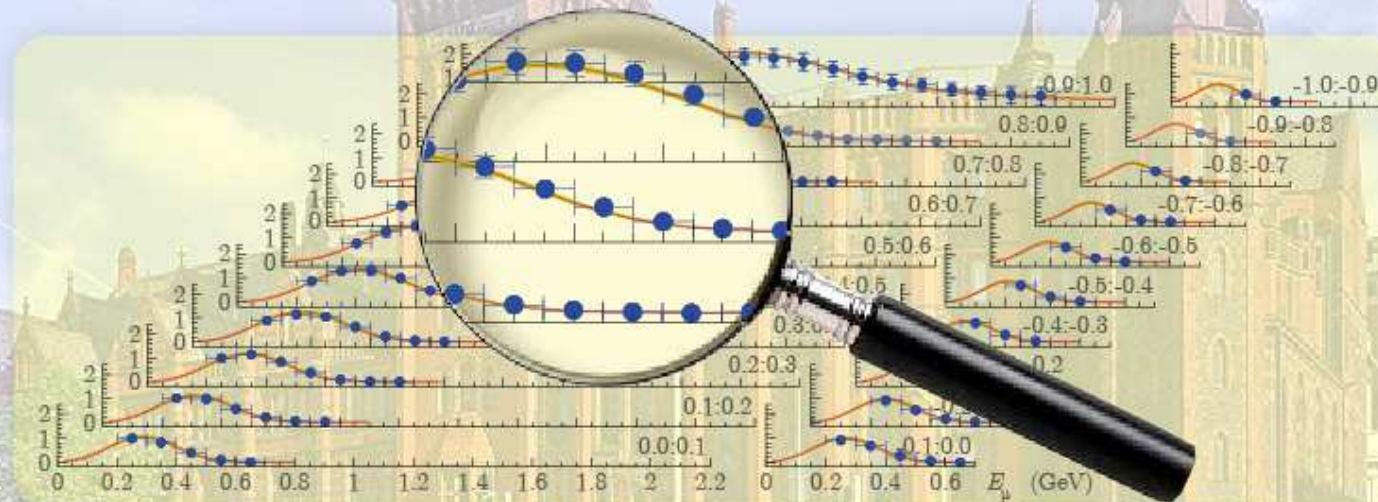


# Running axial mass for quasielastic neutrino-nucleus scattering

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Workshop on Global Fits to Neutrino Scattering Data and Generator Tuning (NuTune2016)  
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The aim of our report is to provide a simple “**cookbook recipe**” for a precise description of the CCQE neutrino and antineutrino interactions with nuclei using only **two adjustable parameters**.



*Cooking a duck*



*NO! Ducking a cook*

We do not intend discussing the physical meaning of the proposed recipe; instead we wish to demonstrate its usefulness as an **empirical tool** which could be used in the data processing of the neutrino oscillation experiments with accelerator and atmospheric (anti)neutrino beams.

# Outline

- Introduction
- Total Cross Sections – Recent data
- FNAL MiniBooNE 2010-2013
- FNAL MINER $\nu$ A 2013
- T2K ND280 2016
- Parameters of the Smith–Moniz RFG Model
- “Golden” dataset for the global fits
- Comparison with the early golden data
- Comparison with certain datasets unclaimed in the global fit
- Conclusions

# Introduction

One of the main sources of ambiguity in predicted neutrino event rates in the accelerator experiments at low and intermediate energies is caused by **nuclear effects** for the charged-current quasielastic (CCQE)  $\nu/\bar{\nu}$  interactions with various detection targets.

Another closely-related problem arises from the experimental uncertainty of the **nucleon axial mass parameter,  $M_A$** , in the dipole model of the nucleon axial form factor

$$F_A(Q^2) = F_A(0) \left(1 + \frac{Q^2}{M_A^2}\right)^{-2}.$$

The most familiar Relativistic Fermi Gas (RFG) model cannot describe the up-to-date CCQE data (NOMAD, MiniBooNE, MINER $\nu$ A, T2K, etc.) with a unique value of  $M_A$ .

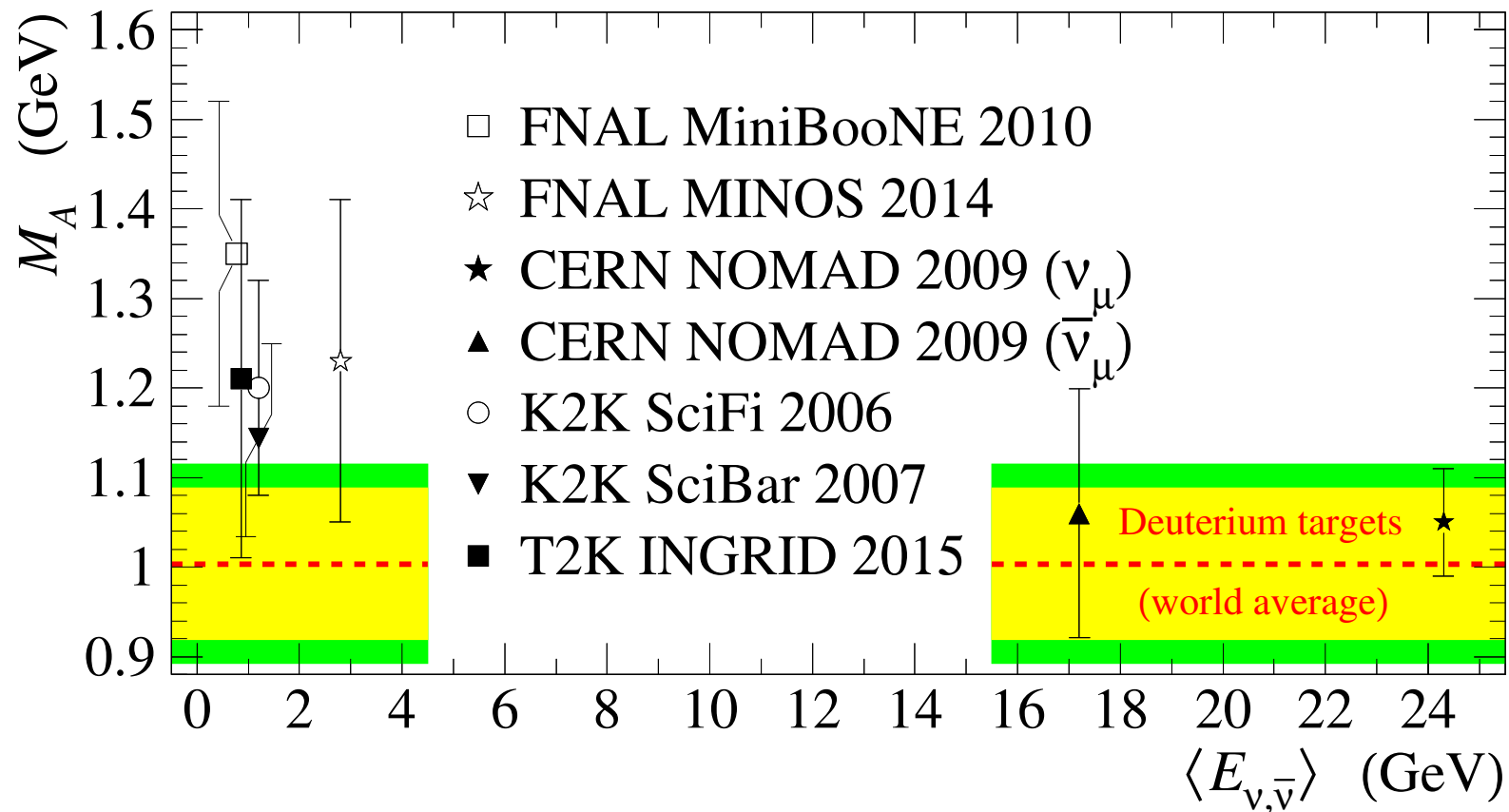
In fact, there are a lot of different models based on the Fermi-gas approach. One of the most popular models is the model by Smith and Moniz (SM).<sup>a</sup> Below, the term “RFG” is used for the slightly-modified SM model,<sup>b</sup> with an updated set of the input parameters.

R.A. Smith and E.J. Moniz, “Neutrino reactions on nuclear targets,” Nucl. Phys. B **43**, 605 (1972); erratum – *ibid.* **101**, 547 (1975).

K.S. Kuzmin, V.V. Lyubushkin and V.A. Naumov, “Quasielastic axial-vector mass from experiments on neutrino-nucleus scattering,” Eur. Phys. J. C **54**, 517 (2008) [arXiv:0712.4384 [hep-ph]].

## Some nuclear models currently on the market.

- **RFG Extensions:**
  - **Local Fermi Gas (LFG)** or **Local Density Approximation (LDA)**  
[Fermi momenta and separation energies are functions of the position in the nucleus]
  - **Spectral Function (SF)**  
[correlations between removal energies and initial state momenta of the nucleons, Gaussian spectrum of each shell]
- **Scaling and Superscaling (SuSA) models.**
- **Relativistic Mean Field (RMF) models.**
- **Relativistic Green's Function (RGF) models.**
- **Transverse Enhancement Model (TEM).**
- **Many-body Currents:**
  - **Short-range correlations (SRC),**
  - **Meson Exchange Currents (MEC, 2p2h),**  $\Leftarrow$  (most popular and successful)
  - **Random Phase Approximation (RPA) and Quasiparticle RPA (QRPA).**
- **Hotchpotchs** (e.g., SuSA+MEC, MEC+RPA, etc.).



The nucleon axial mass vs. neutrino energy, obtained in several recent experiments. The dashed line and surrounding shaded double band represent the value of  $M_A$  and its  $1\sigma$  and  $2\sigma$  uncertainties extracted from the deuterium experiments ( $M_A^D$  hereafter).

This effect is now explained to an extent by the advanced nuclear models.

Besides, the much simpler RFG-based calculations can be somewhat fine-tuned by using a larger value of  $M_A$  in the lower energy range.<sup>a</sup>

This observation can be used for an **empirical** description of the experimental data within the RFG model. Namely, it is shown<sup>b</sup> that disadvantages of the RFG model can be **effectively compensated** by introducing the **energy-dependent (“running”) axial mass**

$$M_A^{\text{run}} = M_0 \left( 1 + \frac{E_0}{E_\nu} \right)$$

with the parameters values

$$M_0 = 1.006 \pm 0.025 (0.030) \text{ GeV} \quad \text{and} \quad E_0 = 0.334_{-0.054}^{+0.058} (0.070)_{(0.064)} \text{ GeV}$$

obtained from a global fit to all available data on neutrino-nucleus and antineutrino-nucleus CCQE scattering with  $\chi^2/\text{NDF} = 300.4/(435-19) = 0.72$ .

This allows to phenomenologically account for the intricate nuclear effects beyond RFG, significant at low and medium neutrino energies.

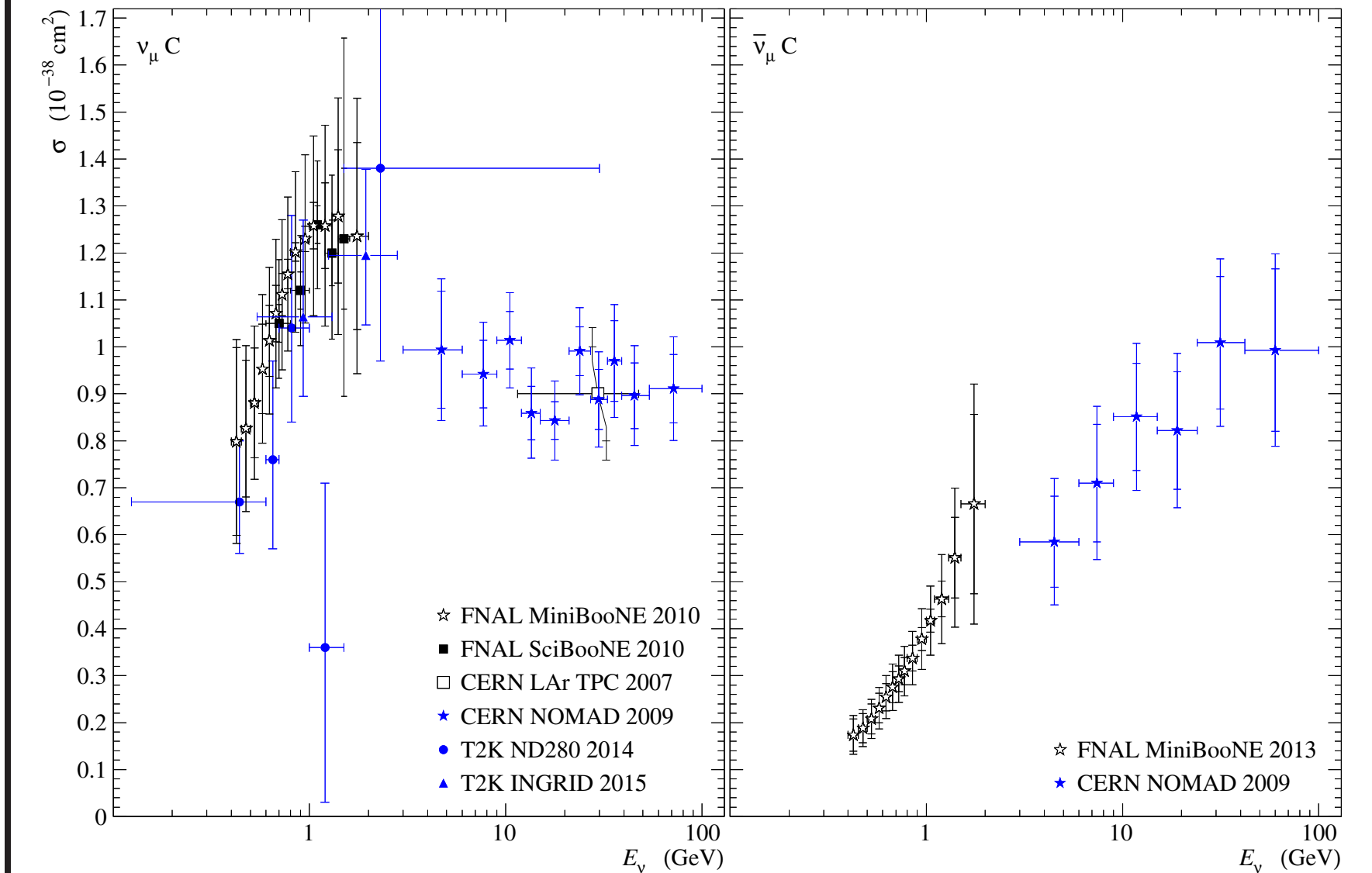
<sup>a</sup>A lot of papers; see, e.g., A.M. Ankowski, O. Benhar, C. Mariani and E. Vagnoni, Phys. Rev. D **93**, 113004 (2016) [arXiv:1603.01072 [hep-ph]] and references therein.

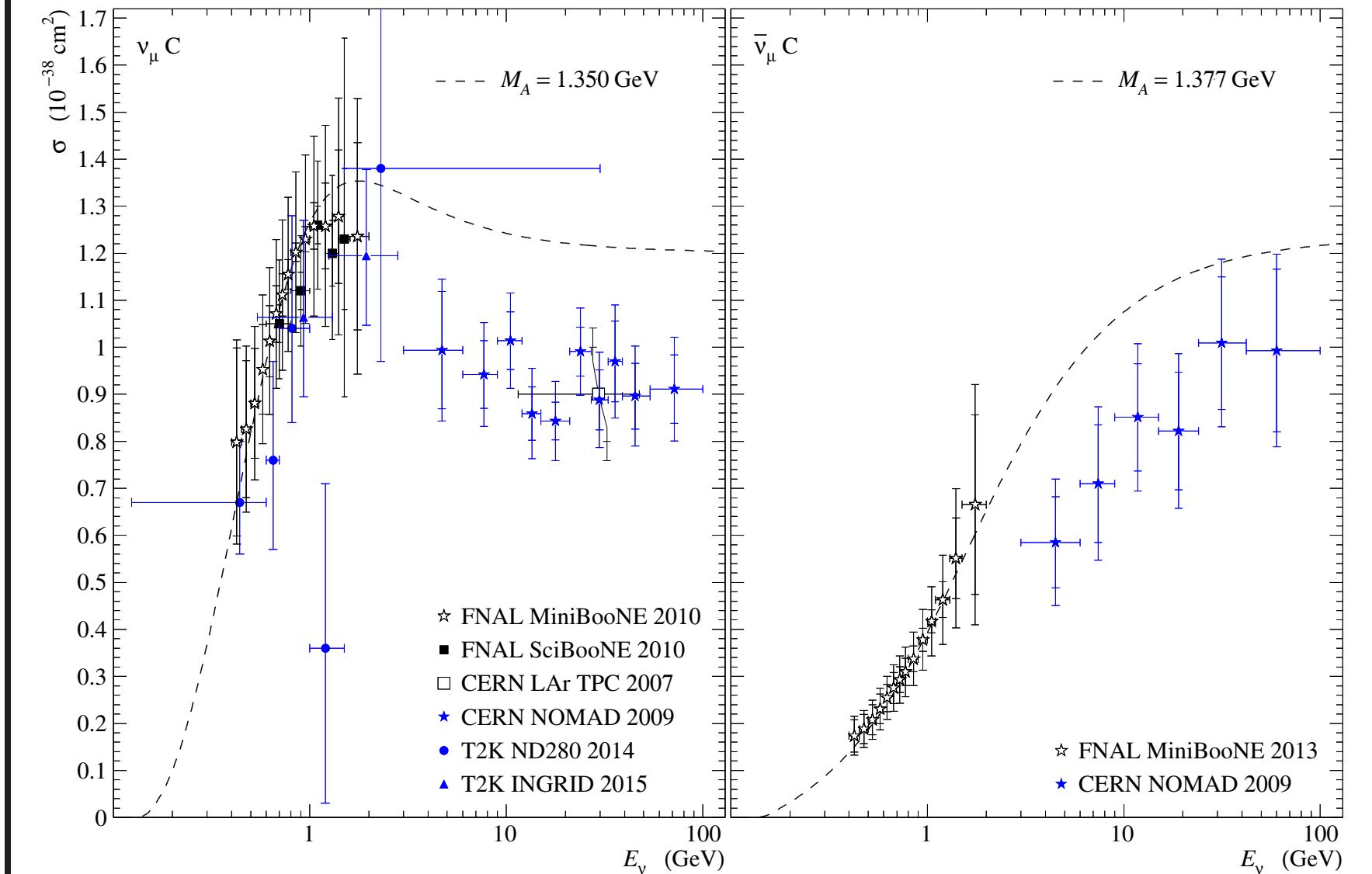
<sup>b</sup>K. S. Kuzmin and V. A. Naumov, in preparation. See also L. D. Kolupaeva, K. S. Kuzmin, O. N. Petrova and I. M. Shandrov, Mod. Phys. Lett. A **31**, 1650077 (2016) [arXiv:1603.07451 [hep-ph]]; K. S. Kuzmin, V. A. Naumov, O. N. Petrova, Accepted by Phys. Part. Nucl. (in press); K. S. Kuzmin, V. A. Naumov, O. N. Petrova, submitted to Acta Phys. Polon. B.

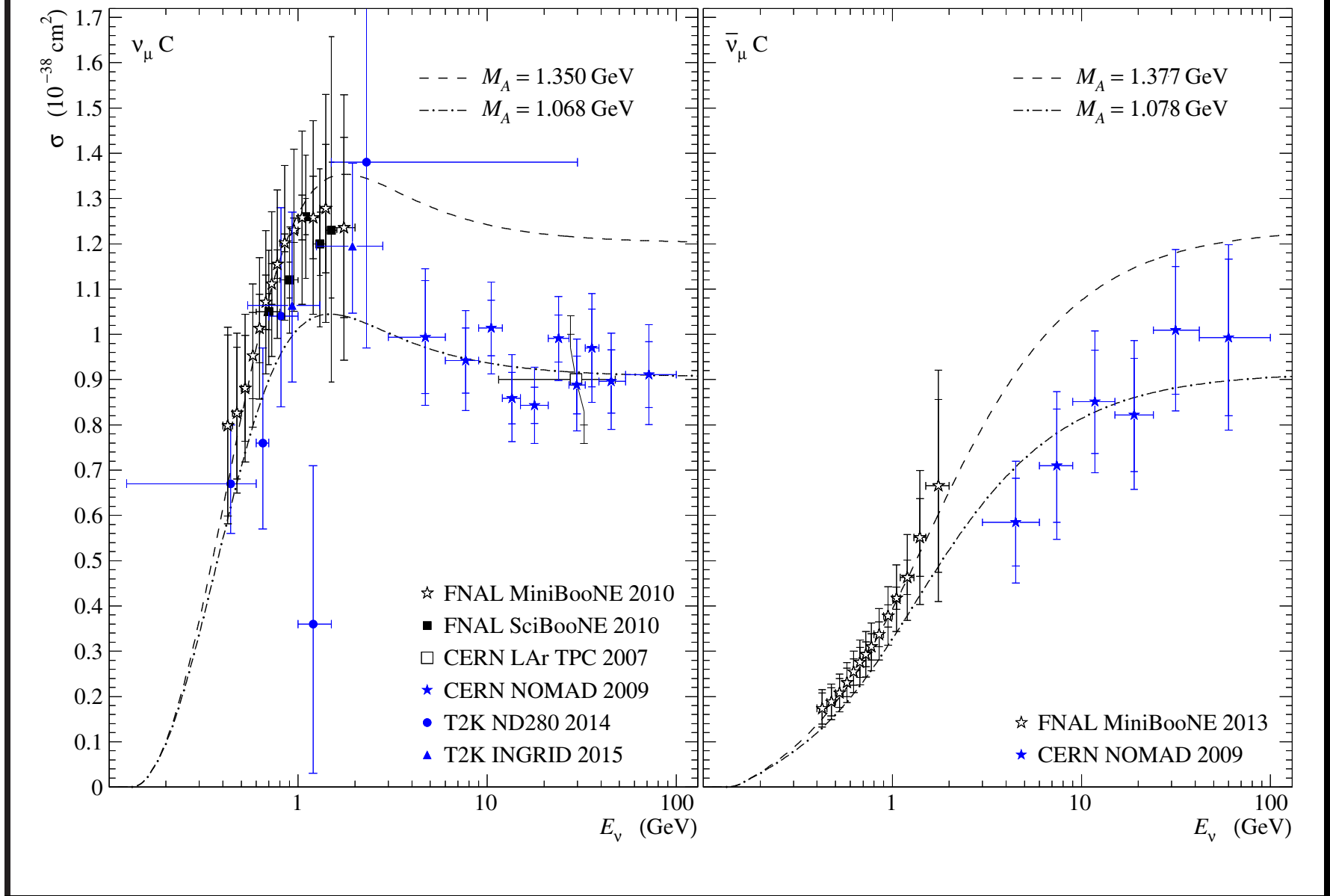


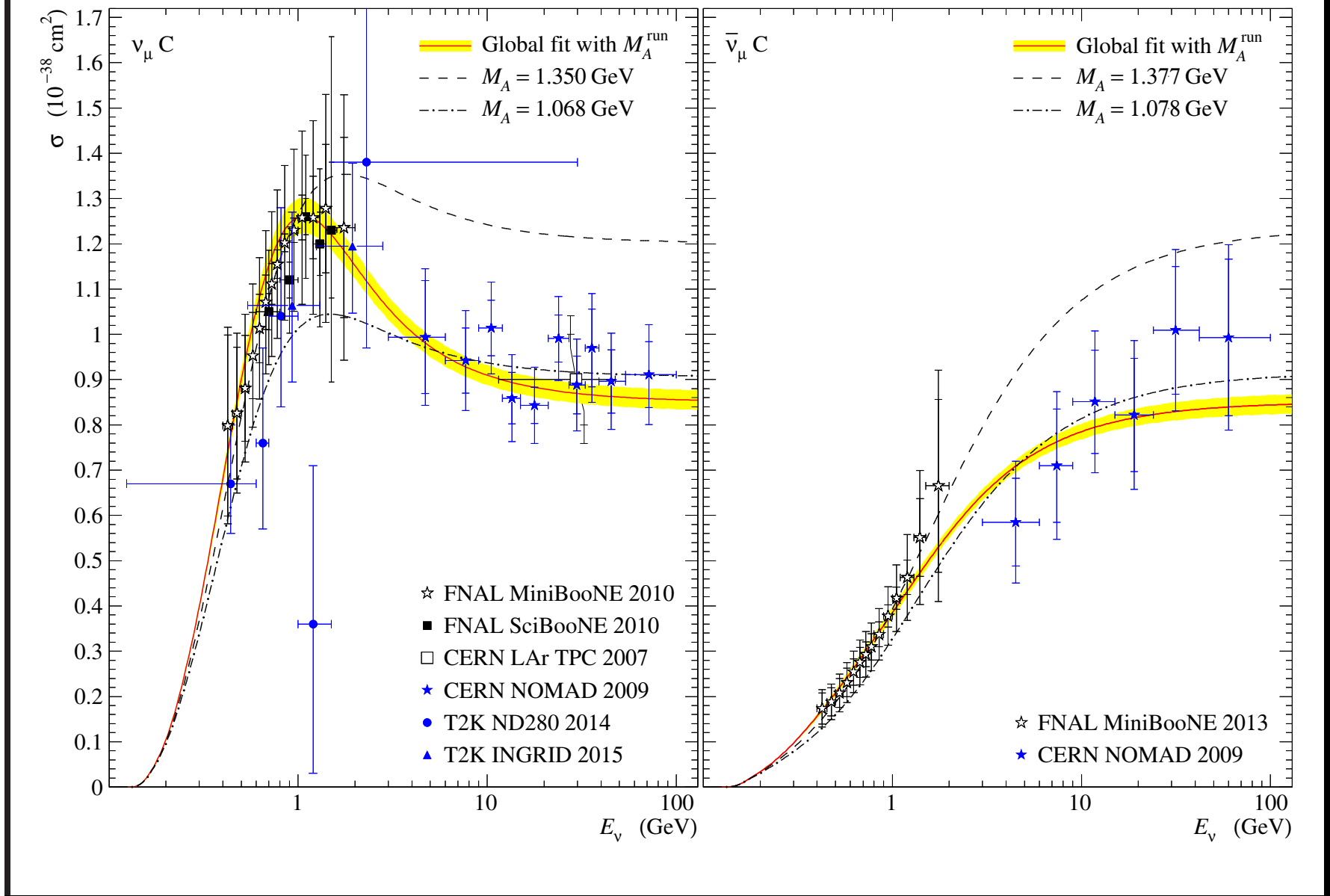
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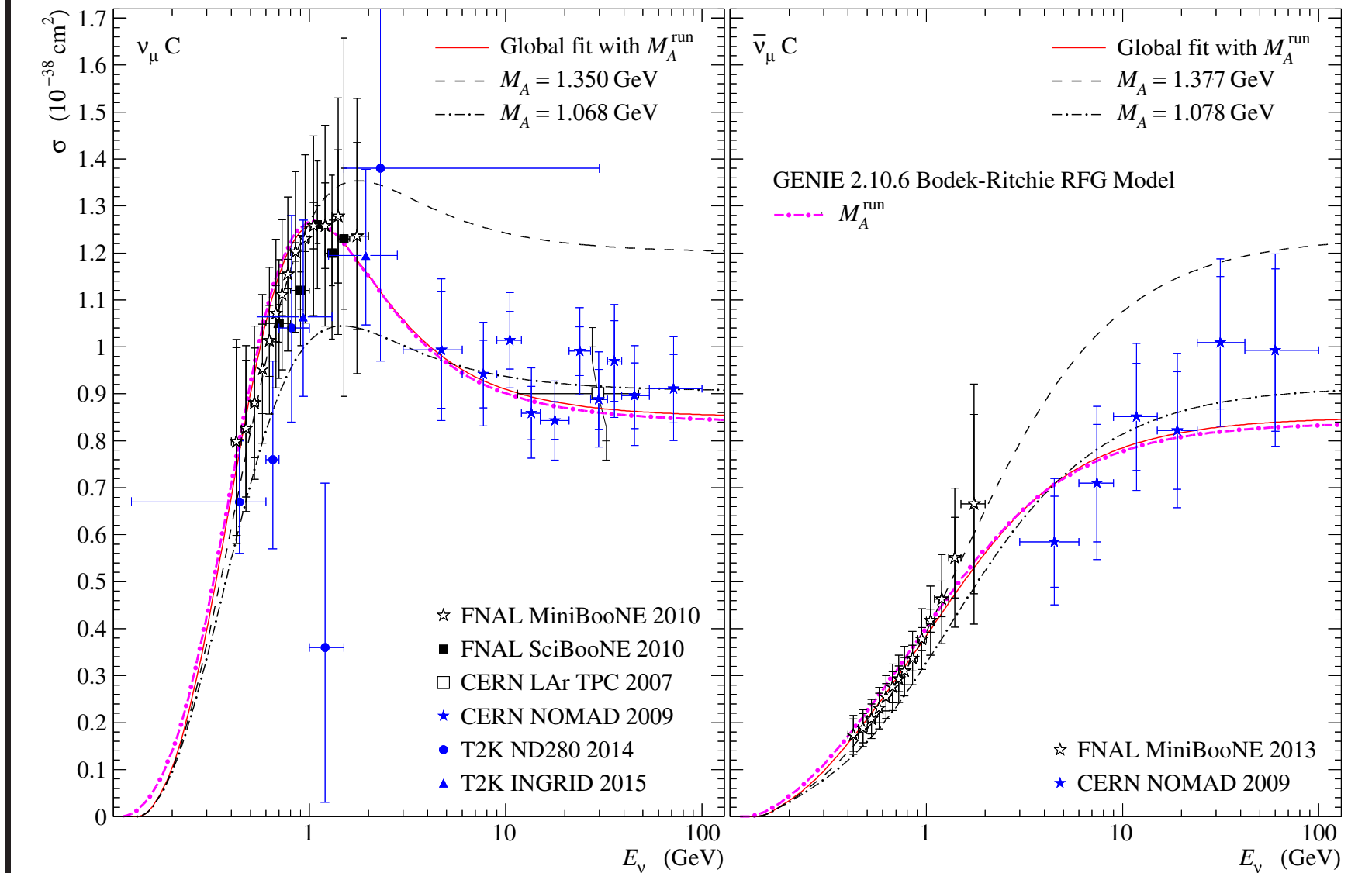
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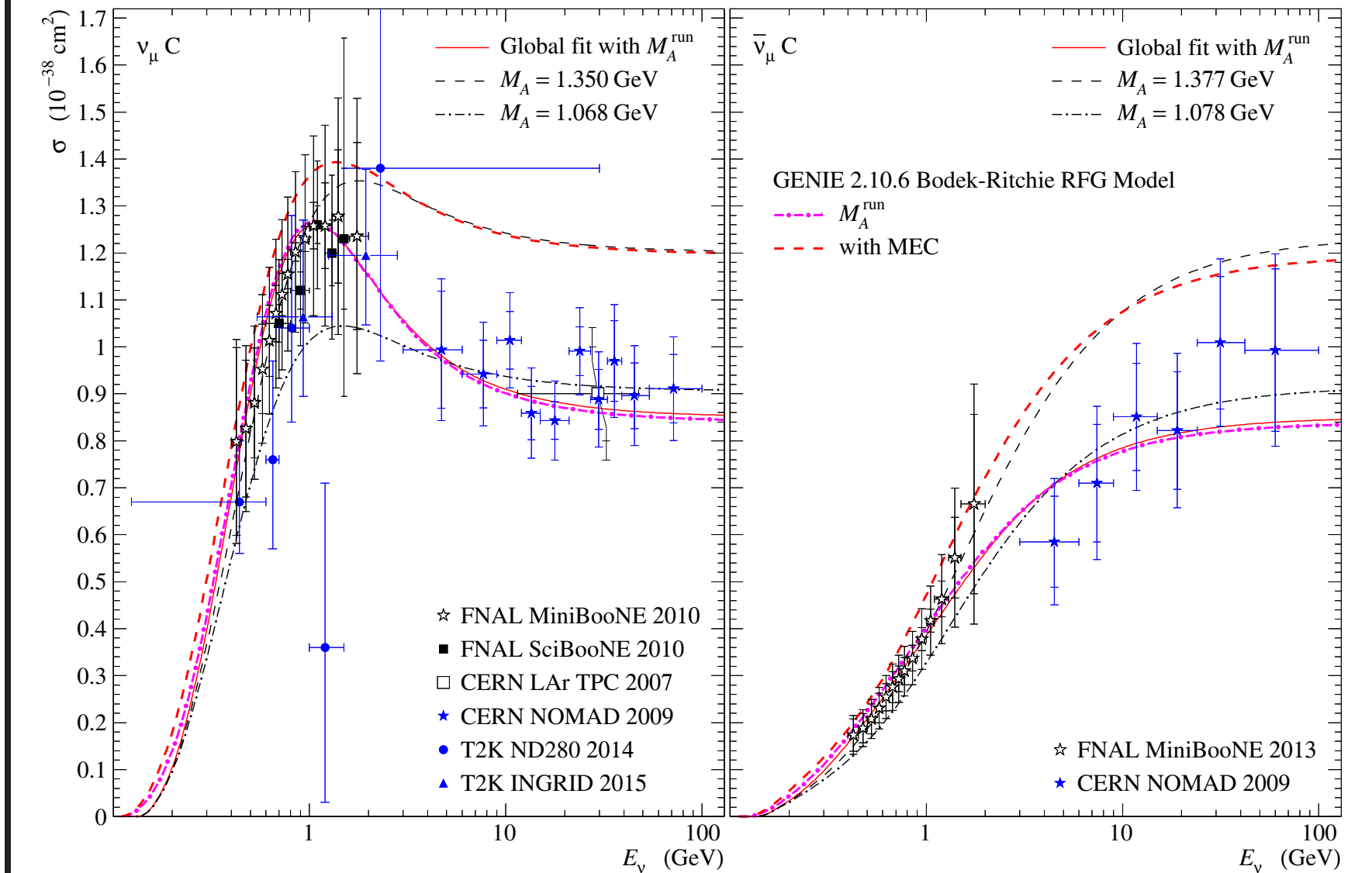


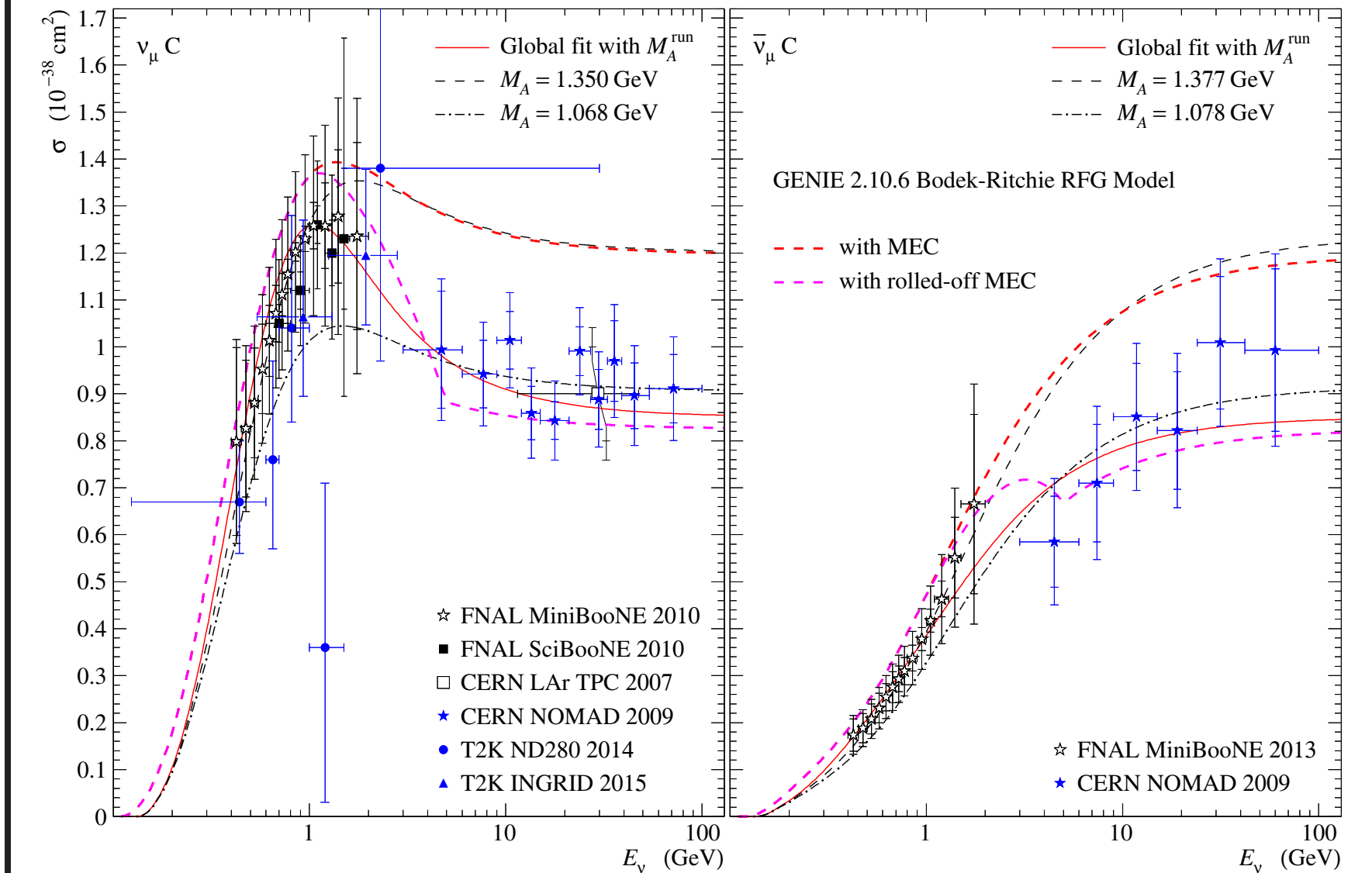






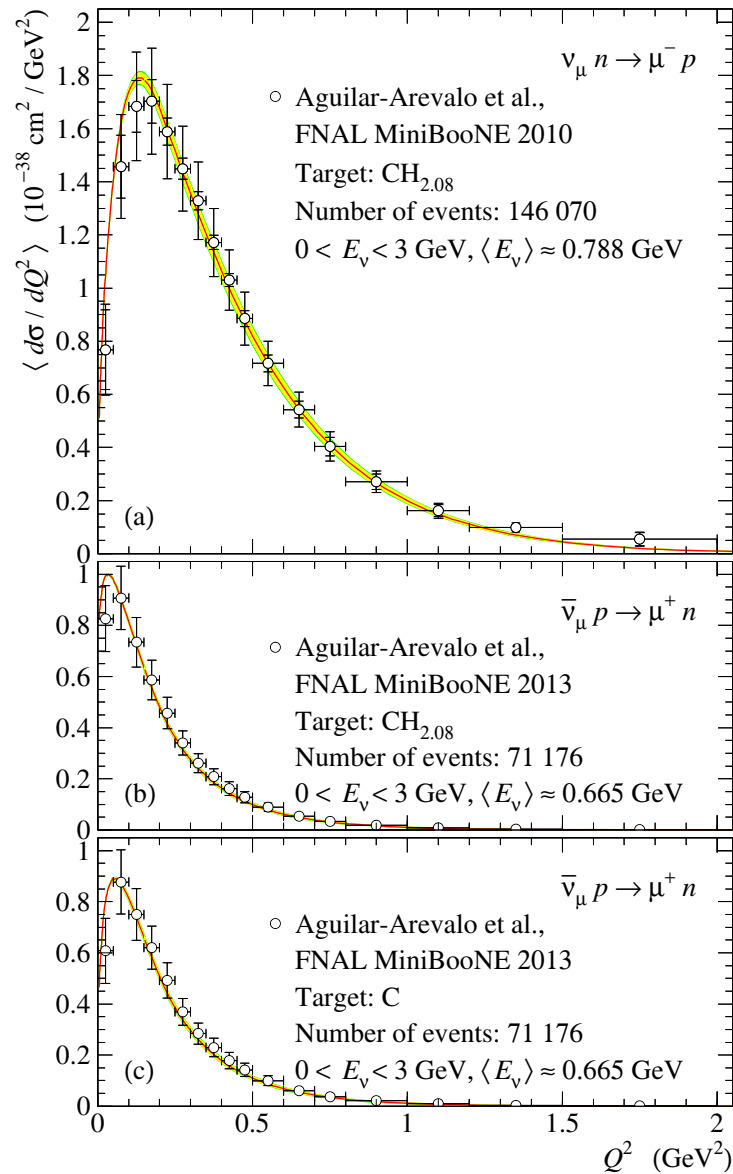




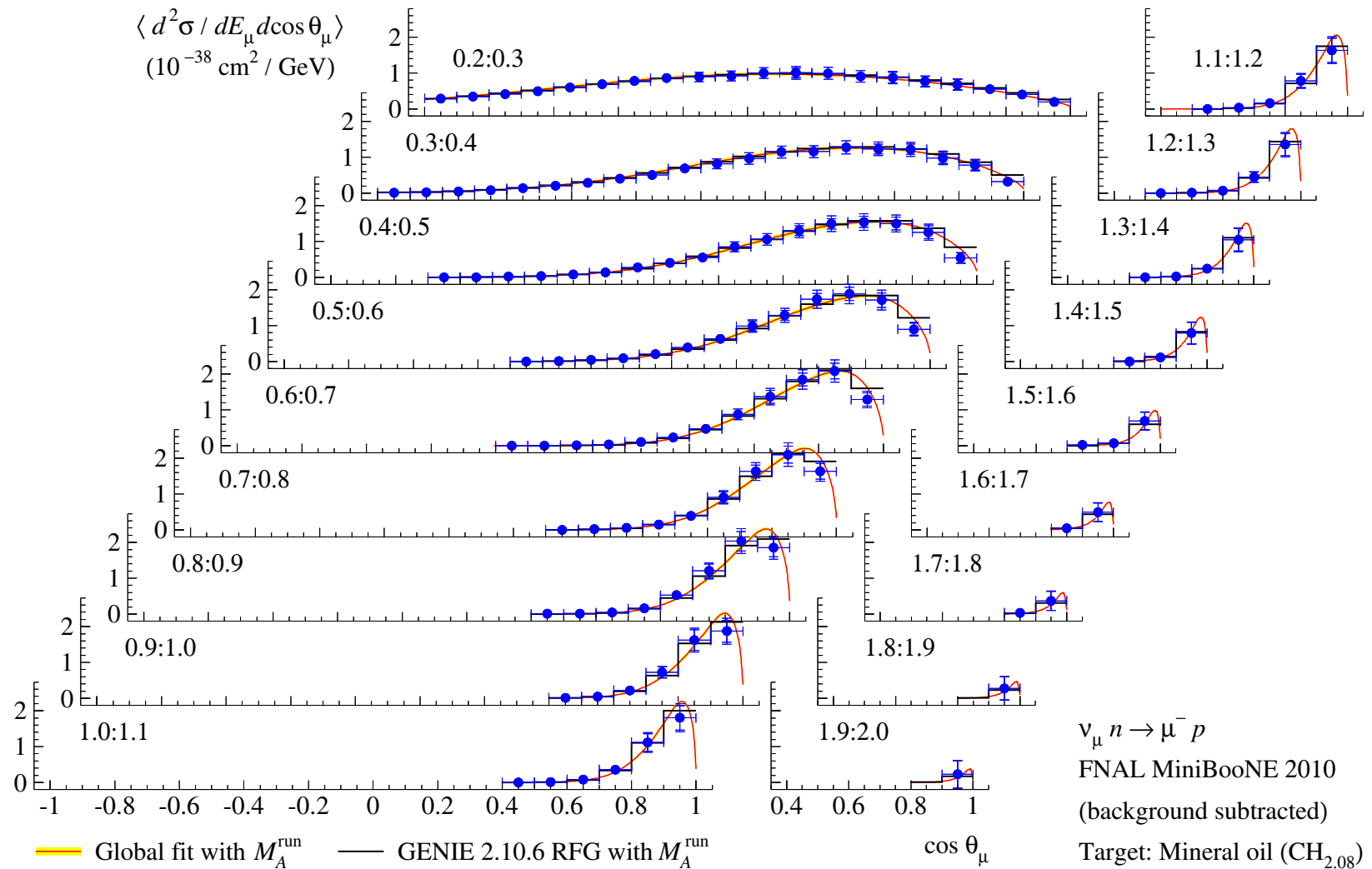


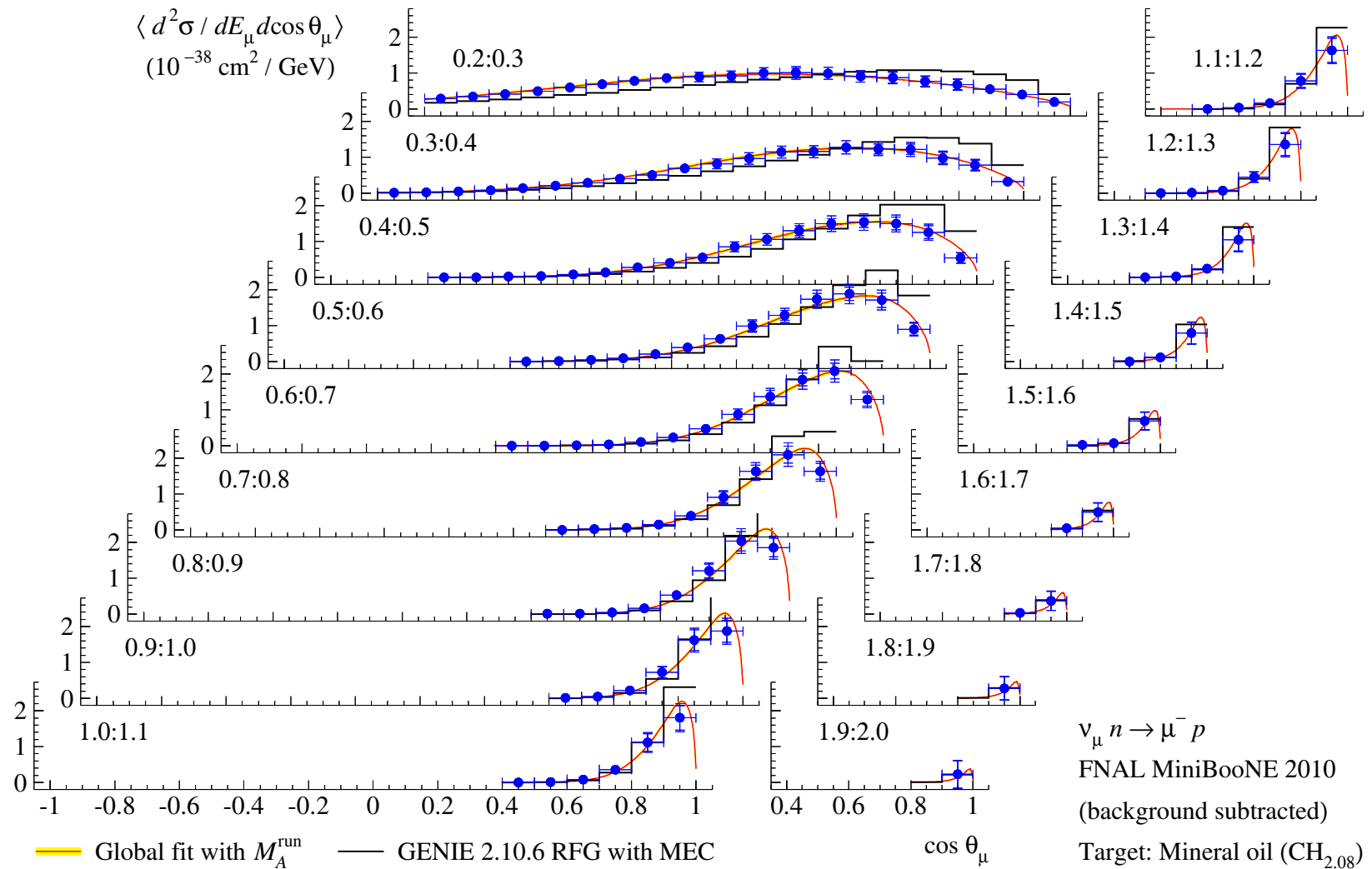


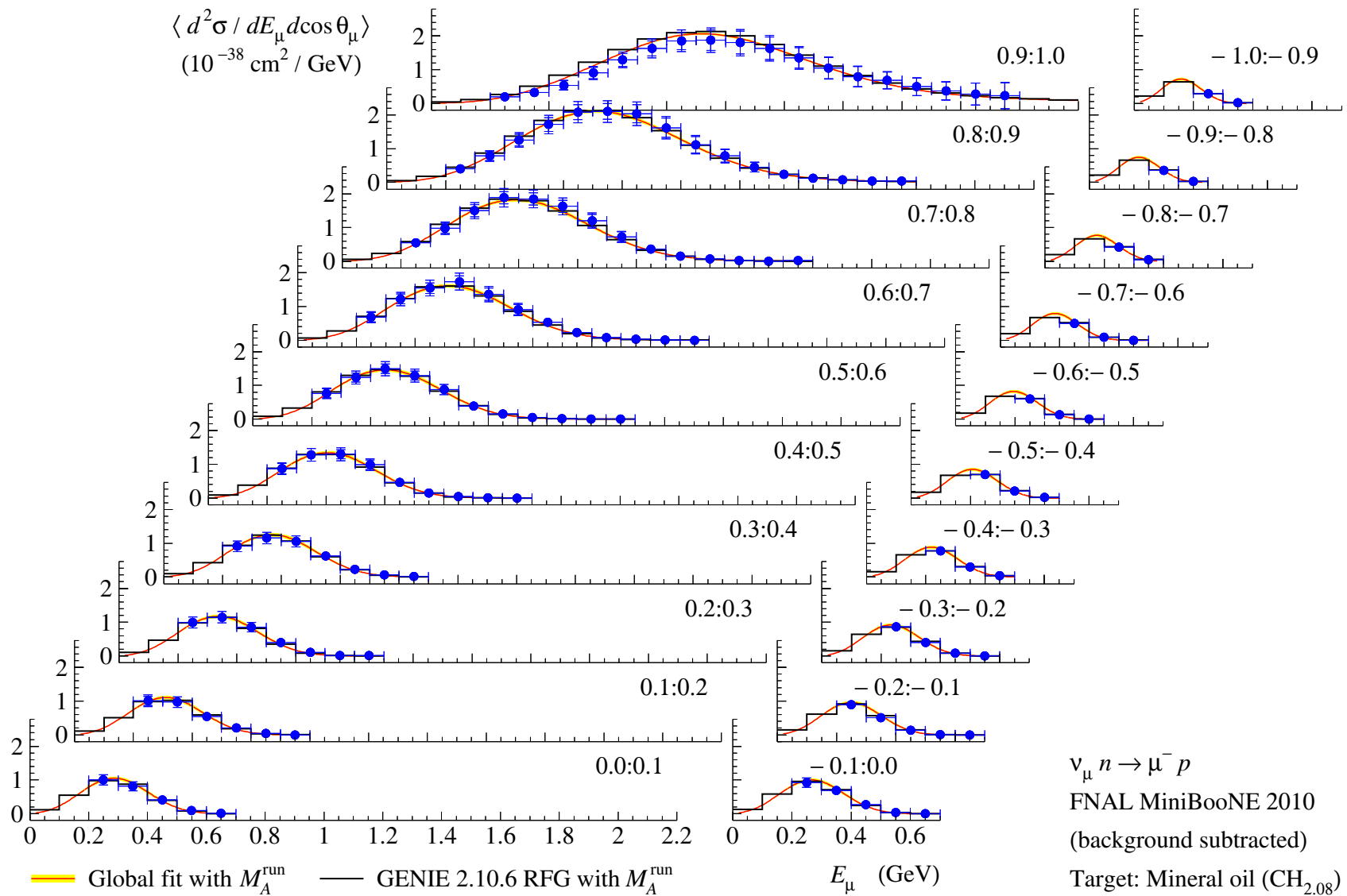
# **FNAL MiniBooNE 2010 – 2013**

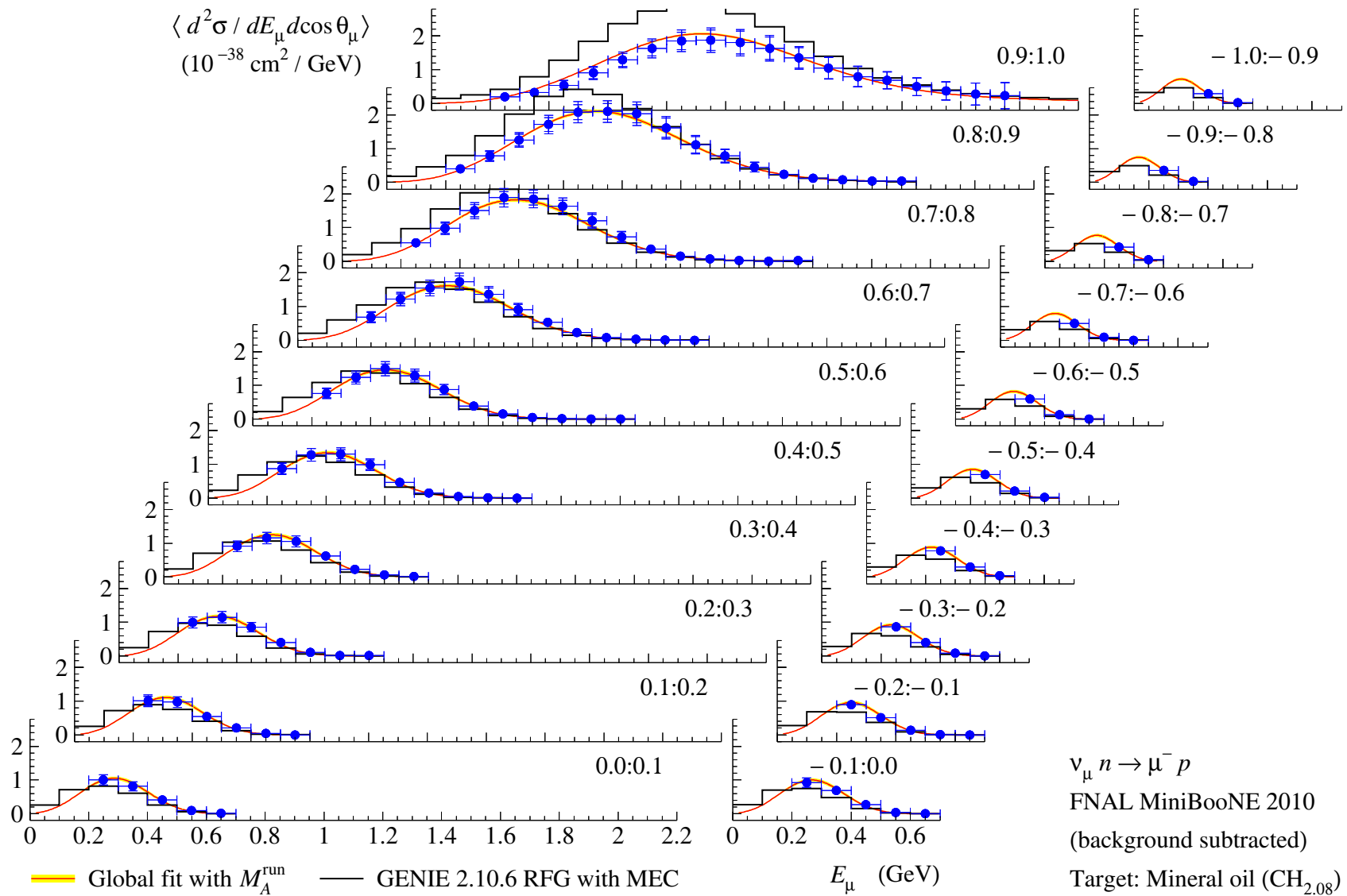


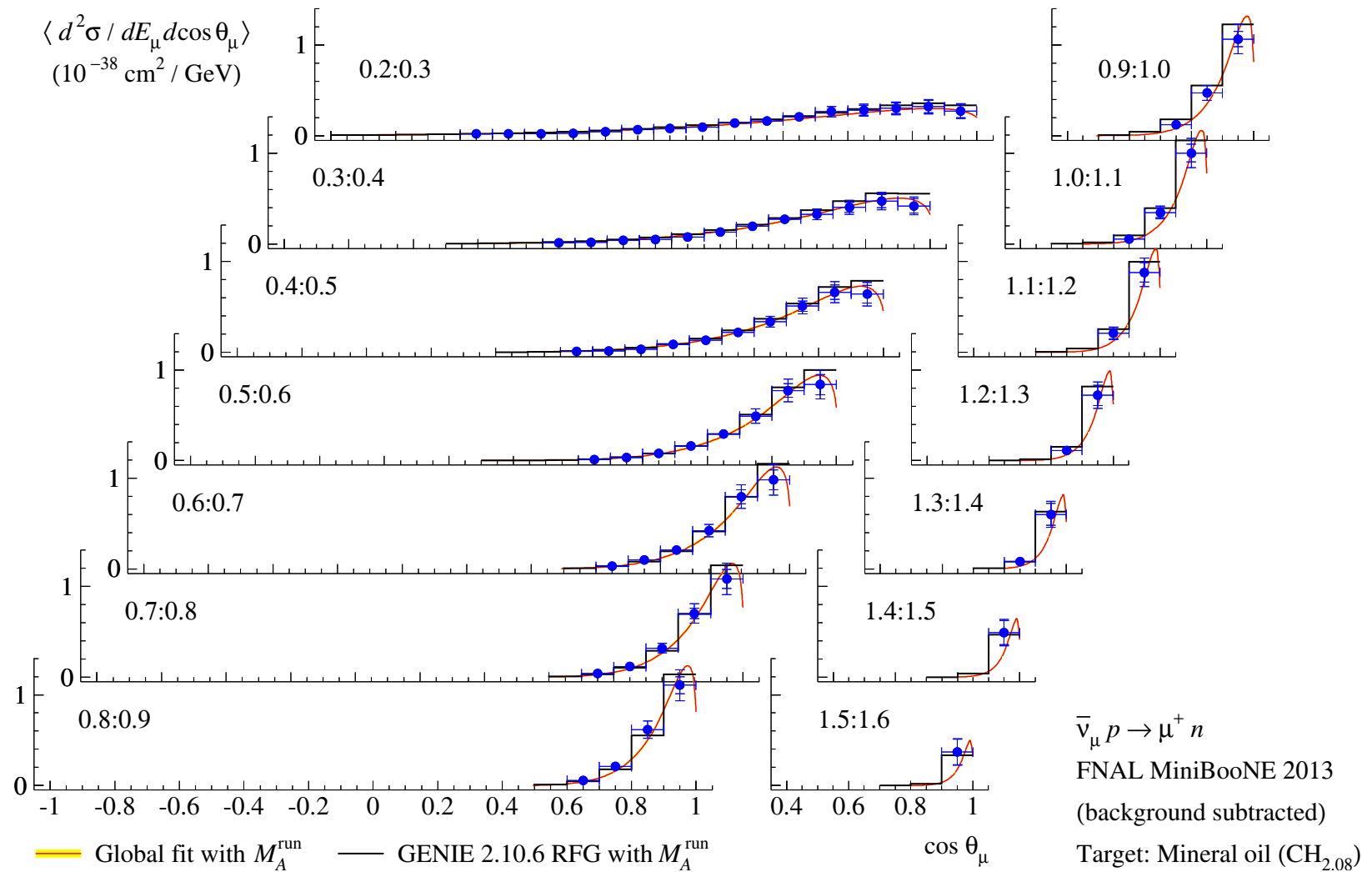
← Recalculated by the MiniBooNE Co. from the “raw” data on the CH<sub>2.08</sub> target. The agreement is a little bit worse and we know why...

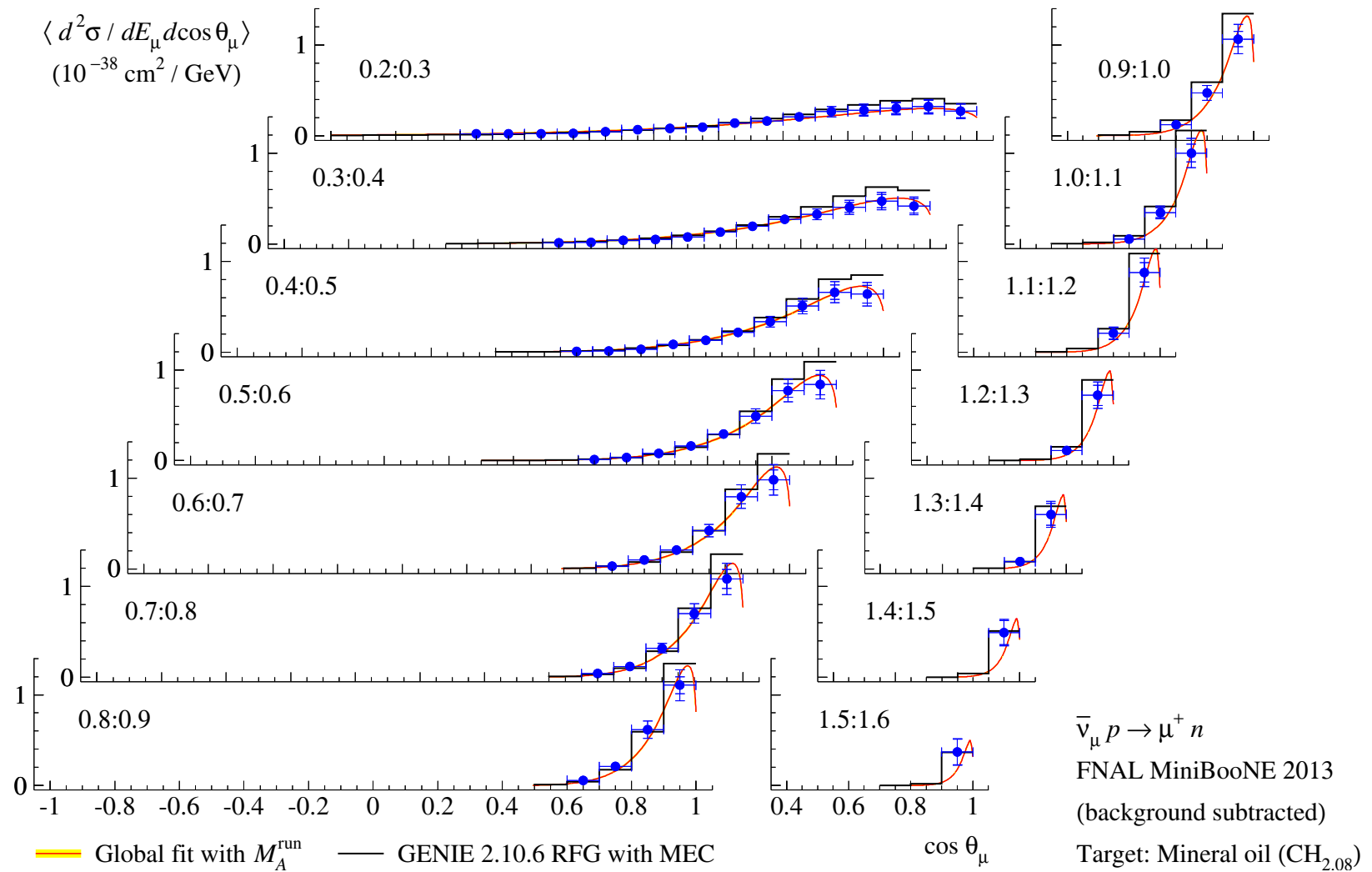




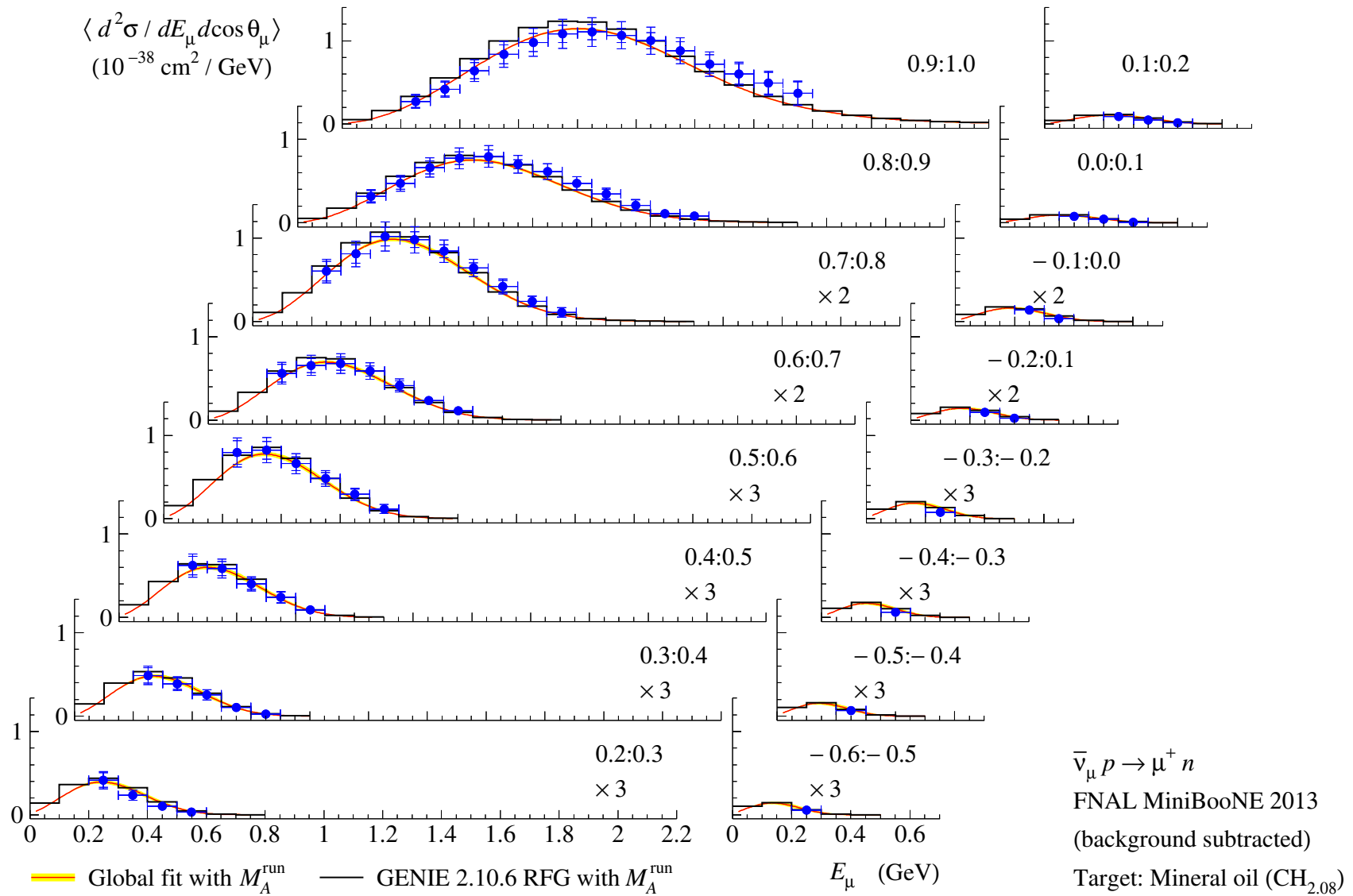


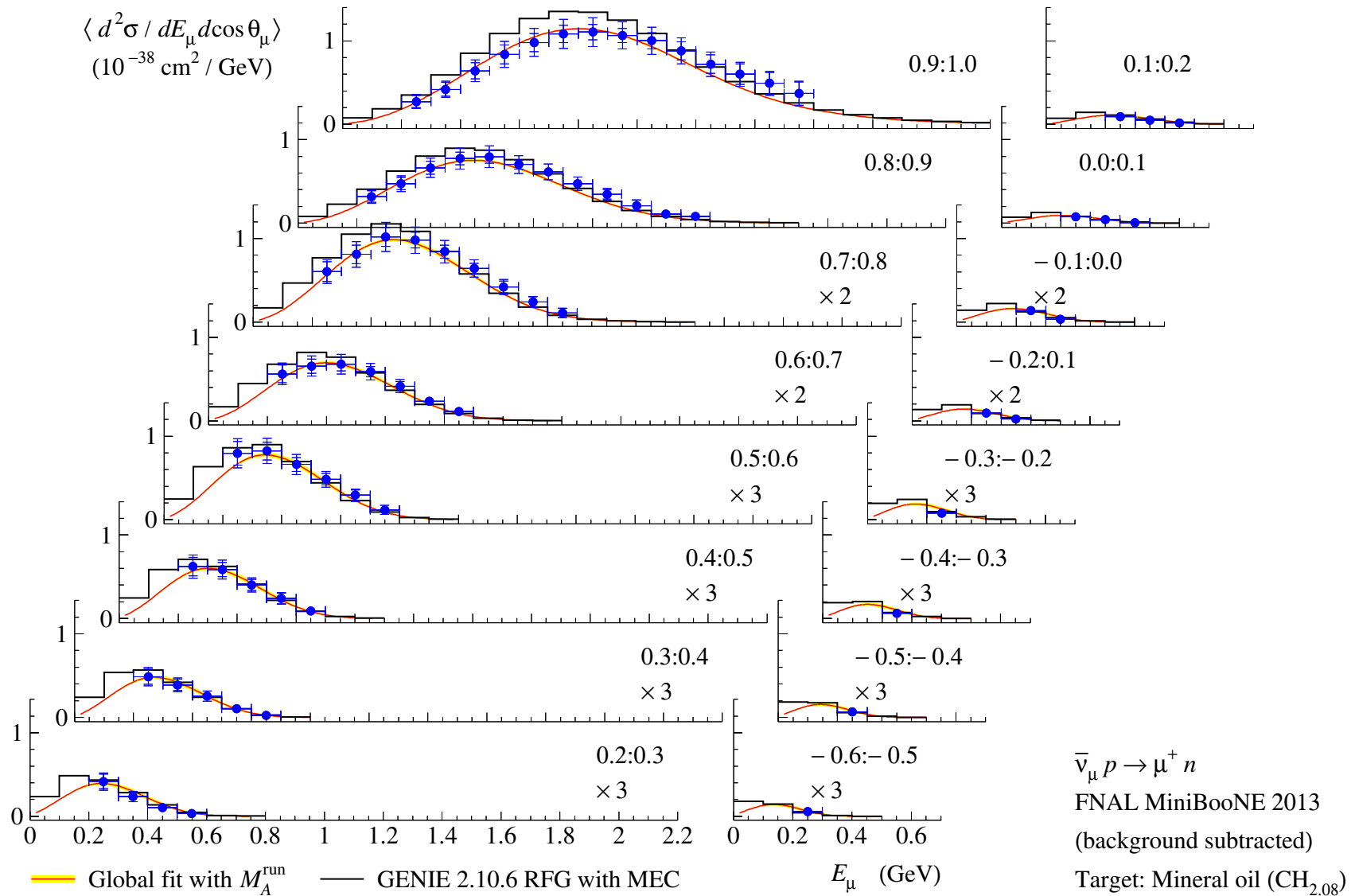




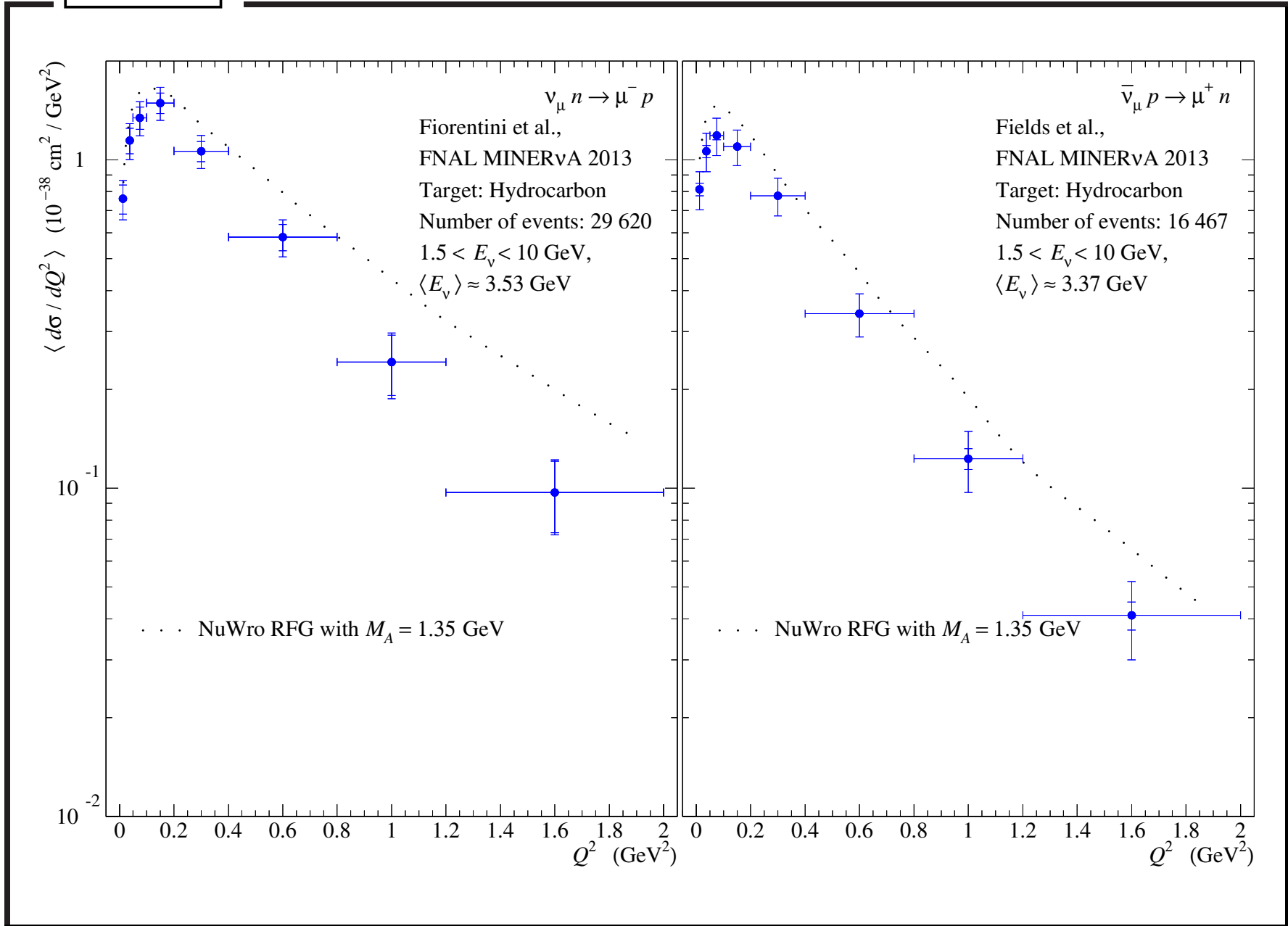


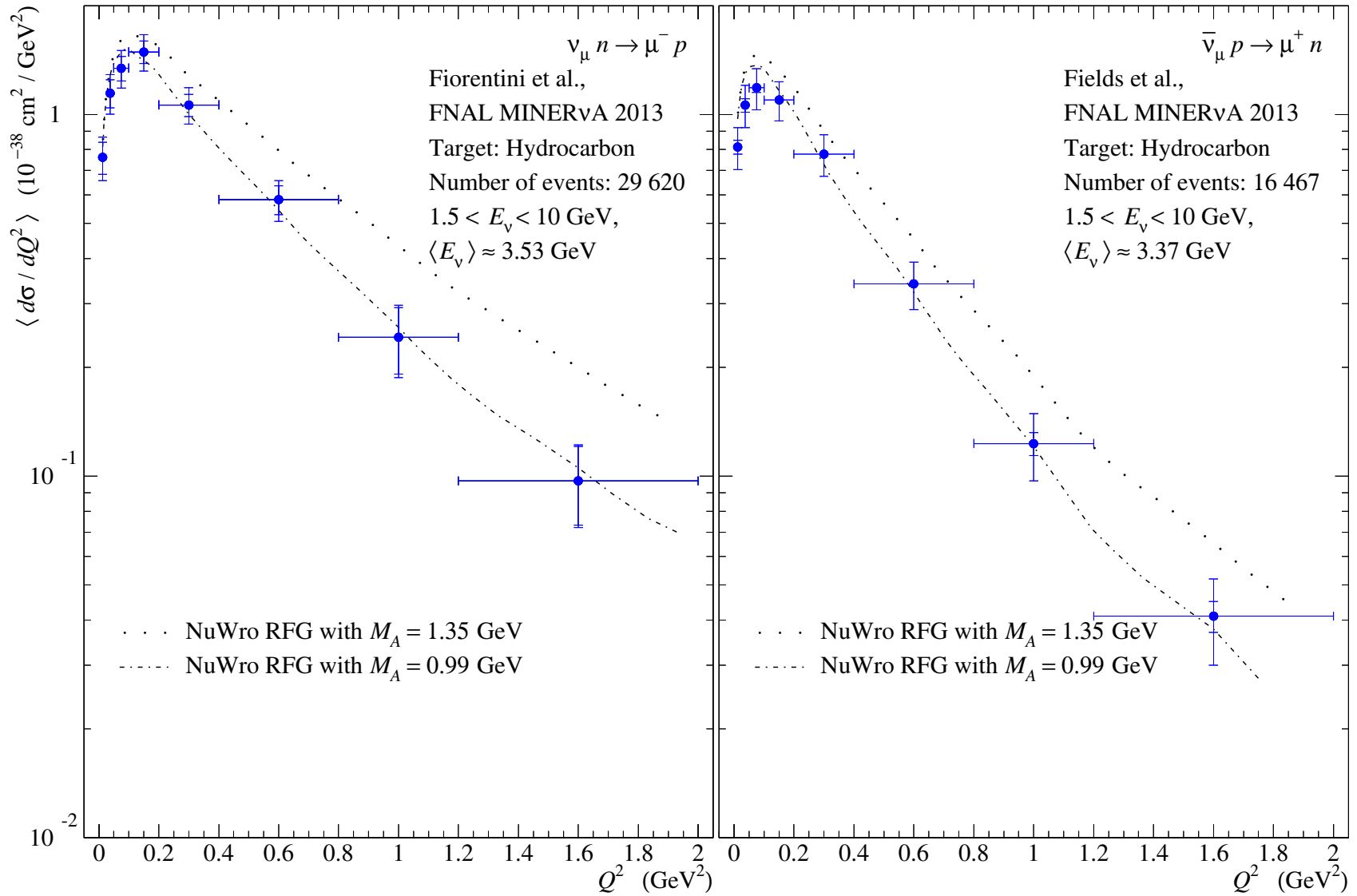


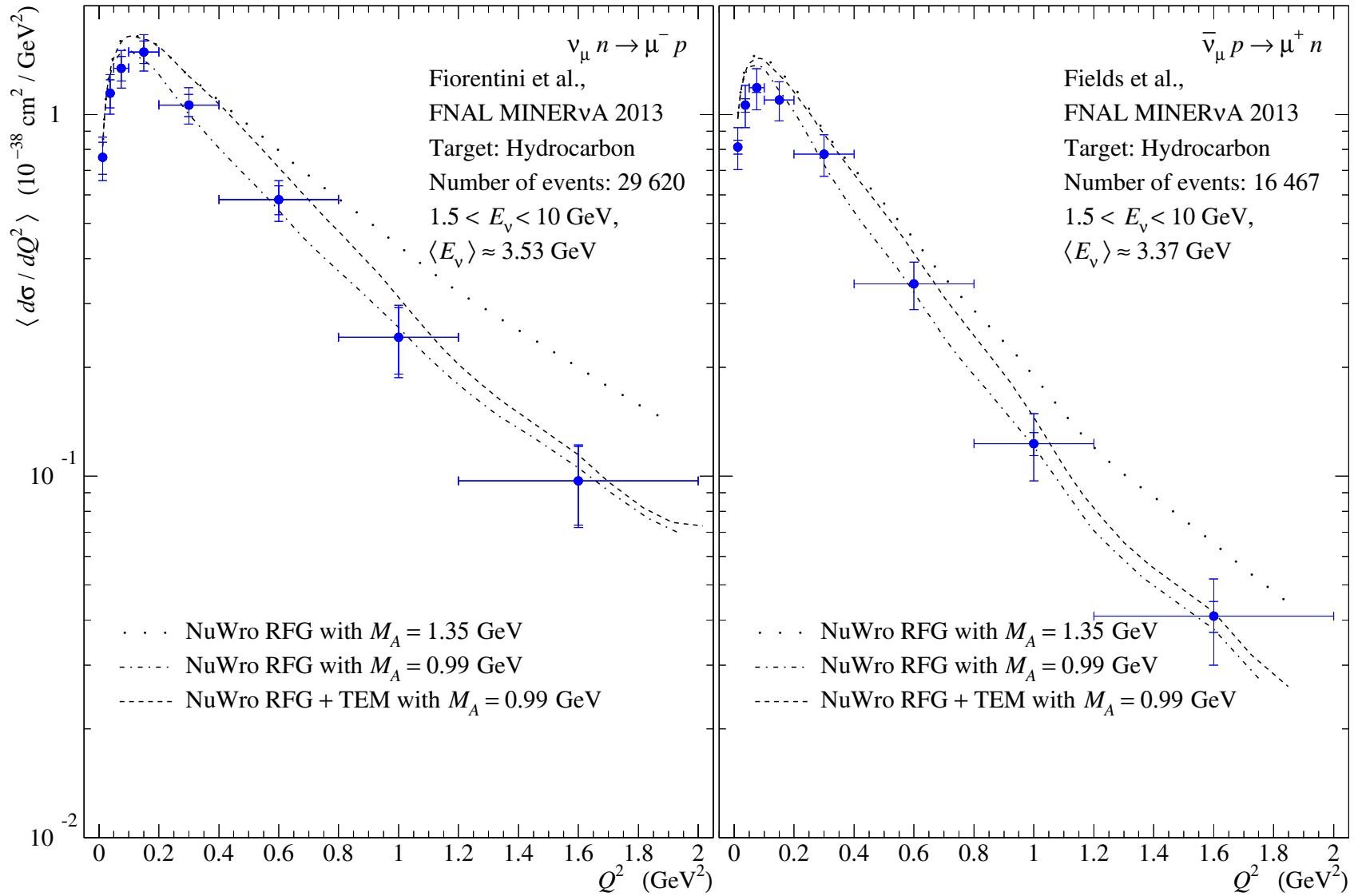


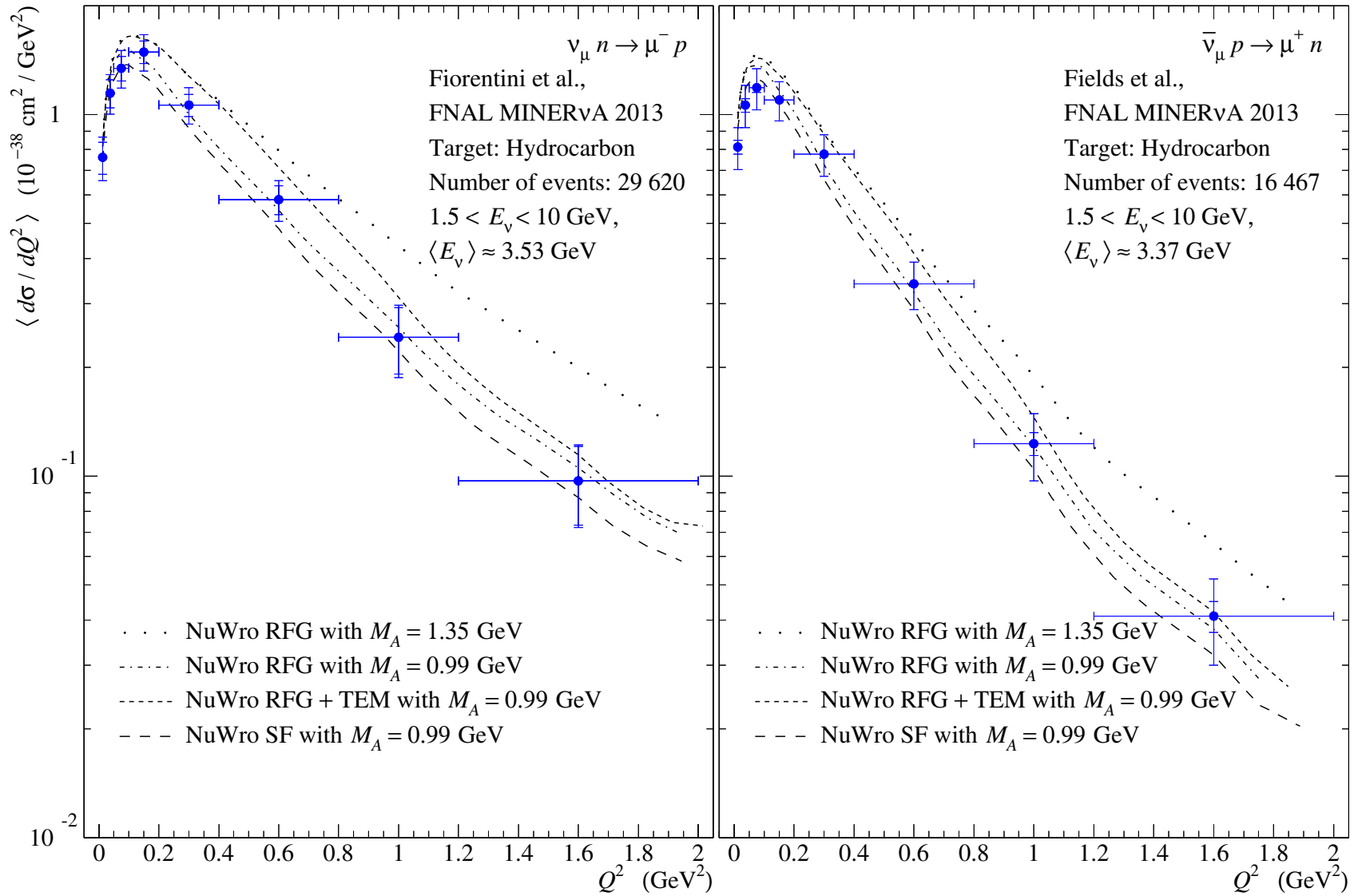


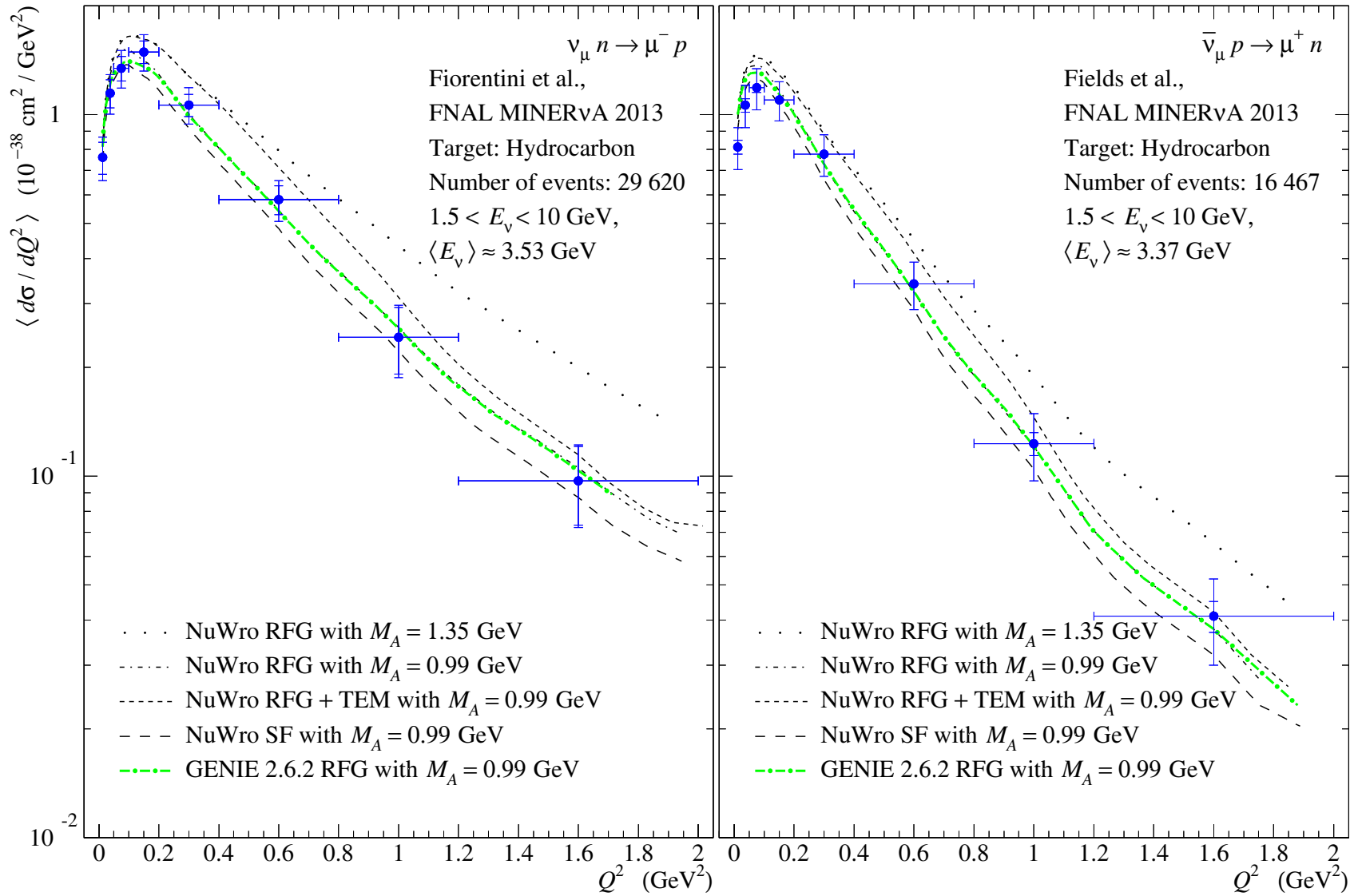
# FNAL MINER $\nu$ A 2013



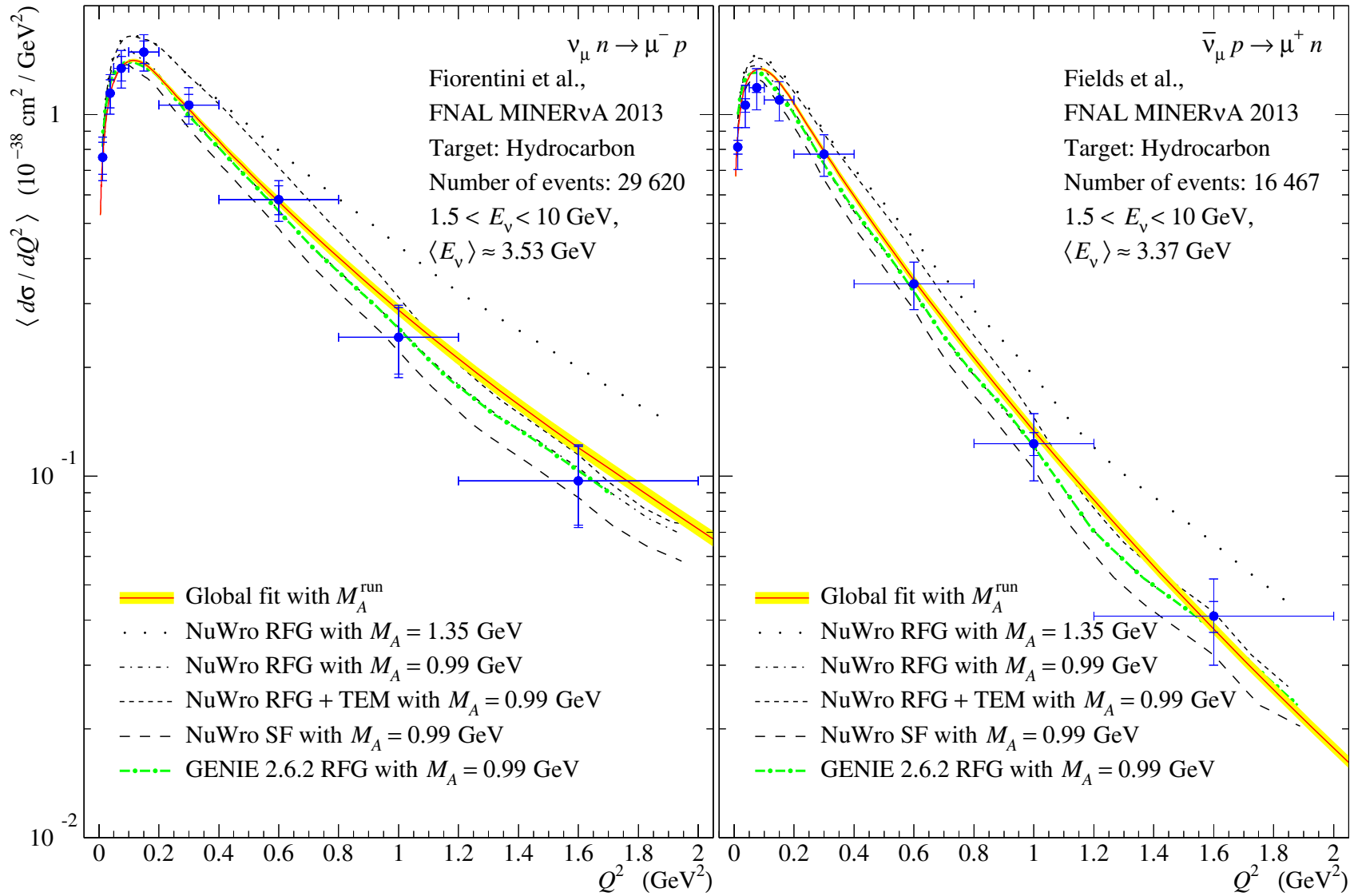


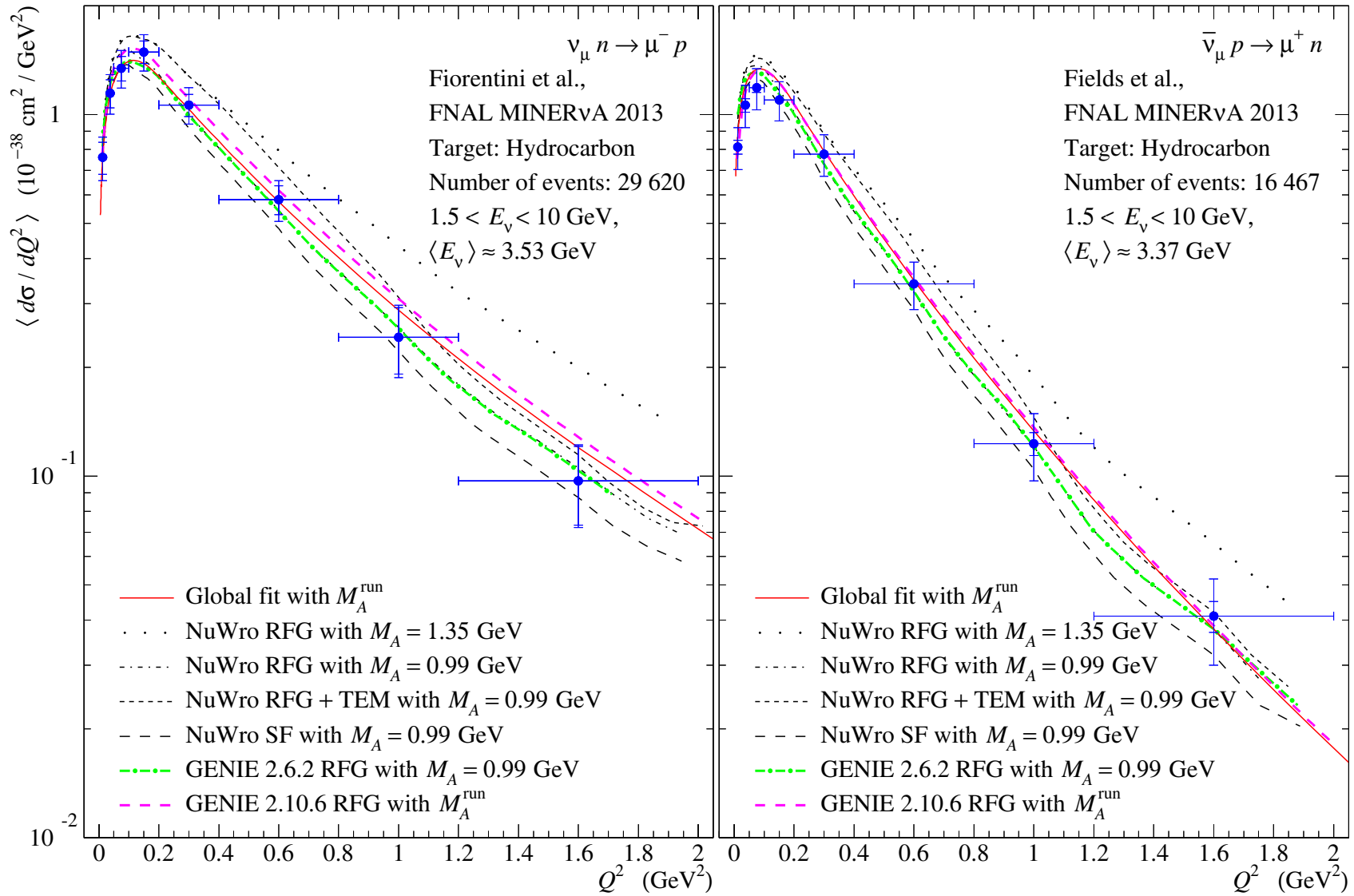






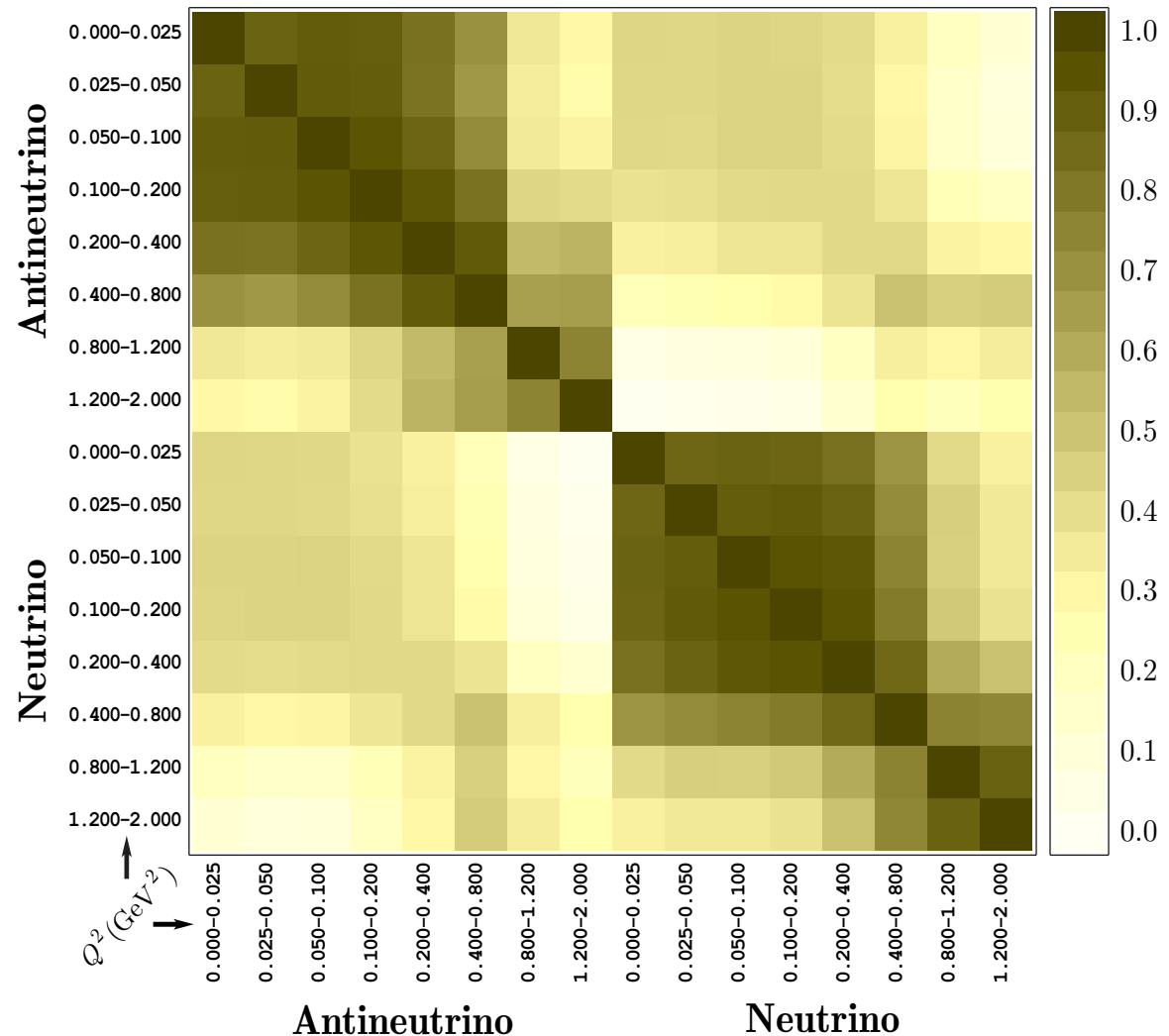






As it is seen from the correlation matrix, the MINER $\nu$ A data are highly correlated. Therefore the *by-eye* comparison between the data and the model predictions is not very conclusive.

The values of  $\chi^2/N$  (where  $N$  is the number of data-points) listed in next slide allow us to conclude that our model provides the best description of the MINER $\nu$ A results among the models shown in slides 28–34.



| Model  | Author's estimation<br>(by bin medians) | Our estimation<br>(by bin medians) | Our estimation<br>(by bin averages) |
|--|---|------------------------------------|-------------------------------------|
| $\nu_\mu$ FNAL MINER $\nu$ A 2013 [27], $N = 8$                          |   |                                    |                                     |
| NuWro RFG with $M_A = 1.35$  | 3.7                                     | 29.26/8 = 3.66                     | 31.13/8 = 3.89                      |
| NuWro RFG with $M_A = 0.99 + \text{TEM}$                                 | 2.4                                     | 18.40/8 = 2.30                     | 20.93/8 = 2.62                      |
| NuWro RFG with $M_A = 0.99$  | 3.5                                     | 28.19/8 = 3.52                     | 23.24/8 = 2.91                      |
| NuWro SF with $M_A = 0.99$   | 2.8                                     | 22.66/8 = 2.83                     | 18.69/8 = 2.34                      |
| GENIE 2.6.2 RFG with $M_A = 0.99$  | —                                       | 16.04/8 = 2.00                     | 15.27/8 = 1.91                      |
| Global fit with $M_A^{\text{run}}$                                       | —                                       | 7.19/8 = 0.90                      | 8.04/8 = 1.00                       |
| GENIE 2.10.6 with $M_A^{\text{run}}$                                     | —                                       | 11.16/8 = 1.39                     | 10.77/8 = 1.35                      |
| $\bar{\nu}_\mu$ FNAL MINER $\nu$ A 2013 [28], $N = 8$                    |   |                                    |                                     |
| NuWro RFG with $M_A = 1.35$  | 2.90                                    | 23.50/8 = 2.94                     | 23.20/8 = 2.90                      |
| NuWro RFG with $M_A = 0.99 + \text{TEM}$                                 | 1.06                                    | 8.35/8 = 1.04                      | 9.50/8 = 1.19                       |
| NuWro RFG with $M_A = 0.99$  | 2.64                                    | 21.39/8 = 2.67                     | 17.80/8 = 2.23                      |
| NuWro SF with $M_A = 0.99$   | 2.14                                    | 17.37/8 = 2.17                     | 14.59/8 = 1.82                      |
| GENIE 2.6.2 RFG with $M_A = 0.99$  | —                                       | 19.96/8 = 2.50                     | 18.23/8 = 2.28                      |
| Global fit with $M_A^{\text{run}}$                                       | —                                       | 7.29/8 = 0.91                      | 5.51/8 = 0.69                       |
| GENIE 2.10.6 with $M_A^{\text{run}}$                                     | —                                       | 19.31/8 = 2.41                     | 15.92/8 = 1.99                      |
| $\nu_\mu \oplus \bar{\nu}_\mu$ FNAL MINER $\nu$ A 2013 [27,28], $N = 16$ |   |                                    |                                     |
| NuWro RFG with $M_A = 1.35$  | —                                       | 59.15/16 = 3.70                    | 59.34/16 = 3.71                     |
| NuWro RFG with $M_A = 0.99 + \text{TEM}$                                 | —                                       | 28.33/16 = 1.77                    | 30.74/16 = 1.92                     |
| NuWro RFG with $M_A = 0.99$  | —                                       | 48.11/16 = 3.01                    | 40.59/16 = 2.54                     |
| NuWro SF with $M_A = 0.99$   | —                                       | 38.12/16 = 2.38                    | 32.39/16 = 2.02                     |
| GENIE 2.6.2 RFG with $M_A = 0.99$  | —                                       | 34.87/16 = 2.18                    | 33.04/16 = 2.07                     |
| Global fit with $M_A^{\text{run}}$                                       | —                                       | 19.93/16 = 1.25                    | 17.90/16 = 1.10                     |
| GENIE 2.10.6 with $M_A^{\text{run}}$                                     | —                                       | 34.77/16 = 2.17                    | 30.02/16 = 1.88                     |

# **T2K ND280 2016**

## A caveat

We have not yet included the new (2016) T2K ND280 data [69] into the global fit due to the following reasons:

- We don't know the actual  $\nu_\mu$  spectrum since in several runs of the experiment, the three magnetic horns were excited by different current pulses.
- We don't know which of two data analyses (or both with certain weights) should be used in the fit.
- There is some small additional uncertainty related to complicated elemental composition of the detector target (fine-grained detector FGD1).<sup>a</sup> The measurement has been effectively considered as if on the FGD1 scintillator ( $C_8H_8$ ). But our experience suggests that such simplification is not quite harmless...

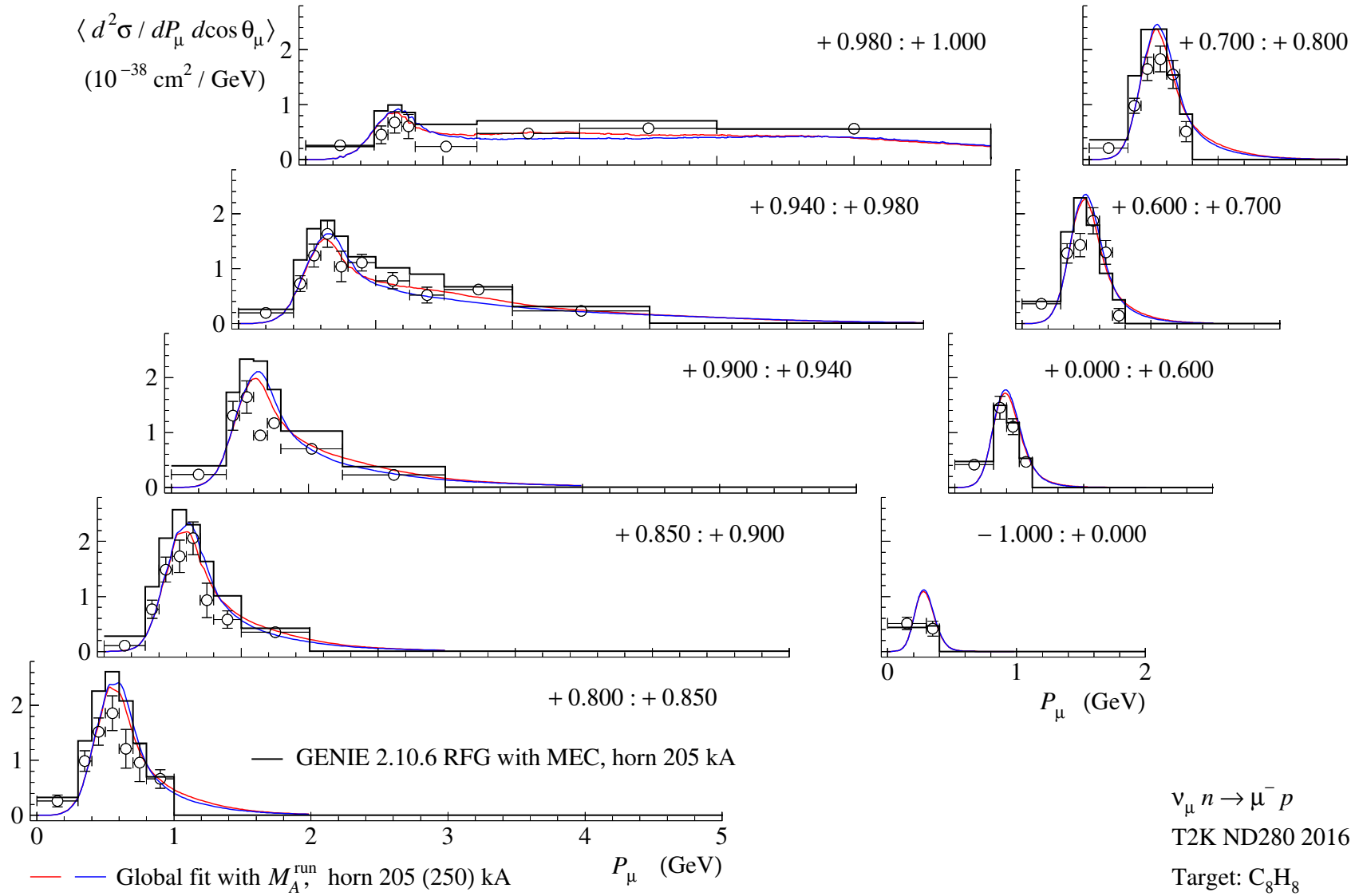
Here we only compare our calculations with the data from [Analyses 1 & 2](#) [69] by using the two  $\nu_\mu$  spectra corresponding to the horn currents of 250 and 205 kA.

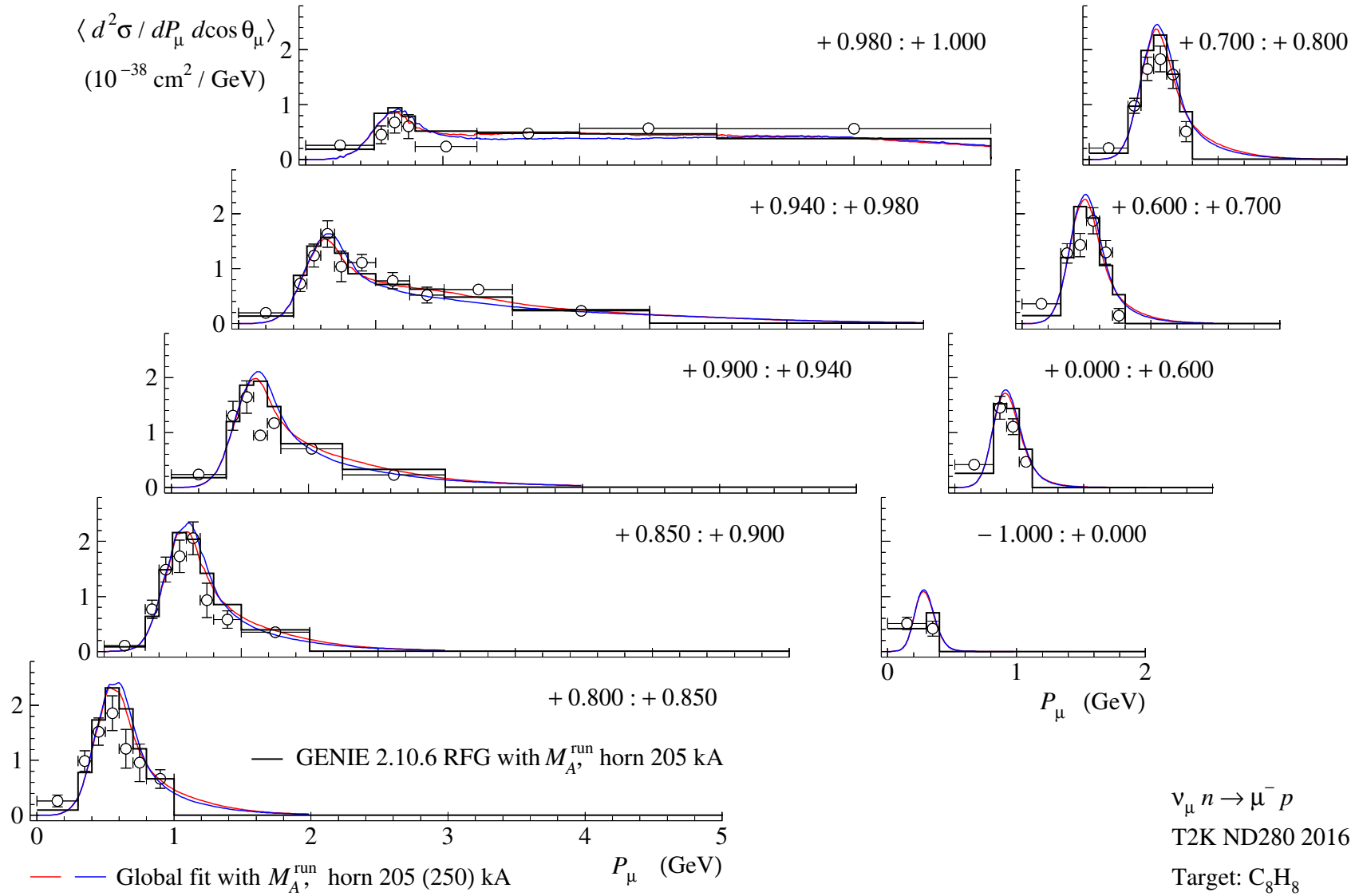
We find a reasonable (within the uncertainties) agreement between the ND280 results and those predicted by our model.

Therefore, it can be preliminary concluded that adding the T2K ND280 2016 data into the global fit would not drastically change the result.

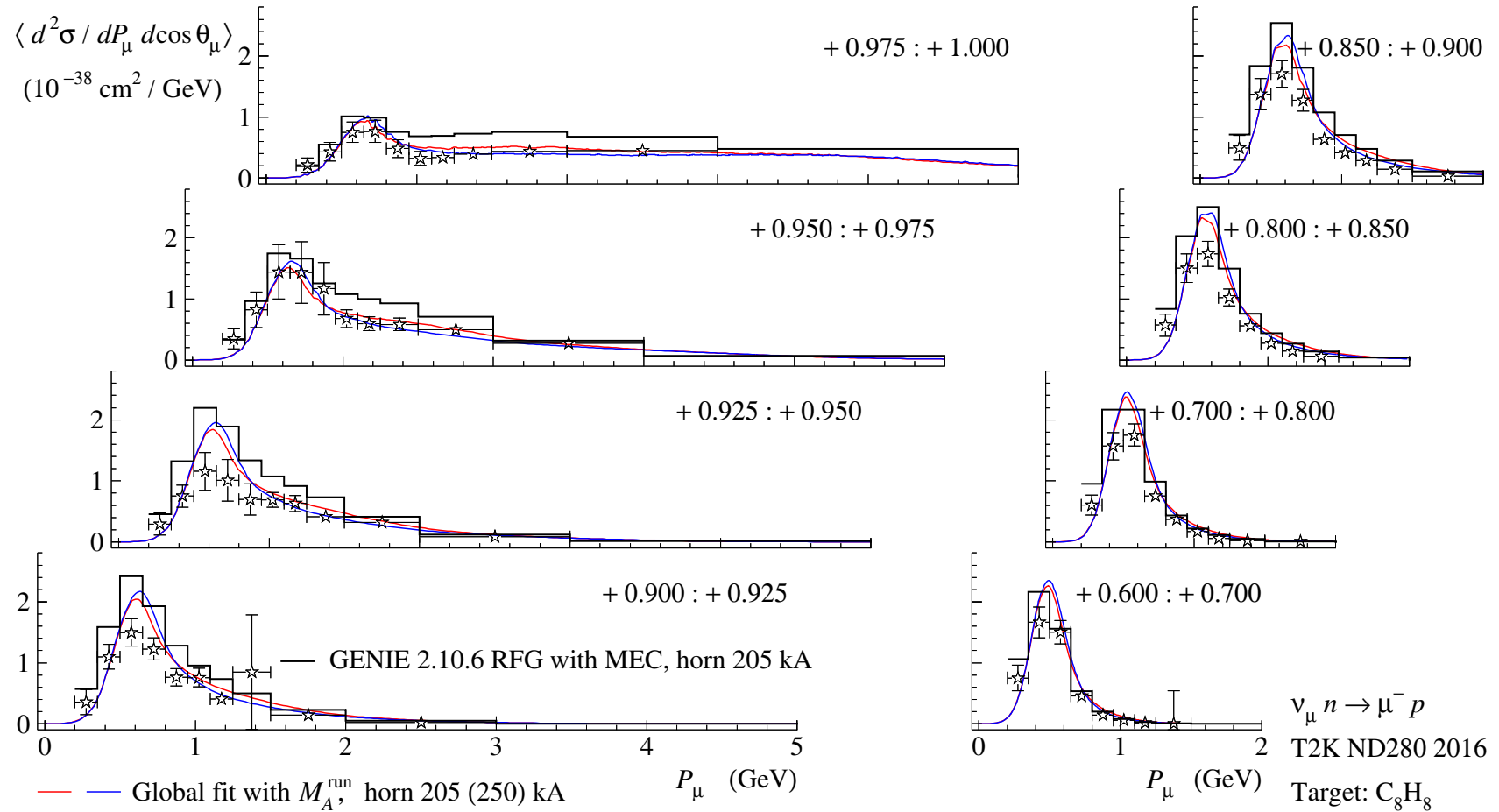
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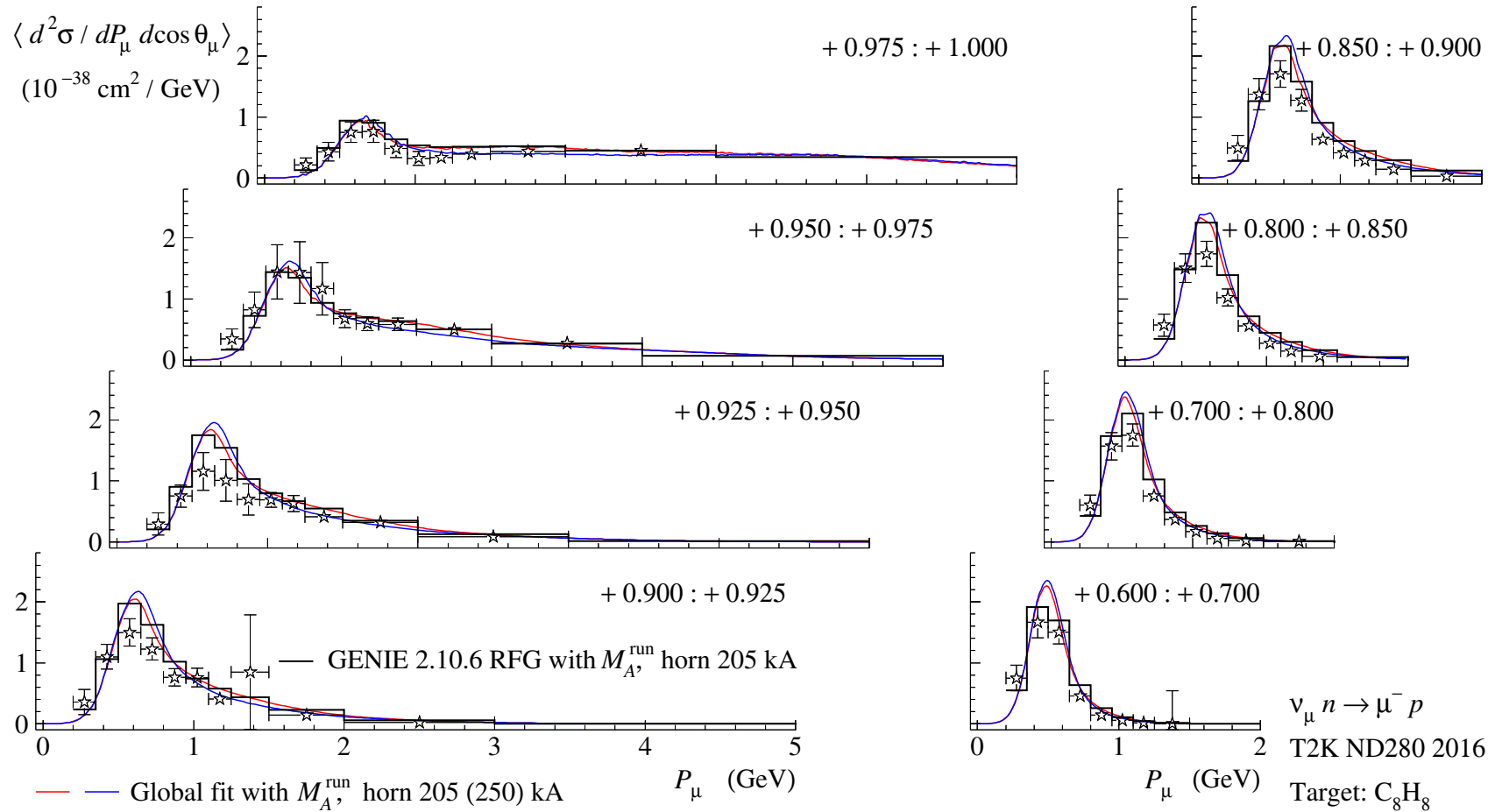
<sup>a</sup>The FGD1 composition: carbon (86.1%), hydrogen (7.35%), oxygen (3.70%), and small quantities of other elements such as titanium, silicon, and nitrogen (2.85%) [69].











# Parameters of the SM RFG Model

## Summary

The binding (separation) energies and Fermi momenta for isoscalar nuclei are calculated by using the parameterizations

$$p_F = 270[1 - 4.2/A + (6.0/A)^2 - (5.3/A)^3]$$

and

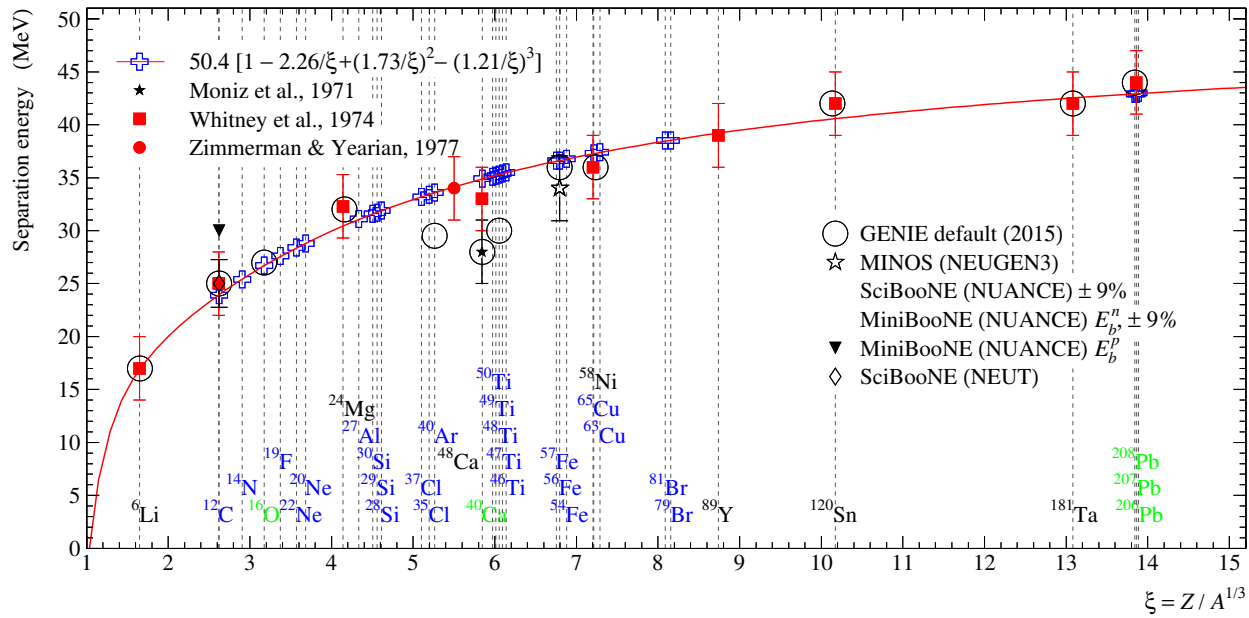
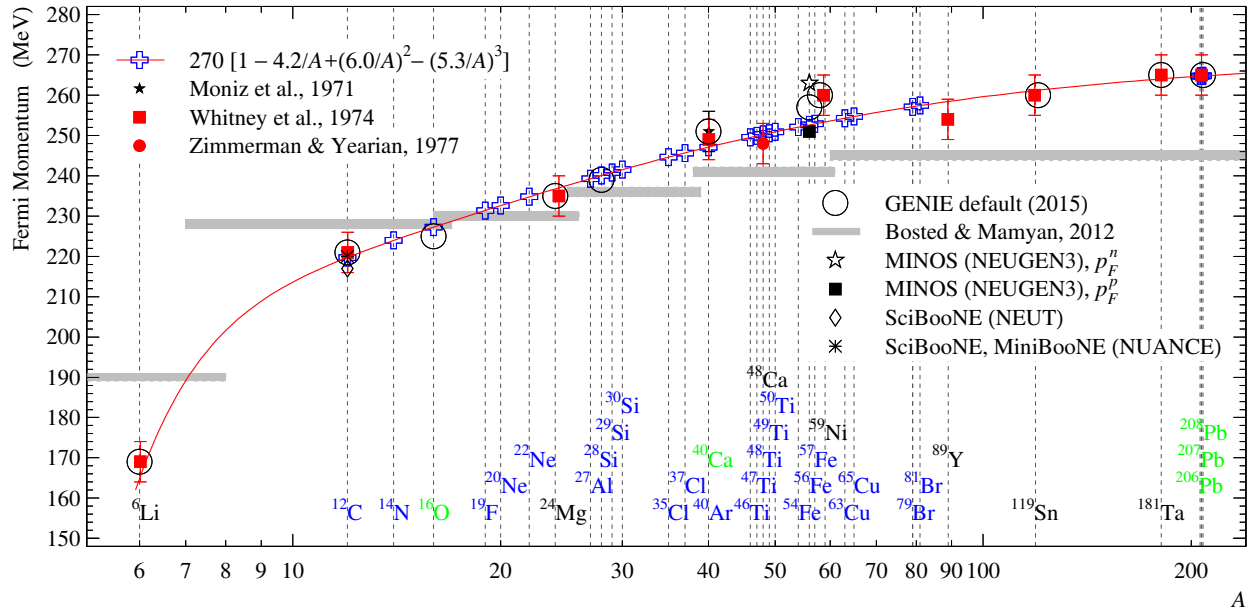
$$\varepsilon_b = 50.4 [1 - 2.26/\xi + (1.73/\xi)^2 - (1.21/\xi)^3], \quad \xi = Z/A^{1/3},$$

obtained from available (unfortunately, rather old) data on electron-nucleus scattering.<sup>a</sup> The proton and neutron Fermi momenta ( $p_F^p$  and  $p_F^n$ , respectively) are then calculated in the conventional way:<sup>b</sup>

$$p_F^p = (2Z/A)^{1/3} p_F \quad \text{and} \quad p_F^n = (2N/A)^{1/3} p_F \quad (N = A - Z).$$

<sup>a</sup> E.J. Moniz, I. Sick, R.R. Whitney, J.R. Ficenece, R.D. Kephart, and W.P. Trower, "Nuclear Fermi momenta from quasielastic electron scattering," Phys. Rev. Lett. **26**, 445 (1971); R.R. Whitney, I. Sick, J.R. Ficenece, R.D. Kephart and W.P. Trower, "Quasielastic electron Scattering," Phys. Rev. C **9**, 2230 (1974); P.D. Zimmerman and M.R. Yearian, "Fermi momenta and separation energies obtained from the quasi-elastic scattering of electrons from <sup>48</sup>Ca and <sup>40</sup>Ca," Z. Phys. A **278**, 291 (1976).

<sup>b</sup>These relations are based on the simplest assumption that the density of nuclear matter is approximately constant irrespective of the proton-to-neutron ratio  $Z/N$ .



**“Golden” dataset  
for the global fits**

## Accountancy & taxonomy

### Experiments with **deuterium** and **hydrogen** targets:

ANL 1977 [5] ( $\sigma_\nu$ , 8 data points),  
 ANL 1982 [7] ( $\langle dN_\nu/dQ^2 \rangle$ , 39 data points),  
 BNL 1980 [10] ( $\sigma_{\bar{\nu}}$ , 5 data point),  
 BNL 1990 [14, 15] ( $\langle dN_\nu/dQ^2 \rangle$ , 37 data points),  
 FNAL 1983 [16] ( $\langle dN_\nu/dQ^2 \rangle$ , 20 data points), and  
 CERN BEBC 1990 [44] ( $\langle d\sigma_\nu/dQ^2 \rangle$ , 8 data points).

### Experiment with **Ne-H<sub>2</sub>** mixture

FNAL 1984 [17] ( $dN_{\bar{\nu}}/dQ^2$ , 14 data points).

### Experiment with **aluminium** target

IHEP-ITEP 1985 [51] ( $\sigma_\nu$ ,  $\sigma_{\bar{\nu}}$ ,  $\langle d\sigma_{\nu+\bar{\nu}}/dQ^2 \rangle$ , 8 data points in each set).

### Experiments with **carbonaceous** targets

FNAL MiniBooNE 2010 [24] ( $\langle d^2\sigma_\nu/dE_\mu d\cos\theta_\mu \rangle$ , 137 data points),  
 FNAL MiniBooNE 2013 [25] ( $\langle d^2\sigma_{\bar{\nu}}/dE_\mu d\cos\theta_\mu \rangle$ , 75 data points),  
 CERN NOMAD 2009 [47, 48] ( $\sigma_\nu$  and  $\sigma_{\bar{\nu}}$ , 10 and 6 data points, respectively), and  
 FNAL MINER $\nu$ A 2013 [27, 28] ( $\langle d\sigma_\nu/dQ^2 \rangle$ ,  $\langle d\sigma_{\bar{\nu}}/dQ^2 \rangle$ , 8 data points in each set).

### Experiments with **carbon-containing mixtures**

CERN GGM 1979 [43, 45] ( $\langle dN_{\bar{\nu}}/dQ^2 \rangle$ , 13 data points),  
 CERN LAr-TPC 2007 [46] ( $\sigma_\nu$ , 1 data point),  
 IHEP SKAT 1990 [56] ( $\langle d\sigma_\nu/dQ^2 \rangle$  and  $\langle d\sigma_{\bar{\nu}}/dQ^2 \rangle$ , 8 and 7 data points, respectively).

Experiment with **iron** target T2K INGRID 2015 [63] ( $\sigma_\nu$ , 2 data points).

- Thus the fitted dataset consists of **430** data points with **286** data points for the  $\nu_\mu$  cross sections (66.5% of the full dataset), **136** data points for the  $\bar{\nu}_\mu$  cross sections (31.6%), and **8** data points for the summarized cross sections of  $\nu_\mu + \bar{\nu}_\mu$  (1.9%).
- The experimental data are presented as follows:
  - 212** data points for the flux-averaged double differential cross sections  $\langle d^2\sigma/dE_\mu d\cos\theta_\mu \rangle$  (49.3% of the full dataset, FNAL MiniBooNE),
  - 47** data points for the flux-averaged differential cross sections  $\langle d\sigma/dQ^2 \rangle$  (10.9%),
  - 128** data points for the flux-averaged  $Q^2$  distributions  $\langle dN/dQ^2 \rangle$  (29.8%), and
  - 43** data points for the total cross sections (10.0%).
- The dataset for extracting the current axial mass  $M_A$  (deuterium and hydrogen targets) contains **85** data points (27.2% of the full dataset): ANL 1977 [5], ANL 1982 [7], BNL 1980 [10], BNL 1990 [14, 15], FNAL 1983 [16], CERN 1990 [44].

### Criteria for exclusion of data from the global fit:

- Obsolete data (superseded or reconsidered due to increased statistics, revised normalization, etc. in the posterior reports of the same experimental groups).
- Data from the experiments with inactive targets (suggested by Arie Bodek).
- Data with  $\chi^2/\text{NDF} > 5$  and  $|N - 1| > 2\sigma_N$  (where  $N$  is the flux normalization).

**Convention:** below, the **solid** symbols in the figures represent the data included into the global fit while **open** symbols represent the data excluded from the fit.

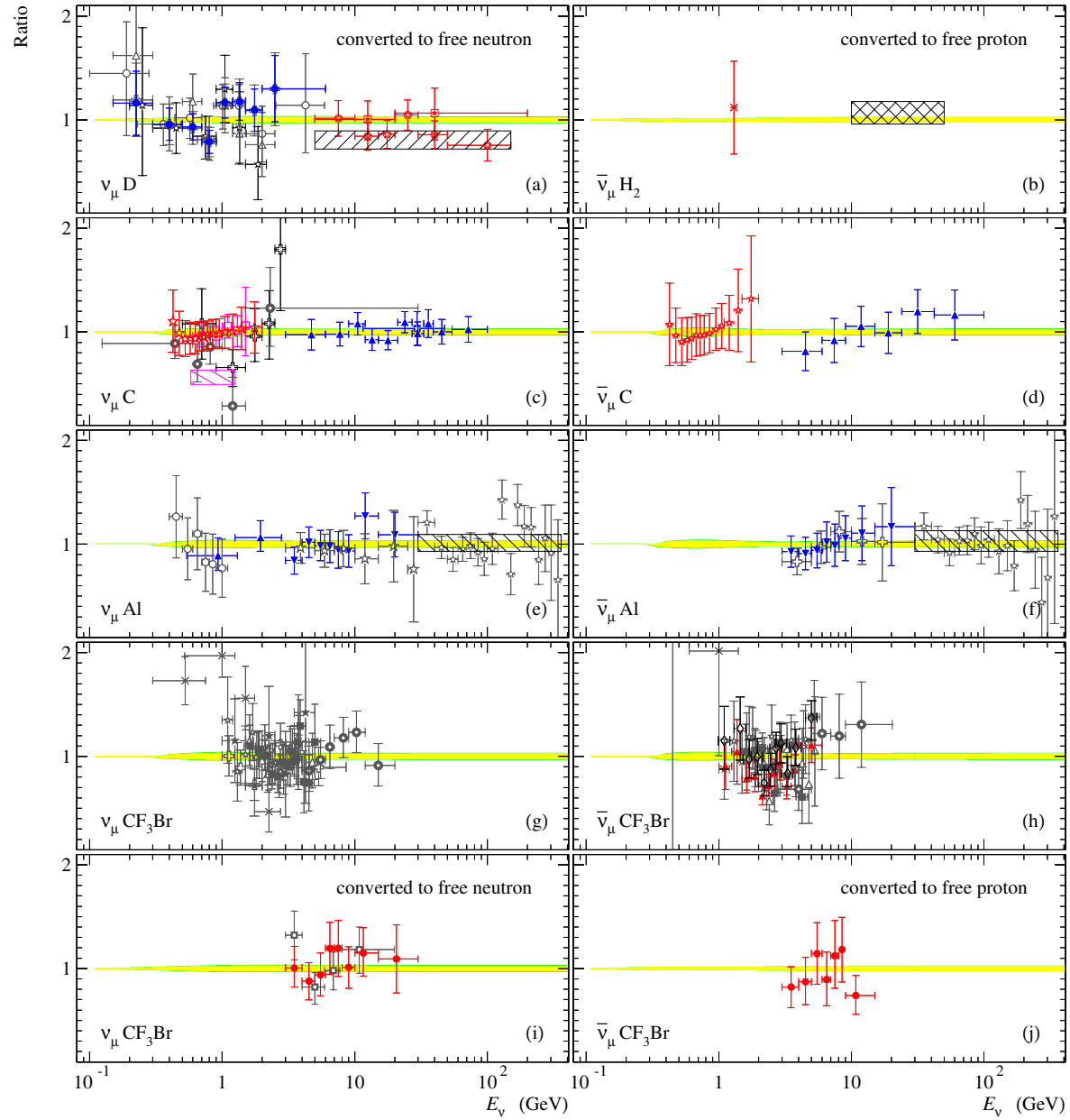


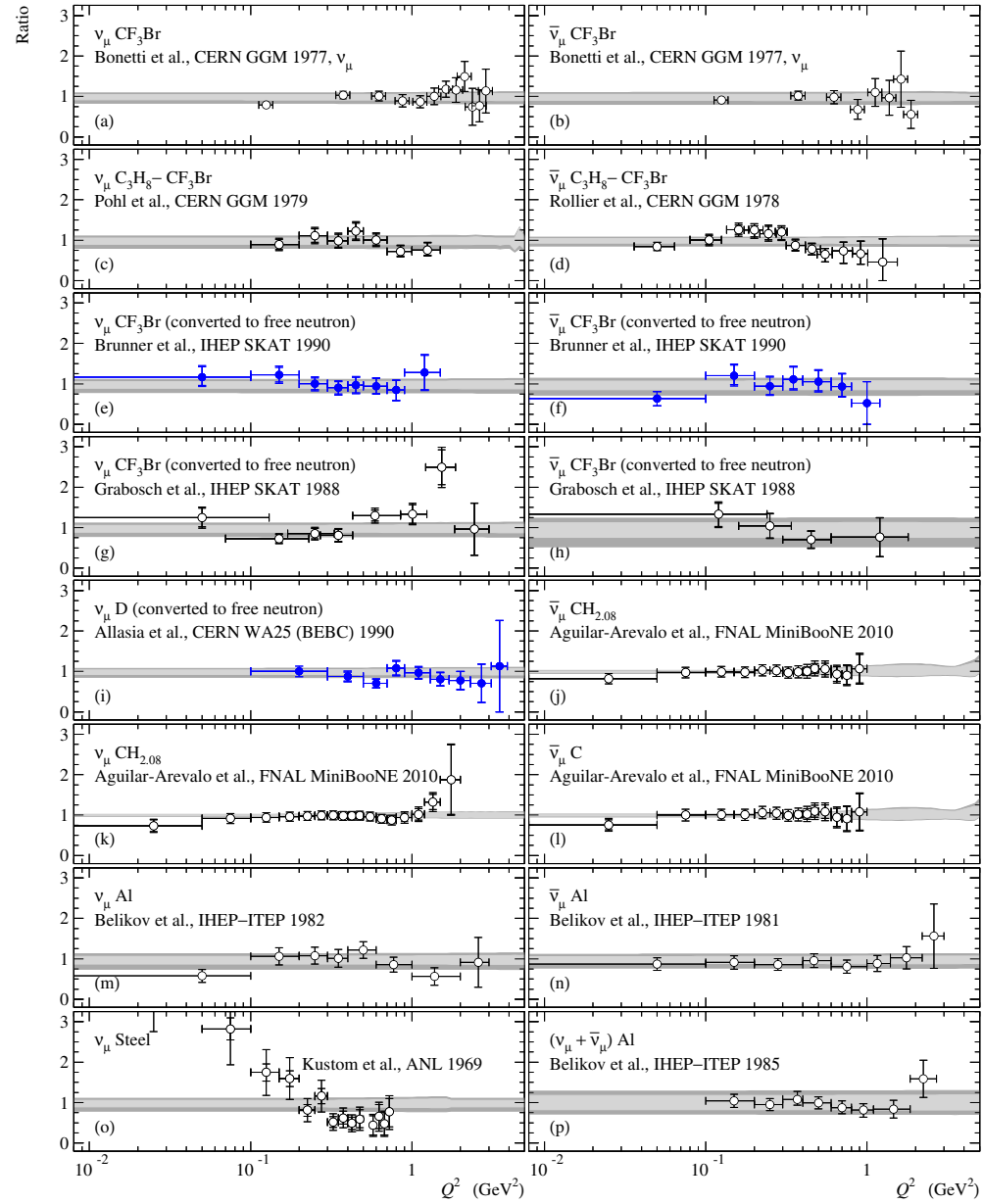
## A few more details

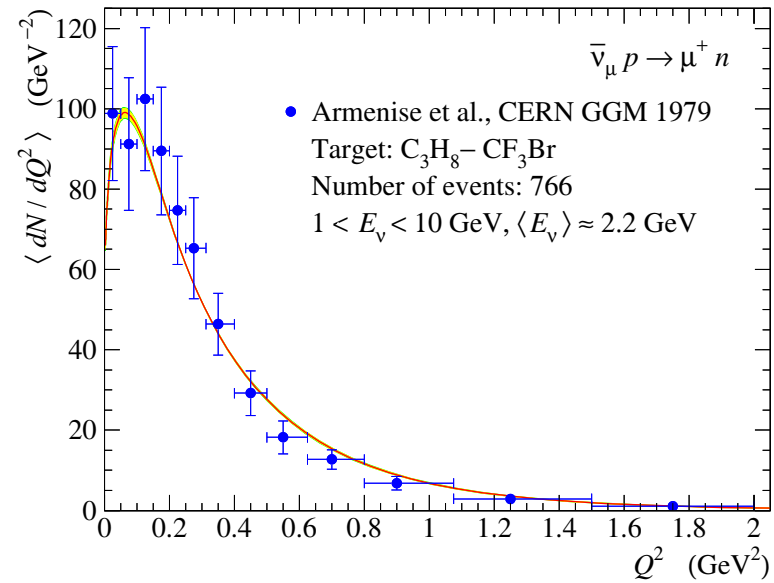
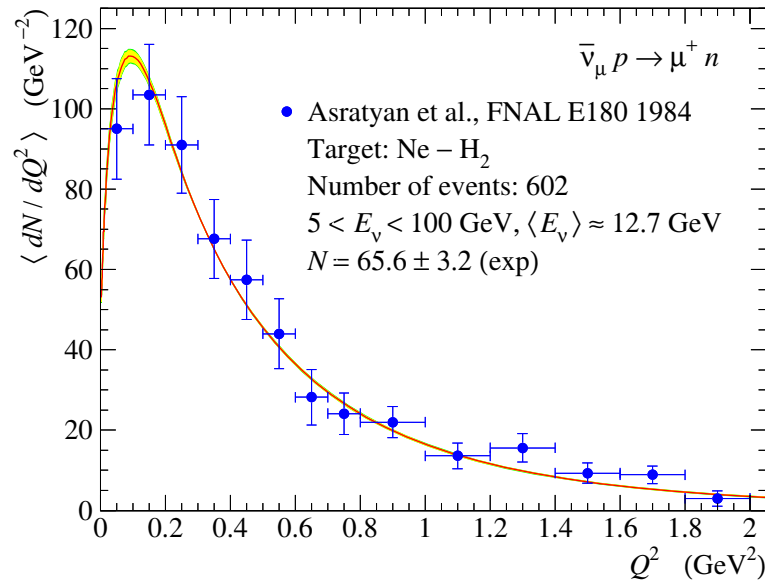
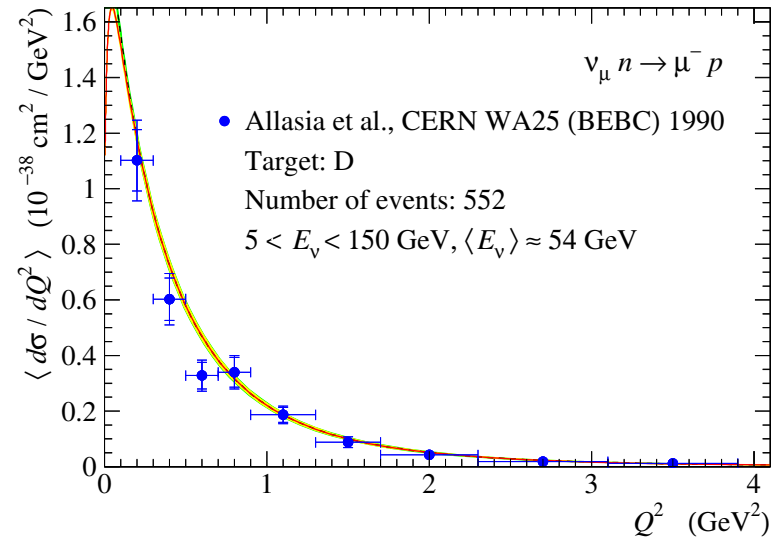
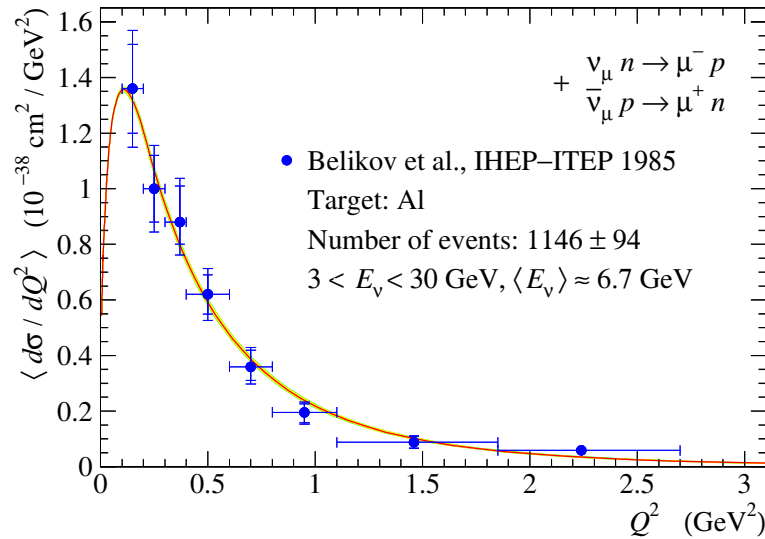
| Experiment   | Data type  | $N$                       | $\chi^2/NP$        |
|--|--|---------------------------|--------------------|
| Deuterium & hydrogen data and data for nuclear targets recalculated to free nucleons |  |                           |                    |
| Barish <i>et al.</i> , ANL 1977 [5]  | $\sigma_\nu^*$   | $0.893 \pm 0.034 (0.068)$ | $15.97/8 = 2.00$   |
| Miller <i>et al.</i> , ANL 1982 [7]  | $\langle dN_\nu/dQ^2 \rangle^*$                                  | $99.7 \times 0.979$       | $73.90/38 = 1.95$  |
| Fanourakis <i>et al.</i> , BNL 1980 [10]   | $\langle dN_{\bar{\nu}}/dQ^2 \rangle^*$                          | $2.4 \times 0.901$        | $1.37/4 = 0.34$    |
|  | $\sigma_{\bar{\nu}}$   | –                         | $0.05/1 = 0.05$    |
| Kitagaki <i>et al.</i> , BNL 1990 (2003 update) [14]                                 | $\langle dN_\nu/dQ^2 \rangle^*$                                  | $245.4 \times 0.991$      | $21.49/37 = 0.58$  |
| Kitagaki <i>et al.</i> , FNAL 1983 [16]  | $\langle dN_\nu/dQ^2 \rangle^*$                                  | $41.5 \times 0.945$       | $8.53/20 = 0.43$   |
|  | $\sigma_\nu$   | $1.020 \pm 0.077 (0.152)$ | $0.18/2 = 0.09$    |
| Allasia <i>et al.</i> , CERN BEBC 1990 [44]  | $\langle d\sigma_\nu/dQ^2 \rangle^*$                             | $0.935 \pm 0.035 (0.070)$ | $4.77/8 = 0.68$    |
|  | $\sigma_\nu$   | $0.873 \pm 0.032 (0.064)$ | $15.11/6 = 3.02$   |
| Brunner <i>et al.</i> , IHEP SKAT 1990 [55,56]                                       | $\langle d\sigma_\nu/dQ^2 \rangle^*$                             | $1.077 \pm 0.059 (0.117)$ | $3.89/8 = 0.49$    |
|  | $\sigma_\nu$   | $1.056 \pm 0.058 (0.115)$ | $3.14/8 = 0.39$    |
|  | $\langle d\sigma_{\bar{\nu}}/dQ^2 \rangle^*$                     | $0.902 \pm 0.064 (0.126)$ | $7.31/7 = 1.04$    |
|  | $\sigma_{\bar{\nu}}$   | $0.823 \pm 0.056 (0.111)$ | $8.10/7 = 1.15$    |
| Heavy targets data   |  |                           |                    |
| Asratyan <i>et al.</i> , FNAL 1984 [17]  | $\langle dN_{\bar{\nu}}/dQ^2 \rangle^*$                          | $68.7 \times 0.993$       | $7.93/14 = 0.57$   |
|  | $\sigma_{\bar{\nu}}$   | –                         | $0.06/1 = 0.06$    |
| Aguilar-Arevalo <i>et al.</i> , FNAL MiniBooNE 2010 [24]                             | $\langle d^2\sigma_\nu/dE_\mu d\cos\theta_\mu \rangle^*$         | $1.000 \pm 0.008 (0.015)$ | $35.76/137 = 0.26$ |
| Aguilar-Arevalo <i>et al.</i> , FNAL MiniBooNE 2013 [25]                             | $\langle d^2\sigma_{\bar{\nu}}/dE_\mu d\cos\theta_\mu \rangle^*$ | $1.017 \pm 0.012 (0.023)$ | $46.86/75 = 0.63$  |
| Fiorentini <i>et al.</i> , FNAL MINER $\nu$ A 2013 [27]                              | $\langle d\sigma_\nu/dQ^2 \rangle^*$                             | –                         | $8.04/8 = 1.01$    |
| Fields <i>et al.</i> , FNAL MINER $\nu$ A 2013 [28]                                  | $\langle d\sigma_{\bar{\nu}}/dQ^2 \rangle^*$                     | –                         | $5.51/8 = 0.67$    |
| Armenise <i>et al.</i> , CERN GGM 1979 [43,45]                                       | $\sigma_{\bar{\nu}}$   | $0.781 \pm 0.031 (0.062)$ | $20.57/11 = 1.87$  |
| Martinez de la Ossa Romero <i>et al.</i> , CERN LAr-TPC 2007 [46]                    | $\sigma_\nu^*$   | –                         | $0.09/1 = 0.09$    |
| Lyubushkin <i>et al.</i> , CERN NOMAD 2009 [47,48]                                   | $\sigma_\nu^*$   | $1.036 \pm 0.026 (0.052)$ | $5.91/10 = 0.59$   |
|  | $\sigma_{\bar{\nu}}^*$   | $1.016 \pm 0.058 (0.114)$ | $3.09/6 = 0.52$    |
| Belikov <i>et al.</i> , IHEP–ITEP 1985 [51]  | $\langle d\sigma_{\nu+\bar{\nu}}/dQ^2 \rangle^*$                 | $0.975 \pm 0.114 (0.184)$ | $5.68/8 = 0.71$    |
|  |  | $0.988 \pm 0.131 (0.212)$ |                    |
|  | $\sigma_\nu^*$   | $0.873 \pm 0.036 (0.071)$ | $7.02/8 = 0.88$    |
|  | $\sigma_{\bar{\nu}}^*$   | $0.897 \pm 0.045 (0.089)$ | $2.66/8 = 0.33$    |
| Abe <i>et al.</i> , T2K ND280 2014 [61]  | $\sigma_\nu^*$   | –                         | $9.57/5 = 2.39$    |
| Abe <i>et al.</i> , T2K INGRID 2015 [63]   | $\sigma_\nu^*$   | $0.977 \pm 0.086 (0.170)$ | $0.55/2 = 0.28$    |

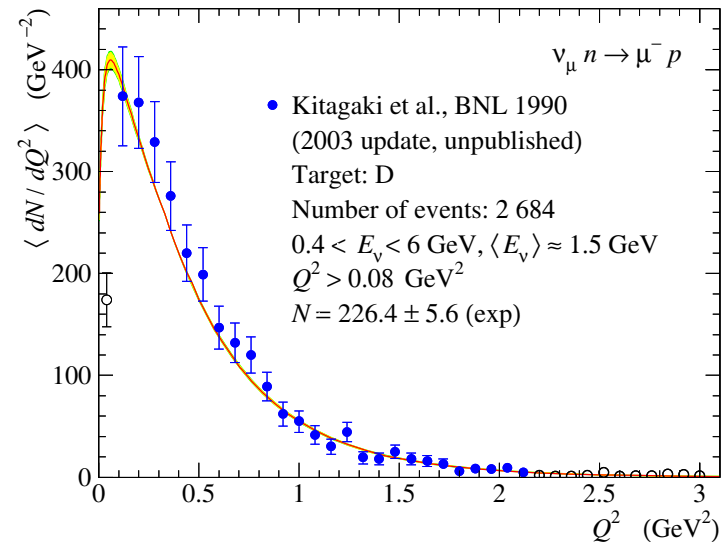
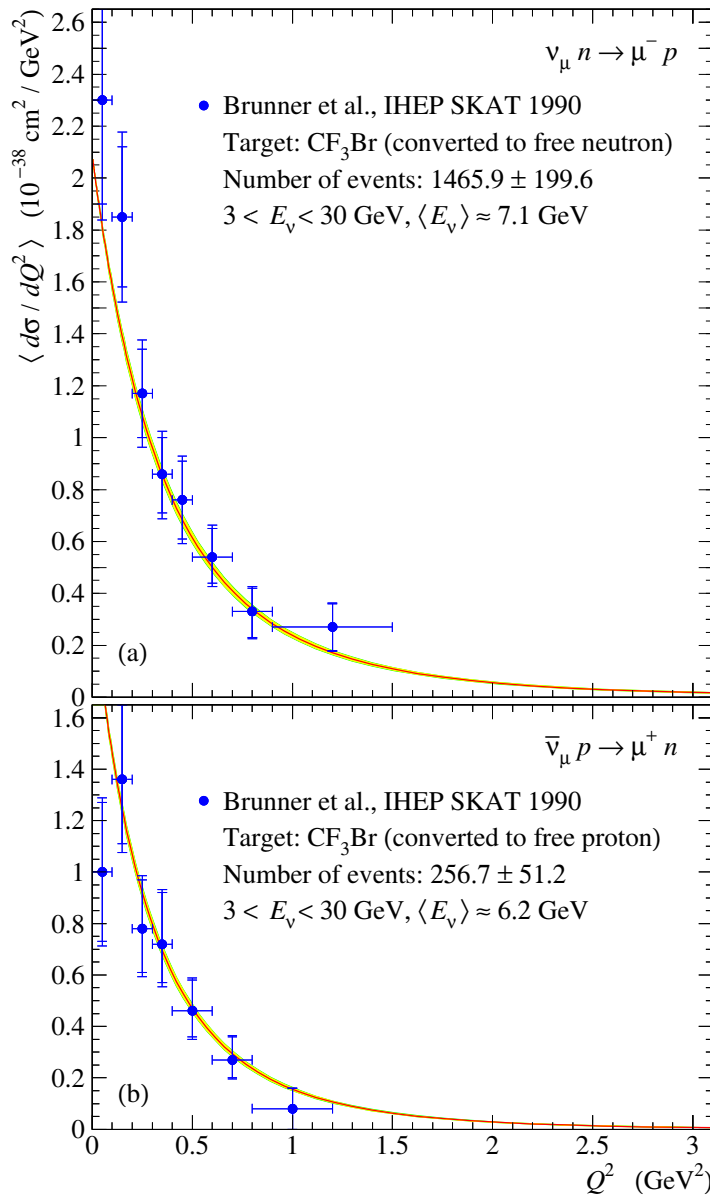
\* The data included into the global fits.

# Comparison with the early golden data

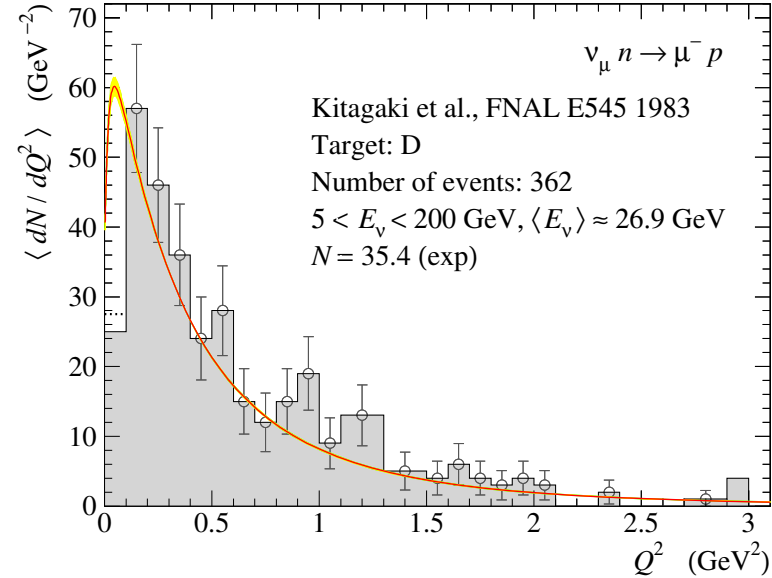
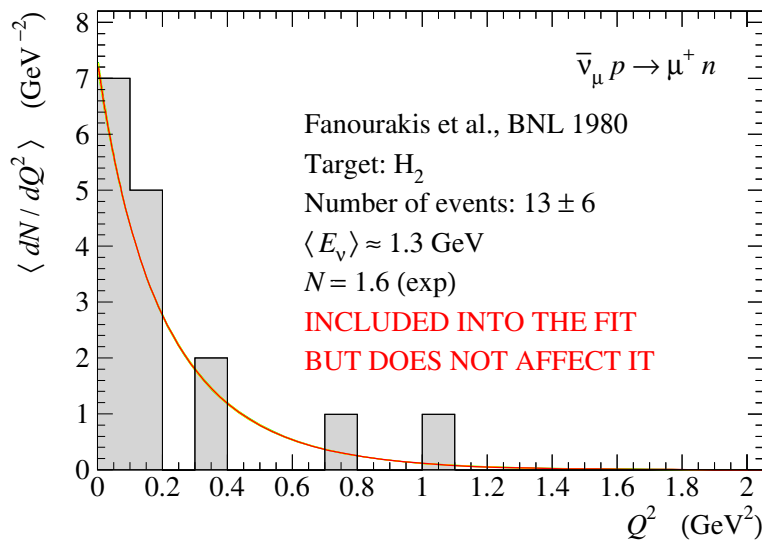
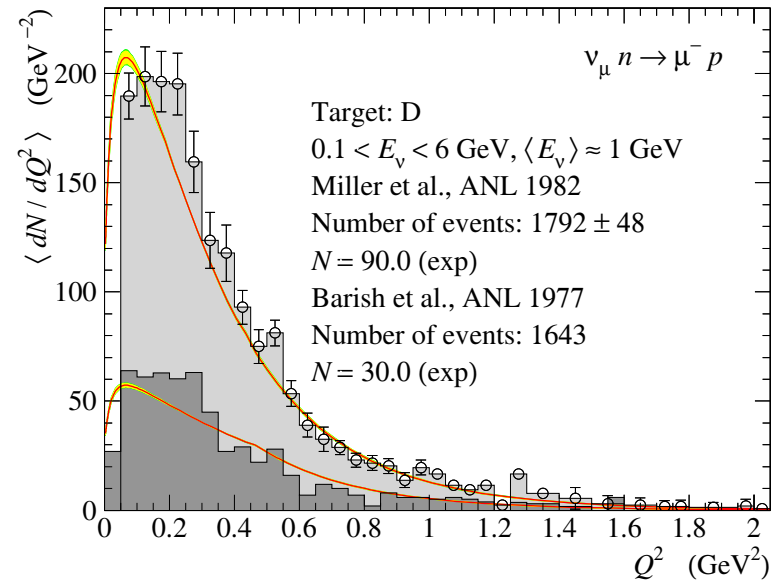
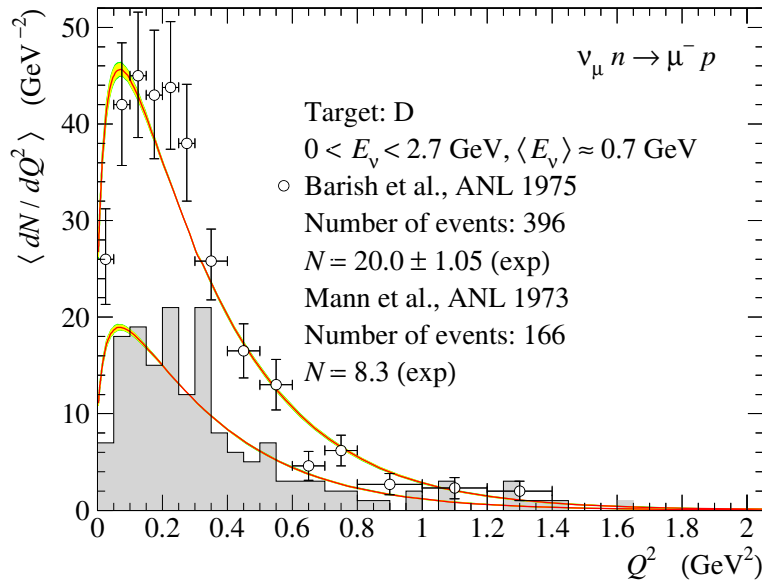




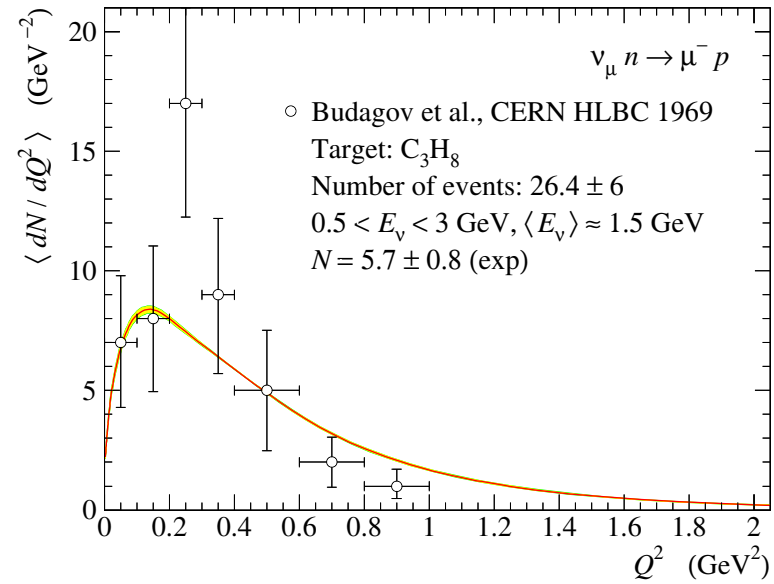
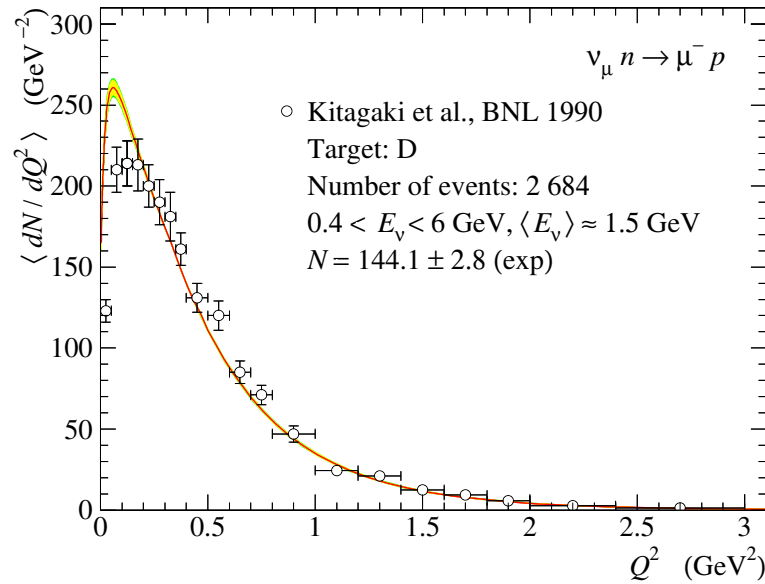
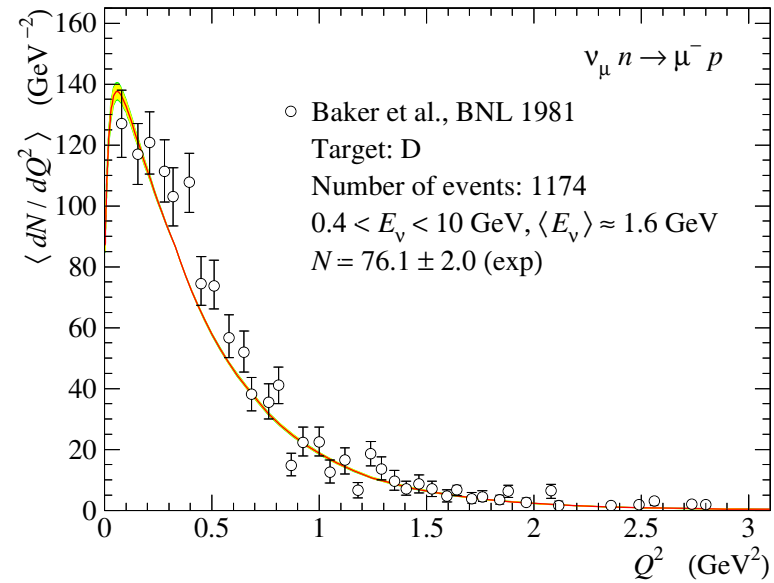
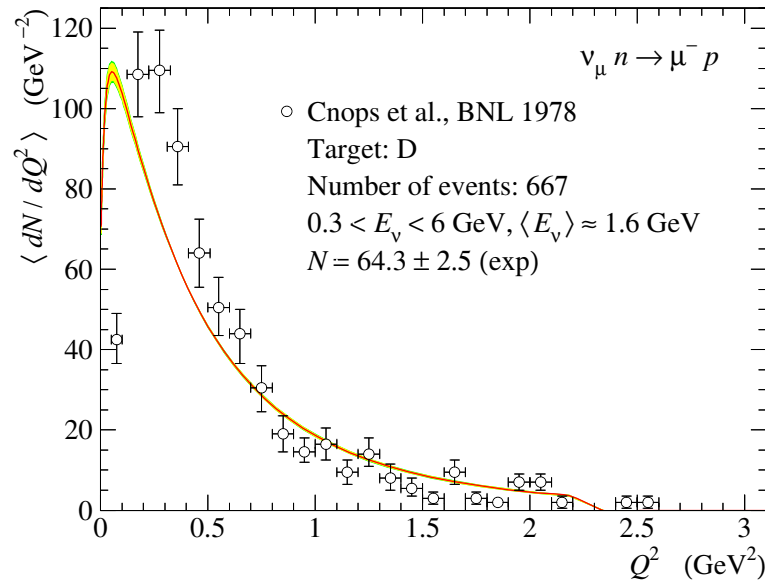


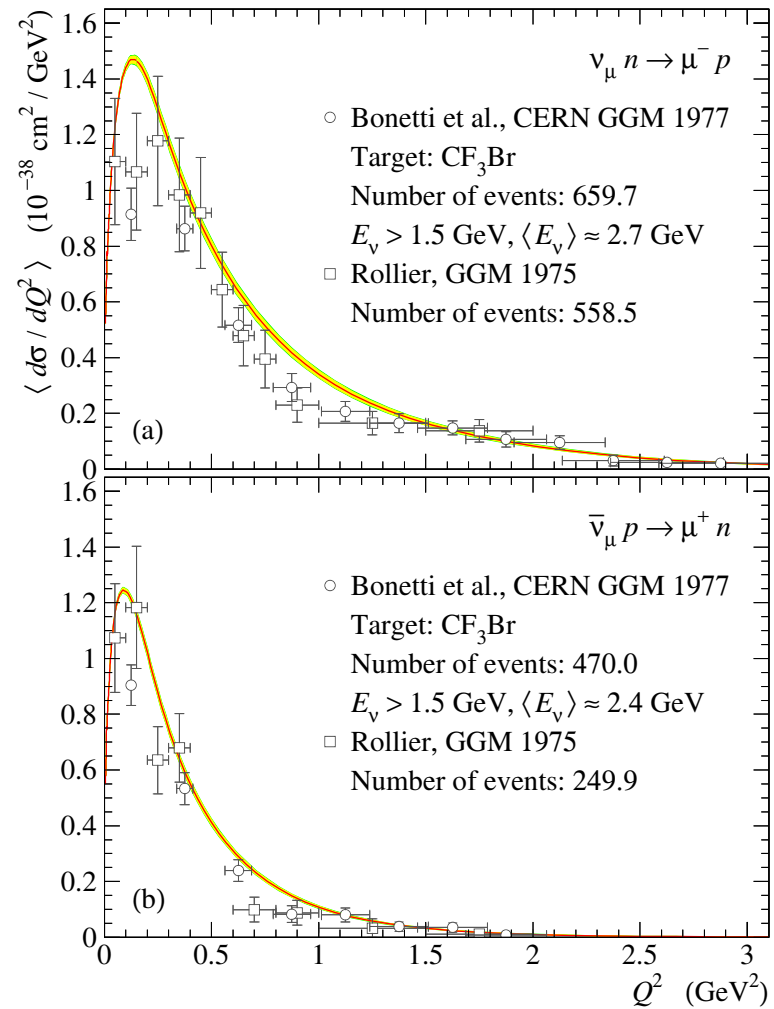
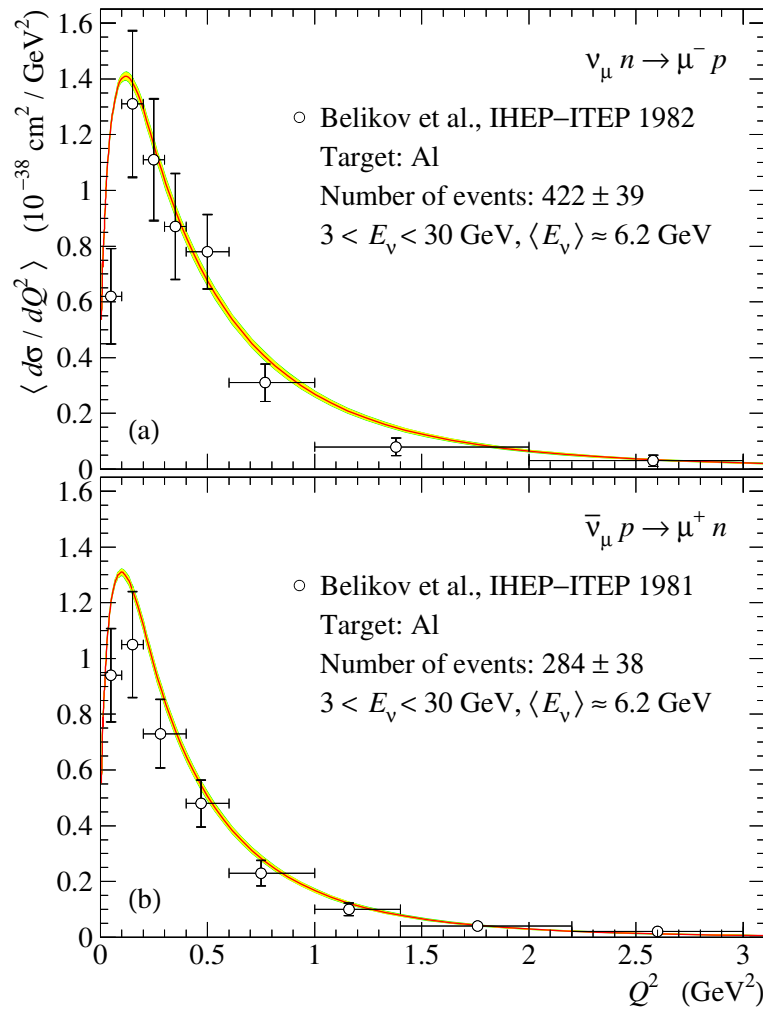


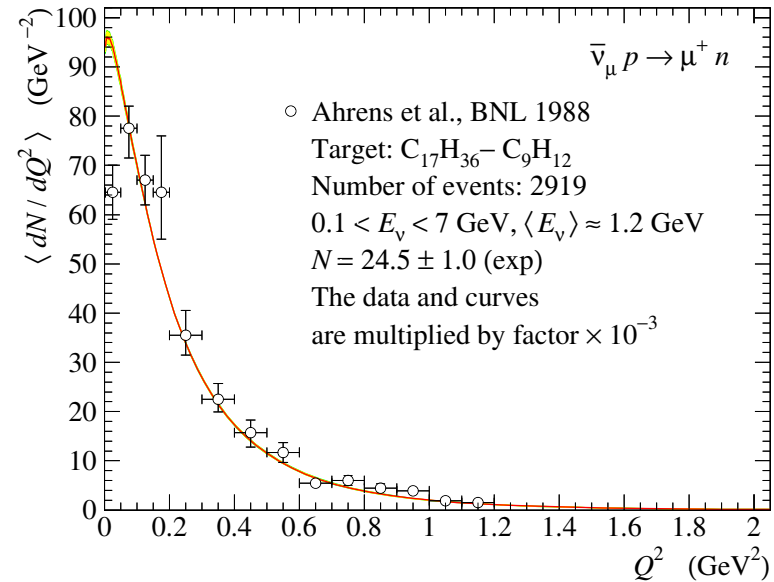
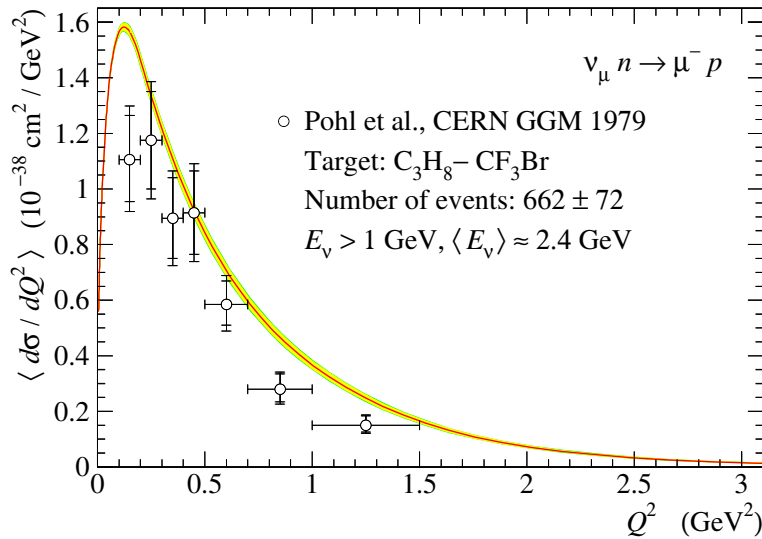
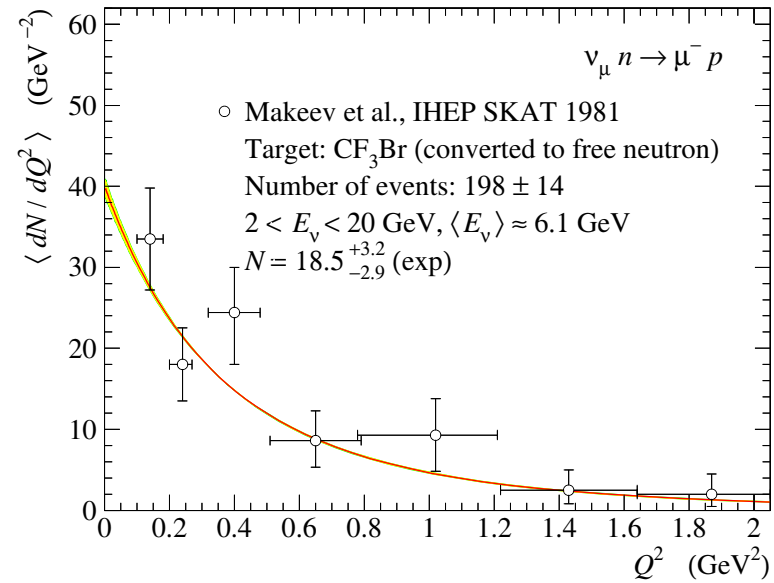
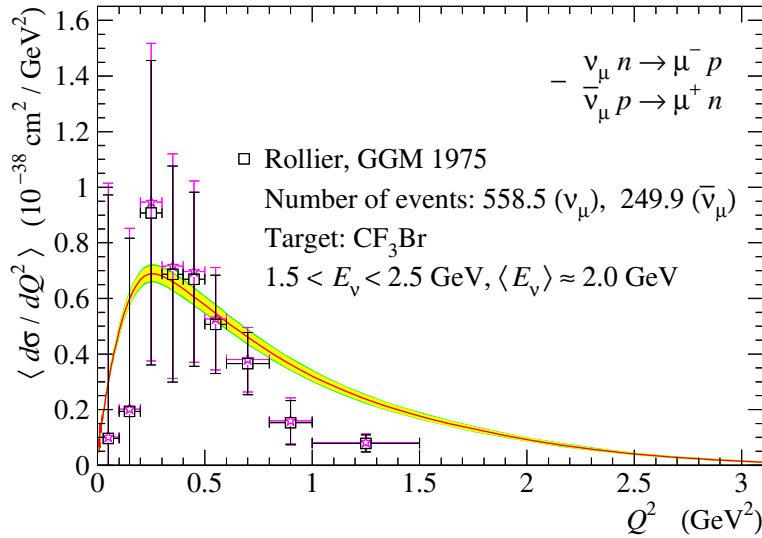
**Comparison with certain datasets  
unclaimed in the global fit**

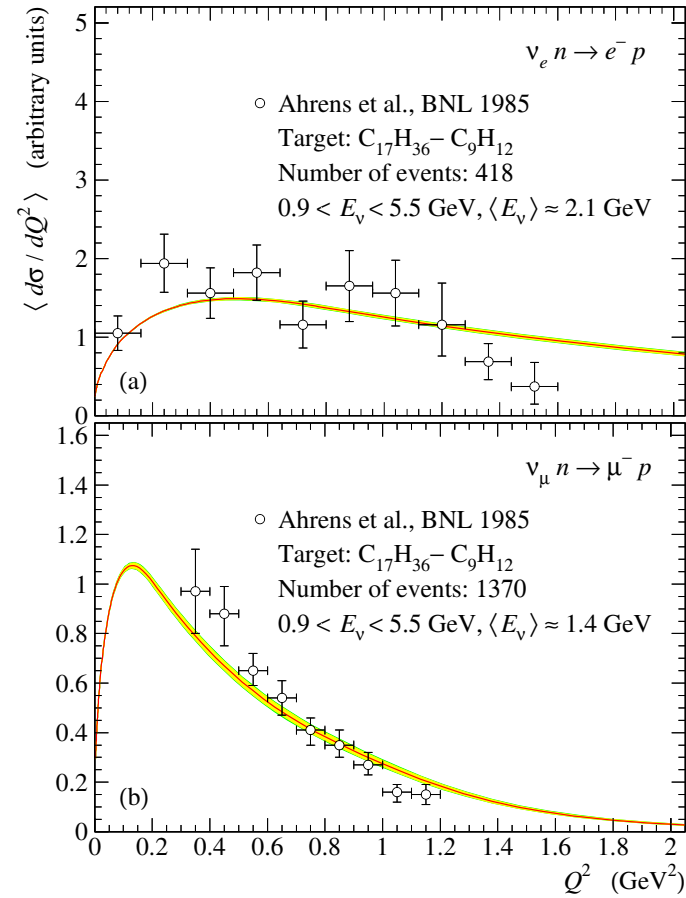
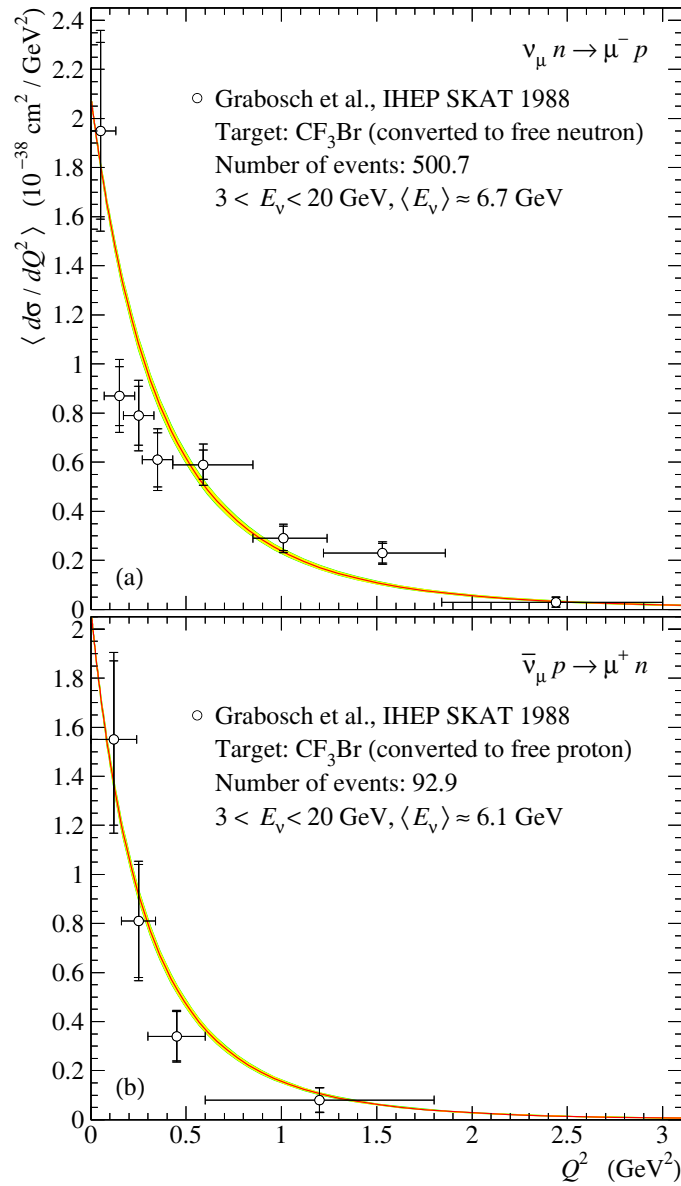


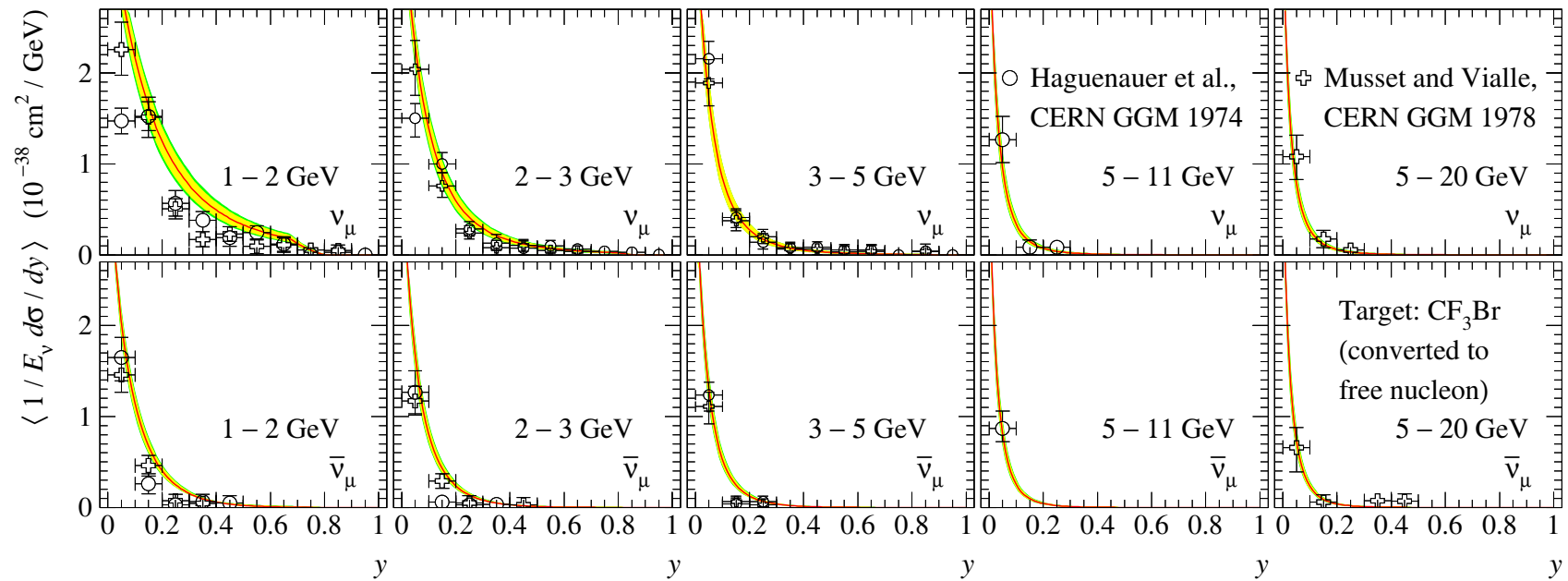










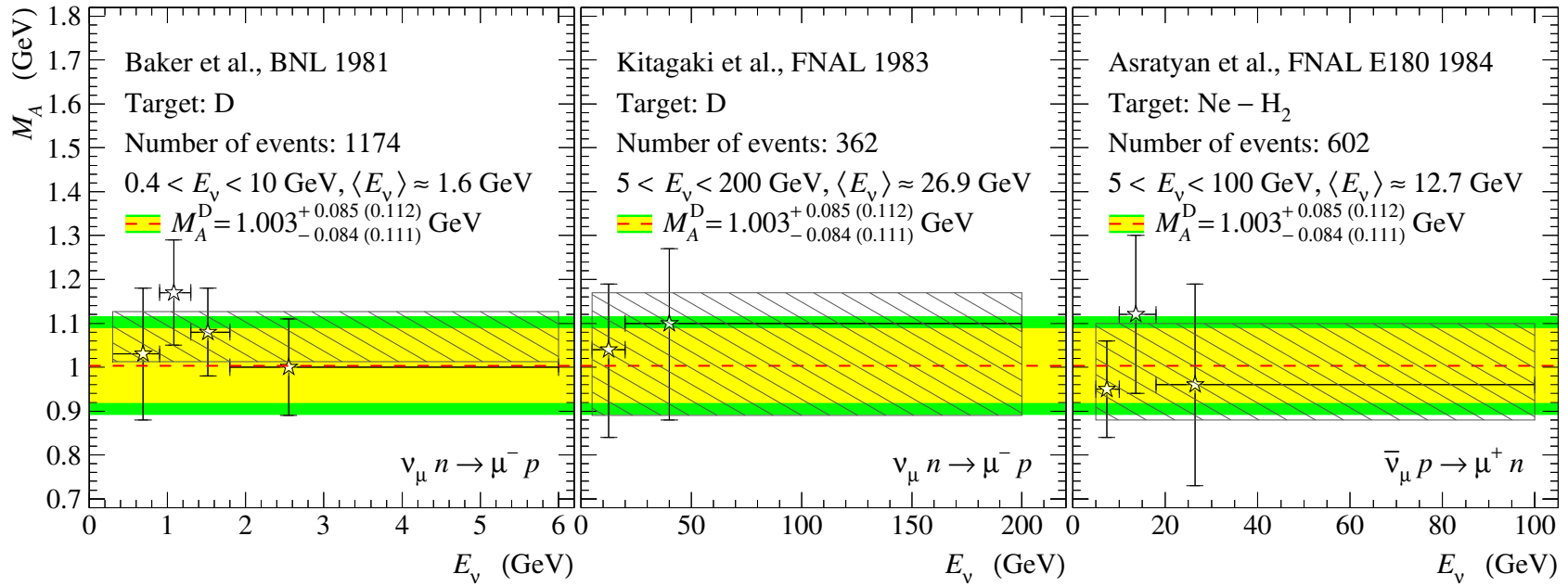


Flux-weighted differential cross sections for  $\nu_\mu n \rightarrow \mu^- p$  (top) and  $\bar{\nu}_\mu p \rightarrow \mu^+ n$  (bottom). The cross sections were measured in the Gargamelle bubble chamber filled with heavy liquid freon and exposed to the CERN-PS  $\nu_\mu$  and  $\bar{\nu}_\mu$  beams.

The shaded regions around the solid curves represent  $1\sigma$  and  $2\sigma$  deviations from the best-fitted values of  $M_A^{\text{run}}$  obtained in the global fit. Only the events with  $E_{\nu, \bar{\nu}} > 1.5$  GeV were accepted.

The energy range for each experiment is shown in the legends.

Obsolete GGM 1975 data are also shown for completeness.

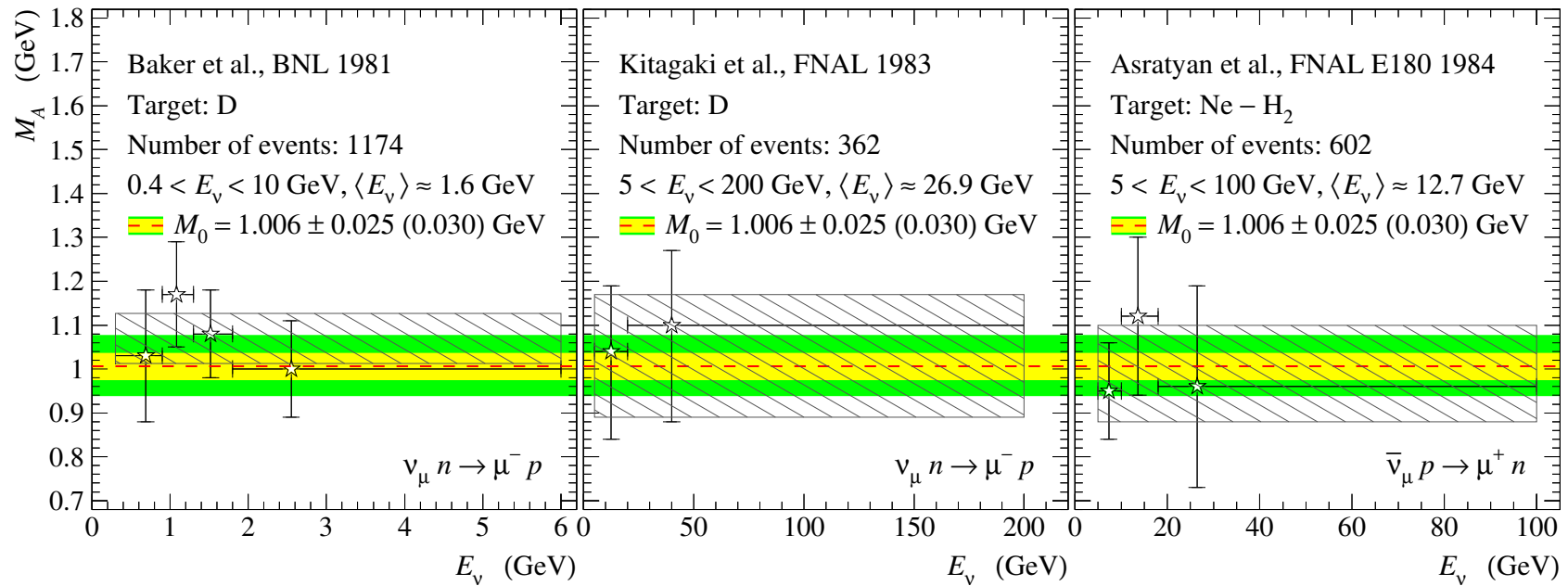


Current axial mass,  $M_A$  vs. neutrino energy measured for the  $\nu_\mu n \rightarrow \mu^- p$  reaction with the deuterium filled detector at BNL [11]. Rectangle denotes the average value of  $M_A$  obtained by the authors of the experiment for the mean neutrino energy  $\langle E_\nu \rangle = 1.48$  GeV.

Current axial mass,  $M_A$  vs. neutrino energy measured for the  $\nu_\mu n \rightarrow \mu^- p$  reaction with the deuterium filled detector at BNL [16]. Rectangle denotes the average value of  $M_A$  obtained by the authors of the experiment for the whole neutrino energy range at the mean energy of  $\langle E_\nu \rangle = 26.9$  GeV.

Current axial mass,  $M_A$  vs. antineutrino energy measured for the  $\bar{\nu}_\mu p \rightarrow \mu^+ n$  reaction with the Ne- $H_2$  filled detector in the FNAL experiment E180 [17, 18]. Dashed rectangle denotes the average value of  $M_A$  obtained by the authors of the experiment for the whole energy range at the mean energy of  $\langle E_\nu \rangle = 15.4$  GeV.

Dashed line and shaded bands in each panel represent the best-fit value of  $M_A^D$  and corresponding  $1\sigma$  and  $2\sigma$  uncertainties, obtained in our global fit to the deuterium and hydrogen QECC data.



The same as in the previous slide except the lines and bands which now represent the best-fit value of the parameter  $M_0$  (which can be identified with the current axial mass of the nucleon,  $M_A$ ) and corresponding  $1\sigma$  and  $2\sigma$  uncertainties, obtained in our global fit to all CCQE data described in slide 47.

In all our fits, the nuclear corrections for the  $\nu - d$  CCQE scattering were accounted for within the impulse approximation and by using the closure over the final dinucleon states and meson exchange current correction at low  $Q^2$ ; the Reid hard core potential and Hulthen wave function for the deuteron were applied.<sup>a</sup>

<sup>a</sup>See S.K. Singh and H. Arenhovel, "Pion exchange current effects in  $\nu_\mu + d \rightarrow \mu^- + p + p$ ," Z. Phys. A **324**, 347 (1986) and references therein.

## Conclusions

- The Smith-Moniz RFG model supplemented by the energy-dependent (running) axial mass,  $M_A^{\text{run}}$ , defined by only **two universal parameters**  $M_0$  and  $E_0$ , describes well the earlier and modern CCQE data on various nuclear targets.
- The best-fit values of  $M_0$  and  $E_0$  **are sensitive to the details of the RFG model** (including the input Fermi momenta and separation energies) and to the model of the nucleon electromagnetic form factors. However the fit can almost automatically be repeated with other versions of RFG (e.g., Bodek–Ritchie), or its extensions (SF, LFG, etc.), or with more advanced models and improved inputs.
- A more sophisticated parametrization of  $M_A^{\text{run}}$  **is unreasonable** for the present-day level of accuracy of the CCQE (and CCQE-like) data but may be needed in the future.
- Individual parametrizations for different nuclei or nuclear groups **are not warranted**, but mainly because the currently available dataset for the inorganic heavy nuclei is not accurate enough and not quite self-consistent.
- There's no statistically significant difference between the  $M_A^{\text{run}}$  parameters extracted separately from the  $\nu$  and  $\bar{\nu}$  data, but there's **a weak hint to a possible difference**. Any case, the available  $\bar{\nu}$  dataset is not yet sufficient for a more definite statement.

Finally, we suggest incorporating (as an option) the running axial mass into the neutrino MC generators used in the modern neutrino oscillation experiments.



**Thanks for attention**



## $\chi^2$ schematically

$$\chi^2 = \sum_{i \text{ (uncorr)}} \sum_{j \in G_i} \frac{[N_i T_{ij}(\boldsymbol{\lambda}) - E_{ij}]^2}{\sigma_{ij}^2} + \sum_{i \text{ (uncorr)}} \frac{(N_i - 1)^2}{\sigma_i^2} + \sum_{i \text{ (corr)}} [\mathbf{T}_i(\boldsymbol{\lambda}) - \mathbf{E}_i] \mathbf{W}_i^{-1} [\mathbf{T}_i(\boldsymbol{\lambda}) - \mathbf{E}_i]^T .$$

Here

- index  $i$  enumerates the experiments (or data groups)  $G_i$ ,
- index  $j \in G_i$  enumerates the bin-averaged experimental data  $E_{ij}$  from the group  $G_i$  with the errors  $\sigma_{ij}$  which do not include the flux normalization uncertainty,  $\sigma_i$ . The data sets include the total CCQE cross sections, flux-averaged differential cross sections ( $d\sigma_{\nu,\bar{\nu}}/dQ^2$ ,  $d^2\sigma_{\nu,\bar{\nu}}/dE_\mu d\cos\theta_\mu$ , etc.), and unnormalized  $Q^2$  distributions ( $dN_{\nu,\bar{\nu}}/dQ^2$ ).
- $T_{ij}(\boldsymbol{\lambda})$  represent the model predictions dependent on the set of fitting parameters  $\boldsymbol{\lambda}$  (in our particular case,  $\boldsymbol{\lambda} = (M_0, E_0)$  for heavy nuclear targets and  $\boldsymbol{\lambda} = M_0$  for hydrogen and deuterium targets).
- $\mathbf{W}_i$  is the correlation matrix for the  $i$ th group.

Nuclear corrections for the CCQE cross sections on deuterium are, as a rule, already taken into account by the authors of the experiment and are represented as the cross sections or  $Q^2$  distribution for a free nucleon target.

The formal analytic expressions for the normalization factors  $N_i$  can be found by solving the system of equations

$$\partial\chi^2/\partial N_i = 0.$$

The solution is of the form:

$$N_i = \frac{1 + \sigma_i^2 \sum_{j \in G_i} \sigma_{ij}^{-2} T_{ij}(\boldsymbol{\lambda}) E_{ij}}{1 + \sigma_i^2 \sum_{j \in G_i} \sigma_{ij}^{-2} [T_{ij}(\boldsymbol{\lambda})]^2}.$$

The penalty terms  $(N_i - 1)^2 / \sigma_i^2$  contain the flux uncertainties ( $\sigma_i$ ) common to the data group  $G_i$ .

The  $\chi^2$  for the final fit to all data also includes another penalty term

$$\left( \frac{M_0 - M_A^D}{\Delta M_A^D} \right)^2,$$

which provides a “soft anchoring” of the parameter  $M_0$  to the current axial mass  $M_A^D \pm \Delta M_A^D$  obtained from the fitting the robust deuterium data only.

## Including the “mixed” data into the $\chi^2$ sum

The expression for  $\chi^2$  shown in the slide 67 becomes a bit more complicated when the data is a combination of the values obtained in the experiments with the neutrino and antineutrino beams.

Here we consider the case when we have to fit the *uncorrelated* data like  $E_1$ ,  $E_2$  and  $E_3 = E'_1 + E'_2$ , where the pairs  $(E_1, E'_1)$  and  $(E_2, E'_2)$  have the common normalization uncertainty. Let's construct the sum for minimization for this case:

$$\begin{aligned} \chi_{\text{mix}}^2 = & \sum_{j \in G_1} \frac{[N_1 T_{1j}(\boldsymbol{\lambda}) - E_{1j}]^2}{\sigma_{1j}^2} + \sum_{j \in G_2} \frac{[N_2 T_{2j}(\boldsymbol{\lambda}) - E_{1j}]^2}{\sigma_{2j}^2} \\ & + \sum_{j \in G_3} \frac{[N_1 T_{1'j}(\boldsymbol{\lambda}) + N_2 T_{2'j}(\boldsymbol{\lambda}) - E_{3j}]^2}{\sigma_{3j}^2} \\ & + \left( \frac{N_1 - 1}{\sigma_1} \right)^2 + \left( \frac{N_2 - 1}{\sigma_2} \right)^2 \end{aligned}$$

(the notation is obvious). From the minimization conditions

$$\frac{\partial \chi_{\text{mix}}^2}{\partial N_1} = \frac{\partial \chi_{\text{mix}}^2}{\partial N_2} = 0$$

we obtain the expressions for the normalization factors:

$$N_1 = \frac{1}{d} (C_1 A_{22} - C_2 A_{12}), \quad N_2 = \frac{1}{d} (C_2 A_{11} - C_1 A_{12}),$$

where

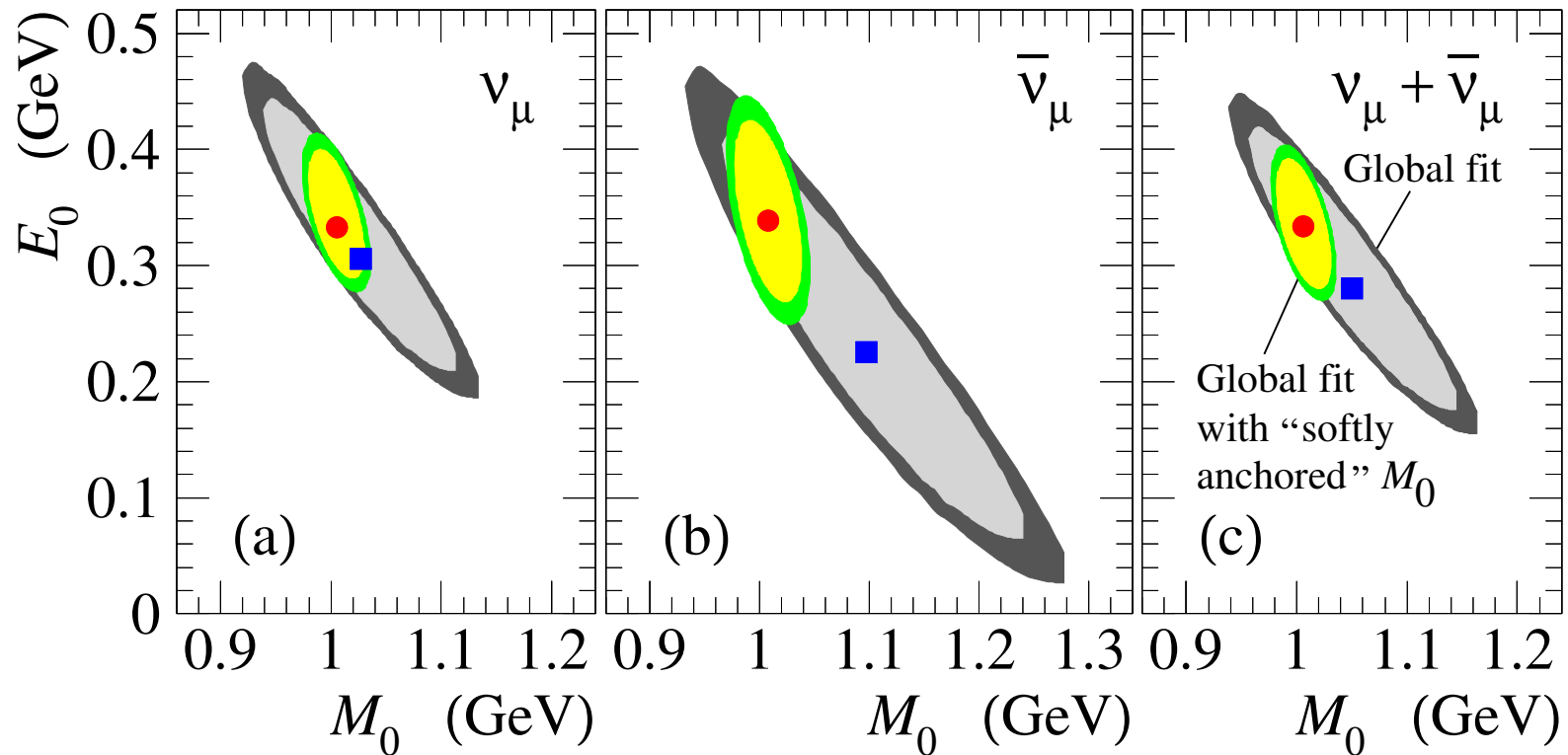
$$d = A_{11} A_{22} - A_{12}^2,$$

$$A_{11} = \sum_{j \in G_1} \frac{[T_{1j}(\lambda)]^2}{\sigma_{1j}^2} + \sum_{j \in G_3} \frac{[T_{1'j}(\lambda)]^2}{\sigma_{3j}^2} + \frac{1}{\sigma_1^2},$$

$$A_{22} = \sum_{j \in G_2} \frac{[T_{2j}(\lambda)]^2}{\sigma_{2j}^2} + \sum_{j \in G_3} \frac{[T_{2'j}(\lambda)]^2}{\sigma_{3j}^2} + \frac{1}{\sigma_2^2},$$

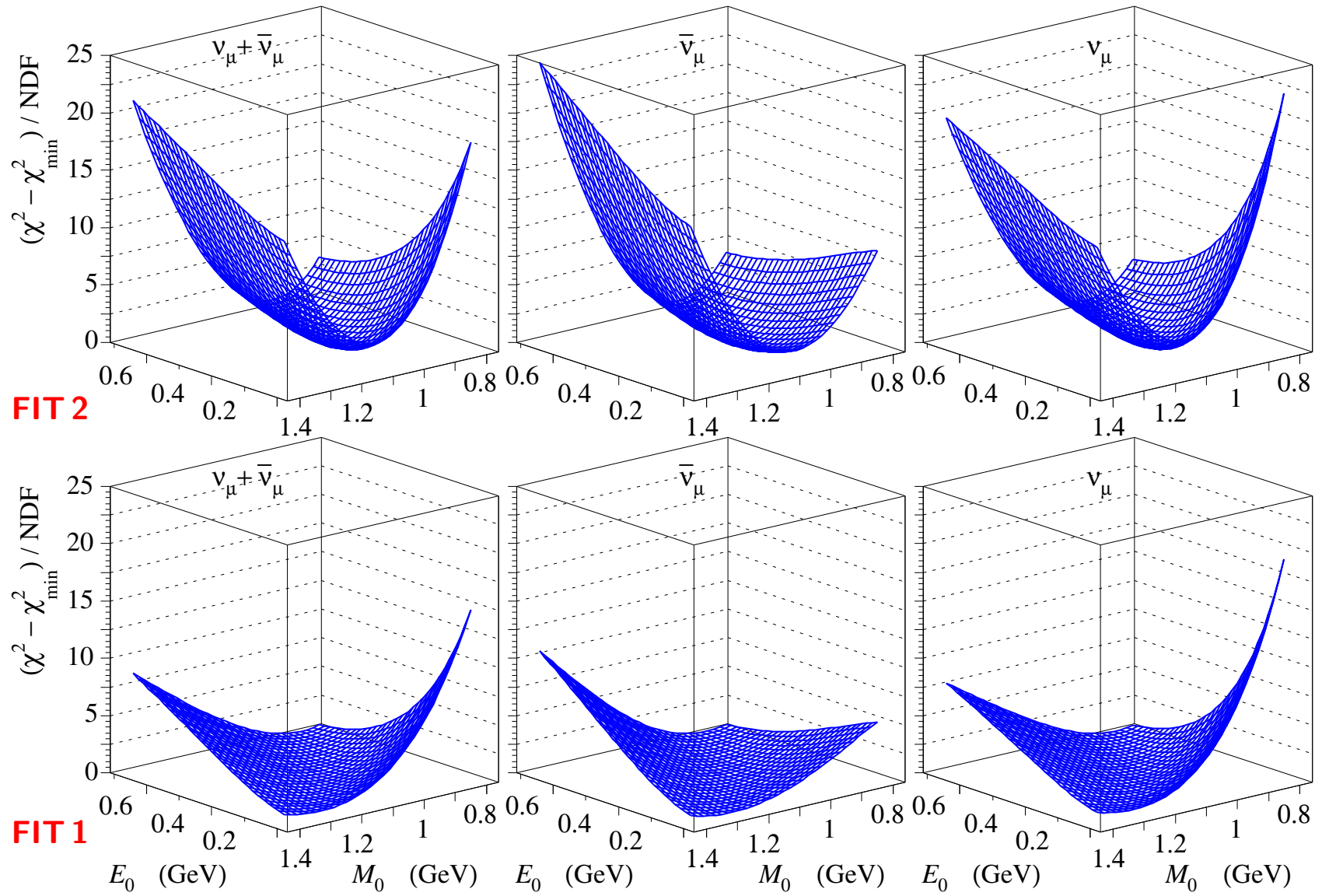
$$A_{12} = \sum_{j \in G_3} \frac{T_{1'j}(\lambda) T_{2'j}(\lambda)}{\sigma_{3j}^2},$$

$$C_i = \sum_{j \in G_2} \frac{T_{ij}(\lambda) E_{ij}(\lambda)}{\sigma_{ij}^2} + \sum_{j \in G_3} \frac{T_{i'j}(\lambda) E_{3j}(\lambda)}{\sigma_{3j}^2} + \frac{1}{\sigma_2^2}.$$

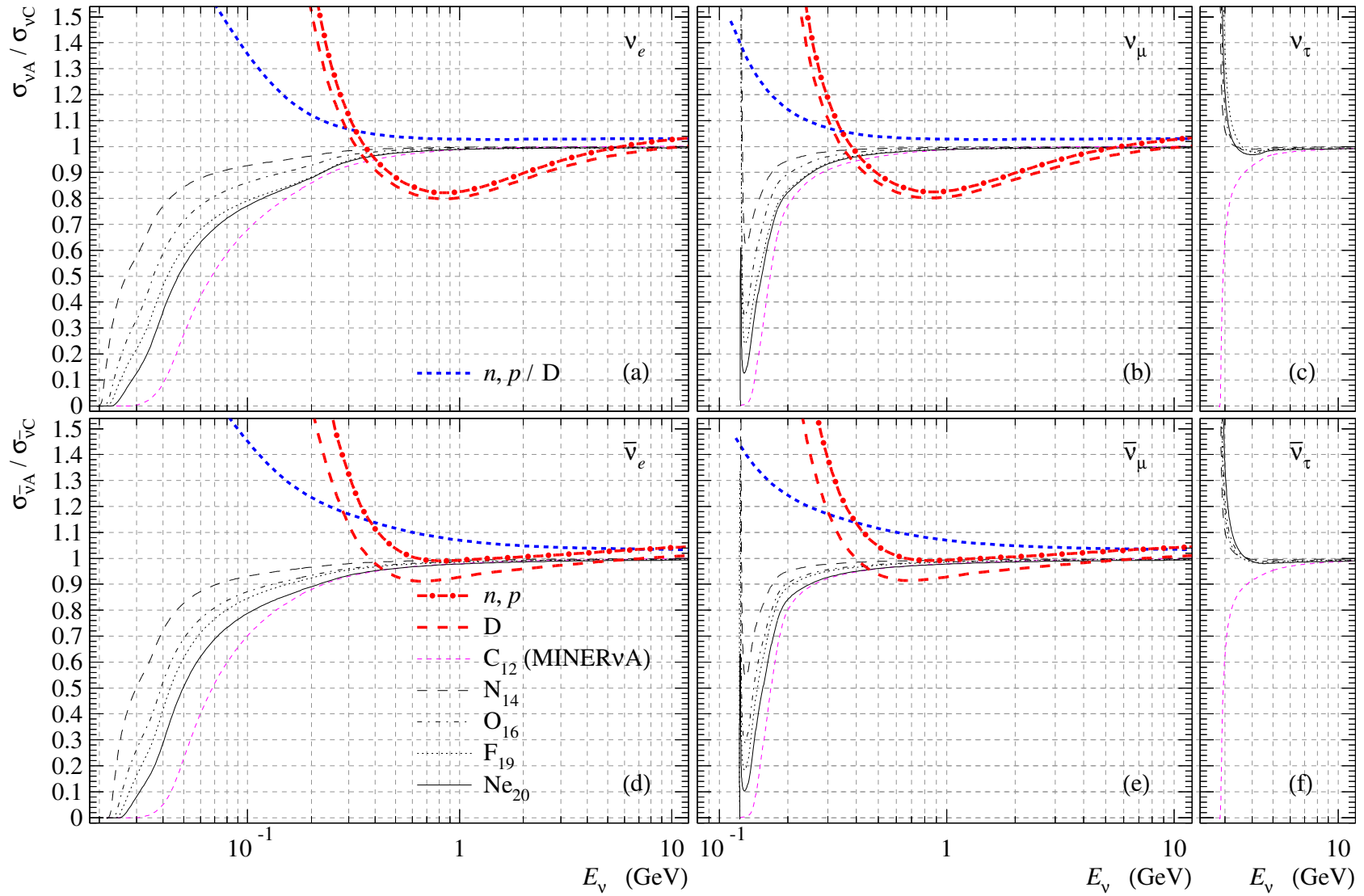


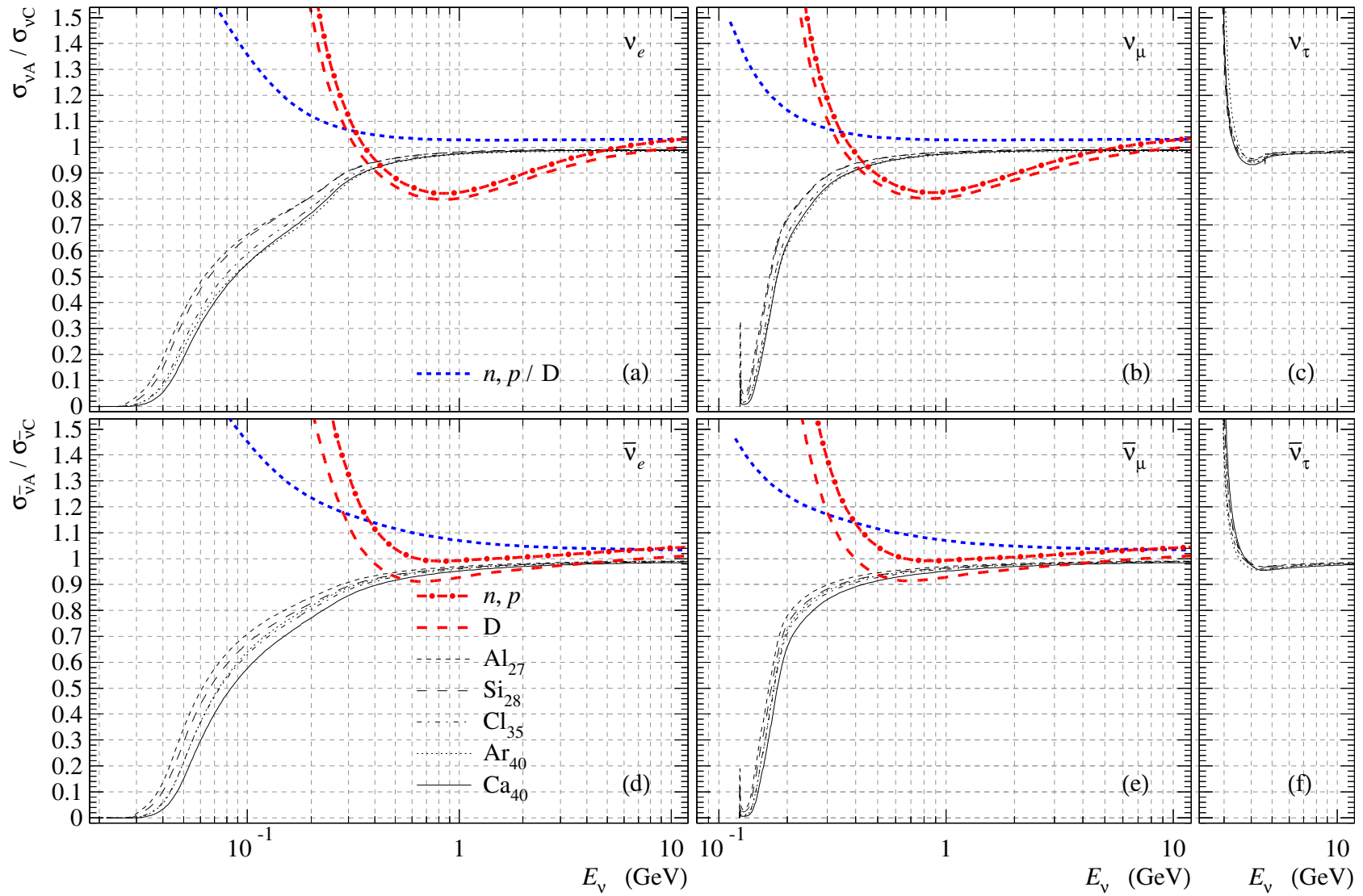
Error contours for the parameters  $M_0$  and  $E_0$  obtained from the global fits for  $\nu$  (a),  $\bar{\nu}$  (b), and  $\nu + \bar{\nu}$  (c) CCQE data, performed *without anchoring of  $M_0$*  (FIT 1, greater areas) and *with anchoring of  $M_0$*  (FIT 2, smaller areas); see slide 68. The inner and outer contours for these fits indicate the 68% and 95% C.L. regions, respectively.

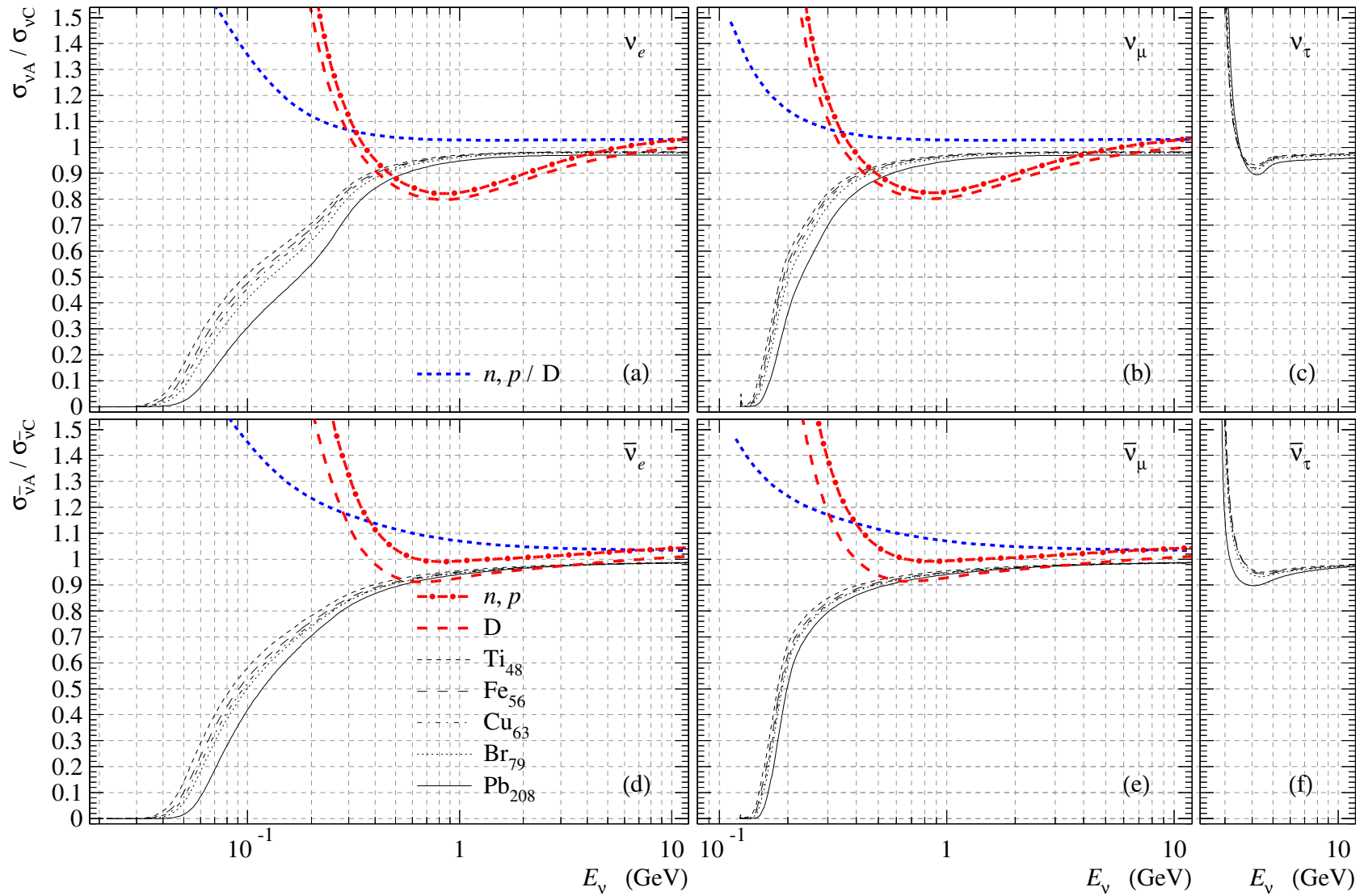
The shapes of  $(\chi^2 - \chi_{\min}^2)/\text{NDF}$  for FIT 1 and FIT 2 are shown in the next slide. ↓

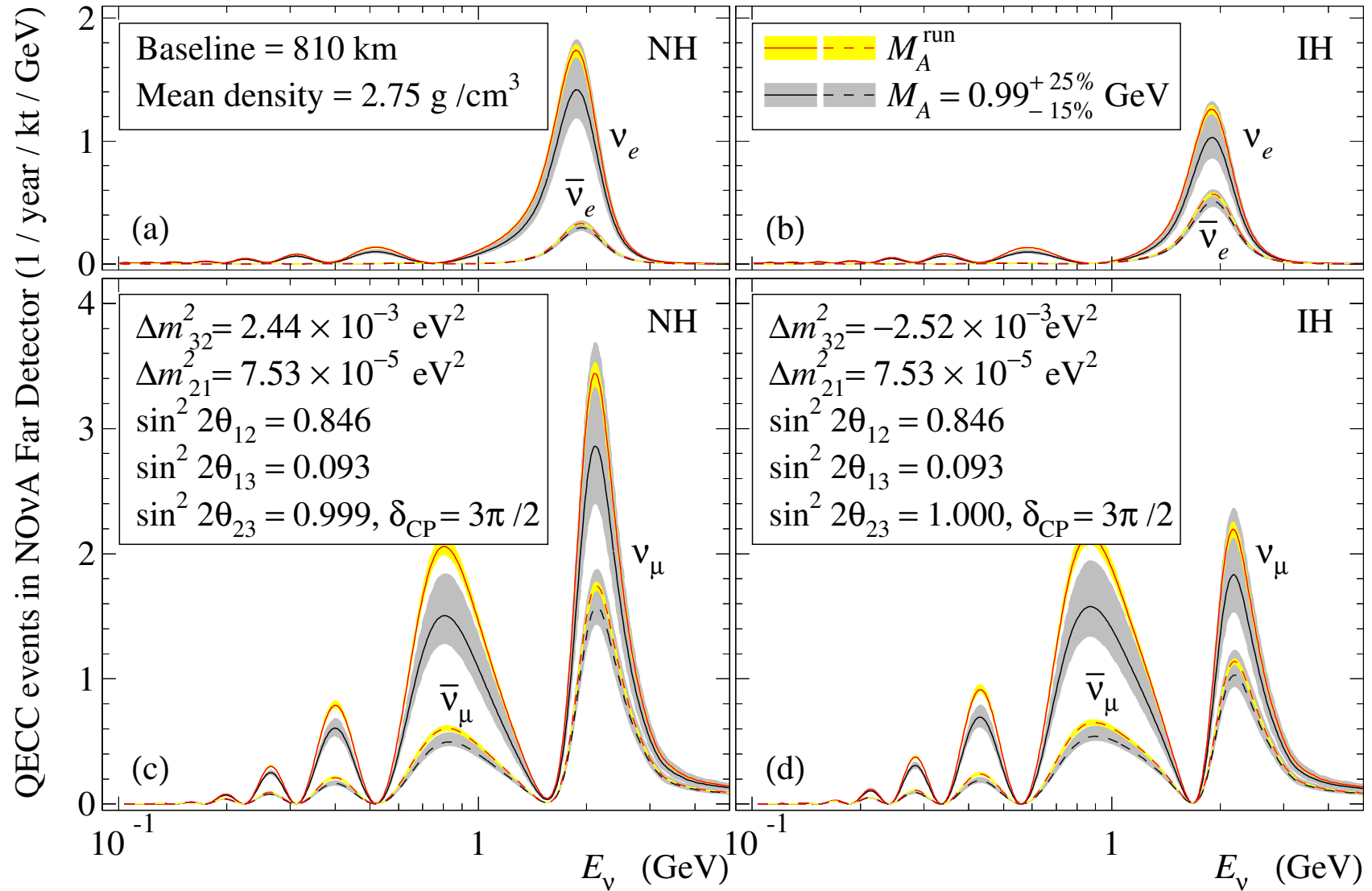




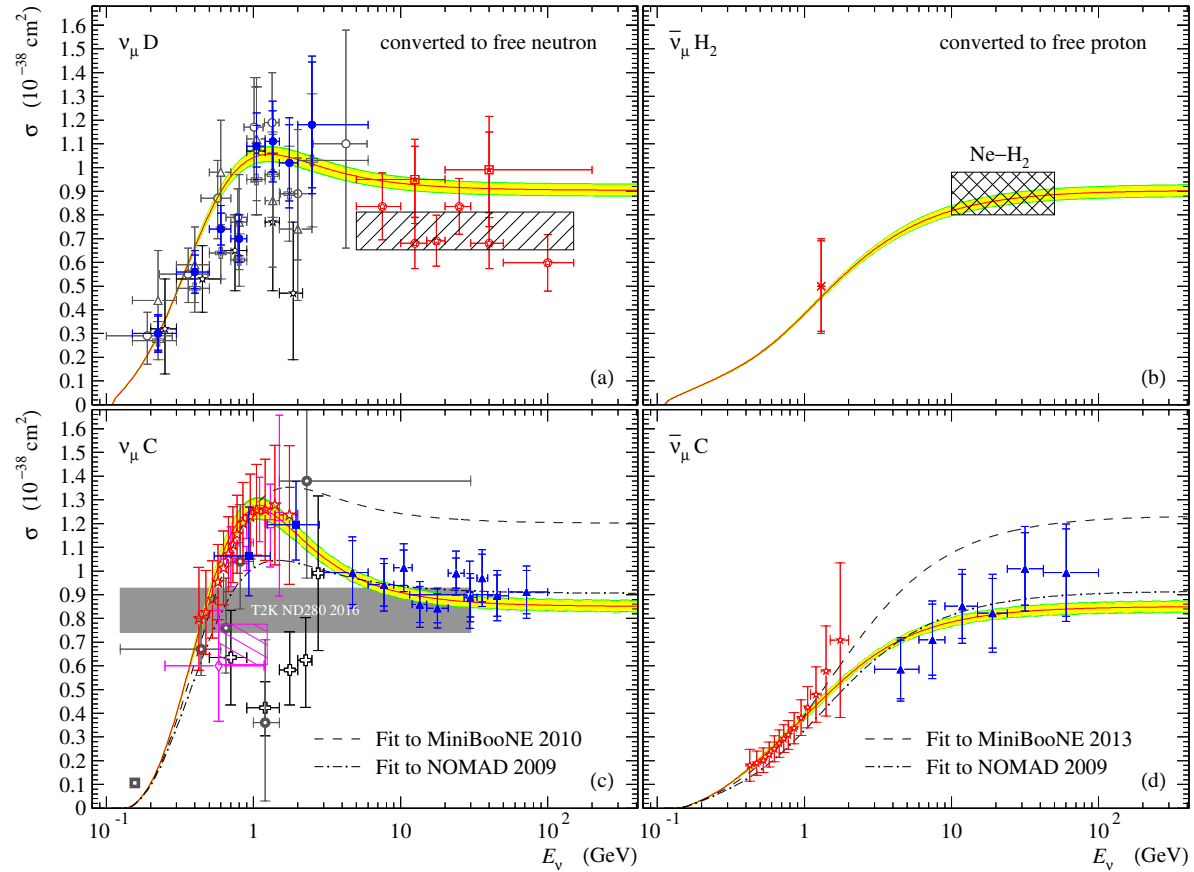




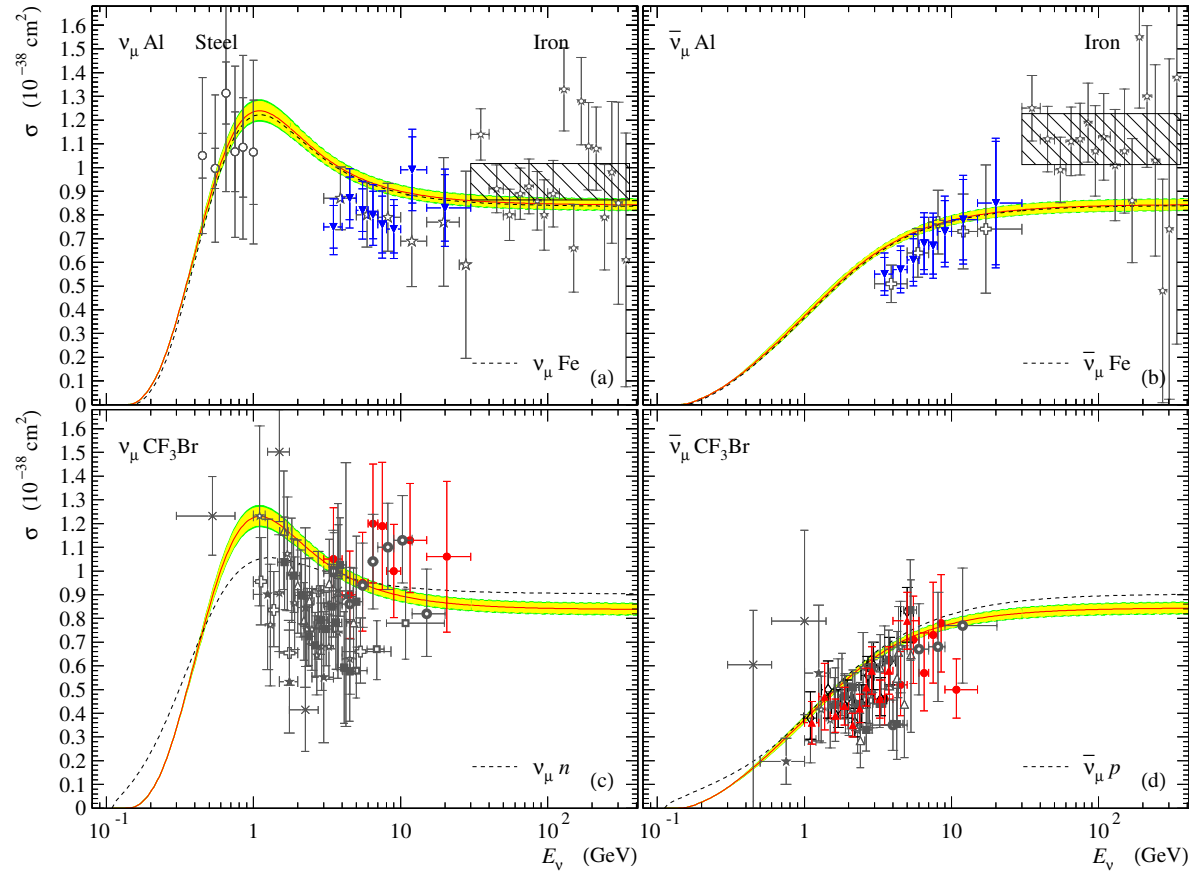




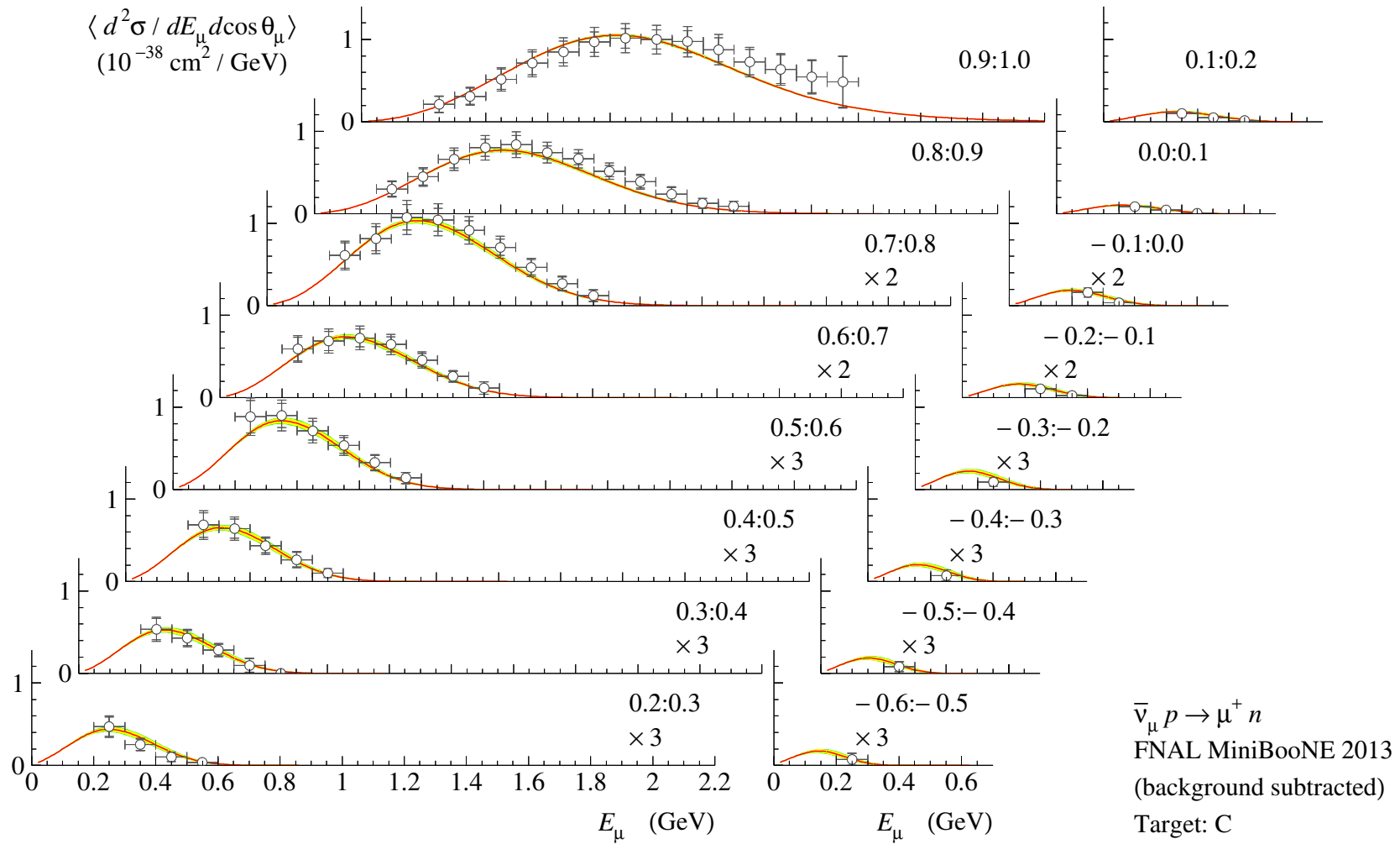
An application to NOνA experiment (see Refs. [70–72] for details).



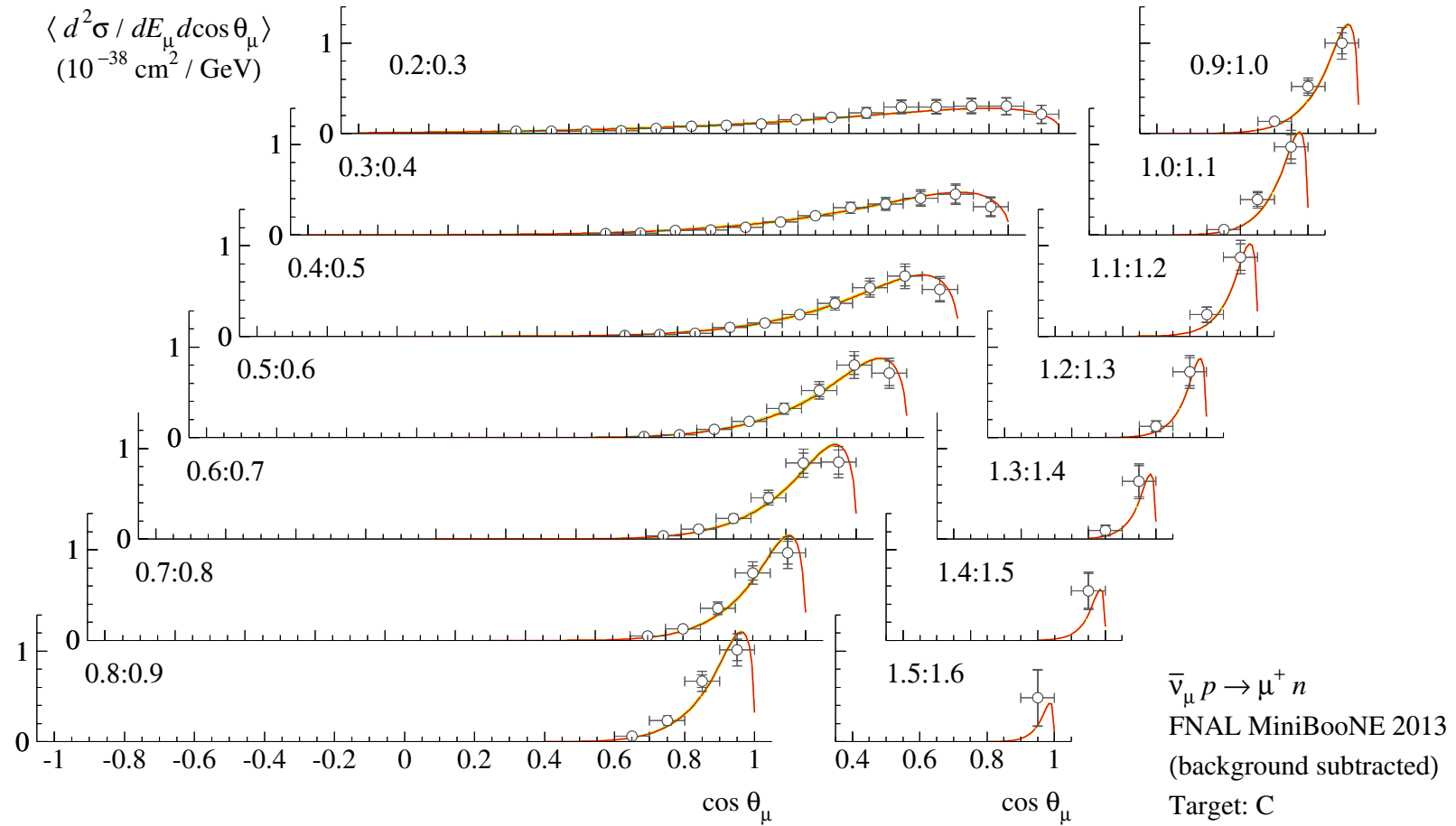
|  |  |
|--|--|
| ☆ Mann et al., ANL 1972 (D)                      | ★ Aguilar-Arevalo et al., FNAL MiniBooNE 2010 ( $\text{CH}_{2.08}$ )   |
| △ Mann et al., ANL 1973 (D)                      | ★ Aguilar-Arevalo et al., FNAL MiniBooNE 2013 ( $\text{CH}_{2.08}$ )   |
| ○ Barish et al., ANL 1975 (D)                    | ◻ Alcaraz-Aunion et al., FNAL SciBooNE 2010 ( $\text{C}_8\text{H}_8$ ) |
| ⊕ Singer et al., ANL 1977 (D)                    | ◊ Betancourt, FNAL NOvA PD 2013 (Carbonaceous targeted)                |
| ● Barish et al., ANL 1977 (D)                    | ⊕ Budagov et al., CERN HLBC 1969 ( $\text{C}_3\text{H}_8$ )            |
| ✱ Fanourakis et al., BNL 1980 ( $\text{H}_2$ )   | ⊗ Allasia et al., CERN BEBC 1990 (D)                                   |
| ■ Auerbach et al., LSND 2002 (C)                 | ⊠ Martinez de la Ossa Romero, CERN LAr TPC 2007 (Ar)                   |
| ⊠ Kitagaki et al., FNAL 1983 (D)                 | ▲ Lyubushkin et al., CERN NOMAD 2009 (carbonaceous target)             |
| ⊠ Asratyan et al., FNAL 1984 ( $\text{Ne-H}_2$ ) | ● Abe, T2K ND280 2014 (carbonaceous target)                            |
|  | ◻ Abe et al., T2K ND280 2013 (carbonaceous target)                     |
|  | ■ Abe et al., T2K INGRID 2015 (carbonaceous target)                    |



- |   |  |
|---|--|
| ○ Kustom et al., ANL 1969 (Steel)   | ▲ Armenise et al., CERN GGM 1979 (C <sub>3</sub> H <sub>8</sub> -CF <sub>3</sub> Br) |
| ⊠ Suwonjandee, FNAL NuTeV 2004 (Fe)   | ⊕ Pohl et al., CERN GGM 1979 (C <sub>3</sub> H <sub>8</sub> -CF <sub>3</sub> Br)     |
| * Franzinetti et al., CERN HLBC 1966 (CF <sub>3</sub> Br)                           | ● Makeev et al., IHEP SKAT 1981 (CF <sub>3</sub> Br)*                                |
| ★ Young, CERN HLBC 1967 (CF <sub>3</sub> Br)  | ● Grabosch et al., IHEP SKAT 1988 (CF <sub>3</sub> Br)                               |
| ☆ Eichten et al., CERN GGM 1973 (CF <sub>3</sub> Br)                                | ● Brunner et al., IHEP SKAT 1990 (CF <sub>3</sub> Br)*                               |
| △ Rollier et al., CERN GGM 1975 (CF <sub>3</sub> Br)                                | ⊕ Belikov et al., IHEP-ITEP 1981 (Al)  |
| ■ Bonetti et al., CERN GGM 1977 (CF <sub>3</sub> Br)                                | ☆ Belikov et al., IHEP-ITEP 1982 (Al)  |
| ◇ Rollier et al., CERN GGM 1978 (C <sub>3</sub> H <sub>8</sub> -CF <sub>3</sub> Br) | ▼ Belikov et al., IHEP-ITEP 1985 (Al)  |
- \* converted to free nucleons



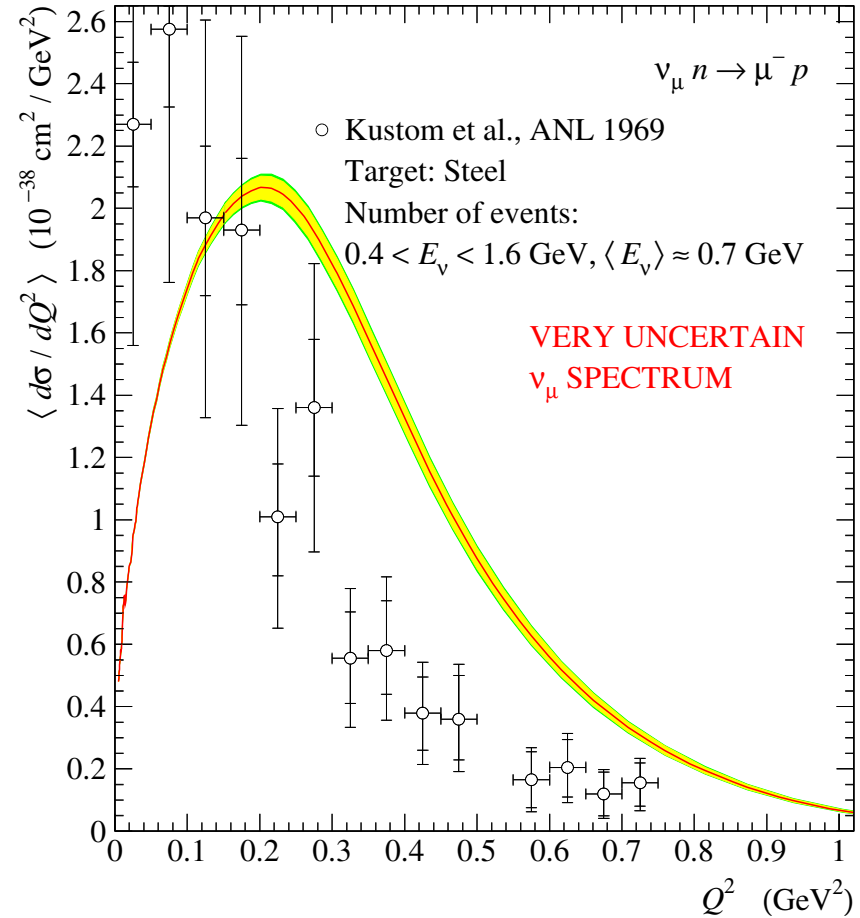
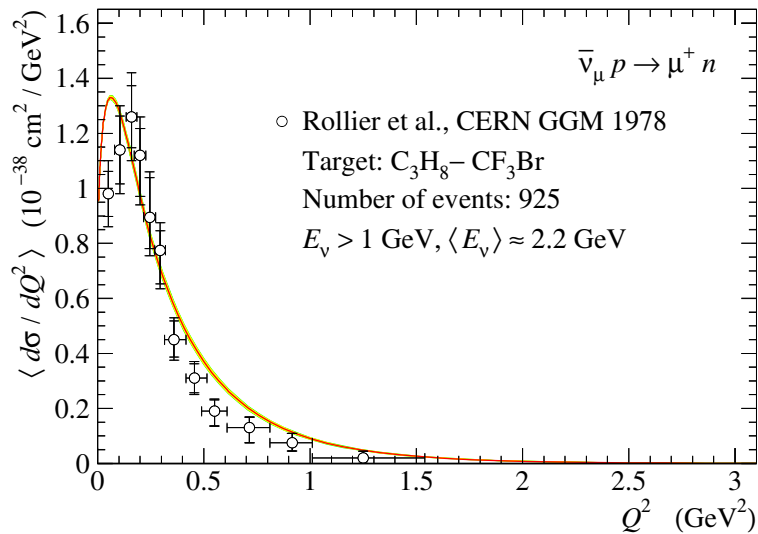
See note in slide 18.

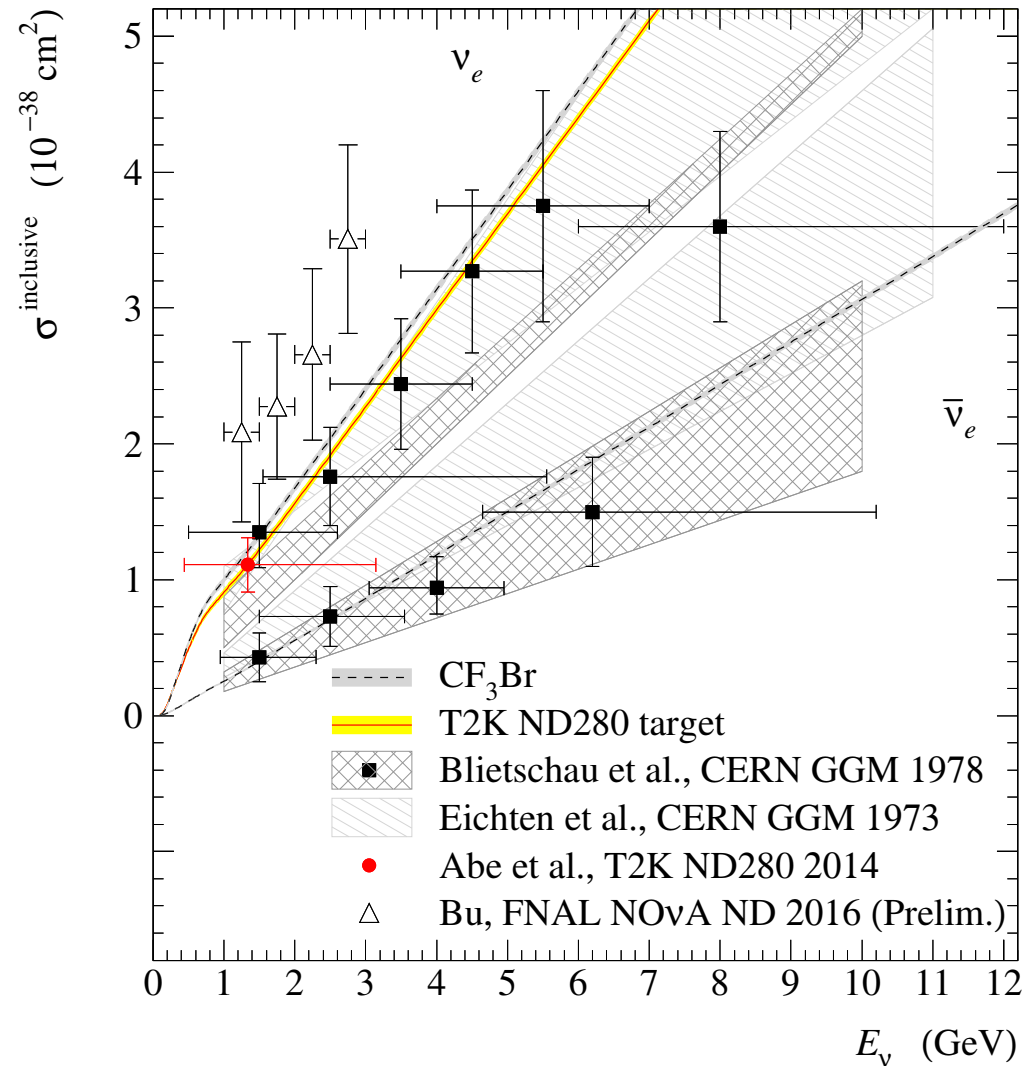


See note in slide 18.



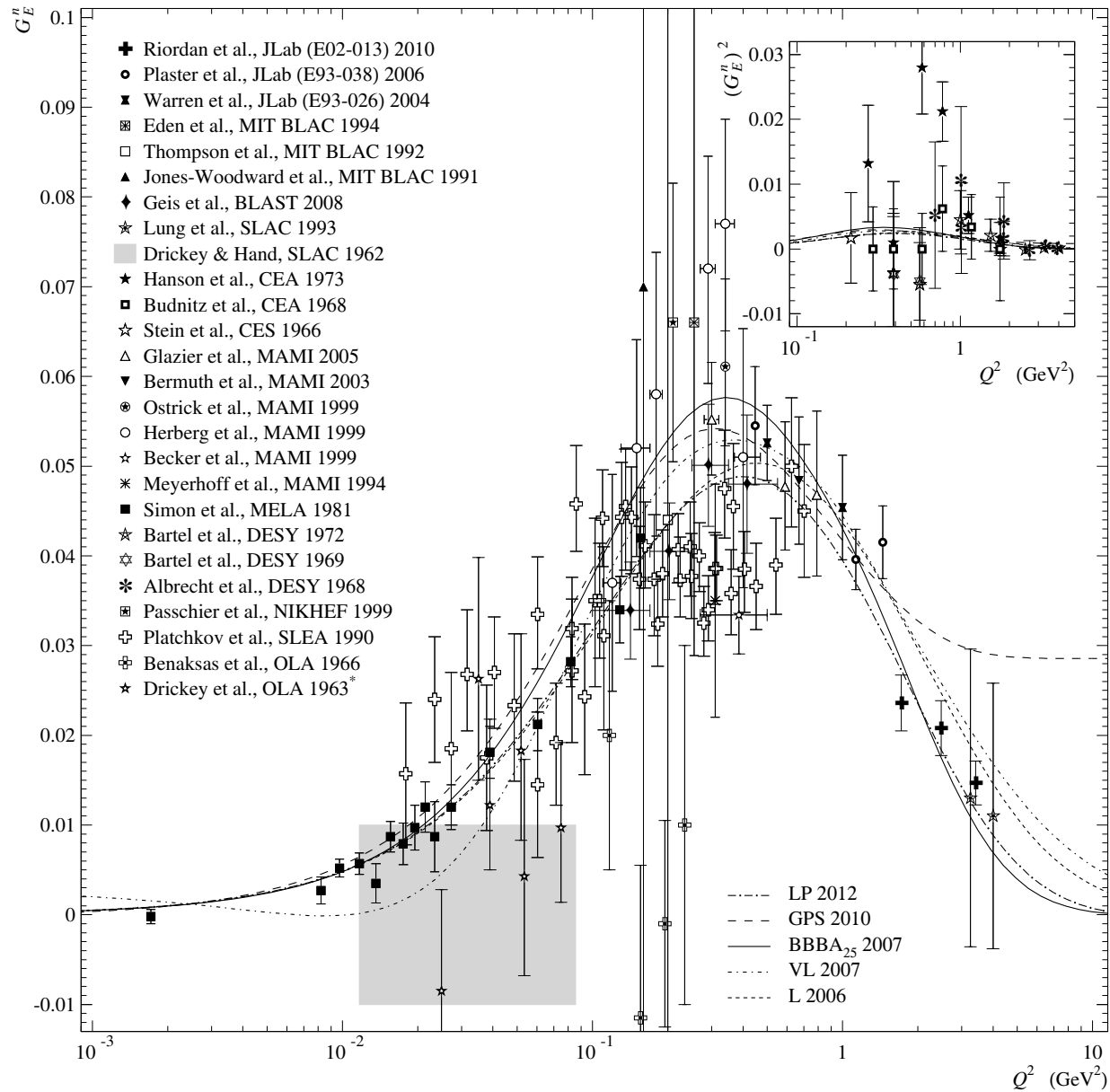
# A fly in the ointment (examples of doubtful agreement)

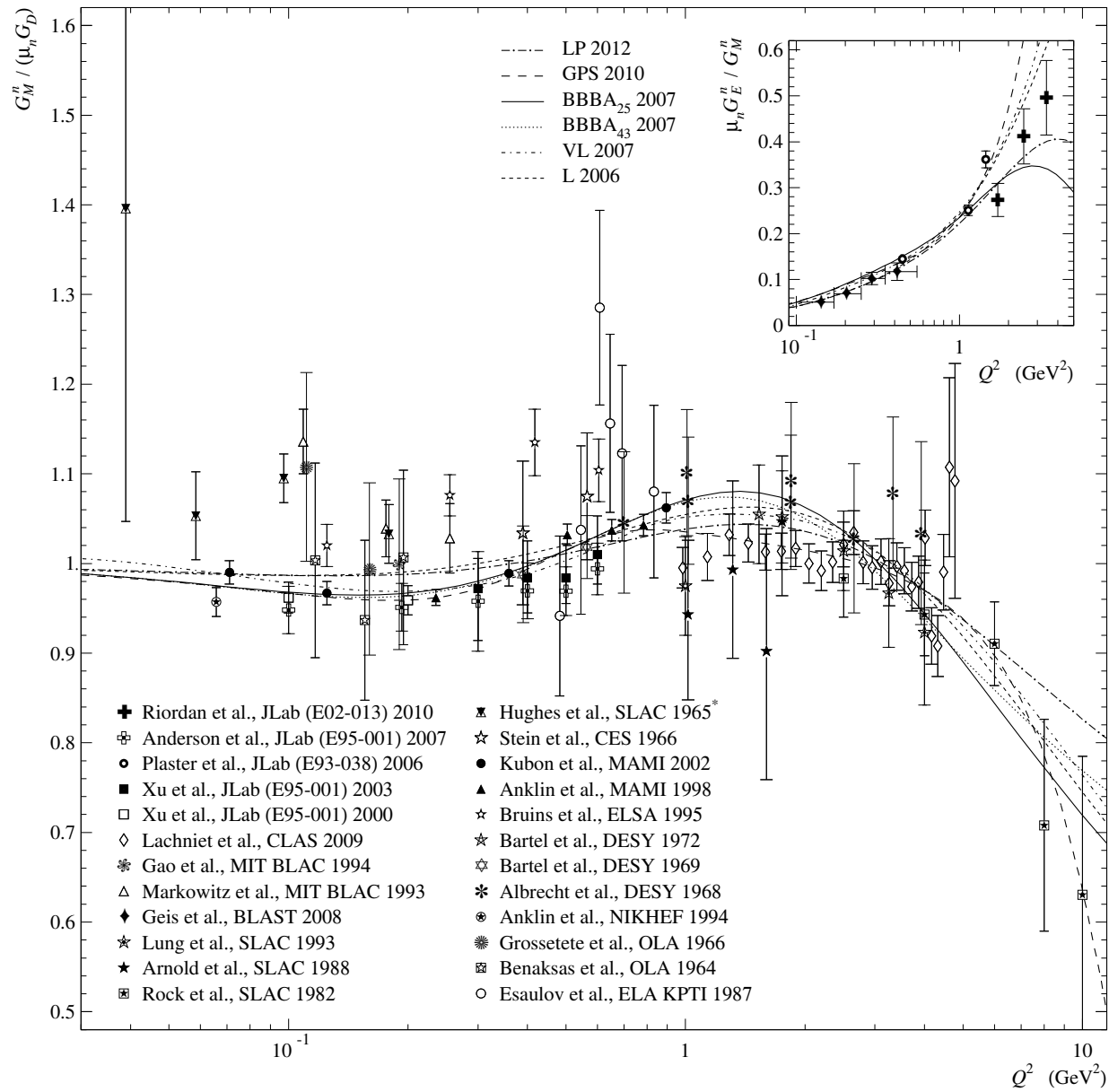


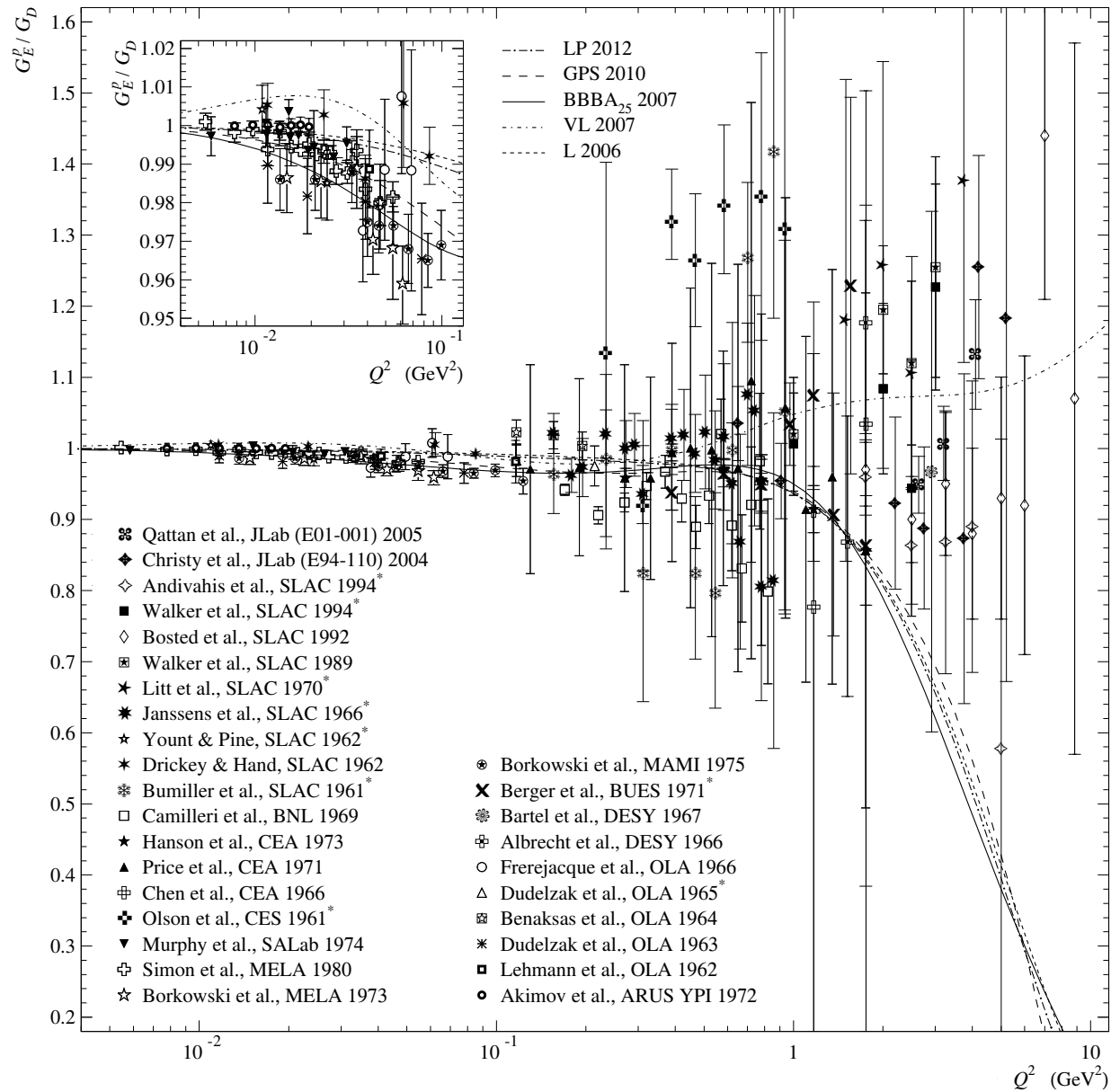


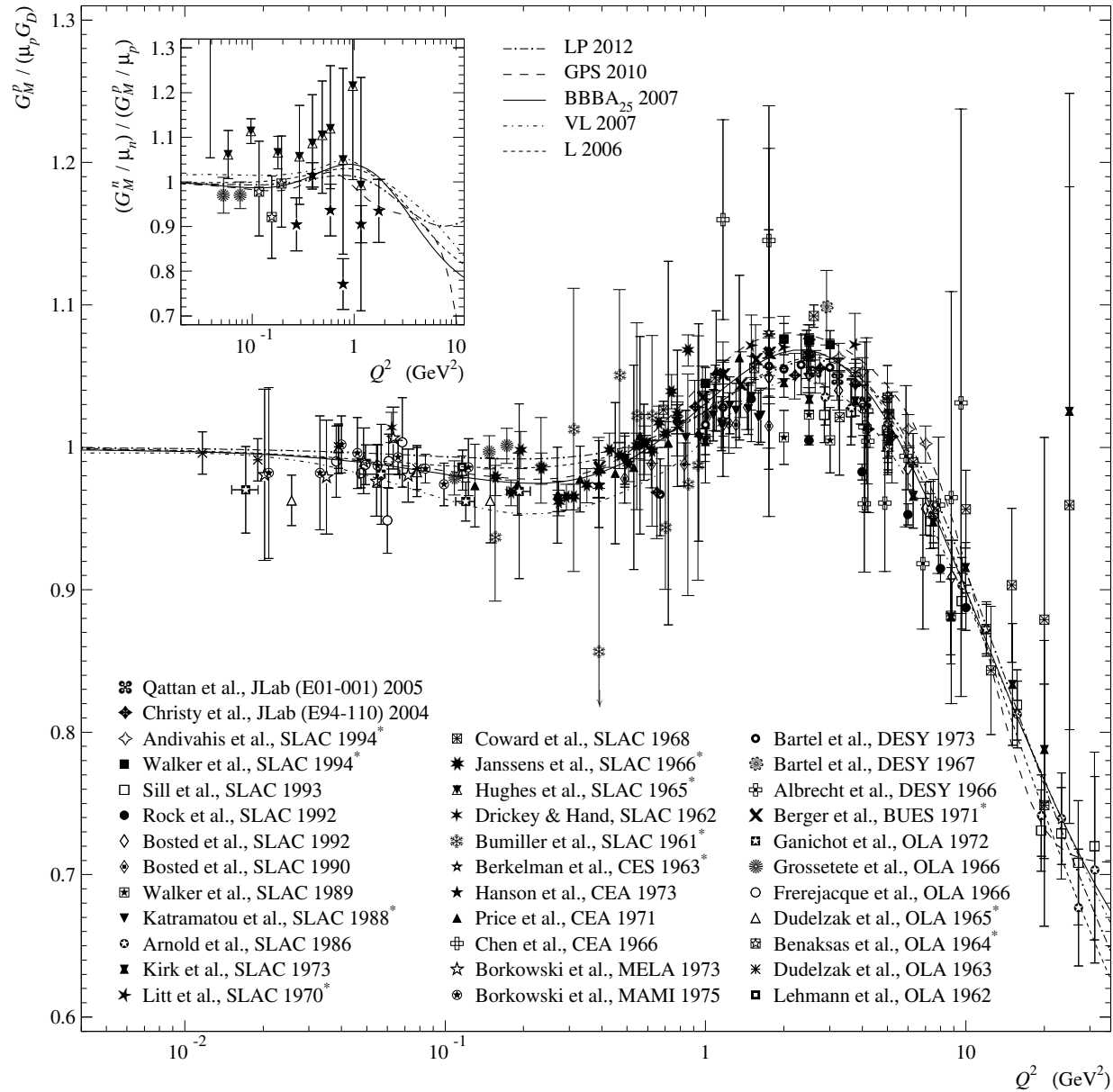
Total inclusive CC cross sections for  $\nu_e$  and  $\bar{\nu}_e$  at low energies with a bit surprising (*preliminary*) data from NO $\nu$ A Near Detector (X. Bu, arXiv:1601.01213 [hep-ex]).

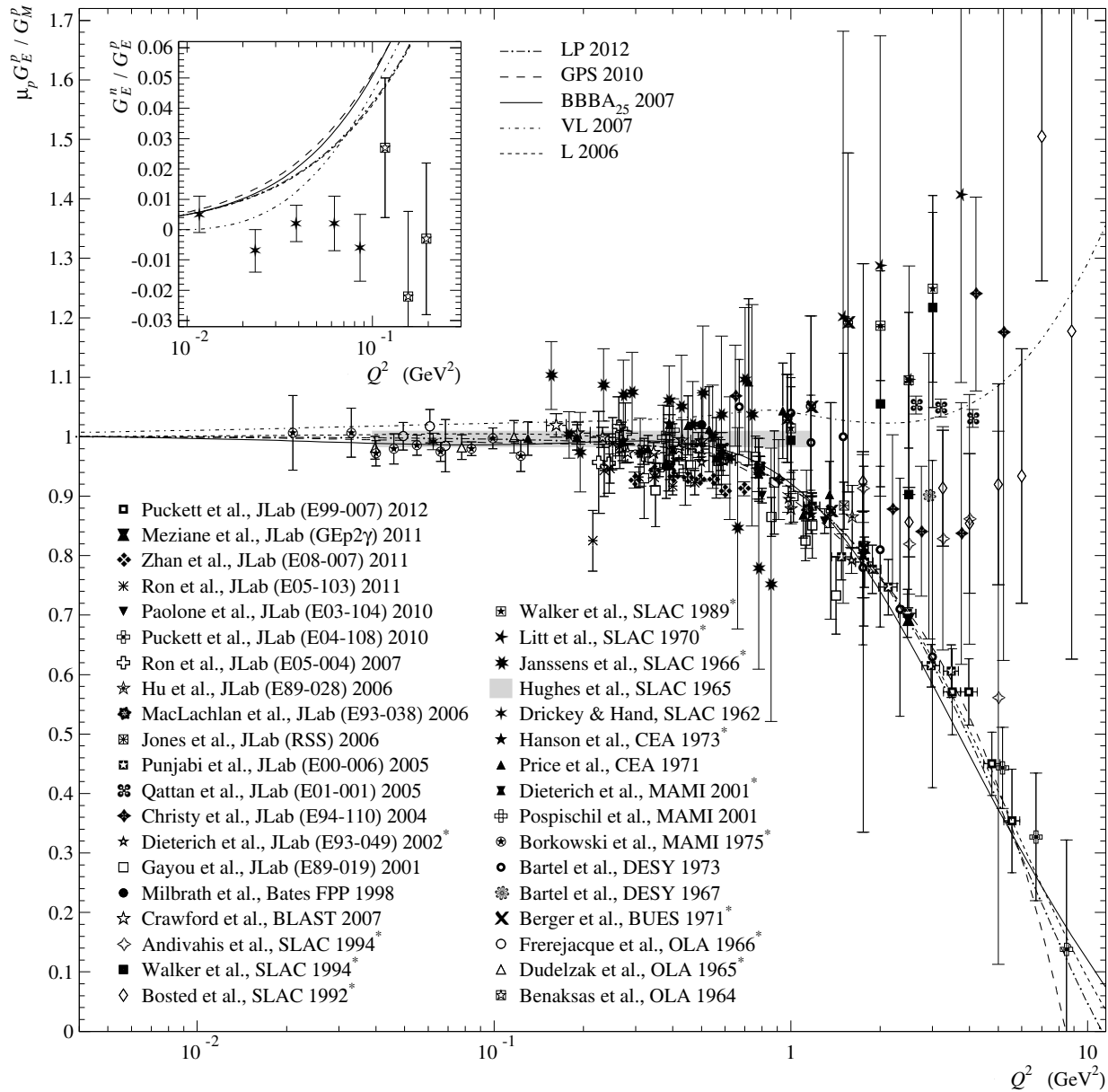
# Comparison of several models for the nucleon electromagnetic form factors with data













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