

# Backward proton production in nuclei-nuclei collisions

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# Plan

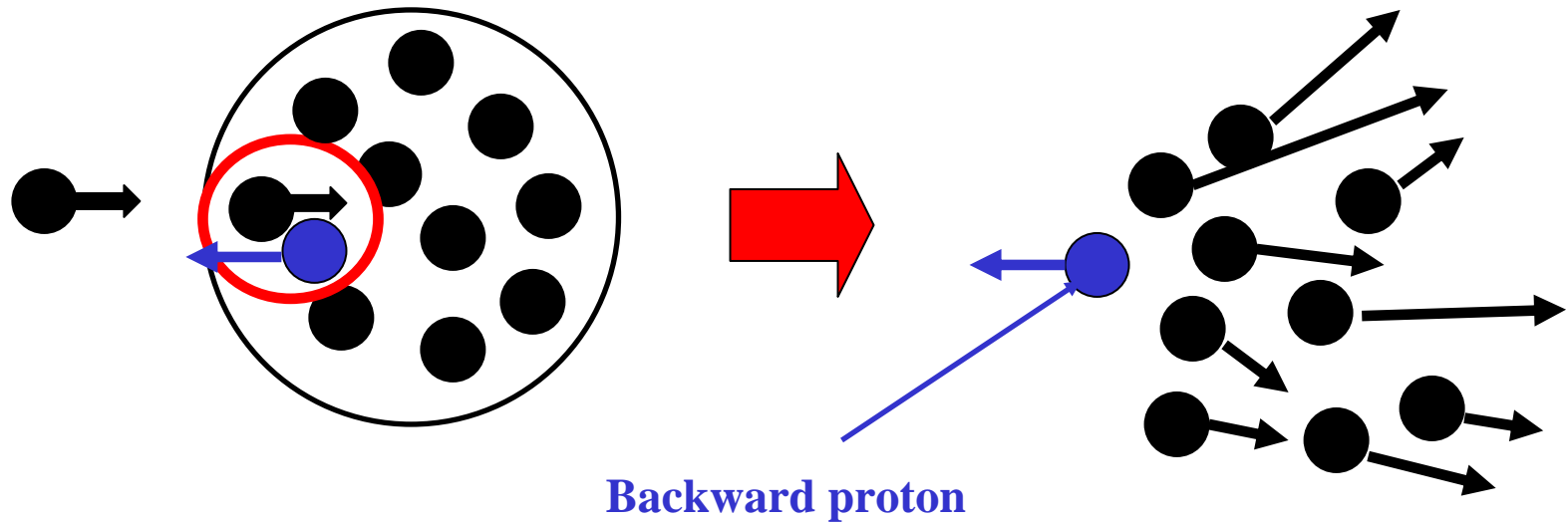
1. Backward proton production in nucleon-nuclei collisions.
2. Backward proton production in nuclei-nuclei collisions.
3. Counting rate for CBM conditions.
4. Possible set-up for backward proton detection on CBM.
5. Conclusions.

## Some experimental features of backward proton production in nucleon-nuclei collisions :

1. The spectra show a weak beam energy dependence starting from momentum 4 GeV/c.
2. The spectrum shape is independent from target mass number
3. There is also a weak spectrum shape dependence from the type of beam particle (for the lepton, proton and light nuclei ( $A < 12$ )).
4. The cross section has a strong atomic mass dependence  $\sim A^n$ ,  $n > 1$ .
  - A. Baldin, *Part. And Nucl.*, V.8(3), p.429 (1977)
  - V. Lukyanov, A. Titov, *Part. And Nucl.*, V.10(4), p.815 (1979)
  - V. Stavinsky, *Part. And Nucl.*, V.10(5), p.979 (1979)
  - L. Frankfurt and M. Strikman, *Phys.Rep.*, V.76, p.215, (1981)
  - A. Efremov, *Part. And Nucl.*, V.13(5), p.613 (1982)

The most sophisticated explanation:

High momentum component in target nucleus  
(different nature for different models)



$$P_{\text{int}} \geq 0.3 \text{ GeV} / c \text{ i.e. } L_{\text{NN}} \leq 1 \text{ Fm.}$$

Why it can be interesting in nuclei-nuclei collisions?

Analogy with high  $P_T$  hadrons production at RHIC.  
Jet Quenching at RHIC.

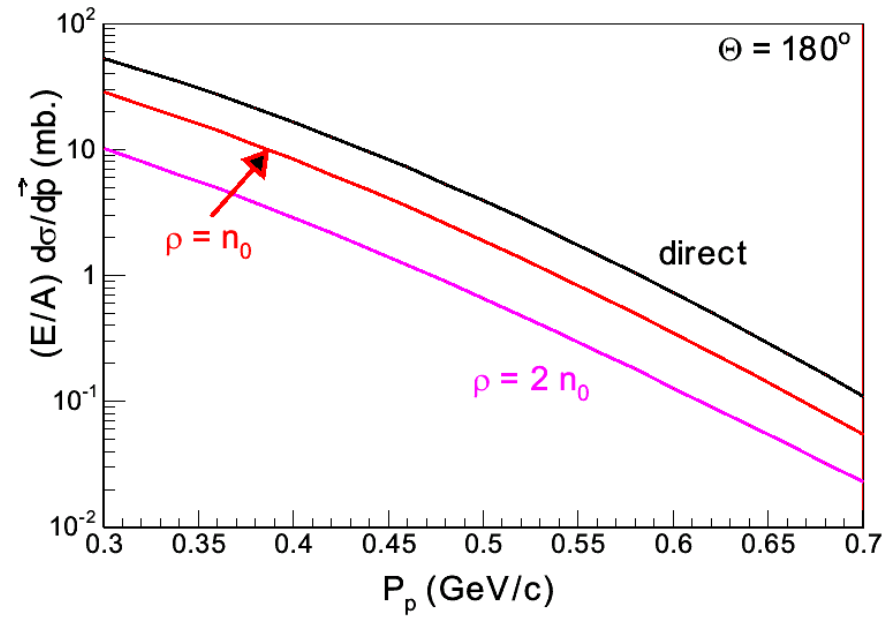
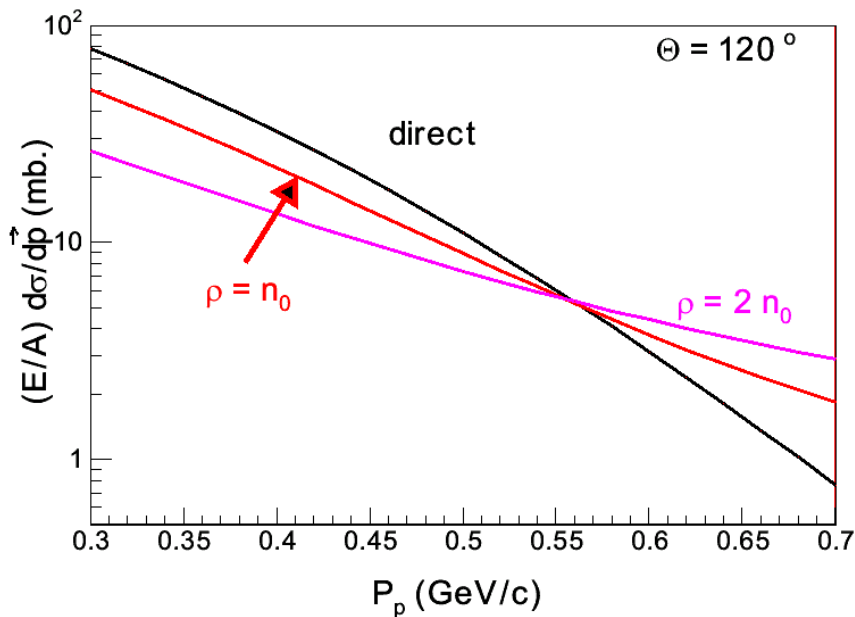
1. Backward protons are produced during the first stage of collision (See spectrum on the next slides) .
2. Spectrum have to be sensitive to the density of the compressed fireball.

The possible mark for initial state.

Spectra can not be explain as a thermal from the fireball with reasonable velocity and temperature (See next slide).

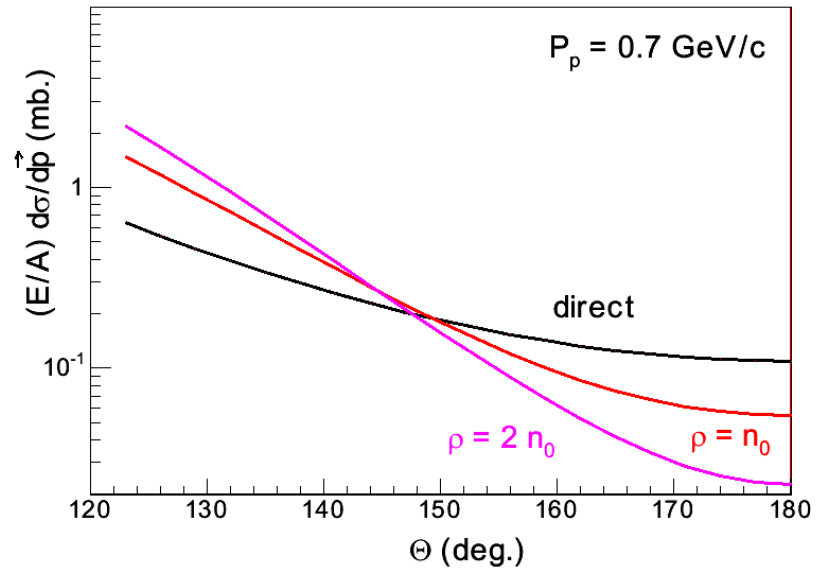
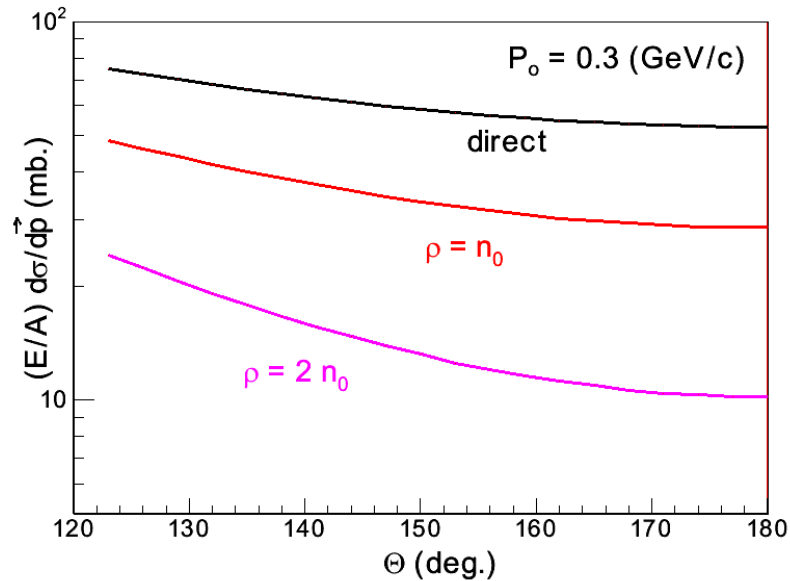
# Medium density effects .

“direct spectrum” and spectrum after proton penetrating through 4 fm hadronic matter with two different densities.



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“direct spectrum” and spectrum after proton penetrating through 4 fm hadronic matter with two different densities.



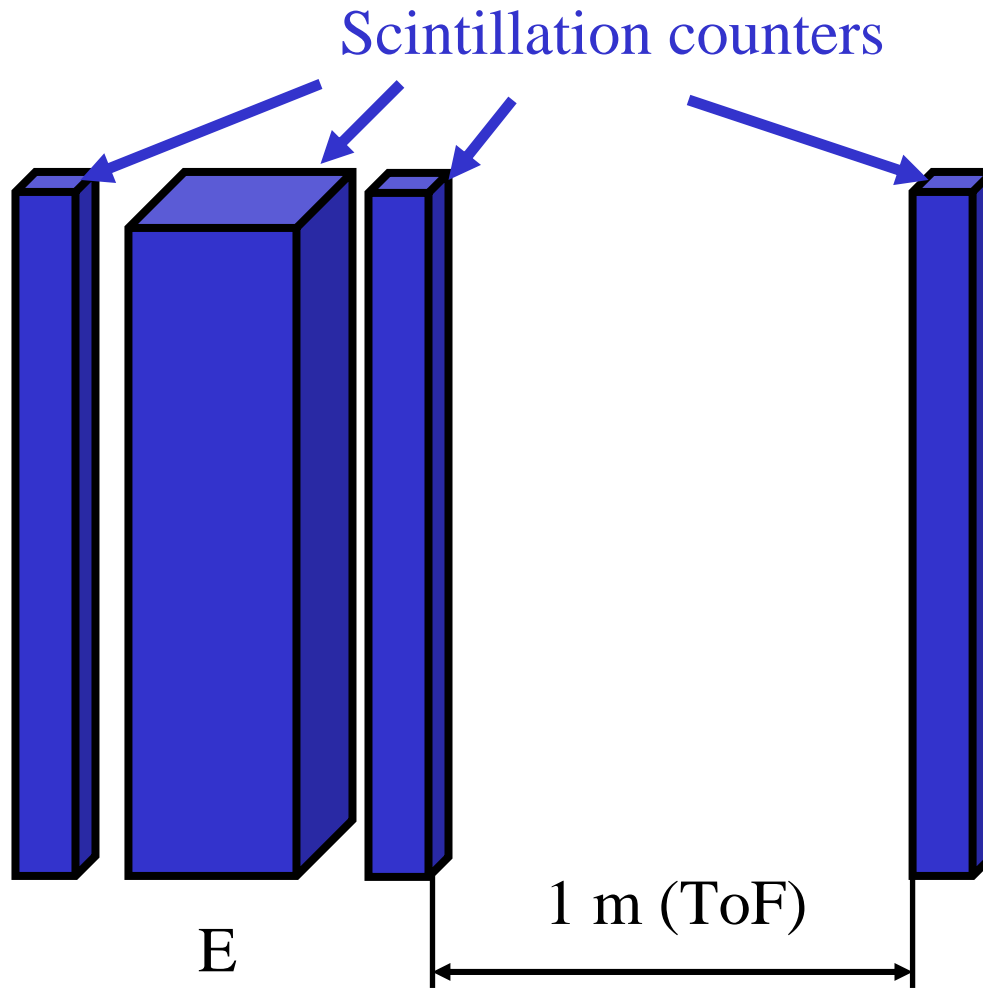
# Counting rates for CBM condition.

The counting rates were estimated for the wall of detectors **1x1 m<sup>2</sup>** situated on the distance **2 m** at the backward direction. The estimations have been done for two different A dependencies ( $A^{(1/3)}$  and  $A^{(1)}$ ) and  **$I_t=0.1$  (g/cm<sup>2</sup>)**.

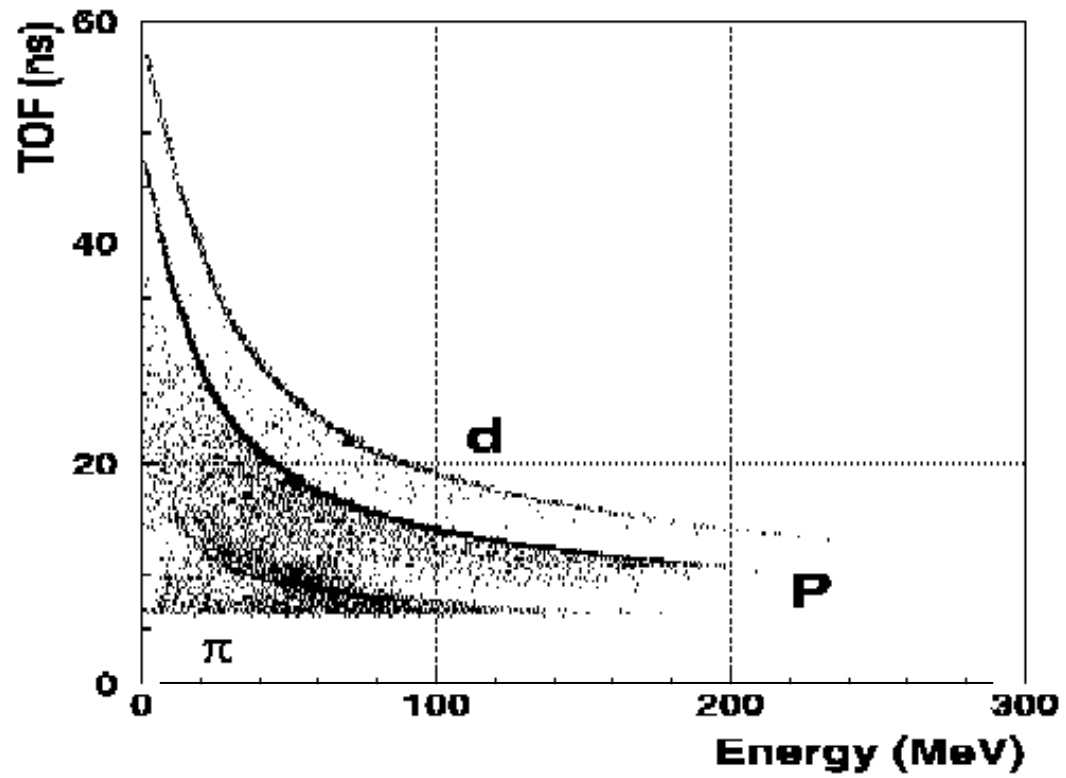
| counting rate ( $10^7$ 1/s.) | $A^{(1/3)}$    | $A^{(1)}$ |
|------------------------------|----------------|-----------|
| One proton                   | $6 \cdot 10^2$ | $10^6$    |
| Two protons                  | $10^1$         | $10^3$    |



# The possible set up .

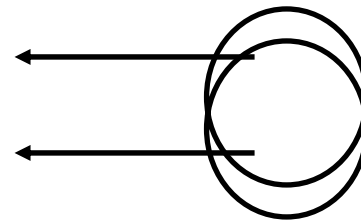
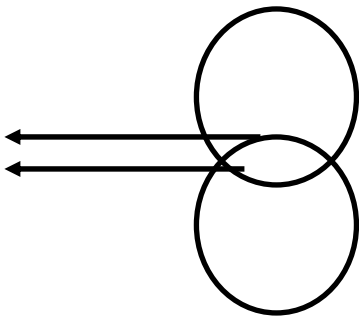


Hadron spectrometer without a magnet

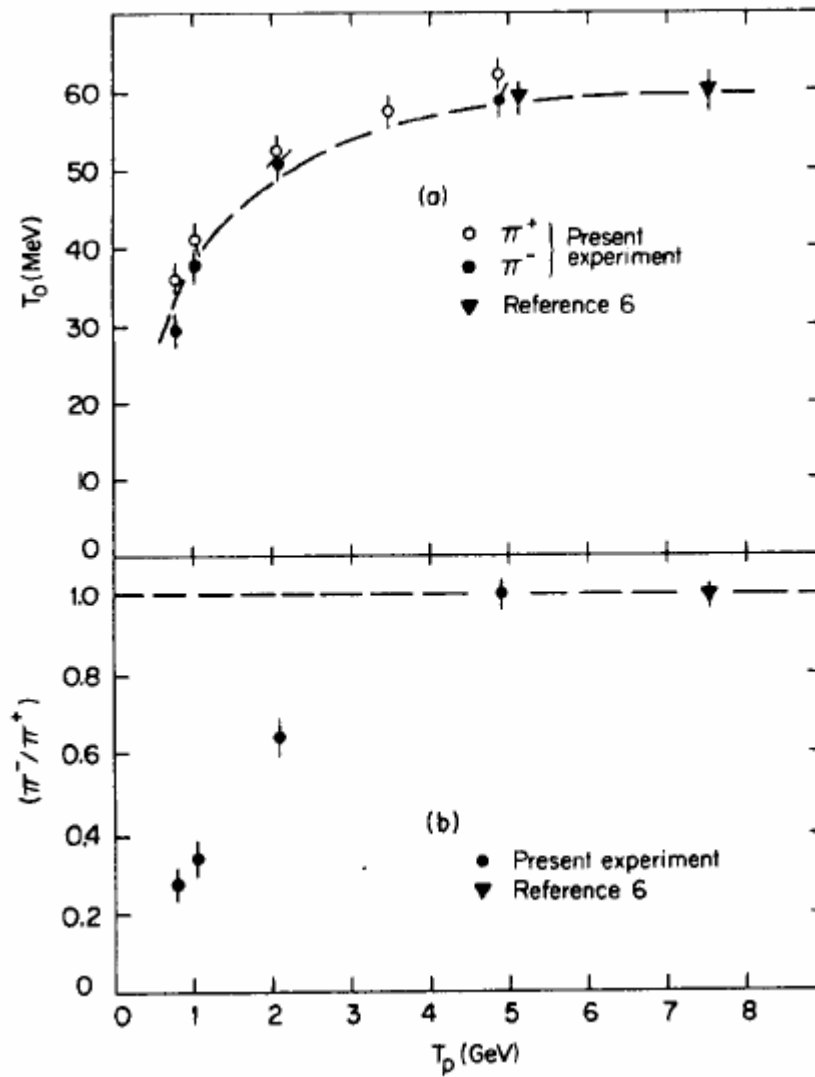


# Conclusions.

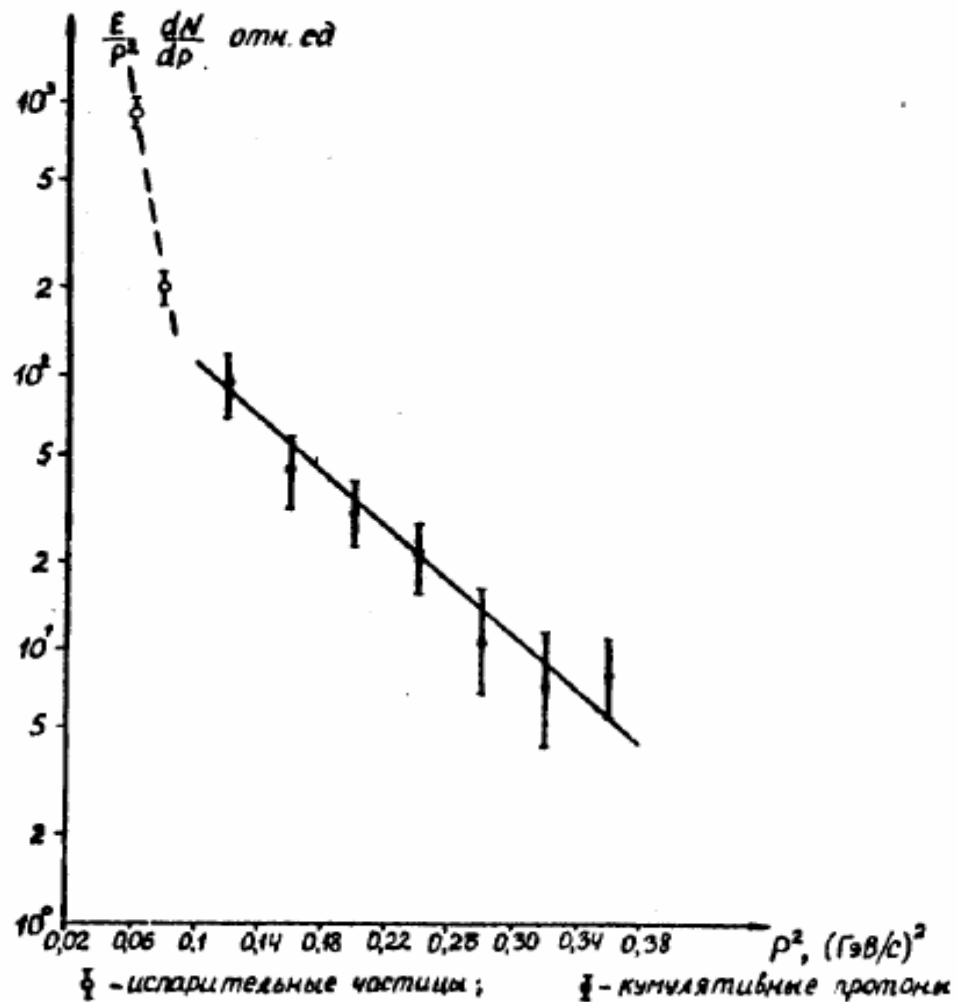
1. With this type of spectrometers it is possible to detect not only protons, but also pions, deuterons and neutrons with 20 % efficiency.
2. The numbers of detected deuterons are expected to be three times smaller than the numbers of protons.
3. The numbers of detected pions is expected to be at about 5% with respect to the protons.
4. The counting rate for two protons detection is high enough for HBT.



# Backup slides

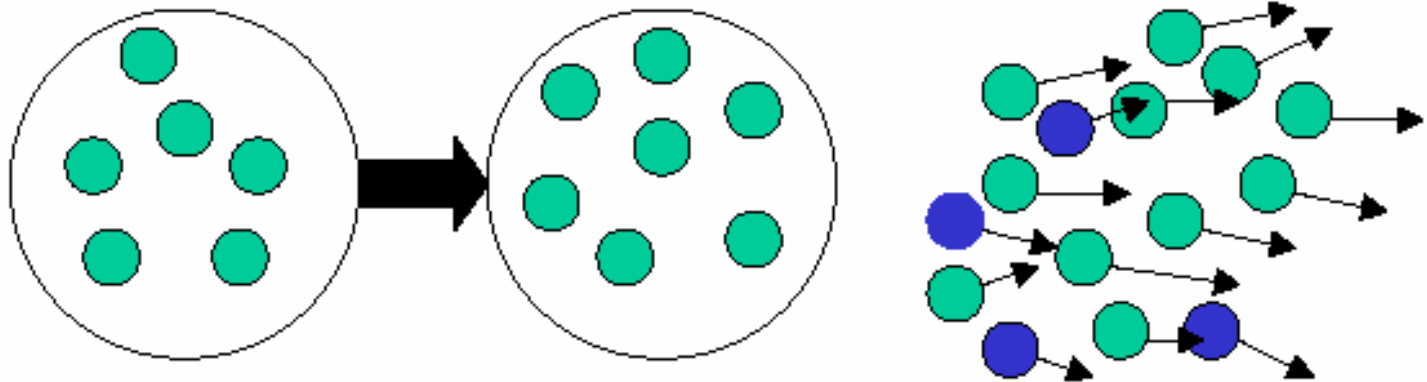


- **A.M.Baldin** *Sov.J.Part and Nucl.*, V8(3),p.429 (1977)
- **L.S.Shreder et al.**, *Phys.Rev.Lett.*, V43, No.24,p.1787 (1979)

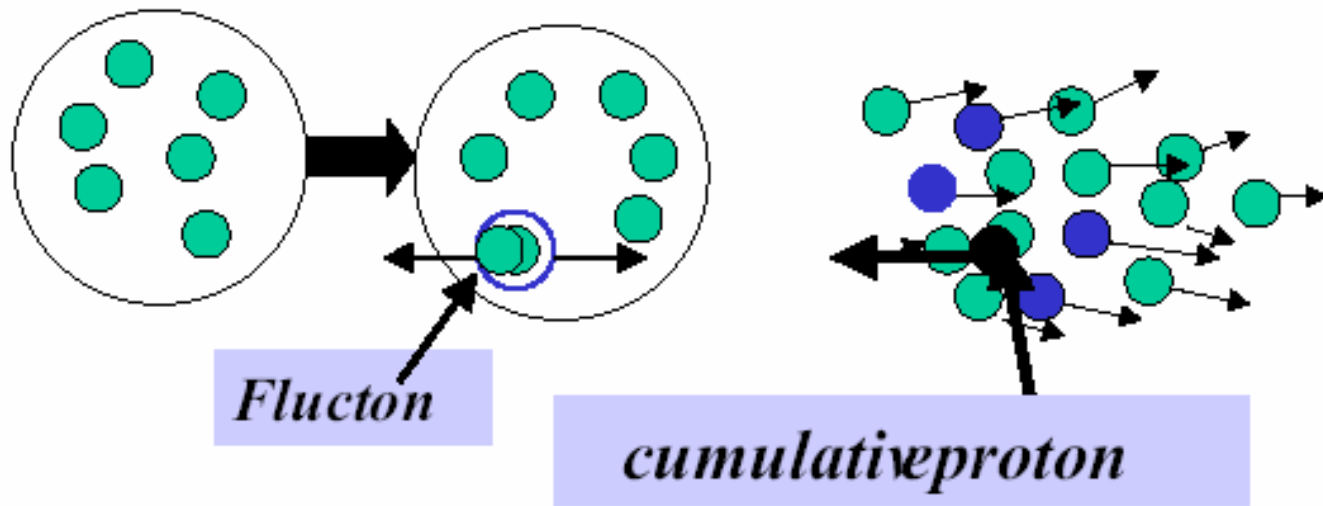


*P.Ammar et al., Pisma v JETF, V.49, No.4, p.189 (1989)*

*without cumulative proton*

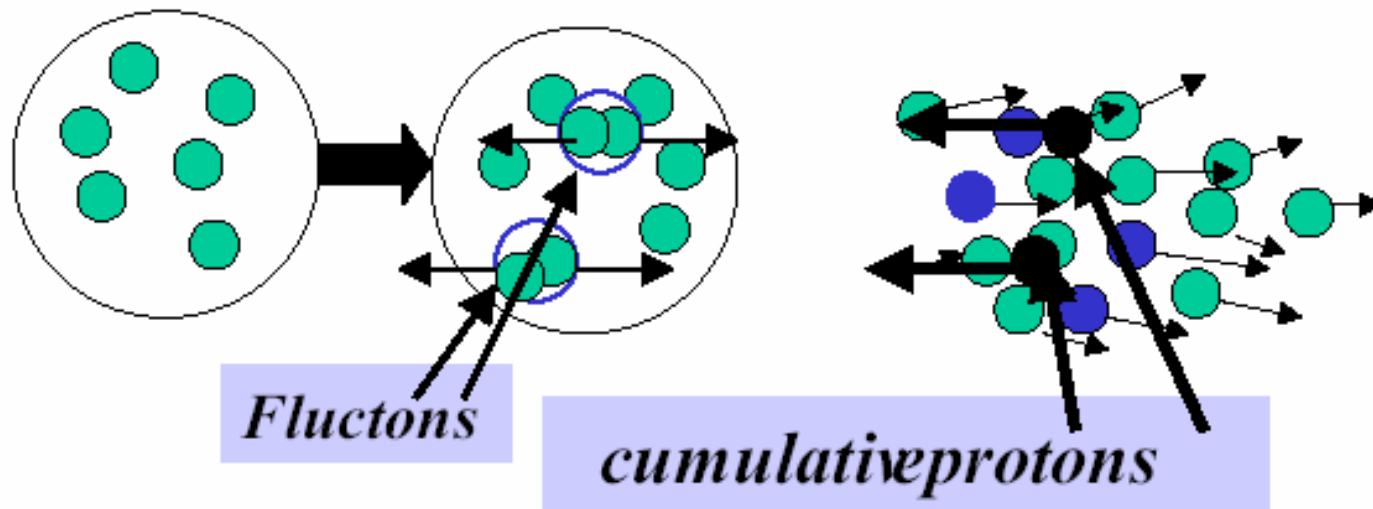


*one cumulative proton*

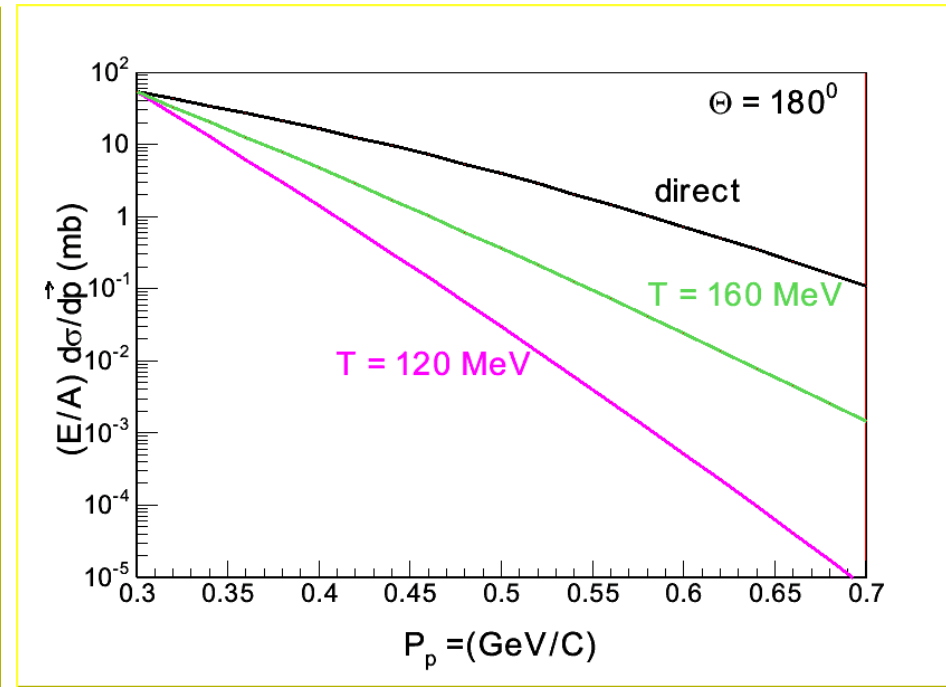
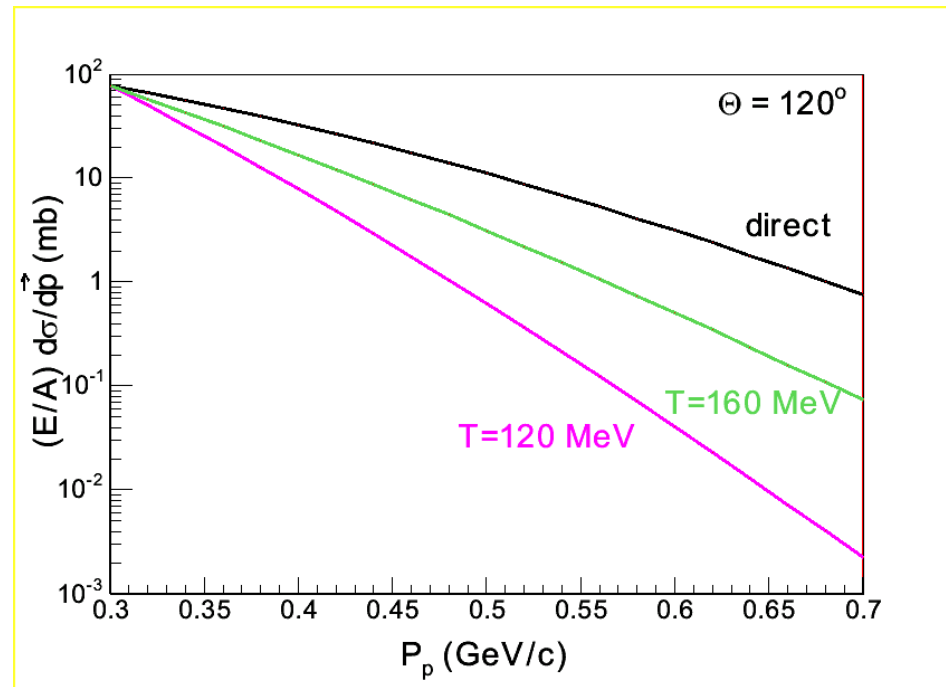




*twocumulative proton*



Cross section as a proton momentum from (pA) collisions and from a thermal fireballs with two temperatures in (AA) collisions  
 (P(beam) = 20 GeV/c).



Cross section as a proton emission angle from (pA) collisions and from a thermal fireballs with two temperatures in (AA) collisions  
( $P(\text{beam}) = 20 \text{ GeV}/c$ ).

