfrak{a}}}; \alpha \perp \beta\}$  be the subset of positive roots of  $\mathfrak{g}$  orthogonal to the root system of  $\hat{\mathfrak{a}}$ .

Denote by  $W_{\mathfrak{b}}$  the subgroup of the Weyl group  $W$  generated by the reflections  $\omega_{\beta}$  corresponding to the roots  $\beta \in \Delta_{\mathfrak{b}}^+$ .

The subsystem  $\Delta_{\mathfrak{b}}$  determines the subalgebra  $\mathfrak{b} = \mathfrak{a}_{\perp} \subset \mathfrak{g}$ .

$\hat{\mathfrak{a}}, \mathfrak{b}$  is the "orthogonal pair" of subalgebras in  $\mathfrak{g}$ .

Cartan subalgebra is decomposed  $\mathfrak{h} = \mathfrak{h}_{\hat{\mathfrak{a}}} + \mathfrak{h}_{\mathfrak{b}}$ .

Introduce

$$\mathcal{D}_{\hat{\mathfrak{a}}} := \rho_{\hat{\mathfrak{a}}} - \pi_{\hat{\mathfrak{a}}}\rho.$$

$$\mathcal{D}_{\mathfrak{b}} := \rho_{\mathfrak{b}} - \pi_{\mathfrak{b}}\rho.$$

# Recurrent relation for branching coefficients

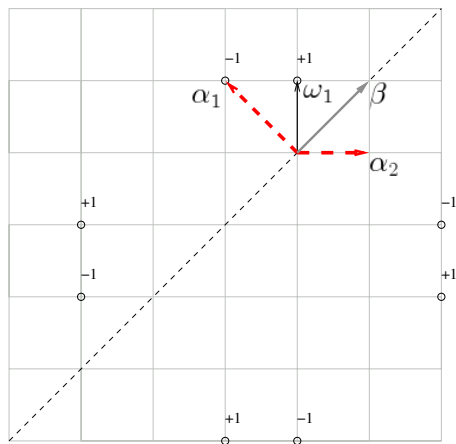
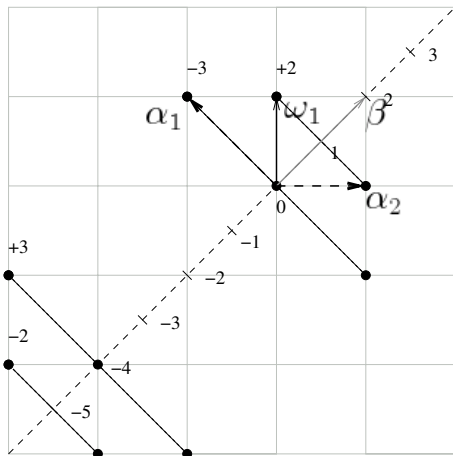
$$k_{\xi}^{(\mu)} = -\frac{1}{s(\gamma_0)} \left( \sum_{u \in W/W_b} \epsilon(u) \dim \left( L_b^{\pi(b)[u(\mu+\rho)-\rho]-\mathcal{D}_b} \right) \right. \\ \left. \delta_{\xi-\gamma_0, \pi(\mathfrak{a} \oplus \mathfrak{h}_{\mathfrak{d}})[u(\mu+\rho)-\rho]+\mathcal{D}_b} + \sum_{\gamma \in \Gamma_{\hat{\mathfrak{a}}\hat{\mathfrak{c}}\hat{\mathfrak{g}}}} s(\gamma + \gamma_0) k_{\xi+\gamma}^{(\mu)} \right).$$

The recursion is governed by the set  $\Gamma_{\hat{\mathfrak{a}}\hat{\mathfrak{c}}\hat{\mathfrak{g}}}$  of the weights  $\{\xi\}$  appearing in the expansion

$$\prod_{\alpha \in \Delta^+ \setminus \Delta_b^+} (1 - e^{-\pi_{\hat{\mathfrak{a}}}\alpha})^{\text{mult}(\alpha) - \text{mult}_{\hat{\mathfrak{a}}}(\pi_{\hat{\mathfrak{a}}}\alpha)} = - \sum_{\gamma \in P_{\hat{\mathfrak{a}}}} s(\gamma) e^{-\gamma}$$

The weights are to be shifted by  $\gamma_0$  – the lowest vector in  $\{\xi\}$  – and the zero element is to be eliminated:

$$\Gamma_{\hat{\mathfrak{a}}\hat{\mathfrak{c}}\hat{\mathfrak{g}}} = \{\xi - \gamma_0\} \setminus \{0\}.$$

Simple example:  $A_1 \subset B_2$ Figure: Roots of  $B_2$ ,  $A_1$  and  $\Psi^{\omega_1}$ Figure: Orthogonal subalgebra  $\mathfrak{b}$  and  $\mathfrak{b}$ -module dimensions



## Surprises

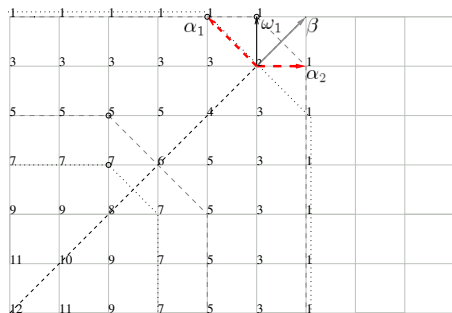


Figure: Embedding  $A_1 \oplus u(1) \subset B_2$  and generalized Verma modules. Dashed – positive  $\epsilon(u)$ , dotted – negative.

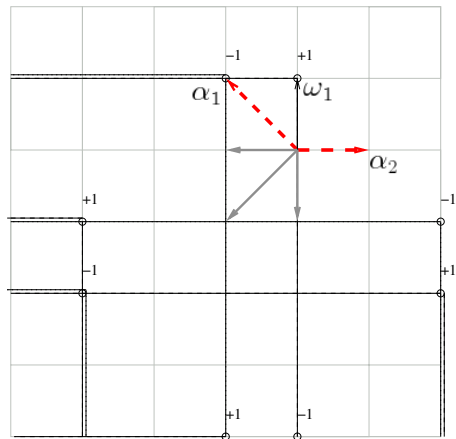
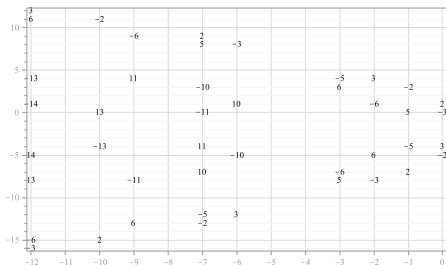
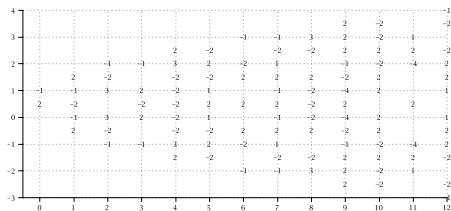
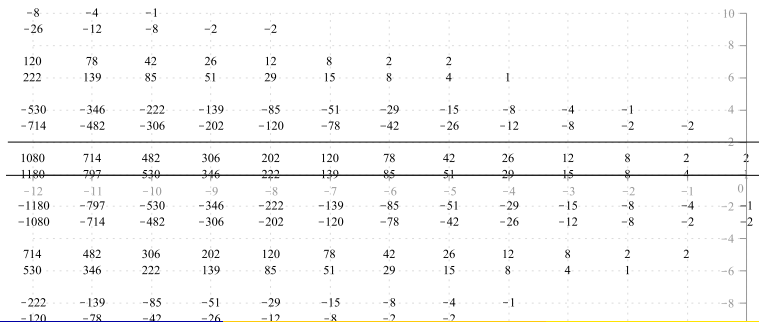


Figure: Embedding  $A_1 \oplus u(1) \subset B_2$  and (deformed)  $A_2$ -modules

Figure:  $\Gamma_{\hat{A}_1 \subset \hat{B}_2}$ Figure:  $\pi_{\hat{A}_1}(\Psi_{\hat{B}_2}^{\omega_1})$ 

## Affine case

- Infinite number of generalized Verma modules – in each grade containing weights of singular element  $\Psi$
- Not so simple relation with multiplicities of another algebra.

$\Gamma_{\hat{a} \subset \hat{g}} \leftrightarrow \prod_{\alpha \in \Delta^+} (1 - e^{-\alpha})^{\text{mult}(\alpha)}$  due to the difference in imaginary roots. E.g.  $\widehat{A_1 \oplus u(1)} \subset \widehat{B_2}$ :

$$\frac{\prod_{n>0} (1 - e^{-n\delta})^2 \prod_{n>0} (1 - e^{-\alpha_1 - n\delta}) \dots}{\prod_{n>0} (1 - e^{-n\delta})^2 \prod_{n>0} (1 - e^{-\alpha_1 - n\delta}) \dots}$$

$$\prod_{n>0} (1 - e^{-\alpha_2 - n\delta}) \prod_{n>0} (1 - e^{-\alpha_1 - \alpha_2 - n\delta}) \prod_{n>0} (1 - e^{-\alpha_1 - 2\alpha_2 - n\delta}) \dots$$

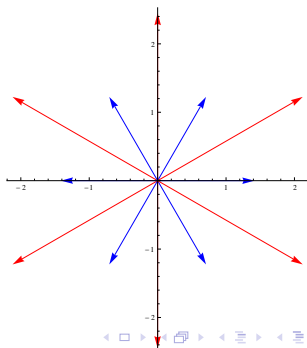
$$\leftrightarrow \prod_{n>0} (1 - e^{-n\delta})^2 \prod_{n>0} (1 - e^{-\alpha_1 - n\delta}) \prod_{n>0} (1 - e^{-\alpha_2 - n\delta})$$

$$\prod_{n>0} (1 - e^{-\alpha_1 - \alpha_2 - n\delta}) \dots$$

# Conclusion

- Recurrent relation for branching coefficients works as in finite-dimensional case
- It can be used for computations
- We can relate recurrent relations for branching coefficients of one pair algebra-subalgebra with recurrent relation of another pair, but we can not interpret singular weights as highest weights of Verma modules.

$$\begin{aligned} \text{E.g.: } & A_1 \oplus \widehat{u(1)} \subset \widehat{B_2} \longleftrightarrow \\ & u(1) \oplus \widehat{u(1)} \subset \widehat{A_2}, \\ \text{or} & \\ & \widehat{A_2} \subset \widehat{G_2} \longleftrightarrow u(1) \oplus \widehat{u(1)} \subset \widehat{A_2} \end{aligned}$$



Thank you for your attention!