

A new perspective on nonrelativistic gravity

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Overview

- Motivation
- GR by gauging the Poincaré algebra
- Newton-Cartan
- Newton-Cartan by gauging the Bargmann algebra
- Extension to strings

Motivation

- A lot of physics in our universe is nonrelativistic
- Problems can become simpler
- Possible applications to AdS/CFT?

General Relativity from Poincaré

- Poincaré algebra:

$$[P, P] = 0$$

$$[J, P] = P$$

$$[J, J] = J$$

Gauge fields Curvatures

$$P_a : e_\mu^a \qquad R_{\mu\nu}^a(P)$$

$$J_{ab} : \omega_\mu^{ab} \qquad R_{\mu\nu}^{ab}(J)$$

- $R_{\mu\nu}^a(P) = 0 \rightarrow$ (1) $\omega_\mu^{ab} = \omega_\mu^{ab}(e, \partial e)$
(2) $P_a \rightarrow gct + local Lorentz$

- Impose vielbein postulate \rightarrow General Relativity

Newton-Cartan : geometrizing Newton

- Spatial metric h and temporal metric τ
- Space and time decouple: $h^{\mu\nu}\tau_{\nu\rho} = 0$
- Introduce 'inverse metrics' $\tilde{h}_{\mu\nu}, \tilde{\tau}^{\mu\nu}$
- Connection via $\begin{matrix} \nabla_{\rho} h^{\mu\nu} = 0 \\ \nabla_{\rho} \tau_{\mu\nu} = 0 \end{matrix} \Rightarrow \Gamma_{\mu\nu}^{\rho}(h, \tilde{h}, \tau, \tilde{\tau}, K)$
- 'Ehlers conditions' to constrain $K \rightarrow K_{\mu\nu} = 2\partial_{[\mu} m_{\nu]}$

Newton-Cartan : EOM

- Poisson eqn. : $R_{\mu\nu}(\Gamma) = 4\pi G\rho\tau_{\mu\nu}$
- Geodesic eqn. : $\ddot{x}^\rho + \Gamma_{\mu\nu}^\rho \dot{x}^\mu \dot{x}^\nu = 0$
- Only nonzero connection: $\Gamma_{00}^i = \partial^i \phi$
- All degrees of freedom in ϕ

Newton-Cartan from Bargmann I

- Bargmann algebra generators: $\{ H, P, G, J, M \}$
Gauge fields Curvatures
- Time transl. $H : \tau_\mu$ $R_{\mu\nu}(H)$
- Space transl. $P_i : e_\mu^i$ $R_{\mu\nu}^i(P)$
- Boosts $G_i : \omega_\mu^{i0}$ $R_{\mu\nu}^{i0}(G)$
- Space rot. $J_{ij} : \omega_\mu^{ij}$ $R_{\mu\nu}^{ij}(J)$
- Central ext. $M : m_\mu$ $R_{\mu\nu}(M)$

- Central extension : $[G, P] = 0 \rightarrow [G, P] \sim M$

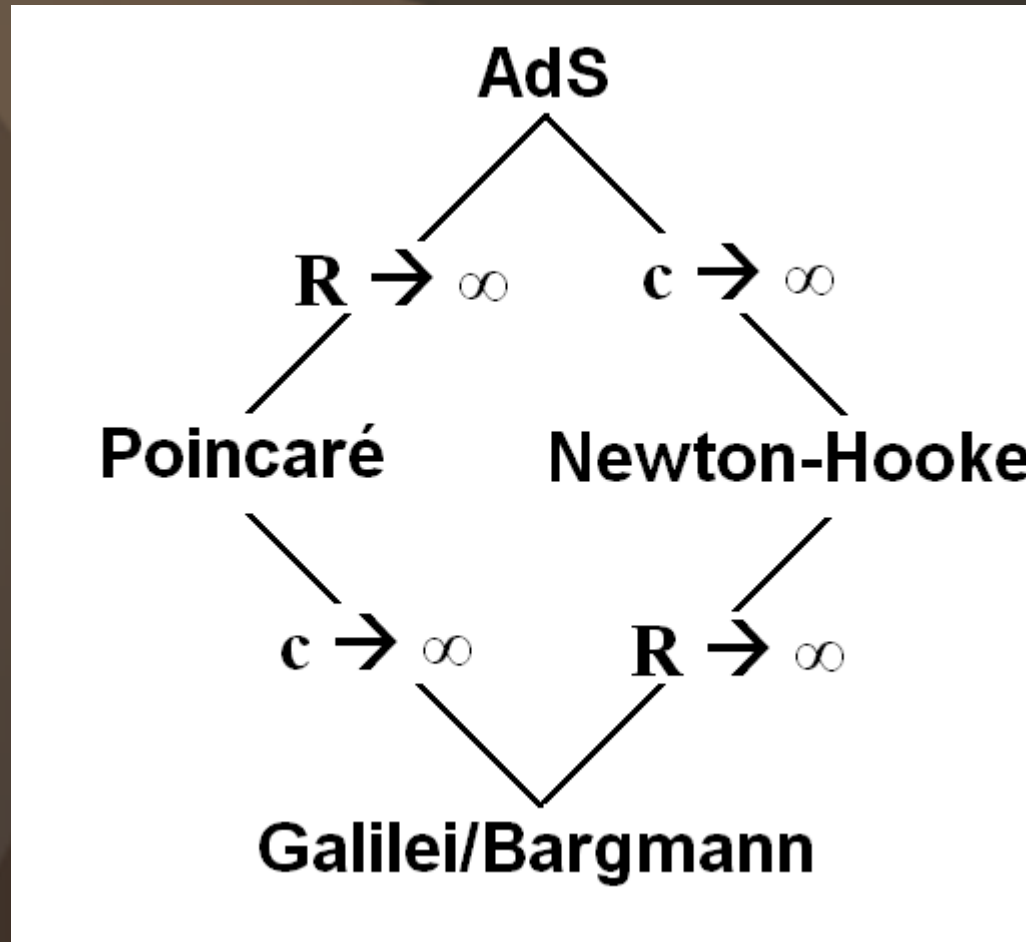
Newton-Cartan from Bargmann II

- $R(P) = R(M) = R(H) = 0$ such that
 $\{H, P\} \rightarrow$ gct + remaining transformations
- $-R(M) = R(P) = 0$: solve for spin connections
 $-R(H) = 0$: $\tau_\mu = \partial_\mu t$
- On-shell constraint: $R(J)=0 \rightarrow$ flat space
- Vielbein postulate
- $K_{\mu\nu} = 2\partial_{[\mu} m_{\nu]} \rightarrow$ Newton-Cartan!

Role of central extension M

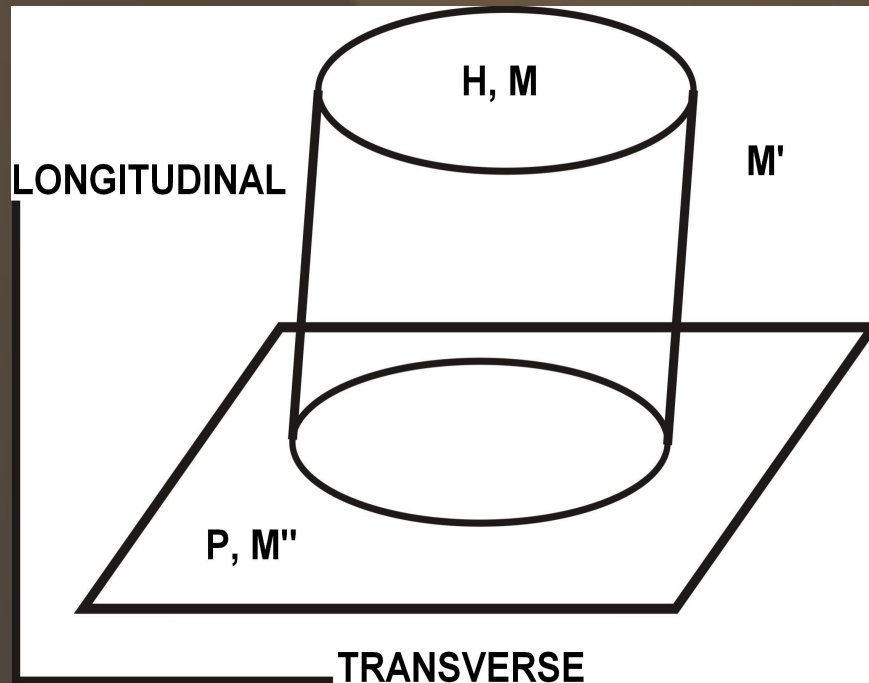
- 1) Without M we cannot solve for spin connections, we need both $R(P) = 0$ *and* $R(M) = 0$!
- 2) Remember that in Newton-Cartan the K in the connection is closed: $K_{\mu\nu} = 2\partial_{[\mu}m_{\nu]}$
→ Curl of the gauge field of M !
- In short: M is crucial

Algebra contractions



- Newton-Hooke: nonrelativistic limit of AdS, cosmological constant in longitudinal direction!

The deformed string Galilei algebra



Generators:

H: long. translations

M: long. boost

P: transv. translation

M'': transv. rotation

M': stringy 'boost'

Z: deformations

- Galilei: $\{ H, M \}$ form 2-dim. Poincaré algebra
- Newton-Hooke: $\{ H, M \}$ form 2-dim. AdS algebra

Stringy Newton-Cartan

- Idea: gauge deformed string Galilei algebra
- Add cosmological constant to equations of motion
- → Stringy Newton-Hooke-Cartan theory!

- Changes:
 - Newton potential: $\phi \rightarrow \phi_{ab}$
 - Deformation: $M \rightarrow \{Z_a, Z_{ab}\}$
 - K in connection: $K_{\mu\nu} = 2\partial_{[\mu}m_{\nu]} \rightarrow K_{\mu\nu}^a = 2\partial_{[\mu}m_{\nu]}^a$

Stringy Newton-Cartan: EOM

- Poisson equation: $\partial^i \partial_i \phi_{ab} = (4\pi G \rho - \Lambda) \tau_{ab}$
- String geodesic equation

Role of deformations $\{Z_a, Z_{ab}\}$

- Without the Z 's we can't solve for the three different spin connections!
- Ambiguity in connection now becomes
$$K_{\mu\nu}{}^a = 2\partial_{[\mu}m_{\nu]}{}^a$$
- $m_{\mu}{}^{ab}$ appears as ambiguity of K , doesn't enter EOM
- In short: the deformation is crucial!

Conclusions

- Bargmann algebra can be gauged to obtain Newton-Cartan theory
- This method can be mimicked for strings, including a cosmological constant
- Algebra deformations play a crucial role
- Stringy Newton-Cartan theory with Newton-Hooke symmetry possibly interesting for AdS/CFT