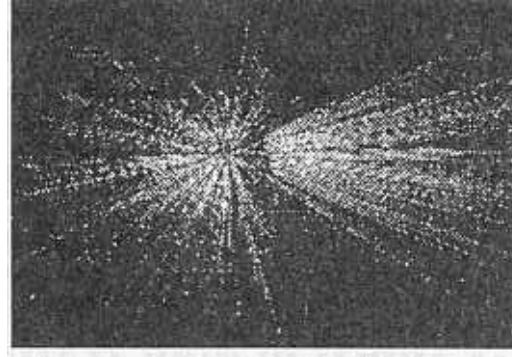


VERY HIGH MULTIPLICITY PHYSICS

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Bose-Einstein correlations in Pythia

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Bose-Einstein Correlations: Enhanced probability for the emission of identical bosons with similar momenta.

Provides information about space-time characteristics of the particle emission region

The two-particles correlation function is given by : $C_2(p_1, p_2) = \frac{P(p_1, p_2)}{P(p_1)P(p_2)}$

$P(p_1, p_2)$ - two-particles probability density function

$P(p_1), P(p_2)$ - single particle probability

$$C_2(Q) = N(1 + \lambda \times e^{-R^2 Q^2})$$

The construction of correlation function is given by:

is the two-particles four-momentum difference distribution
is the corresponding reference sample.

Bose-Einstein effect in Pythia 6.205

There is a possibility to include Bose-Einstein correlations, but it is only a first approximation to BE effect.

First of all the Q_{ij} value is evaluated: $Q_{ij} = \sqrt{(p_i + p_j)^2 - 4m^2}$

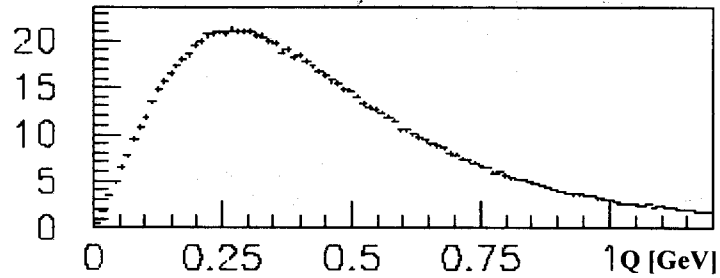
And then the shifted Q'_{ij} is found as the solution to the equation:

$$\int_0^{Q_{ij}} \frac{Q^2 dQ}{\sqrt{Q^2 + 4m^2}} = \int_0^{Q'_{ij}} C_2(Q) \frac{Q^2 dQ}{\sqrt{Q^2 + 4m^2}}$$
$$C_2(Q) = 1 + \lambda \times e^{-\left(\frac{Q}{d}\right)^r}$$

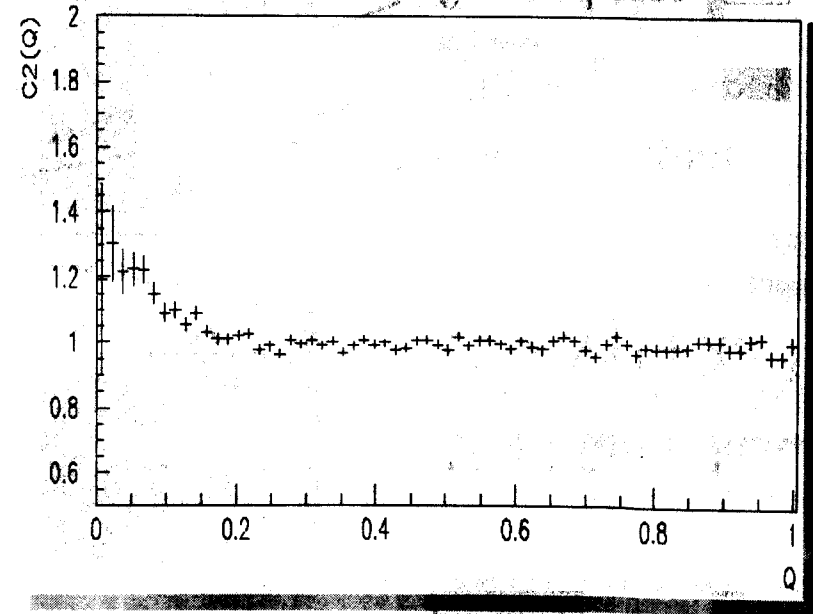
And finally after evaluation of all Q'_{ij} the original momenta are really shifted.

In Pythia collisions at 1.96 TeV (Tevatron CDF) are simulated.

Correlation spectrum $N(Q)$ for π^+ :



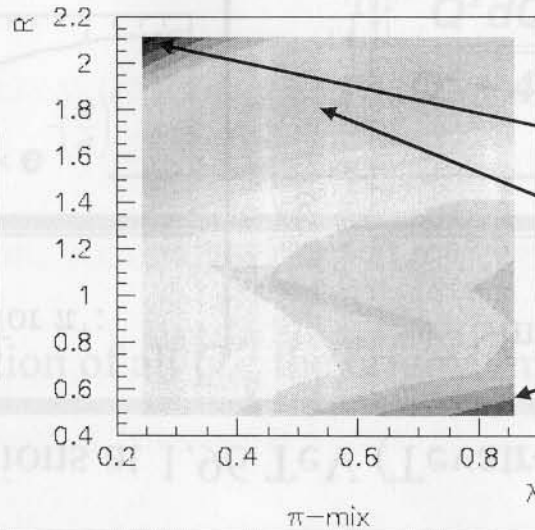
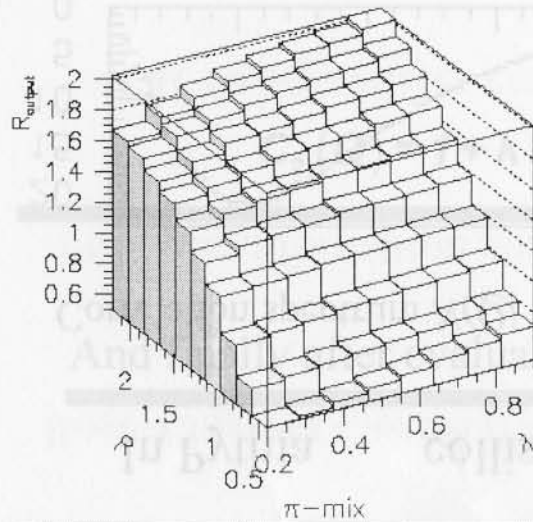
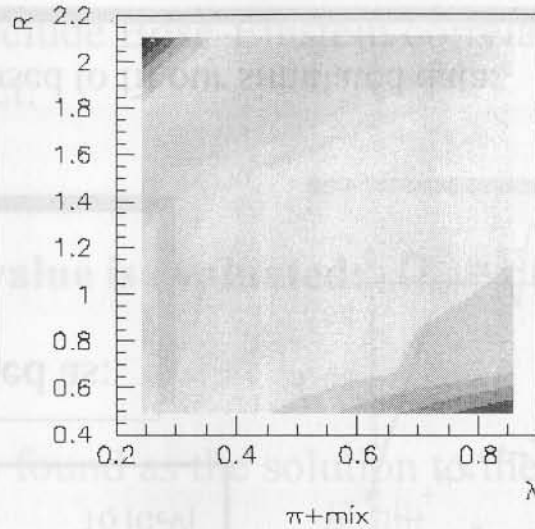
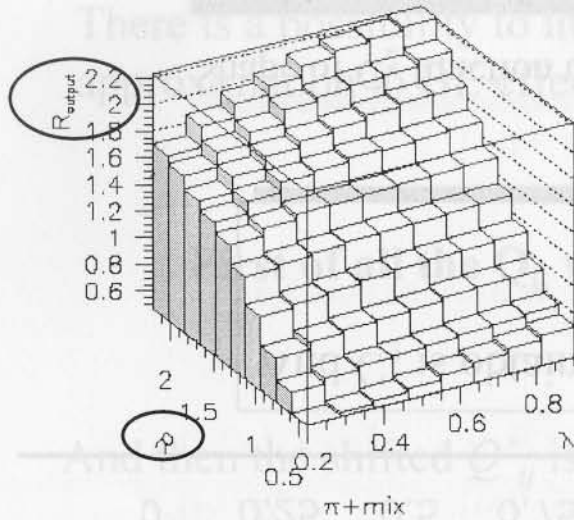
Similar spectrum for $N^{\text{ref}}(Q)$



And C_2 is obtained as:

Shape of C_2 function used to fit our simulated data:

Reproduction of input parameters



Dependence of the output parameter R on the input parameters λ and R

Used formula:

$$C_2 = N \left(1 + \lambda e^{-R^2 x^2} \right)$$

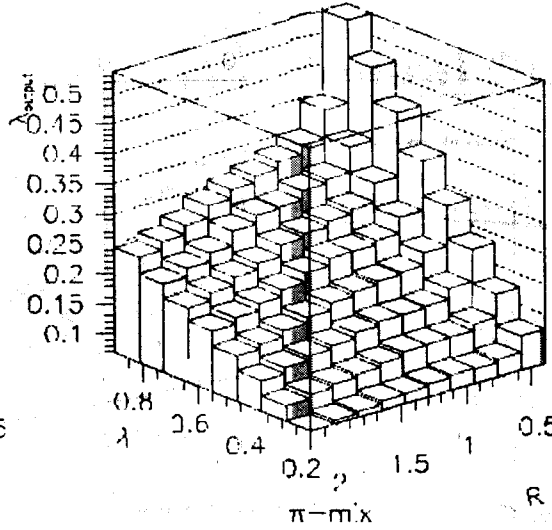
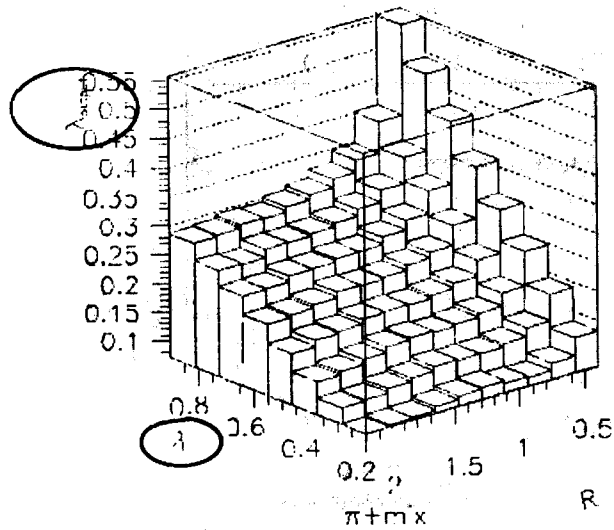
Color spectrum:

- 0.5

0 good area

+ 0.2

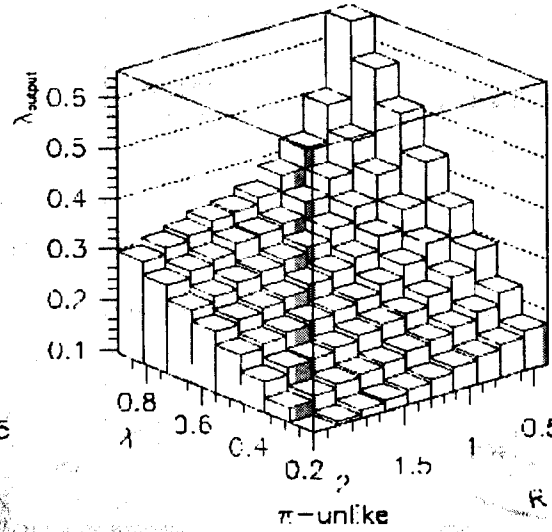
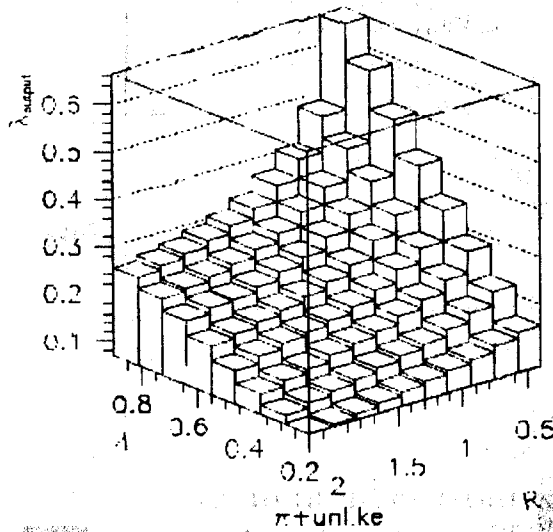
Reproduction of input parameters



Dependence of the output parameter λ on the input parameters λ and R

Used formula:

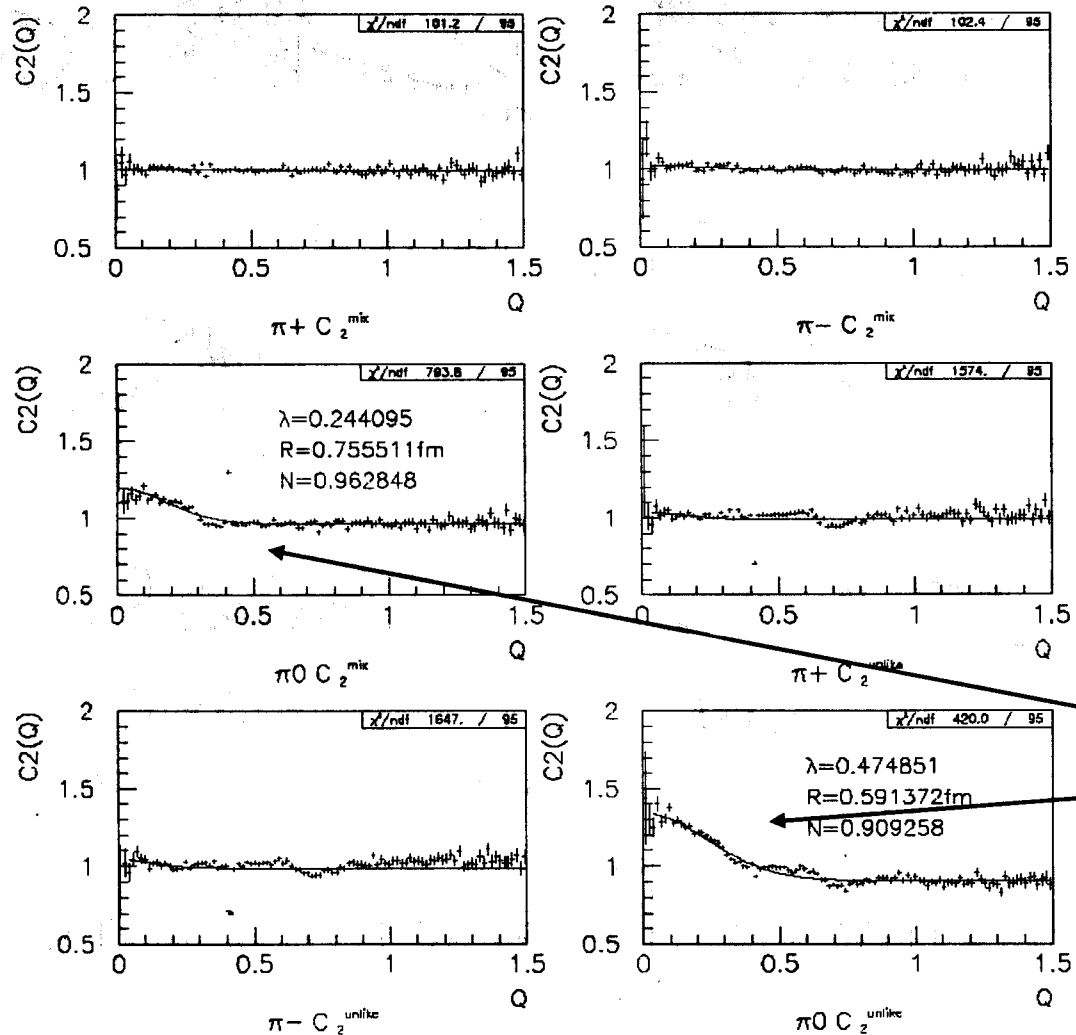
$$C_2 = N \left(1 + \lambda e^{-R^2 x^2} \right)$$



There is no green area!

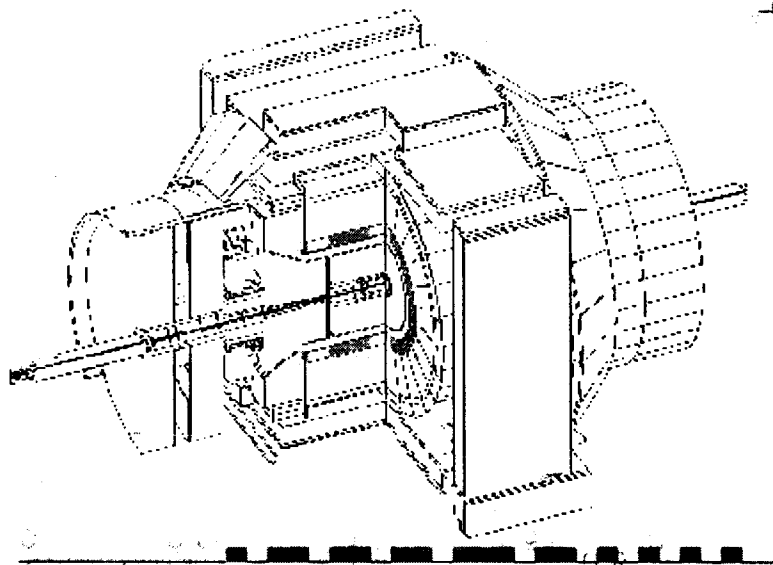
The case when correlations are not included.

BEC off.



It is expected that correlations disappear when they are not included in Pythia

There is a problem with π^0



Central Outer Tracker

Momentum Resolution: $\frac{\delta p_t}{p_t} \approx 0.15\% p_t [\text{GeV}^{-1}]$

Central Calorimeter Hadron

Calorimeter

$$\frac{\delta E}{E} \approx 0.5/\sqrt{E} [\text{GeV}]$$

Energy Resolution:

But the energy is too small for some pions to reach hadron calorimeter.

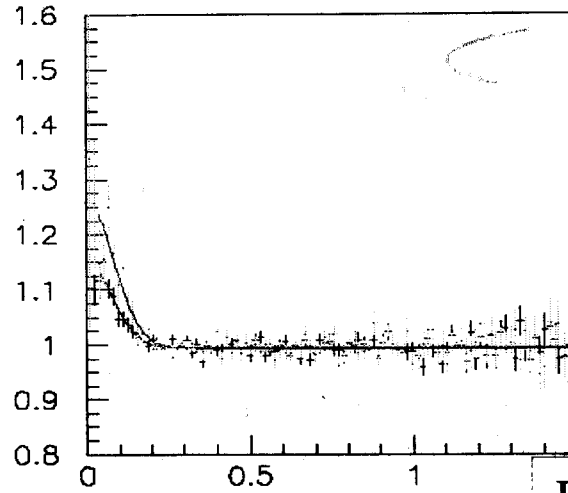
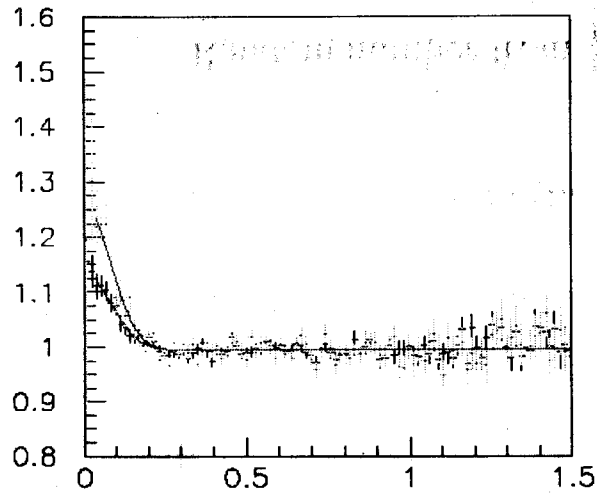
To blur all momenta we used the formulas: $p'_{x,y} = p_{x,y} + rp_t^2 \frac{0.15\%}{\sqrt{2}}$, $p'_z = p_z + rp_z^2 \frac{0.15\%}{\tan(\theta)}$

Then using this momenta the new energy for all pions was calculated:

Random number from
Gaussian distribution

$$E_{\text{new}} = \sqrt{p_x^2 + p_y^2 + p_z^2 + m_\pi^2}$$

More realistic simulation: Momentum blurring



Input parameters:

$R = 1.6\text{fm}$

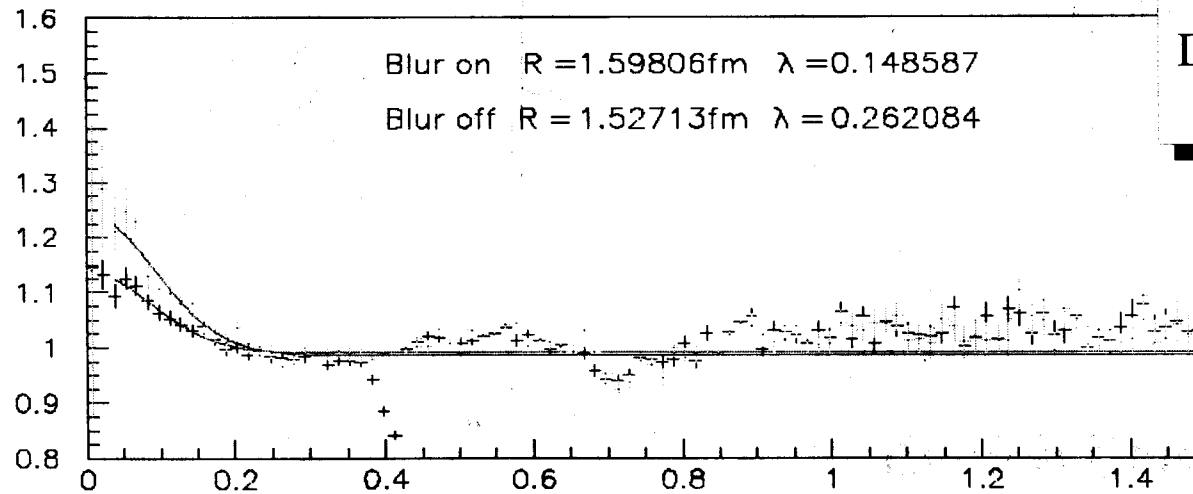
$\lambda = 0.8$

Data simulated with blurring on:

The blue ones

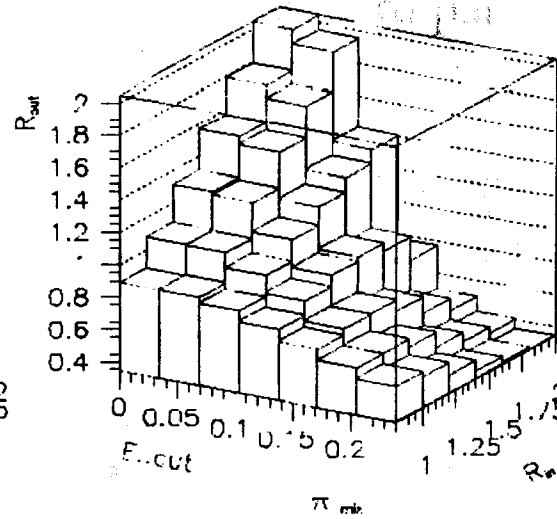
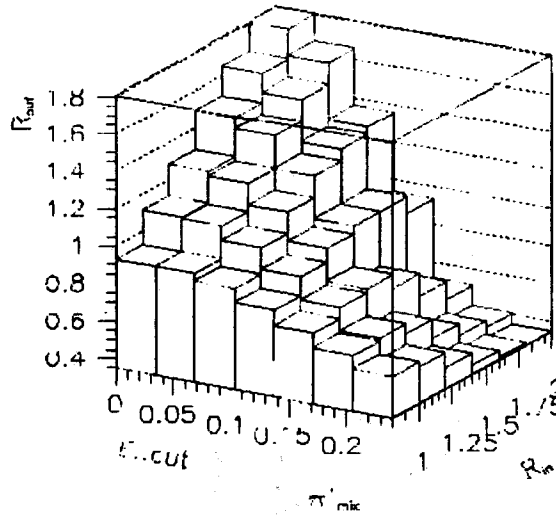
Data simulated with blurring off:

The green ones



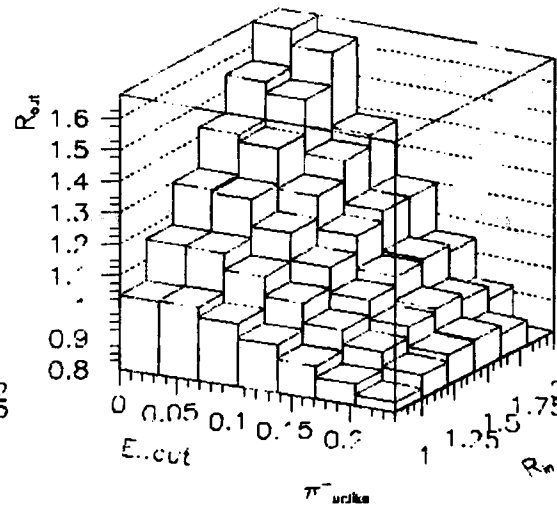
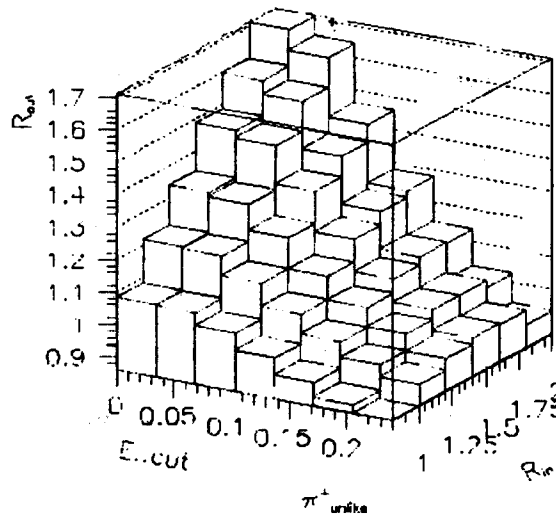
The red curves are fitting lines

More realistic simulation: Energy cut



λ set to 0.8

Dependence of the output parameter R on energy cut and on the input parameter R



We can see the importance to register low energy particles

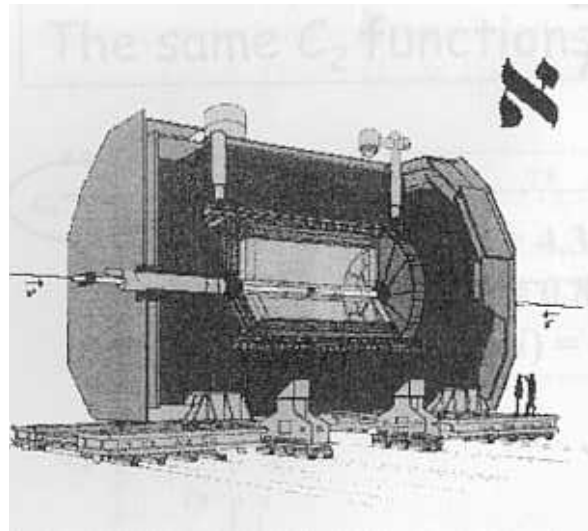
Conclusion of 1st part

•Pythia:

- Parameter R has bounded area of good reproduction
- Parameter λ has no area of good reproduction
- There are non vanishing correlations for π^0 mesons.

- It is important to register low energy particles.

Experiment ALEPH



- Vertex Detector
- Inner Tracking Chamber
- Time Projection Chamber
- Electromagnetic Calorimeter
- Superconducting Magnet Coil
- Hadron Calorimeter
- Muon Chambers
- Luminosity Monitors

Standard parameterization:

$$C_2(Q) = N[1 + \lambda \exp(-(RQ)^2)](1 + \epsilon Q + \dots)$$

- λ coherence strength factor.
- R is related to the size of the boson emission region.

$$R_2(Q) = \frac{\rho_2(p_1, p_2)}{\rho_0(p_1, p_2)} \quad \text{- Two-particle correlation function}$$

$$\rho_2(p_1, p_2) \quad \text{- Two-particle density distribution}$$

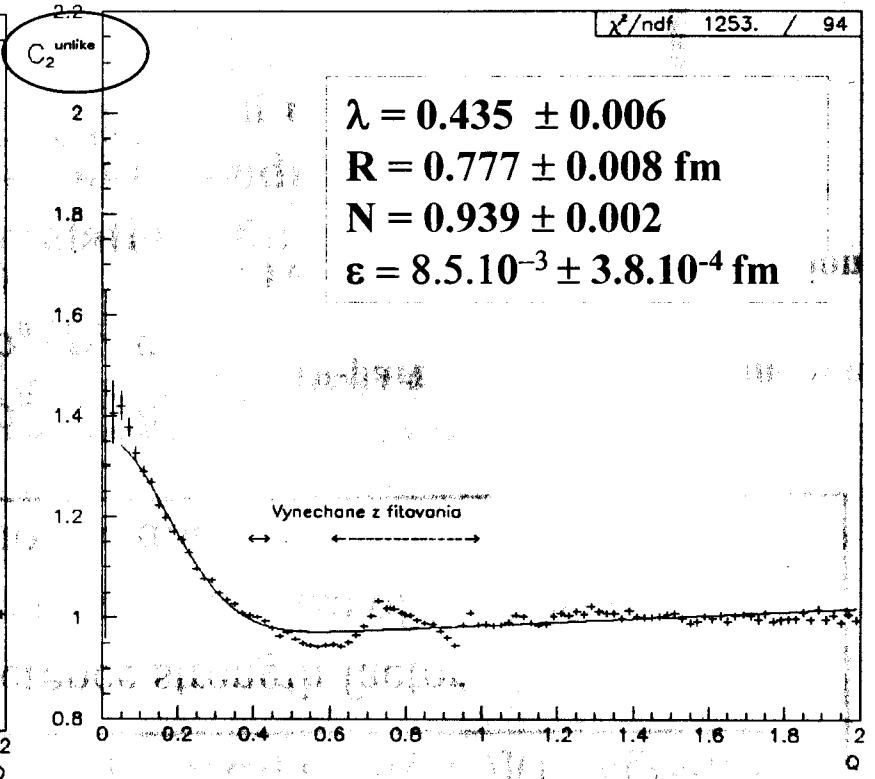
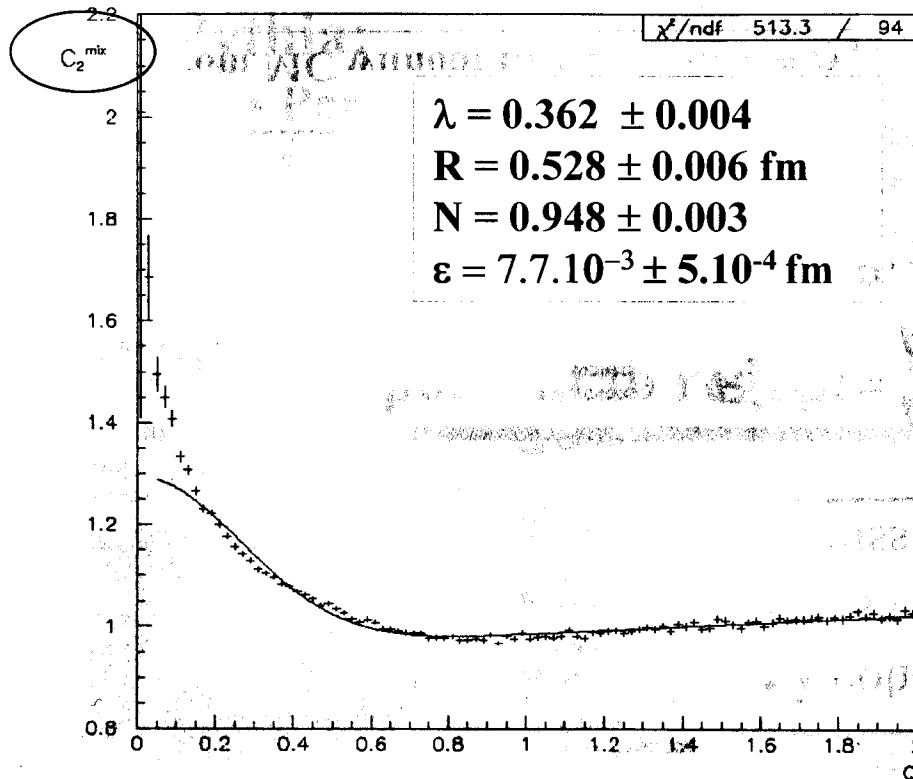
Detector inadequacies, effects introduced by the reference sample are obtained from MC without BEC: $R_2^{MC} = N_{\pm\pm}^{MC} / N_{ref}^{MC}$

The measured correlation function is given by:

More informations in:
Eur. Phys. J. C36: 2004, 147

ALEPH fit

C_2 function for ALEPH data and different reference samples fitted with $c_2(Q) = N \left(1 + \lambda e^{-(QR)^2} \right) [1 + \varepsilon Q]$



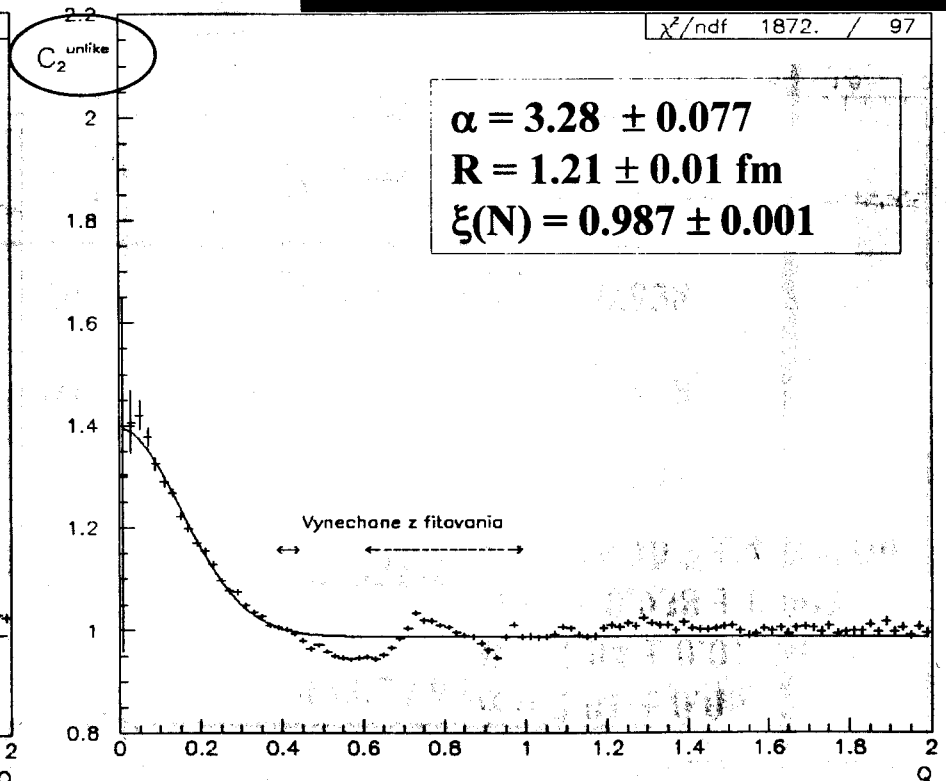
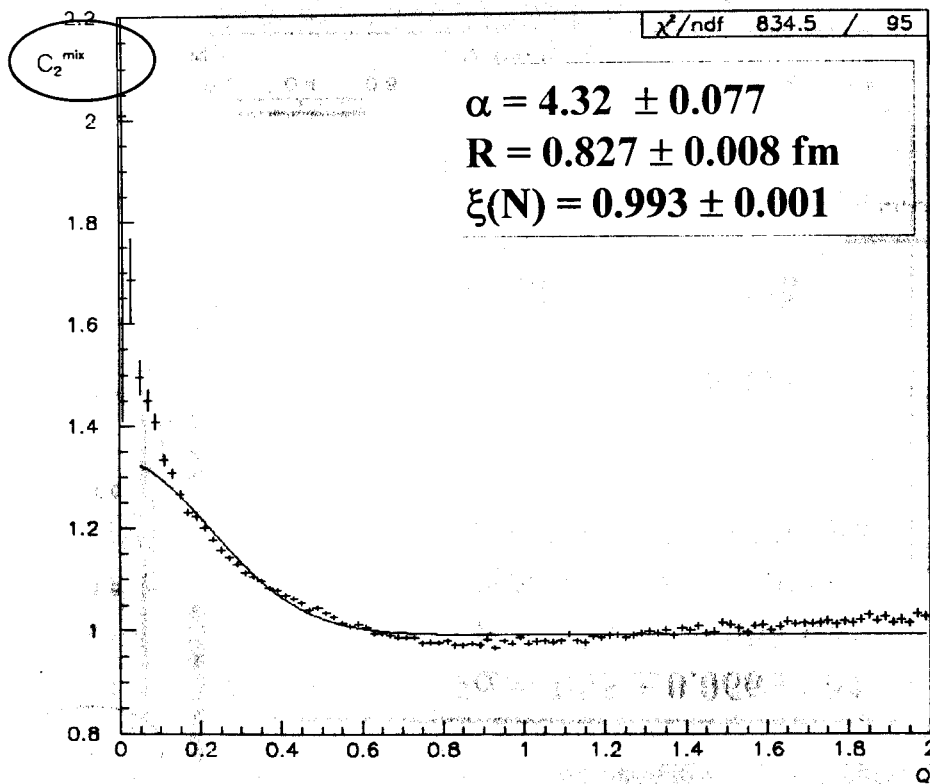
We reproduce their results.

Kozlov fit

The same C_2 functions versus Kozlov theory

QFT based, Langevin evolution equation used.

Ref: arXiv:hep-ph/0304091 v3

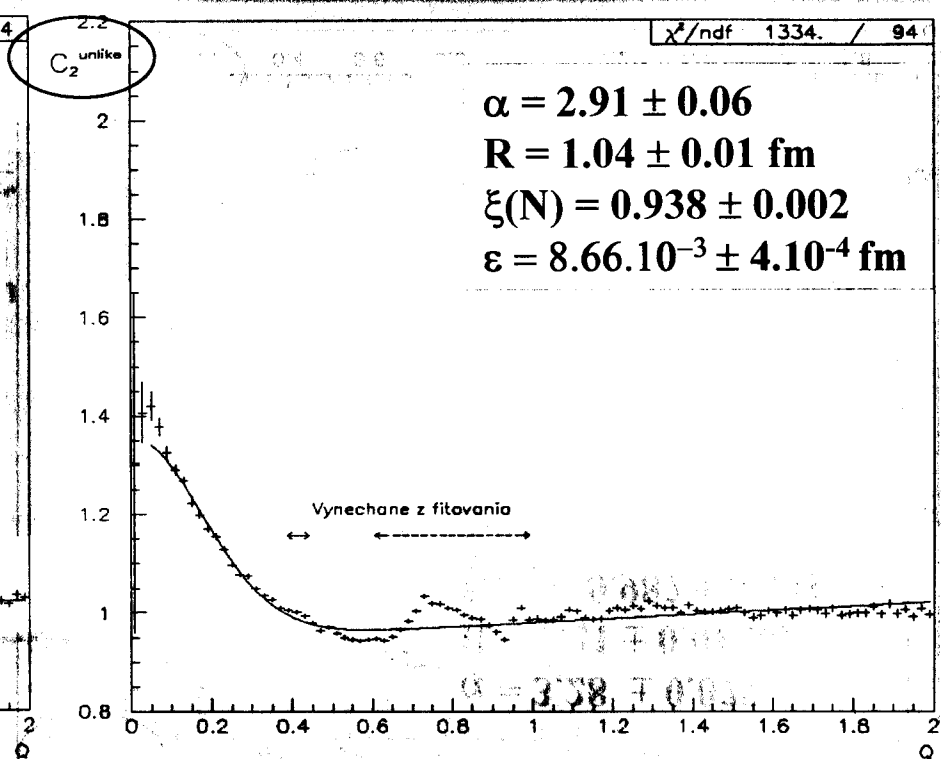
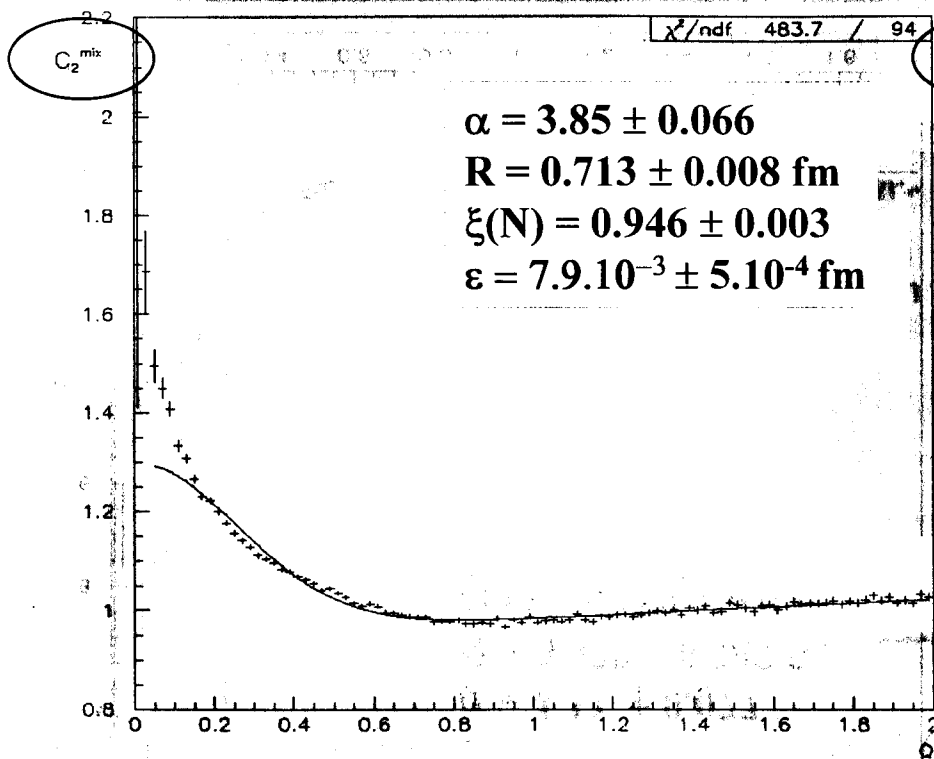


The function used to fit ALEPH data:

$$C_2(Q) = \xi(N) \left(1 + \frac{2\alpha}{(1+\alpha)^2} e^{-\frac{(QR)^2}{2}} + \frac{1}{(1+\alpha)^2} e^{-(QR)^2} \right)$$

Modified Kozlov fit

The same C_2 function versus modified Kozlov C_2 function.



Modified function of Kozlov approach:

$$C_2(Q) = \xi(N) \left(1 + \frac{2\alpha}{(1+\alpha)^2} e^{-\frac{(QR)^2}{2}} + \frac{1}{(1+\alpha)^2} e^{-(QR)^2} \right) (1 + \varepsilon Q)$$

Summary table of our results:

	Aleph C_2		Kozlov C_2		
	$C_2^{\text{mix}}(Q)$	$C_2^{+-}(Q)$		$C_2^{\text{mix}}(Q)$	$C_2^{+-}(Q)$
N	0.948	0.939	$\xi(N)$	0.946	0.938
λ	0.362	0.435	α	3.848	2.918
R (fm)	0.5287	0.7777	R (fm)	0.713	1.043
ε (fm)	0.769e-2	0.852e-2	ε (fm)	0.793e-2	0.867e-2
χ^2 / ndf	513 / 94	1253 / 94	χ^2 / ndf	483.7 / 94	1334 / 94

Conclusions of 2nd part.

- **Correlation analysis of ALEPH data was done in frame of Kozlov model**
- **To explain the increase of $C_2(Q)$, modification factor $(1+\epsilon Q)$ should be edit to the theoretical C_2 function**
- **This factor can be explained by the presence of the non equilibrium processes (?) – In Kozlov approach stationary (equilibrium) state is assumed.**