# Electrodynamics of black holes

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## Plan

- Astrophysical Introduction
- Unipolar Inductor
- Radio Pulsar Braking
- Black Hole Braking
- Conclusion

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# Plan

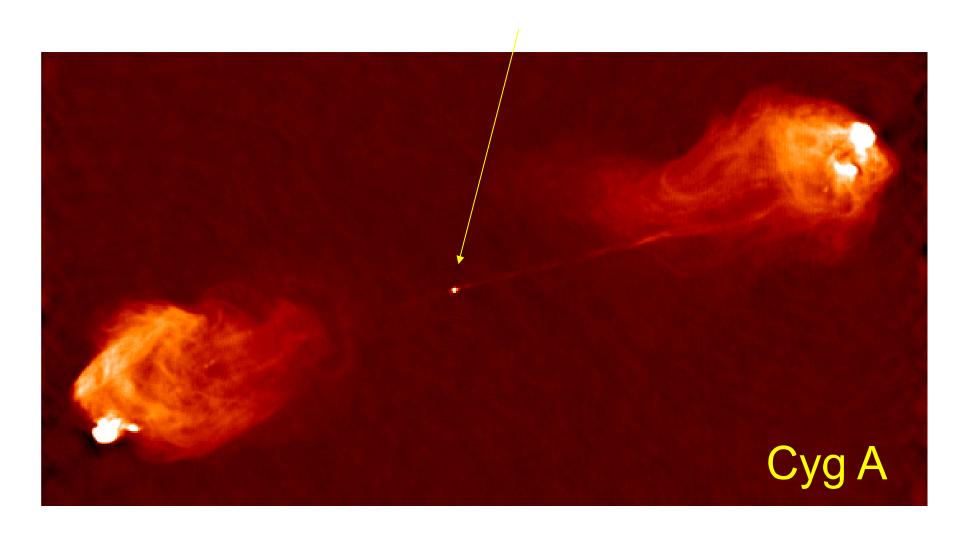
- Astrophysical Introduction
- Black Hole Braking
- Conclusion

# **Astrophysical Introduction**

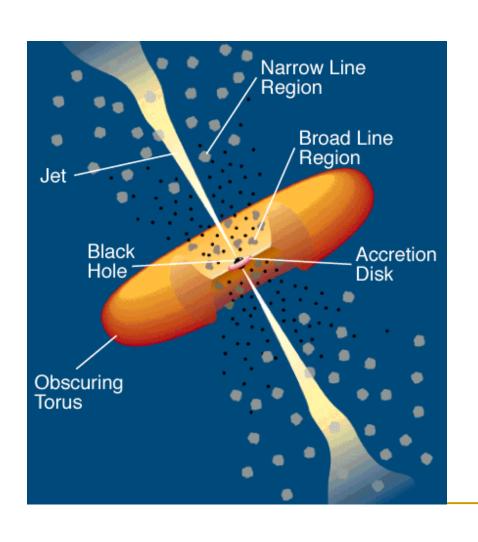
- Active Galactic Nuclei
- Young Stellar Objects
- Radio Pulsars
- Microquasars

Most of them have jets

# Active Galactic Nuclei (AGN) $M \sim (10^8-10^9)M_{\odot}$ , $R \sim 10^{13}$ cm

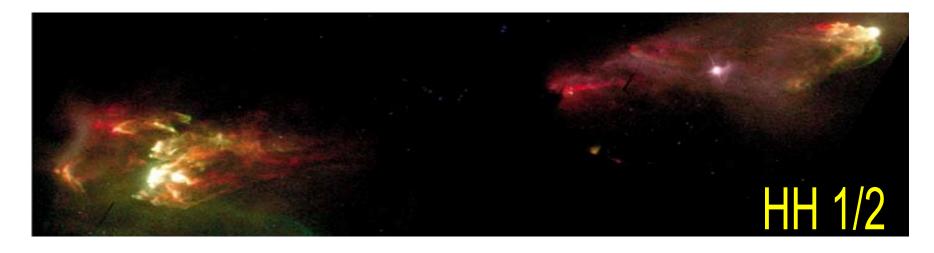


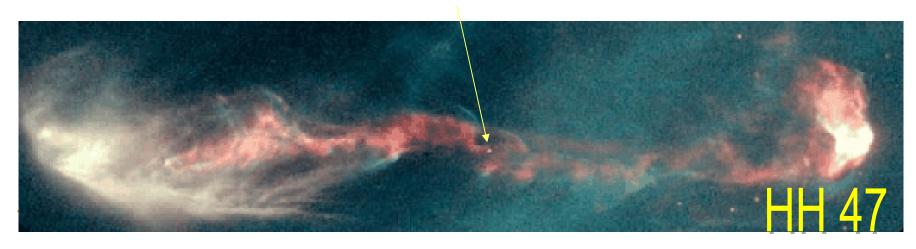
#### Active Galactic Nuclei (model)



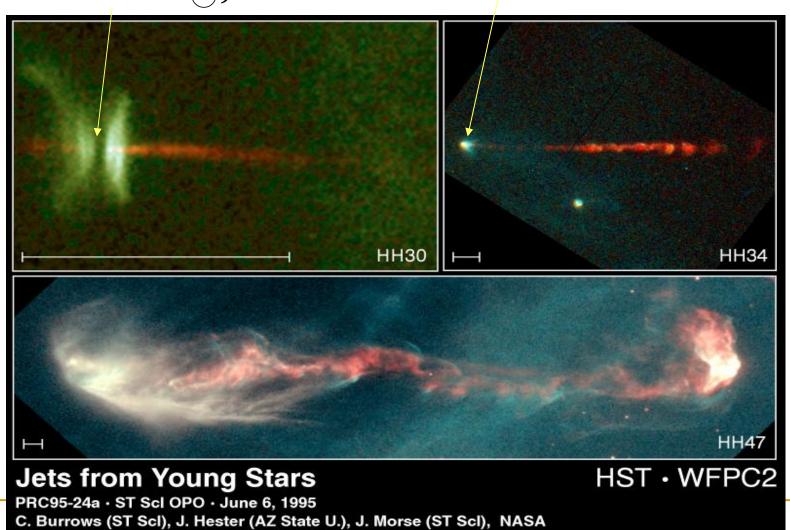


## Young Stellar Objects (YSO) M ~ 10M<sub>•</sub>, R~ 10<sup>10</sup>cm

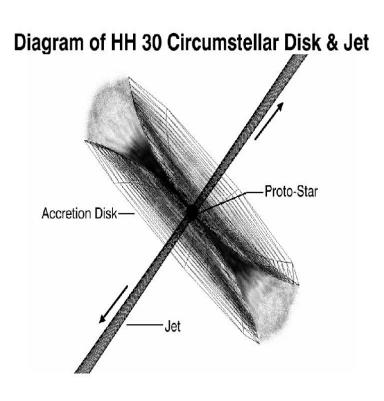




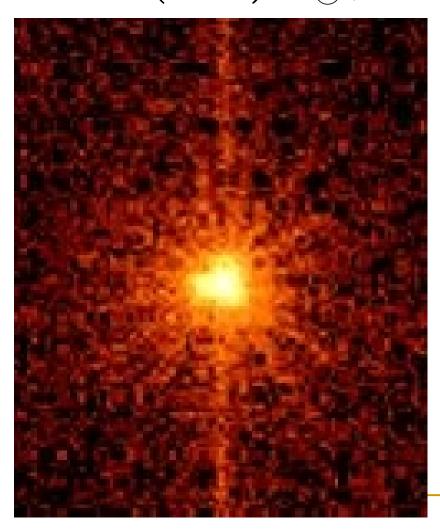
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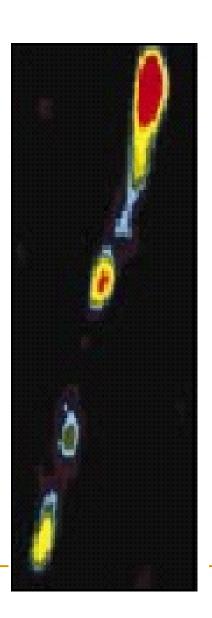


#### Young Stellar Objects (model)

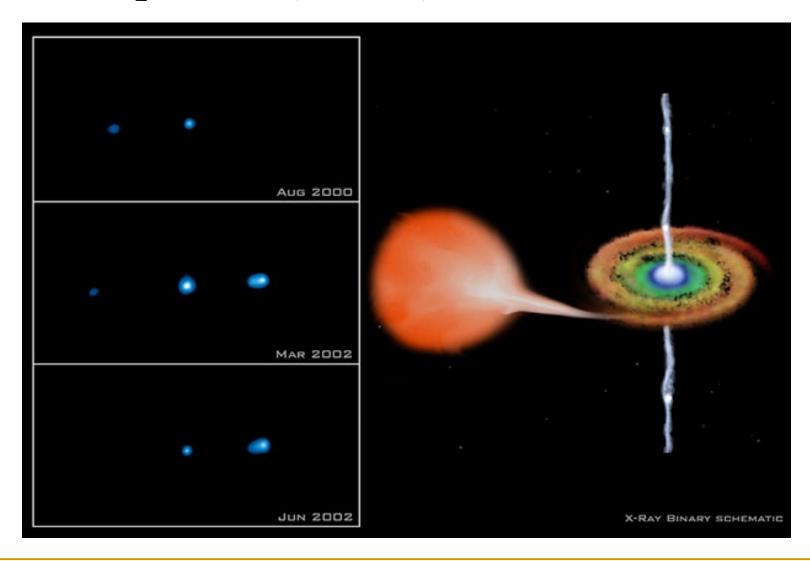


# Microquasars ( $\mu$ QSO) M ~ (3-10)M $_{\odot}$ , R ~ 10 $^6$ cm

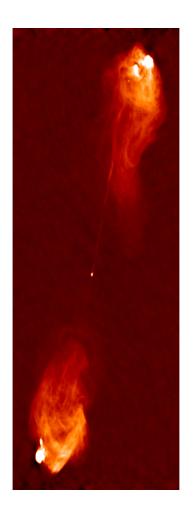


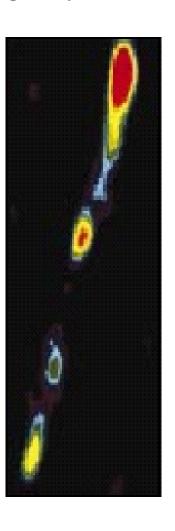


#### Microquasars (model)



#### The same mechanism?







The same mechanism?

- Thermal (gas pressure)?
- Radiative (radiation pressure)?
- Electromagnetic (Ampere force)?

#### Now it is not only the theoretical question

#### Radioastron -

following parameters:

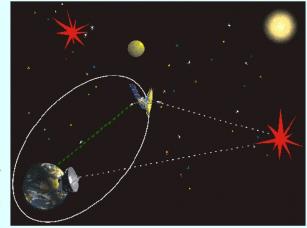
Perigee radius: ≥10,000 km

Initial inclination: 51.6°

Average apogee radius: 350,000 km

Average period of revolution: 9.5 d

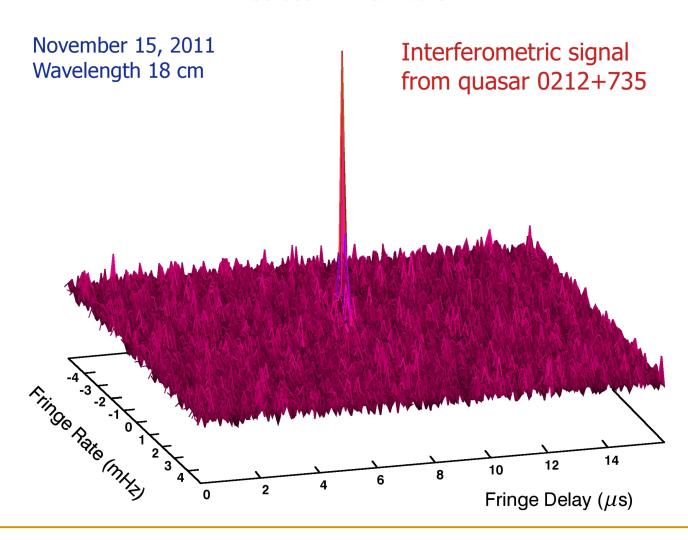
#### The ground-space radio interferometer (artist's view)



Band	Р	L	С	K
Frequencies (MHz) of observations	327	1665	4830	18392-25112
Bandwidth (MHz) for each polarization	4	32	32	32
Fringe size (µas) [base line 350 000 km]	540	106	37	7,1 -10
Min. cor. flux (mJy) [RMS with EVLA, 300 s integration time]	10	1,3	1,4	3,2

Basic parameters of the Space Radio Telescope

#### 100 000 km from Earth

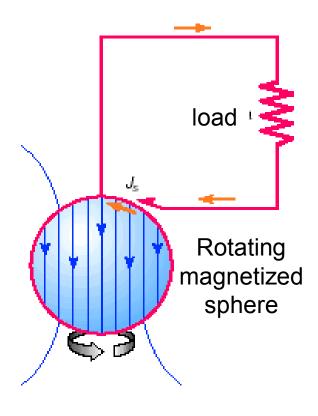


# Unipolar Inductor

Rotating magnetized sphere can work as a battery (the source of the electric current)

#### Unipolar Inductor

- 1. Electric circuit is to be touched to the sphere at different latitudes.
- 2. Electric circuit is to rotate with the angular velocity  $\Omega$  which differs from the angular velocity of a sphere.
- 3. The energy source is the kinetic energy of the rotation.
- 4. EMF does not result from the Faraday effect.



#### Unipolar Inductor (theory)

Freezing-in condition

$$E + V \times B/c = 0$$

$$(E' = j / \sigma = 0)$$

EMF (not connected with the Faraday effect)

$$E \sim (V/c)B, \qquad U = EL$$

Nonzero charge density

$$\rho_{\rm e} \sim E/L$$
  $(E \sim Q/L^2)$ 

$$(div \mathbf{E} = 4\pi \rho_{\rm e})$$

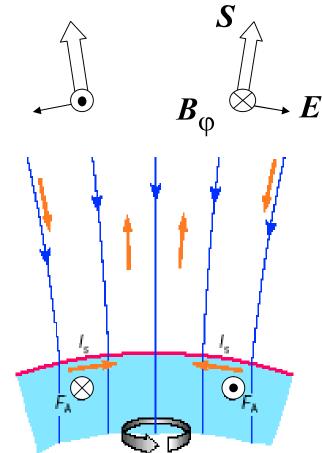
for rotation 
$$V = \Omega r$$

$$\rho_{\rm e} = -\frac{\Omega B}{2\pi c}$$

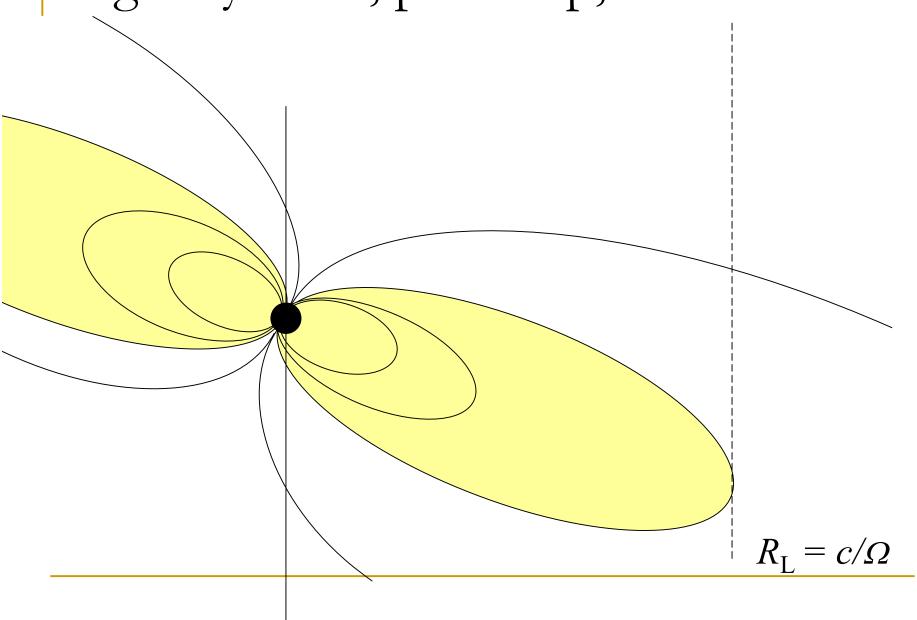
# Radio Pulsar Braking

Rotating neutron star works as a unipolar inductor loosing its rotational energy by electric currents flowing in the magnetosphere.

- 1. Neutron star braking results from Ampere force  $F_A = (1/c) J_s \times B$  of the surface currents  $J_s$  closing volume currents flowing in the magnetosphere.
- 2. Within light cylinder the energy flux is determined by the flux the electromagnetic energy  $S = (c/4\pi) E \times B$



### Light cylinder, polar cap, corotation



- Corotation
- Light cylinder
- Polar cap

$$V = \Omega r$$

$$R_{\rm L} = c / \Omega$$

$$R_0 \sim R \left(\frac{\Omega R}{c}\right)^{1/2}$$

W = III

$$U \sim ER_0 \sim (\Omega R_0 / c) B_0 R_0$$
  
 $I \sim \rho_e c \pi R_0^2 \sim \Omega B_0 R_0^2$ 

$$W \sim B_0^2 \left(\frac{\Omega}{c}\right)^2 R_0^4 c \sim \frac{B_0^2 \Omega^4 R^6}{c^3}$$

$$W \sim B_0^2 \left(\frac{\Omega R_0}{c}\right)^2 R_0^2 c$$

#### For radio pulsars

- magnetic field  $B_0 \sim 10^{12} G$
- neutron star radius  $R \sim 10^6 \, cm$

the total energy loss  $W \sim 10^{31} - 10^{33} \, \text{erg/s}$  corresponds to  $J_r \, \Omega \, \mathrm{d}\Omega/\mathrm{d}t$ 

For this mechanism is necessary to have

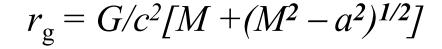
- 1. regular poloidal magnetic field,
- 2. rotation (inductive electric field E, EMF U),
- 3. longitudinal current I (toroidal magnetic field  $B_{\varphi}$ ).

# Black Hole Braking

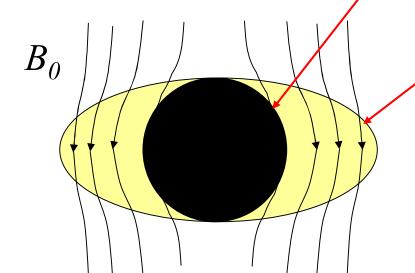
- Rotating black hole
- "No hair" theorem
- Blandford-Znajek process
- A problem
- The answer

#### Rotating Black Hole

#### **Gravitational radius**



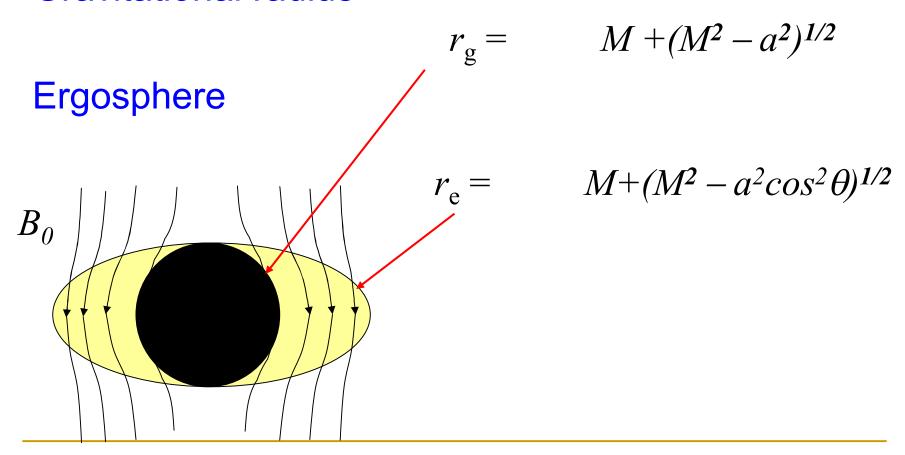




$$r_{\rm e} = G/c^2[M + (M^2 - a^2\cos^2\theta)^{1/2}]$$

## Rotating Black Hole

#### **Gravitational radius**



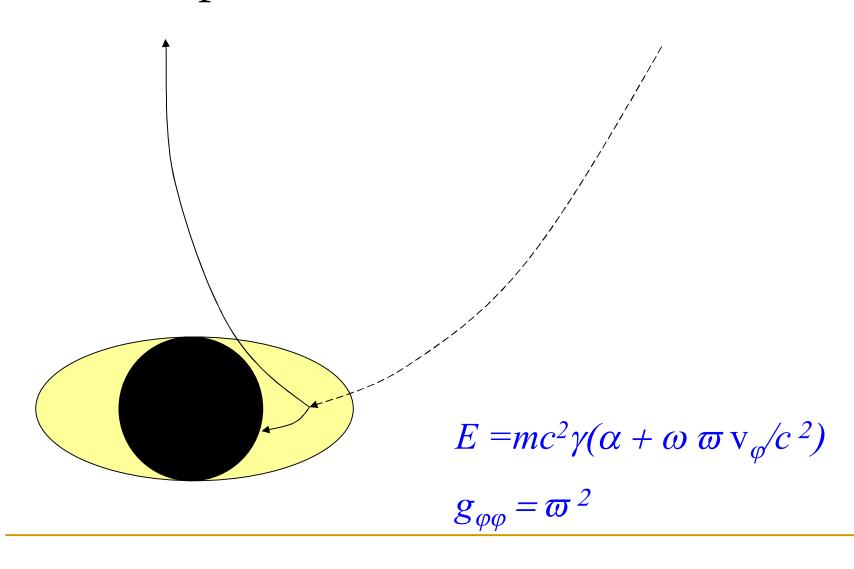
### Frame-dragging (Lense-Thirring effect)

Inside the ergosphere any body (including the light) is to rotate with the black hole. More important. At all distances "the rotation of the space" takes place. Thus, there are a distinguished reference frame rotating with Lense-Thirring angular velocity

$$\omega \sim \Omega_{\rm H} (r_g/r)^3$$

ZAMO (no gyroscope rotation).

#### Penrose process



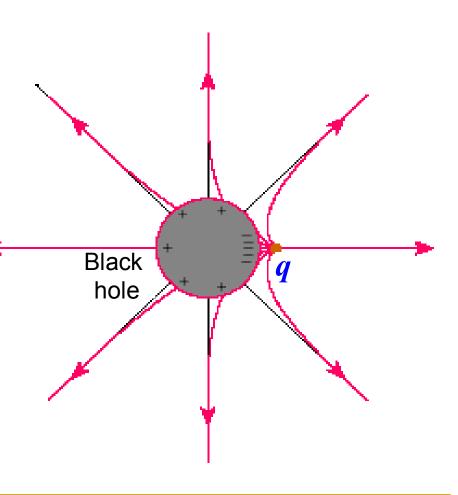
#### "No hair" theorem

Black hole cannot have its own magnetic field.

But it can be submerged into the external magnetic field generated e.g. in the accretion disk.

#### "No hair" theorem

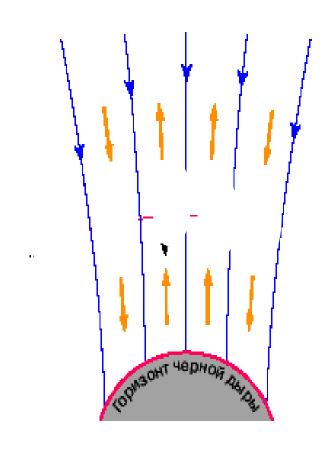
- 1. Electric field of a charge *q* located in the very vicinity of the horizon looks like that of the charge *q* located in the very center of the black hole.
- 2. Thus, the electric field of two different charges (a dipole) vanishes as the charges approach themselves to the horizon.
- 3. The same for the magnetic dipole, quadrupole, etc.



# Blandford – Znajek process (1977) The analogy with radio pulsars

- 1. External poloidal magnetic field (generated in a disk).
- 2. Rotation of a black hole.
- 3. Longitudinal current *I*.

$$W_{\mathrm{BZ}} \sim B_0^2 \left(\frac{\Omega R_g}{c}\right)^2 R_g^2 c$$



## Blandford – Znajek process (1977)

$$W_{\rm BZ} \sim B_0^2 \left(\frac{\Omega R_g}{c}\right)^2 R_g^2 c$$

For the central engine in active galactic nuclei

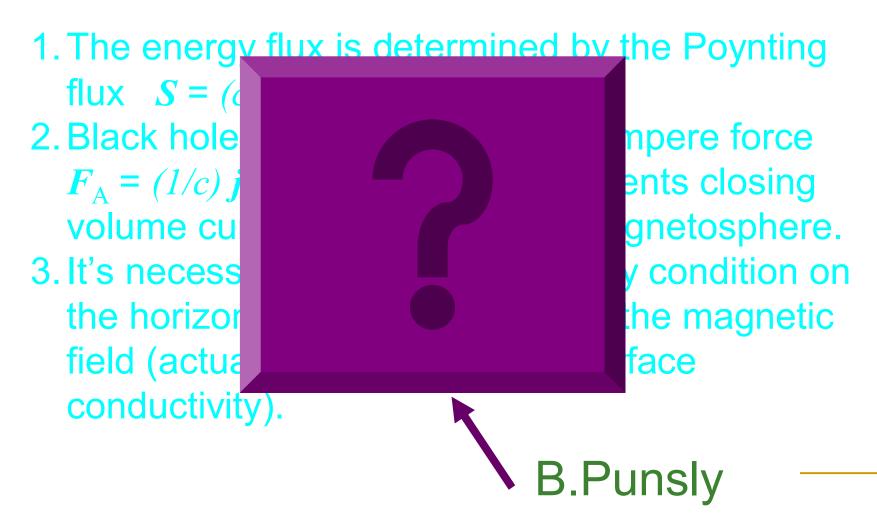
- magnetic field  $B_0 \sim 10^4 G$
- black hole radius  $R_g \sim 10^{13} \, cm$

the total energy loss  $W \sim 10^{45} \, \text{erg/s}$  is in agreement with observations.

# Blandford – Znajek process (1977) The analogy with radio pulsars

- 1. The energy flux is determined by the Poynting flux  $S = (c/4\pi) E \times B$ .
- 2. Black hole braking results from Ampere force  $F_A = (1/c) j \times B$  of the surface currents closing volume currents flowing in the magnetosphere.
- 3. It's necessary to add the boundary condition on the horizon fixing the structure of the magnetic field (actually, to introduce the surface conductivity).

# Blandford – Znajek process (1977) The analogy with radio pulsars



### The problems

The horizon isn't in casual connection with the outer space\*.

The horizon isn't a material boundary of a black hole\*.

Black hole cannot be a battery (the source of the exterior forces)\*.

<sup>\*</sup> They are very serious objections.

### The problems

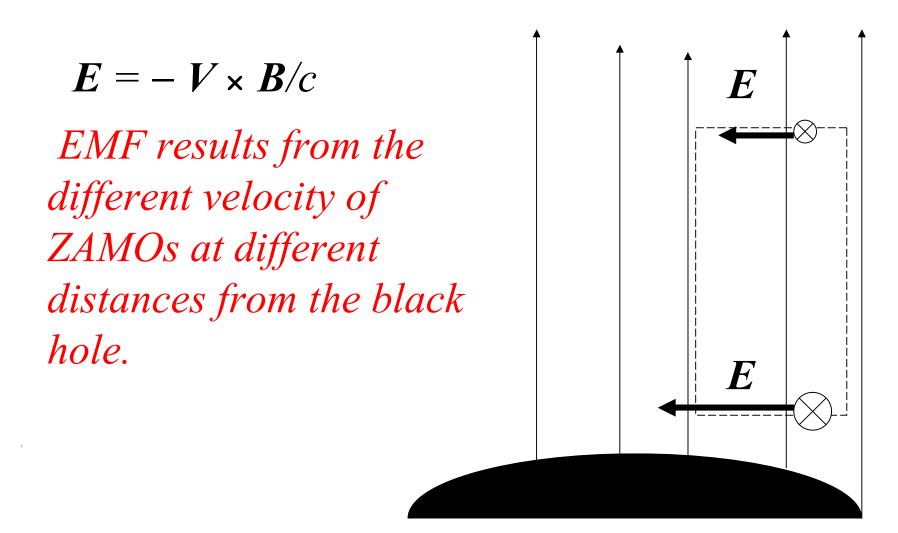
Poynting flux  $S = (c/4\pi) E \times B$  directs downwards in the very vicinity of the horizon.

#### The answer (1)

- EMF results from the Lense-Thirring effect, not the exterior forces inside the black hole.
- Frame dragging results in "the flow of the space" through any circuit which is equivalent to alternative magnetic flux.
- EMF appears over the horizon.

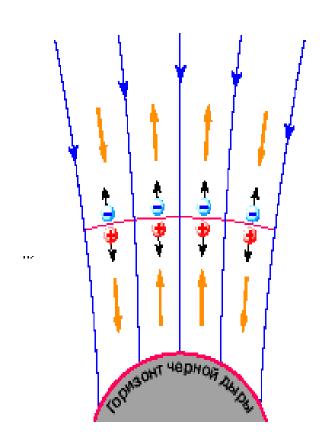
### The answer (1)

The motion relatively ZAMOs results in the generation of the electric field.



# Blandford – Znajek process

- To realize electric current it's necessary to create pairs above the horizon
- Blandford-Znajek process is electromagnetic realization of Penrose process



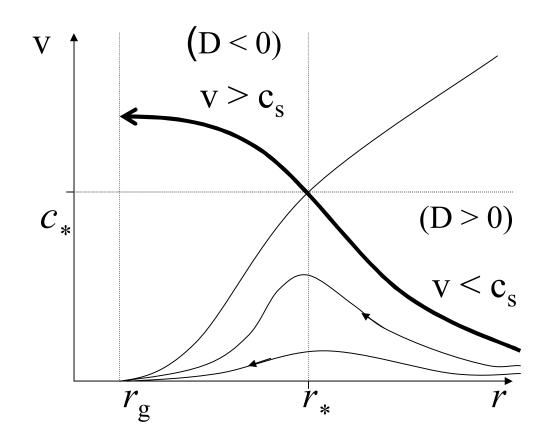
#### Bondi accretion onto black hole

The transonic flow only is possible.

$$v(r_g) > 0$$
 means

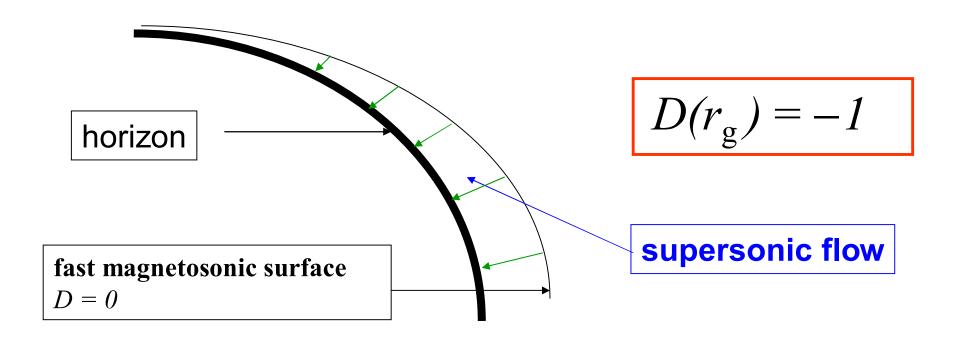
$$v\left(r_{g}\right)>c_{*}.$$

The critical condition fixes the accretion rate.



#### Force-free limit

Force-free limit of the critical condition on the fast magnetosonic surface is "the boundary condition on the horizon"



### Grad-Shafranov equation in Kerr metric

$$A \left[ \frac{1}{\alpha} \nabla_{k} \left( \frac{1}{\alpha \varpi^{2}} \nabla^{k} \Psi \right) + \frac{1}{\alpha^{2} \varpi^{2}} \frac{\nabla^{i} \Psi \cdot \nabla^{k} \Psi \cdot \nabla_{i} \nabla_{k} \Psi}{(\nabla \Psi)^{2} D} \right]$$

$$- \frac{A \nabla_{k}^{'} F \cdot \nabla^{k} \Psi}{2\alpha^{2} \varpi^{2} (\nabla \Psi)^{2} D} + \frac{\nabla_{k}^{'} A \cdot \nabla^{k} \Psi}{\alpha^{2} \varpi^{2}} + \frac{(\Omega_{F} - \omega)}{\alpha^{2}} (\nabla \Psi)^{2} \frac{d\Omega_{F}}{d\Psi}$$

$$+ \frac{32\pi^{4}}{\alpha^{2} \varpi^{2} \mathbf{M}^{2}} \frac{\partial}{\partial \Psi} \left( \frac{G}{A} \right) - 16\pi^{3} \mu n \frac{1}{\eta} \frac{d\eta}{d\Psi} - 16\pi^{3} n T \frac{ds}{d\Psi} = 0$$

$$D = \frac{A}{\mathbf{M}^2} + \frac{\alpha^2}{\mathbf{M}^2} \frac{B_{\phi}^2}{B_{p}^2} + \frac{1}{u_{p}^2} \frac{A}{\mathbf{M}^2} \frac{c_{s}^2}{1 - c_{s}^2}; \mathbf{M}^2 = \frac{4\pi\mu\eta^2}{n}$$

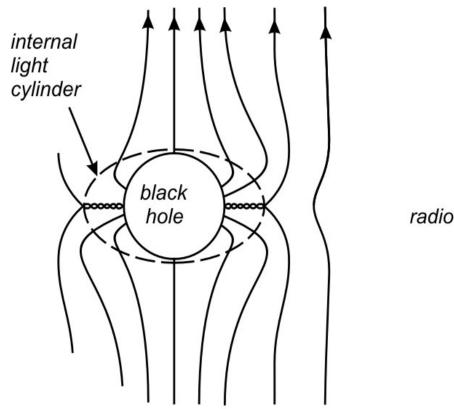
$$F = \frac{64\pi^4}{\mathbf{M}^4} \frac{K}{A^2} - \frac{64\pi^4}{\mathbf{M}^4} \alpha^2 \varpi^2 \eta^2 \mu^2.$$

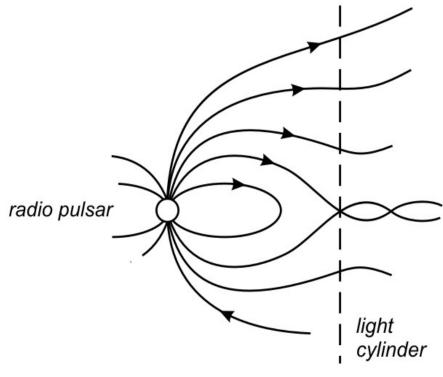
### The answer (2)

- Longitudinal current is determined by the critical condition on the fast magnetosonic surface located outside the horizon.
- For this reason the conditions on the horizon cannot affect the energy loss.

# Black hole magnetosphere

For the magnetosphere filled with plasma all the magnetic flux crossing the internal light surface, intersects the horizon as well

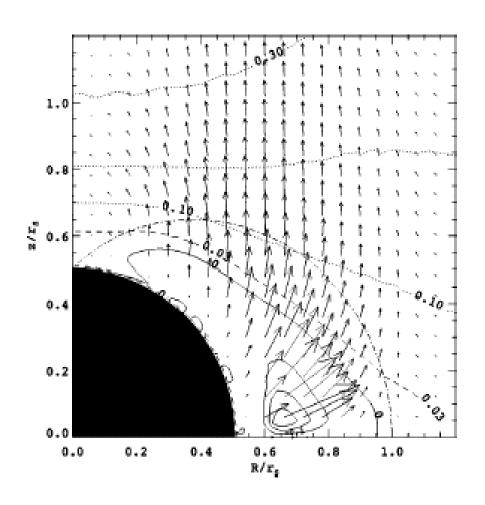




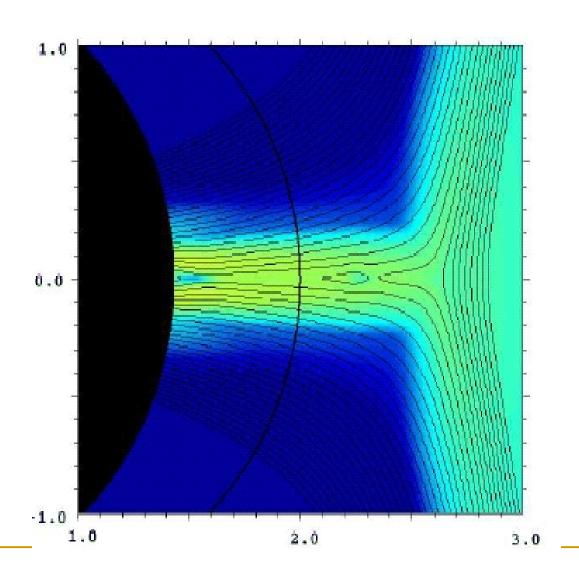
a)

б)

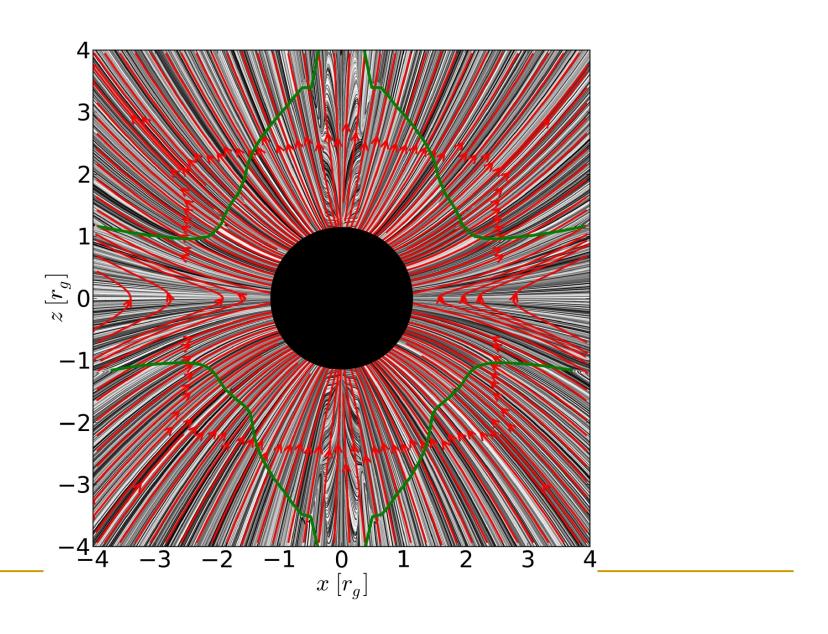
#### S.Koide, Phys. Rev. D., 67, 104010, 2003



#### S.S. Komissarov, MNRAS, **359**, 801, 2005



#### A.Tchekhovskoy, private communication



# Conclusion

Rotating black hole submerged into the external magnetic field can lose its rotating energy by electric currents flowing in the magnetosphere.

# Conclusion

The mechanism of the black hole braking differs from that of radio pulsars.

- It doesn't result from surface currents flowing along black hole horizon.
- EMF results from frame dragging (Lense-Thirring effect) over the horizon, not from exterior forces inside the black hole.

# Conclusion

The mechanism of the black hole braking differs from that of radio pulsars.

• Electric current fixed the energy loss does not depend on the "conductivity of the horizon". It is determined y the critical condition on the fast magnetosonic surface located over the event horizon.