

# Relativistic Heavy Ion Collisions

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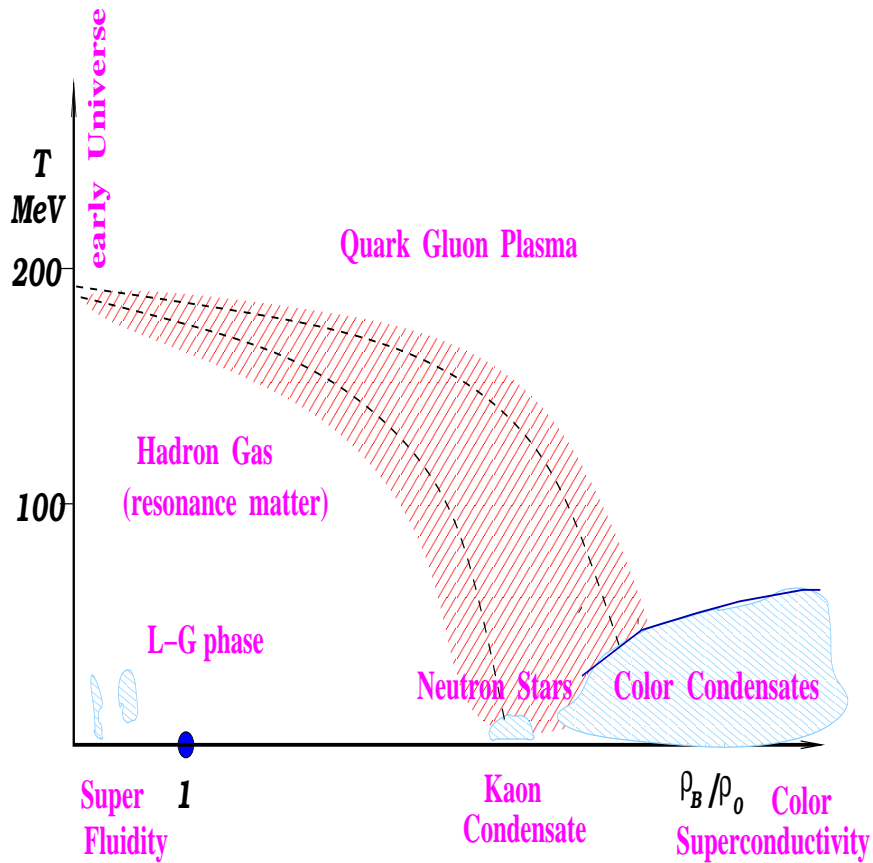
in collaboration with Yu.B. Ivanov<sup>2,3</sup> and V.N. Russkikh<sup>2,3</sup>

*(1)JINR, Dubna; (2)GSI, Darmstadt; (3)Kurchatov Inst., Moscow*

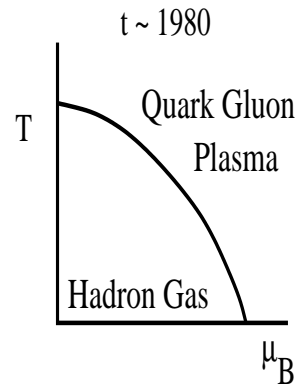
## CONTENT

- ♠ Introductory Remarks
- ♠ Three-Fluid Hydrodynamics: Hadronic Scenario
- ♠ How Well This Model Works
- ♠ Global Space-Time Evolution and Phase Diagram
- ♠ Summary and Outlook
- ♠ Dilepton Production in Expansion Dynamics

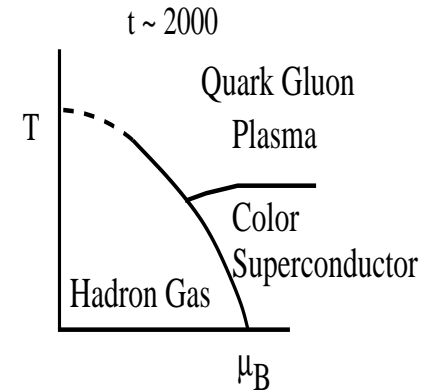
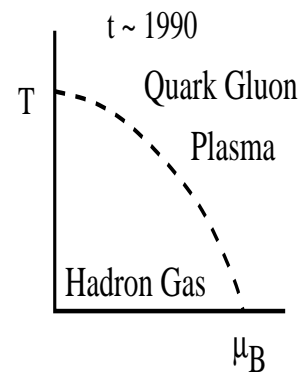
# Motivation (Phase diagram)



## The Evolving QCD Phase Transition



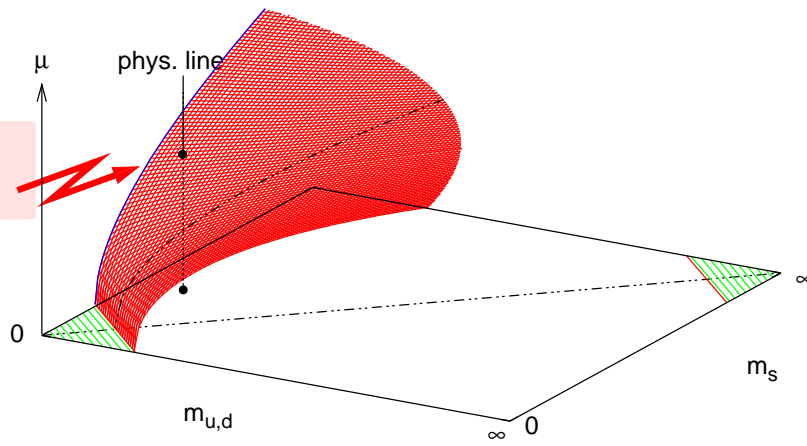
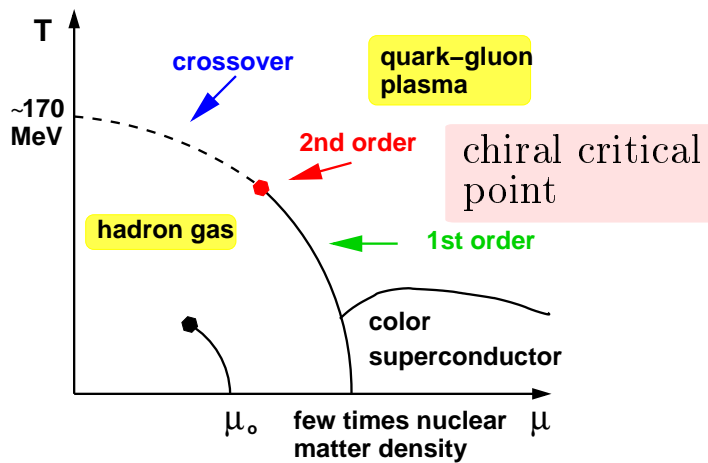
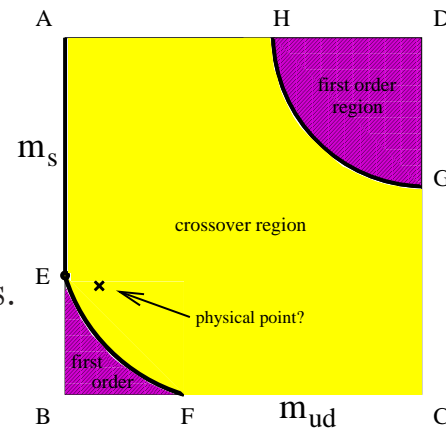
Critical Temperature 150 - 200 MeV ( $\mu_B = 0$ )  
 Critical Density 1/2-2 Baryons/Fm<sup>3</sup> ( $T = 0$ )



# Motivation (The order of phase transition)

$$\mu_B = 0$$

Z.Fodor, Nucl. Phys.  
A271 (2003) 319



# Motivation (Why 3-fluid hydro ?)

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♠ Conservation laws (Gauss theorem)  $\Rightarrow$  Fluid dynamics

$$\partial_\mu J^\mu = 0 \quad \text{net charge conservation} \quad \boxed{4}$$

$$\partial_\mu T^{\mu\nu} = 0 \quad \text{energy momentum conservation} \quad \boxed{10}$$

♠ Tensor decomposition of the charge current  $J^\mu$  and energy-momentum tensor  $T^{\mu\nu}$  with respect to 4-velocity  $u^\mu$

$$\begin{aligned} J_i^\mu &= n_i u^\mu + \dots \\ T^{\mu\nu} &= \underbrace{\varepsilon u^\mu u^\nu - P (g^{\mu\nu} - u^\mu u^\nu)}_{\text{perfect hydro}} + \dots \end{aligned} \quad \boxed{6}$$

- Perfect hydro in local thermodynamical equilibrium + EoS
- First order dissipative corrections (viscosity, heat capacity)  $\Rightarrow$  acausality
- Second order corrections  $\Rightarrow$  + 14 Grad equations

*Spatial-temporal variation of the macro fields have to be SMALL*

# Multi-fluid approximation

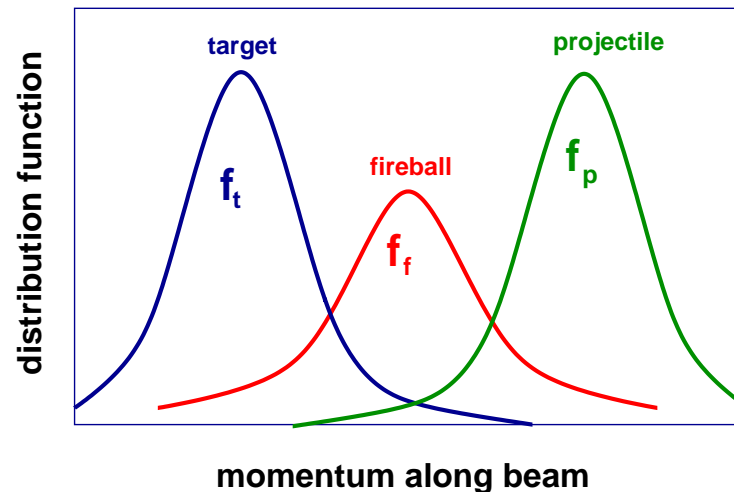
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## ♠ Multi-fluid dynamics

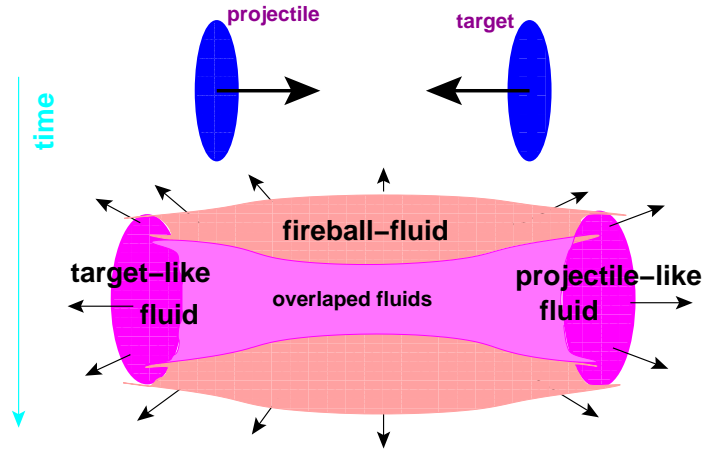
(starting from the relativistic Boltzmann equation to that for moments of distribution function)

$$f(x, p) = \sum_j^M f_j^{(eq)}(x, p)$$

A single fluid may consist of several particle species. Different fluids may be of the same particle species.



# 3-fluid hydro equations



TARGET-LIKE FLUID:  $\partial_\mu J_t^\mu = 0$   $\partial_\mu T_t^{\mu\nu} = -F_{tp}^\nu + F_{ft}^\nu$   
 Lead. particles carry b-charge exchange/emission

PROJECTILE-LIKE FLUID:  $\partial_\mu J_p^\mu = 0,$   $\partial_\mu T_p^{\mu\nu} = -F_{pt}^\nu + F_{fp}^\nu$

FIREBALL FLUID:  $J_f^\mu = 0,$   $\partial_\mu T_f^{\mu\nu} = F_{pt}^\nu + F_{tp}^\nu - F_{fp}^\nu - F_{ft}^\nu$   
 Baryon-free fluid Source term Exchange

The **source term** is delayed due to a formation time  $\tau \sim 1 \text{ fm}/c$

Total energy-momentum conservation:  $\partial_\mu (T_p^{\mu\nu} + T_t^{\mu\nu} + T_f^{\mu\nu}) = 0$

# Friction

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PROJECTIVE-TARGET FRICTION:

$$F_{pt}^\nu = \rho_p \rho_t \left[ (u_p^\nu - u_t^\nu) D_P + (u_p^\nu + u_t^\nu) D_E \right] \xi_h^2(s_{pt})$$

heating

fireball production

EoS dependent

enhancement factor  $\xi_h^2(s_{pt})$

$\rho_\alpha$  - scalar density of  $\alpha$  fluid;  $D_{P/E} = m_N V_{rel}^{pt} \sigma_{P/E}(s_{pt})$ ;  $m_N$  - nucleon mass;

$s_{pt} = m_N^2 (u_p^\nu + u_t^\nu)^2$  - mean invariant energy squared of colliding p- and t-fluids;

$V_{rel}^{pt} = [s_{pt}(s_{pt} - 4m_N^2)]^{1/2} / 2m_N^2$  - mean relative velocity of the p- and t-fluids;

$\sigma_{P/E}(s_{pt})$  - hadron-hadron cross sections integrated with certain weights

$V_{rel}^{pt} < \text{thermal or Fermi velocity}$   $\Rightarrow$  Unification of p and t fluids (equilibrium)

PROJECTIVE(TARGET)-FIREBALL FRICTION:

Absorption of a fireball matter by baryon-rich fluids

$$F_{fp}^\nu = \rho_p \frac{T_f^{(eq)0\nu}}{u_f^0} D_{fp} \quad \text{where} \quad D_{fp} = V_{rel}^{fp} \sigma_{tot}^{N\pi \rightarrow R}(s_{fp}).$$

# Hadronic equation of state

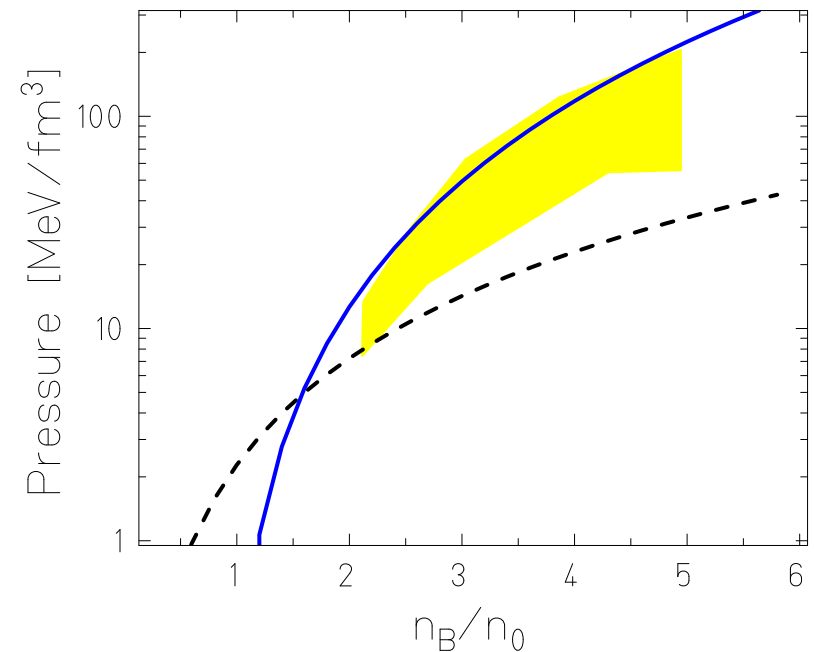
Energy density: 
$$\varepsilon(n_B, T) = \underbrace{\varepsilon_{gas}(n_B, T)}_{\text{gas of free hadrons}} + \underbrace{W(n_B)}_{\text{mean field}}$$

Pressure: 
$$P(n_B, T) = \underbrace{P_{gas}(n_B, T)}_{\text{gas of free hadrons}} + \underbrace{n_B \frac{dW(n_B)}{dn_B} - W}_{\text{mean field}}$$

$T = 0$

$$W(n_B) = n_B m_N \left[ -b \left( \frac{n_B}{n_0} \right) + c \left( \frac{n_B}{n_0} \right)^{\gamma+1} \right]$$

$W(n_B)$  saturates the cold nuclear matter at  $n_0 = 0.15 \text{ fm}^{-3}$  and  $\varepsilon(n_0, T = 0)/n_0 - m_N = -16 \text{ MeV}$ , and provides incompressibility of nuclear matter  $K = 235 \text{ MeV}$ . (P. Danielewicz et al., Science **298**, 1592 (2002))



To preserve causality at high  $n_B$  
$$\varepsilon(n_B, T = 0) = n_0 m_N \left[ A \left( \frac{n_B}{n_0} \right)^2 + C + B \left( \frac{n_B}{n_0} \right)^{-1} \right], \quad n_B > n_c \approx 6n_0$$

Parameters are determined on the condition that  $\varepsilon(n_B, T = 0)$  and its two first derivatives are continuous at  $n_c$ .



# Freeze-out

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- Criteria:  $\underbrace{\textit{Local}}_{\text{(at } x \text{ position)}}$   $\underbrace{\textit{proper}}_{\text{(in local rest frame)}}$   $\underbrace{\textit{energy density of matter}}_{\text{(summed over all fluids)}}$  is less than  $\boxed{\varepsilon_{frz}}$
- Shock-like fr.:  $T_{hydro}$  and  $\mu_{hydro}$  are mapped to  $T_{gas}$  and  $\mu_{gas}$  proceeding from exact baryon, energy and momentum conservation.
- Freeze-out *a la* Milekhin  $E \frac{dN}{d^3p} = \int f_{gas}(x, p) p^\mu d\sigma_\mu$ ,  $d\sigma_\mu = u_\mu(d^3x)_{proper}$   
 $u_\mu = \text{hydro 4-velocity}$  (proper = in the frame, where  $u_\mu = (1, 0, 0, 0)$ )
  - In “space-like regions” it is very similar to Cooper-Frye
  - In “time-like regions” there is no problem with energy conservation, because  $P = 0$  on the system boundary
  - In fact, there is no “time-like freeze-out” in the code. Only tiny fireballs are frozen out.
  - Problem of shadowing still persists
  - Further study of Freeze-out is needed !

# Kinematical variables

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## Notation

Longitudinal rapidity

$$y = \frac{1}{2} \ln \frac{E + p_{\parallel}}{E - p_{\parallel}}$$

♠ additive at the Lorentz transformation

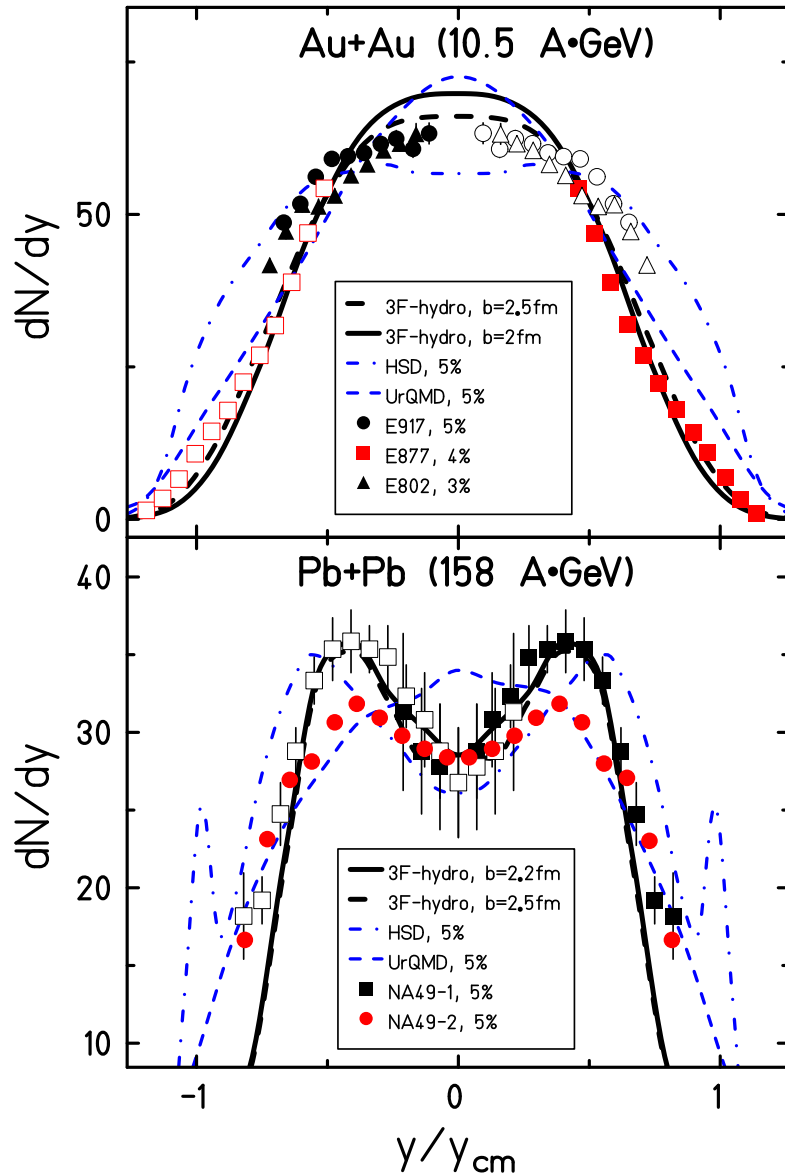
♠ ultrarelativistic limit: if  $E \rightarrow p$  ,  $y \rightarrow \ln \tan \theta/2$

♠ Invariant cross section

$$E \frac{d\sigma}{d^3p} = \frac{d\sigma}{2\pi p_{\perp} dp_{\perp} dy} = \frac{d\sigma}{2\pi m_{\perp} dm_{\perp} dy}$$

with the transverse mass :  $m_{\perp} = \sqrt{m^2 + p_{\perp}^2}$

# Proton and $(p - \bar{p})$ rapidity distributions



## 3-Fluids: hadronic gas EoS

$b = 2\text{ fm}$  for Au+Au(10 AGeV) and  $b = 2.2\text{ fm}$  for Pb+Pb(158 AGeV) are experimental estimates.

Transport models: H. Weber, E.L. Bratkovskaya, W. Cassing and H. Stöcker, Phys. Rev. **C67** (2003) 014904

E802: Phys. Rev. **C60** (1999) 064901

E877: Phys. Rev. **C62** (2000) 024901

E917: Phys. Rev. Lett. **86** (2001) 1970

NA49-1: Phys. Rev. Lett. **82** (1999) 2471

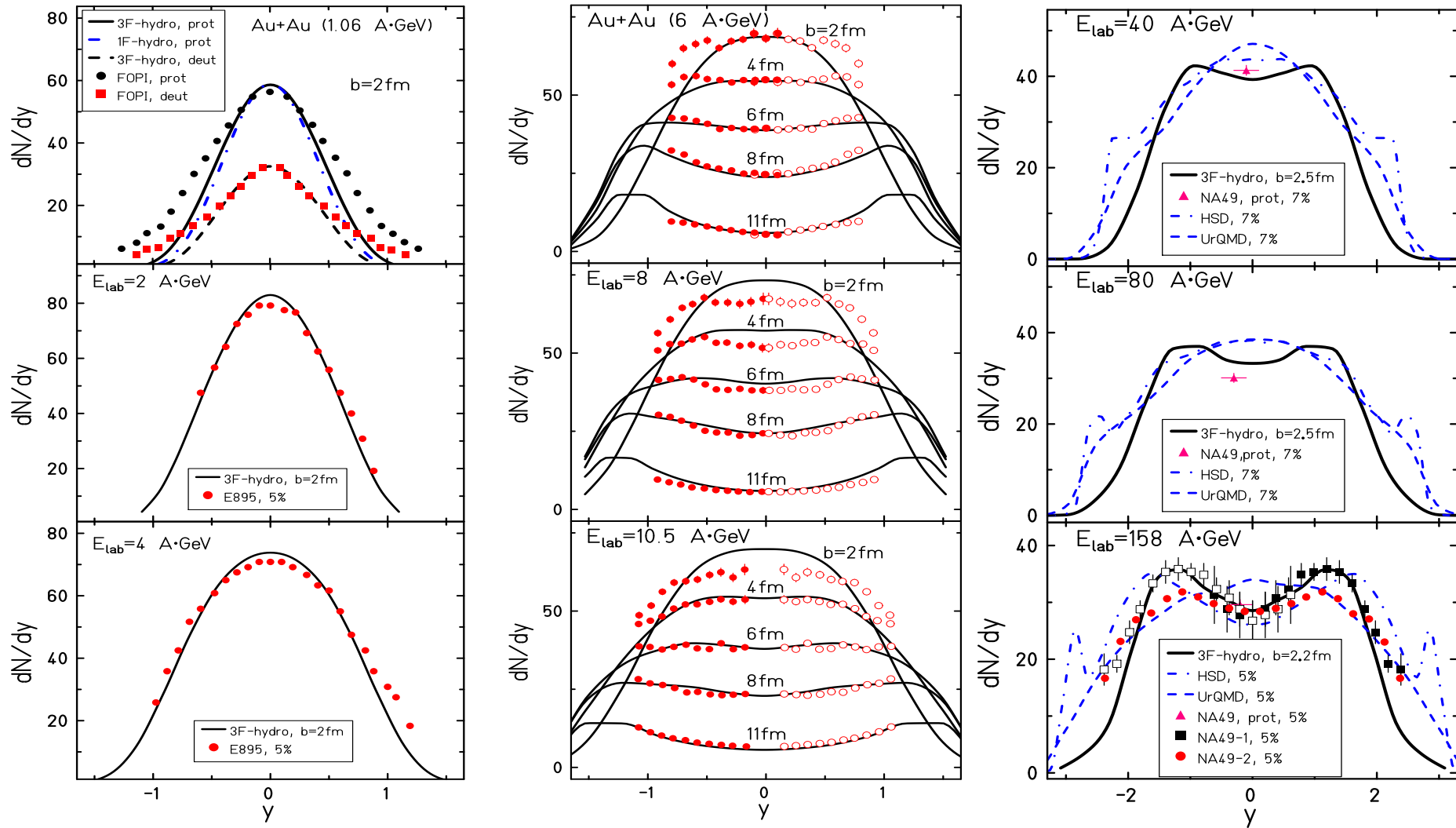
NA49-2(preiminary): Nucl. Phys. **A661** (1999) 362c

# Proton and $(p - \bar{p})$ rapidity distributions

SIS

AGS

SPS



FOPI: N.Herrmann, Nucl. Phys. **A610** (1996) 49c [Au(1.06 GeV/nucleon)+Au]

E895: Phys. Rev. **C68** (2003) 054905 [Au(2 and 4 GeV/nucleon)+Au]

E917: Phys. Rev. Lett. **86** (2001) 1970 [Au(6, 8 and 10.5 GeV/nucleon)+Au]

NA49 (prot.): Phys. Rev. **C69** (2004) 024902

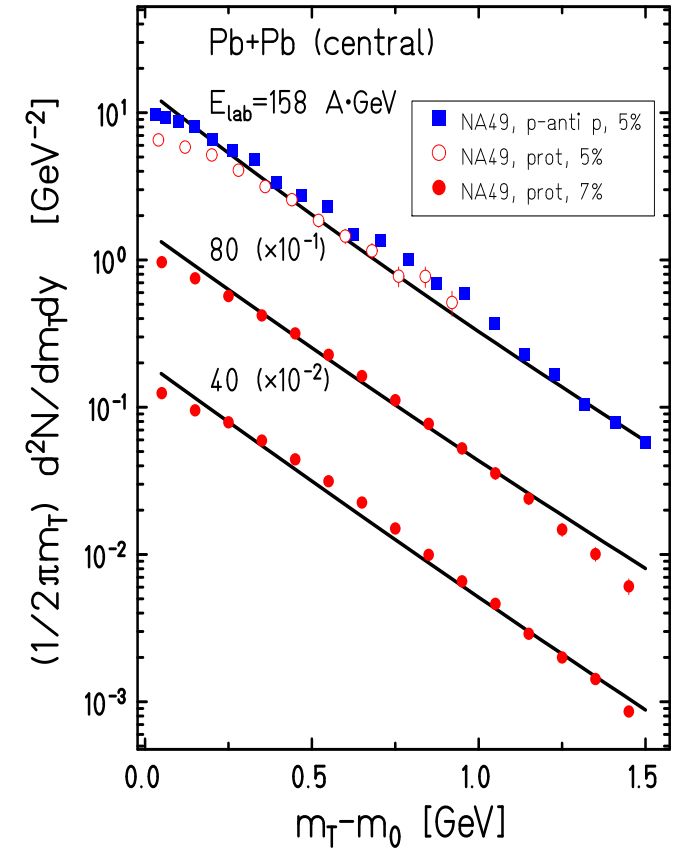
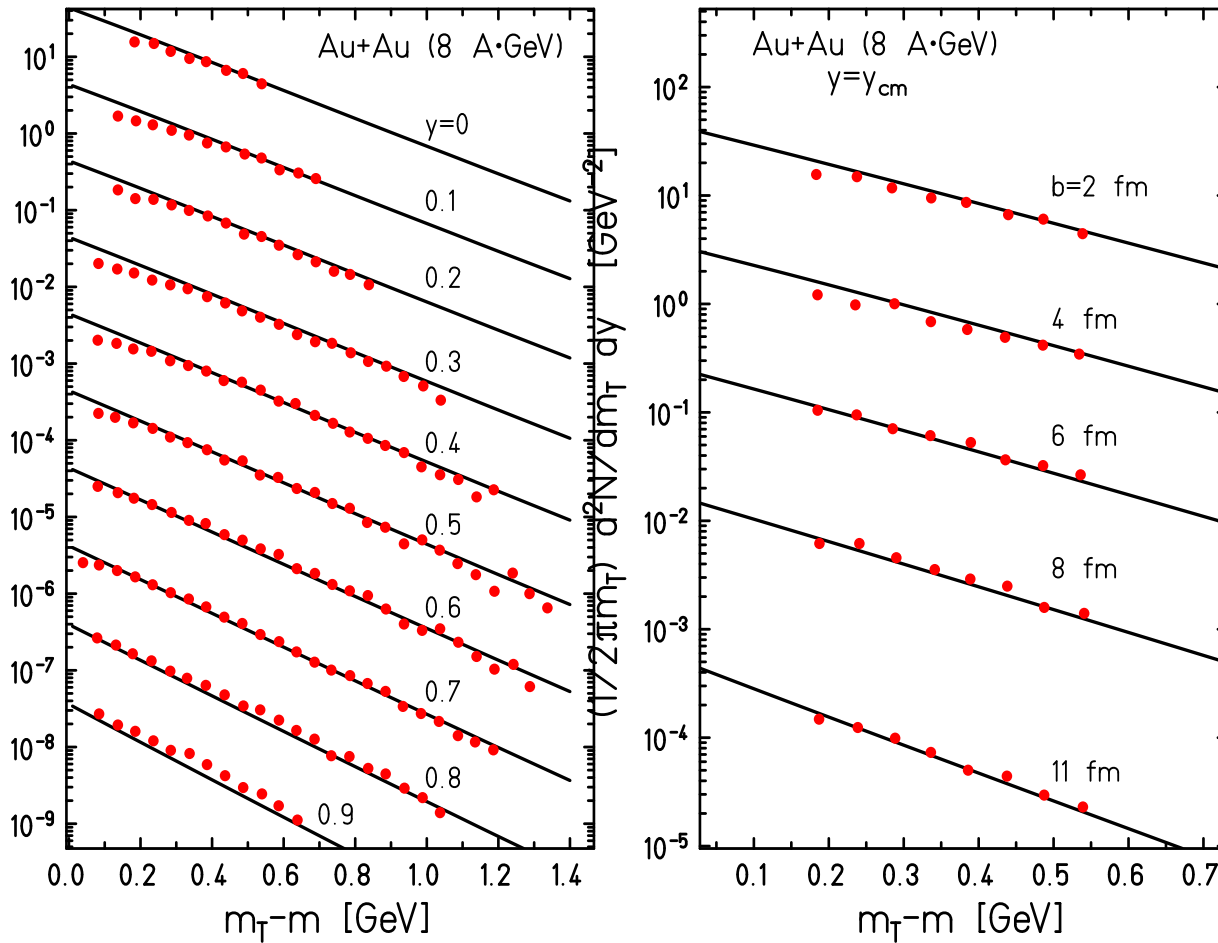
NA49-1: Phys. Rev. Lett. **82** (1999) 2471

NA49-2 (preim.): Nucl. Phys. **A661** (1999) 362c

# Transverse mass distributions of protons

AGS

SPS



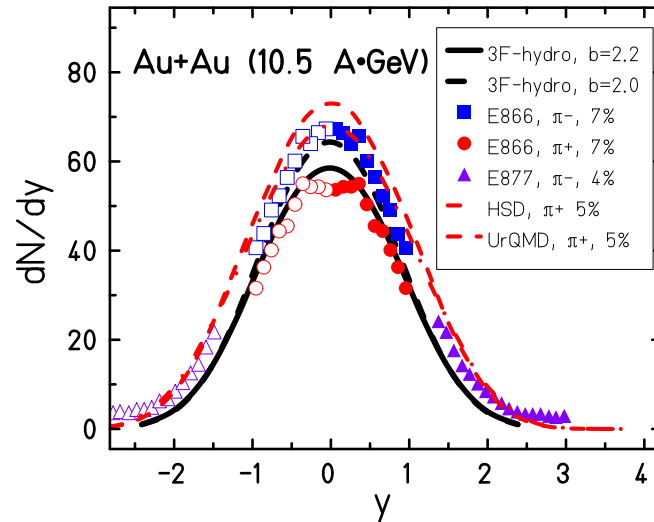
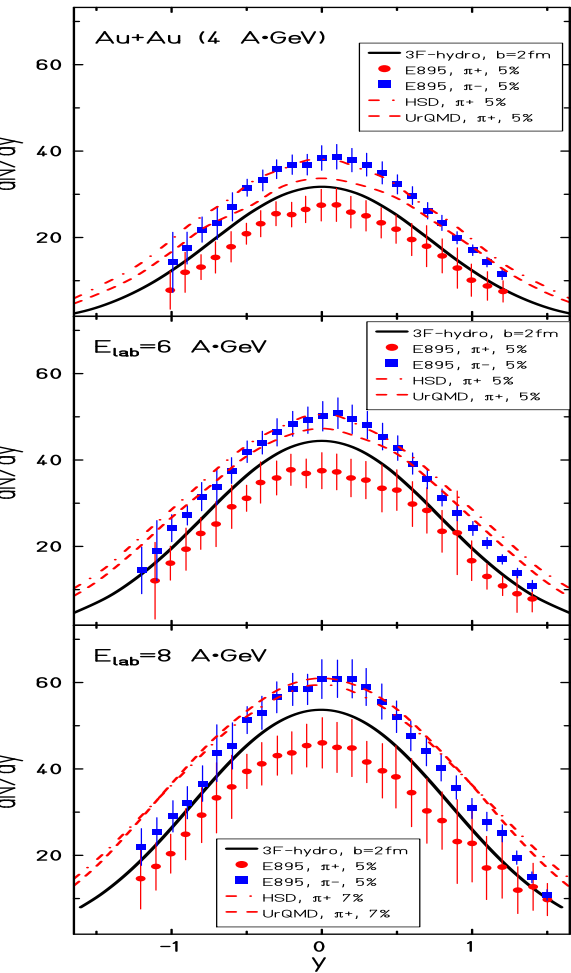
E917: Phys. Rev. Lett. **86** (2001) 1970

$b = 2.2 \text{ fm}$  for 158 AGeV, and  $b = 2.5 \text{ fm}$  for 40 and 80 AGeV are experimental estimates  
 NA49: Phys. Rev. Lett. **82** (1999) 2471  
 NA49: Nucl. Phys. **A715** (2003) 166c

# Pion rapidity spectra

AGS

SPS

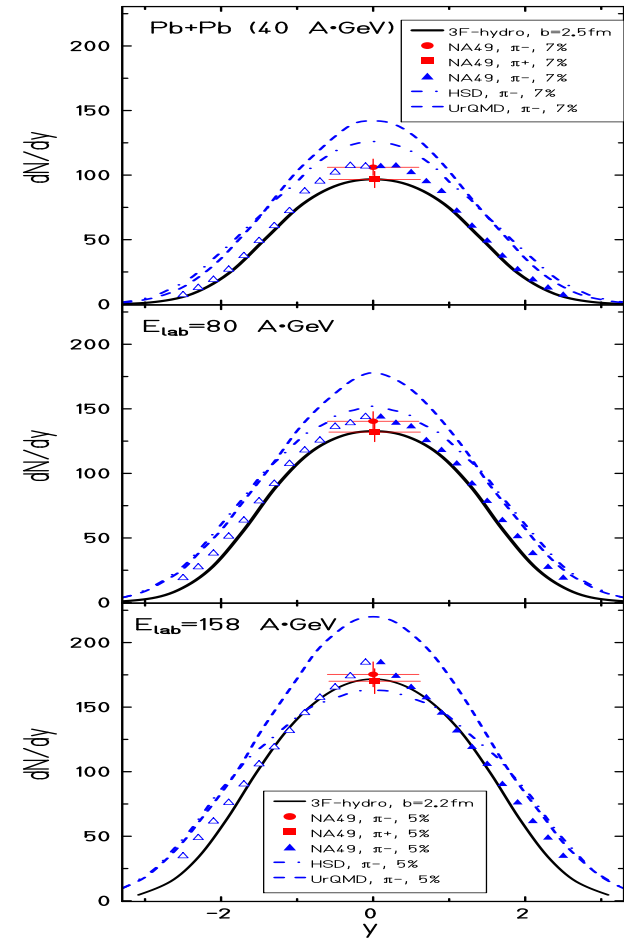


$b = 2.0 \text{ fm}$  for  $7\% \sigma$  and  $b = 1.5 \text{ fm}$  for  $4\% \sigma$  are experimental estimates.

E895: Phys. Rev. **C68** (2003) 054905

E877: Phys. Rev. **C62** (2000) 024901

Transport models: H. Weber,  
E.L. Bratkovskaya, W. Cassing and H. Stöcker,  
Phys. Rev. **C67** (2003) 014904



$b = 2.2 \text{ fm}$  for 158 AGeV, and  $b = 2.5 \text{ fm}$  for 40 and 80 AGeV are experimental estimates

NA49: Phys. Rev. **C66** (2002) 054902

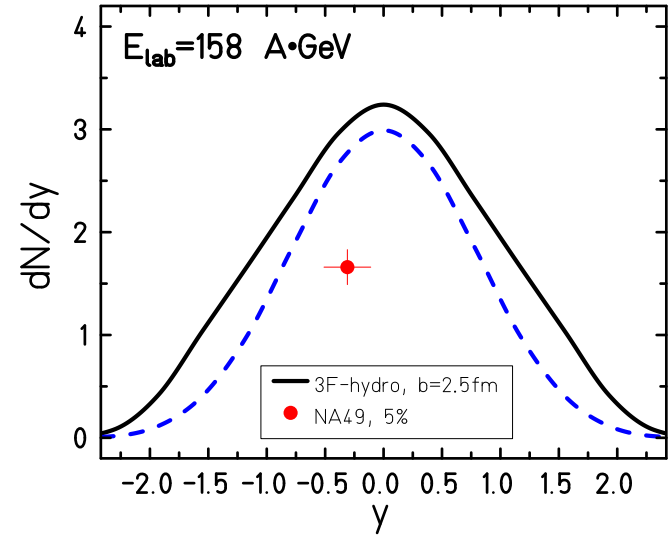
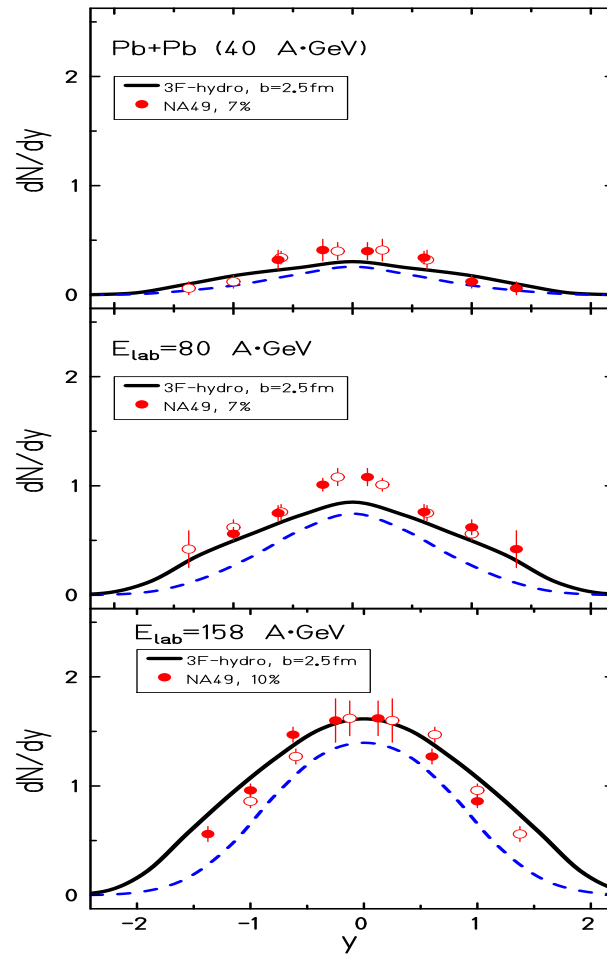
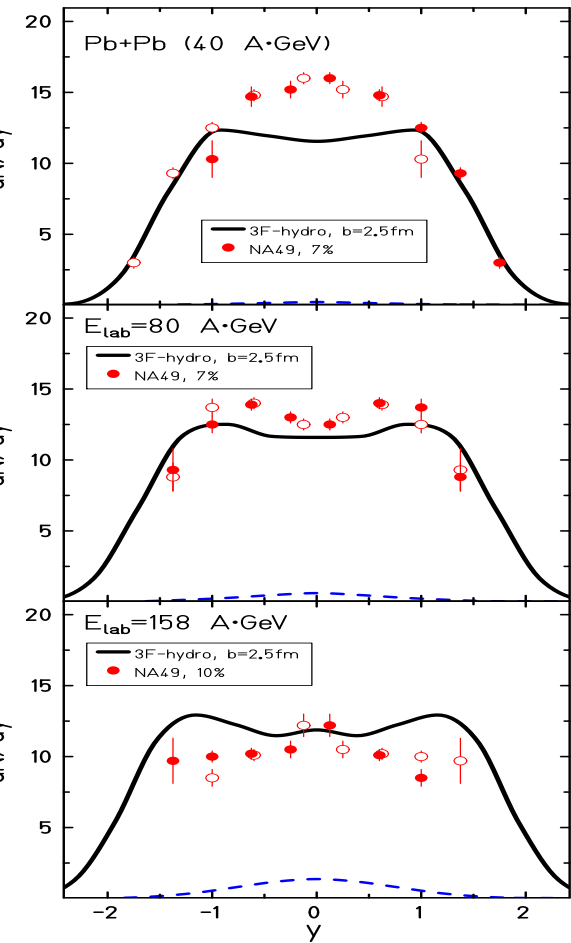
E895: Phys. Rev. **C68** (2003) 054905

# Rapidity distributions for rare particles

$\Lambda + \Sigma^0$

$\bar{\Lambda} + \bar{\Sigma}^0$

$\bar{p}$



NA49: Nucl. Phys. **A661** (1999) 45c

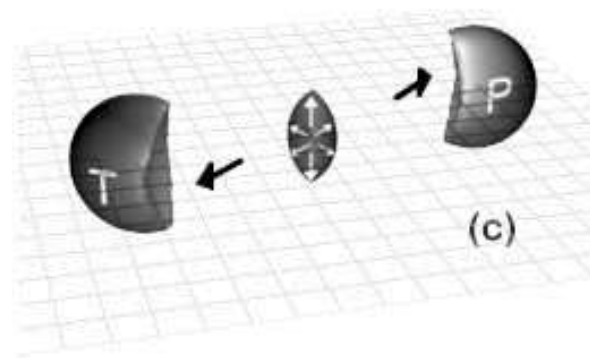
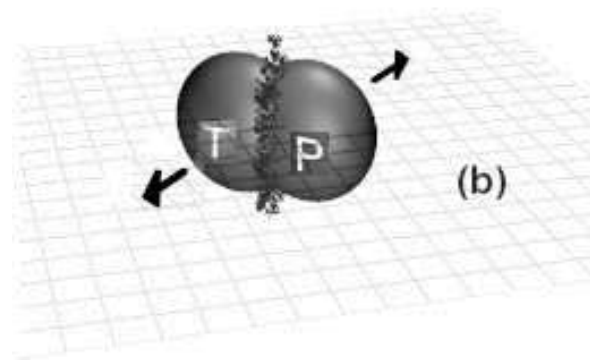
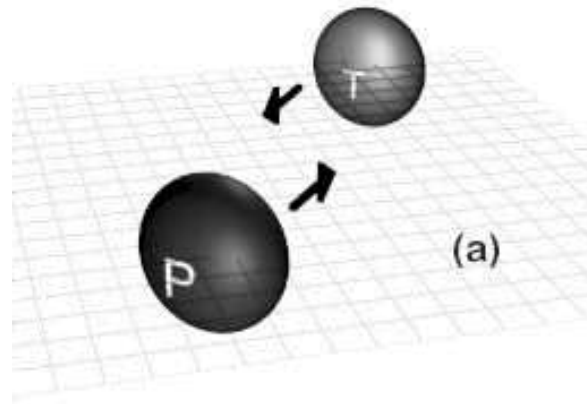
Decays of anti-hyperons are included.

NA49: nucl-ex/0311024

*dashed line = contribution from the fireball fluid*

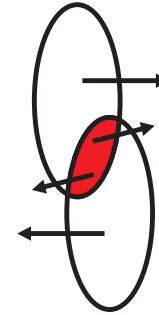
# Nuclear flow

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# Nucleon flow/ AGS data



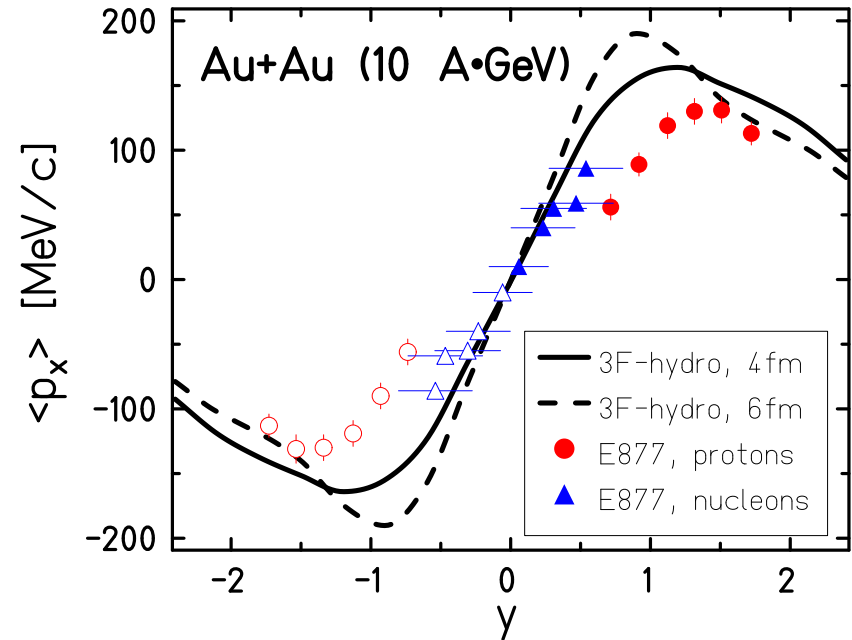
Projection of the hadron transverse momenta on the reaction plane

$$\langle p_x \rangle(y) = \frac{\int d^2 p_T p_x (dN/dy d^2 p_T)}{\int d^2 p_T (dN/dy d^2 p_T)}$$

Resonances are included

E877: Phys. Rev. **C56** (1997)

3254



# Directed and elliptic flow (definitions)

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Directed flow

$$v_1(y) = \int d^2p_T \frac{p_x}{p_T} \frac{dN}{dy d^2p_T} / \int d^2p_T \frac{dN}{dy d^2p_T} \rightarrow \langle \sin \phi \rangle$$

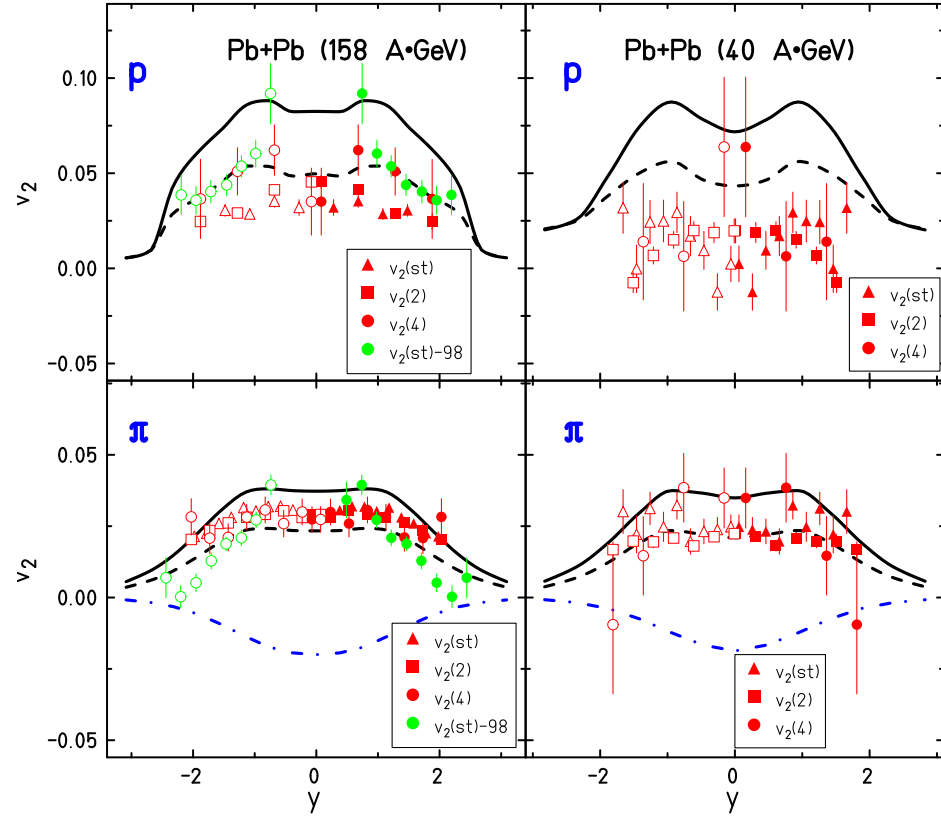
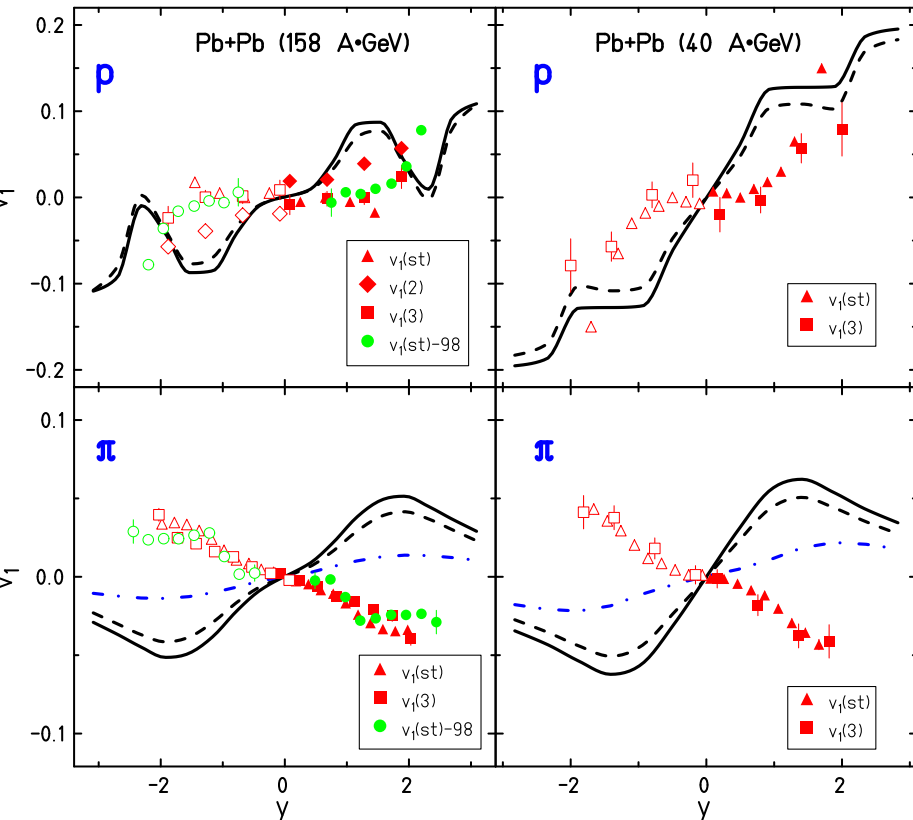
Elliptic flow

$$v_2(y) = \int d^2p_T \frac{p_x^2 - p_y^2}{p_T^2} \frac{dN}{dy d^2p_T} / \int d^2p_T \frac{dN}{dy d^2p_T} \rightarrow \langle \sin 2\phi \rangle$$

# Directed and elliptic flow / SPS data

directed,  $v_1$

elliptic,  $v_2$



3F-hydro:  $b = 5.6$  fm (solid line) and  $b = 4$  fm (dashed line) for mid-central collisions

$v_{1,2}(st)$ : standard method;  $v_{1,2}(n)$ :  $n$ -particle correlation method (N.Borghini, P.M.Dinh, J.-Y.Ollitrault, Phys. Rev. **C93** (2001) 054906)

NA49: Phys. Rev. Lett. **80** (1998) 4136

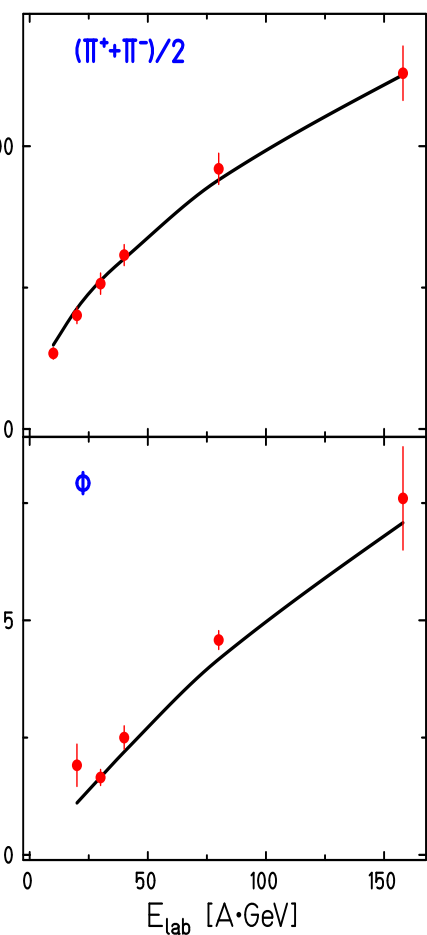
NA49: Phys. Rev. **C68** (2003) 034903

3-Fluids: Hadron gas EoS is too hard

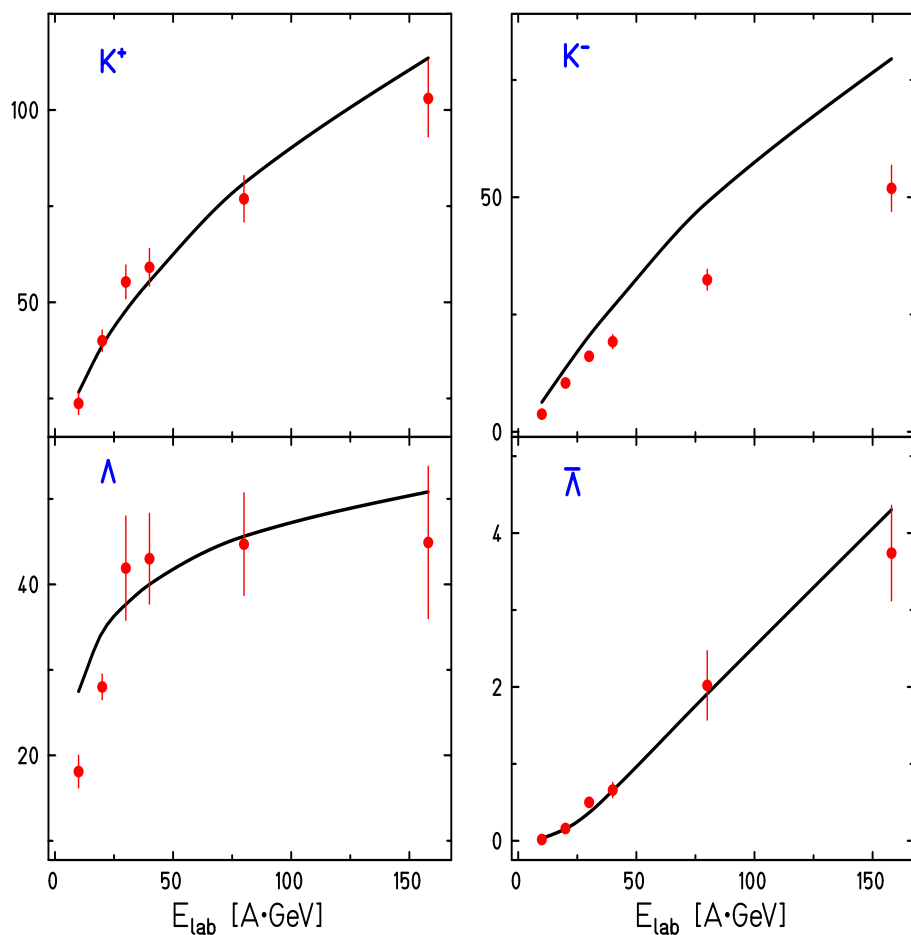
Pion shadowing (?)

# Multiplicity

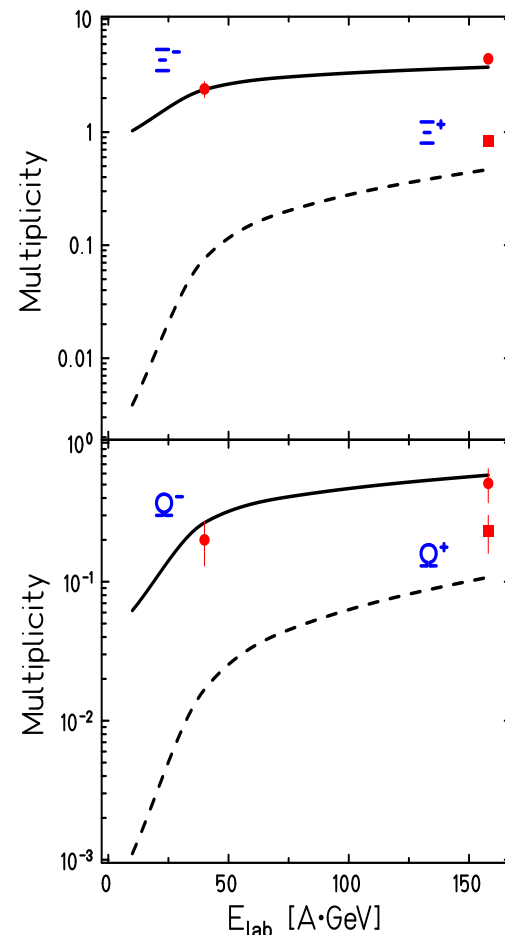
non-strange



strange



multi-strange



as summarized by F.Becattini, M.Gazdzicki et al., hep-ph/0310049

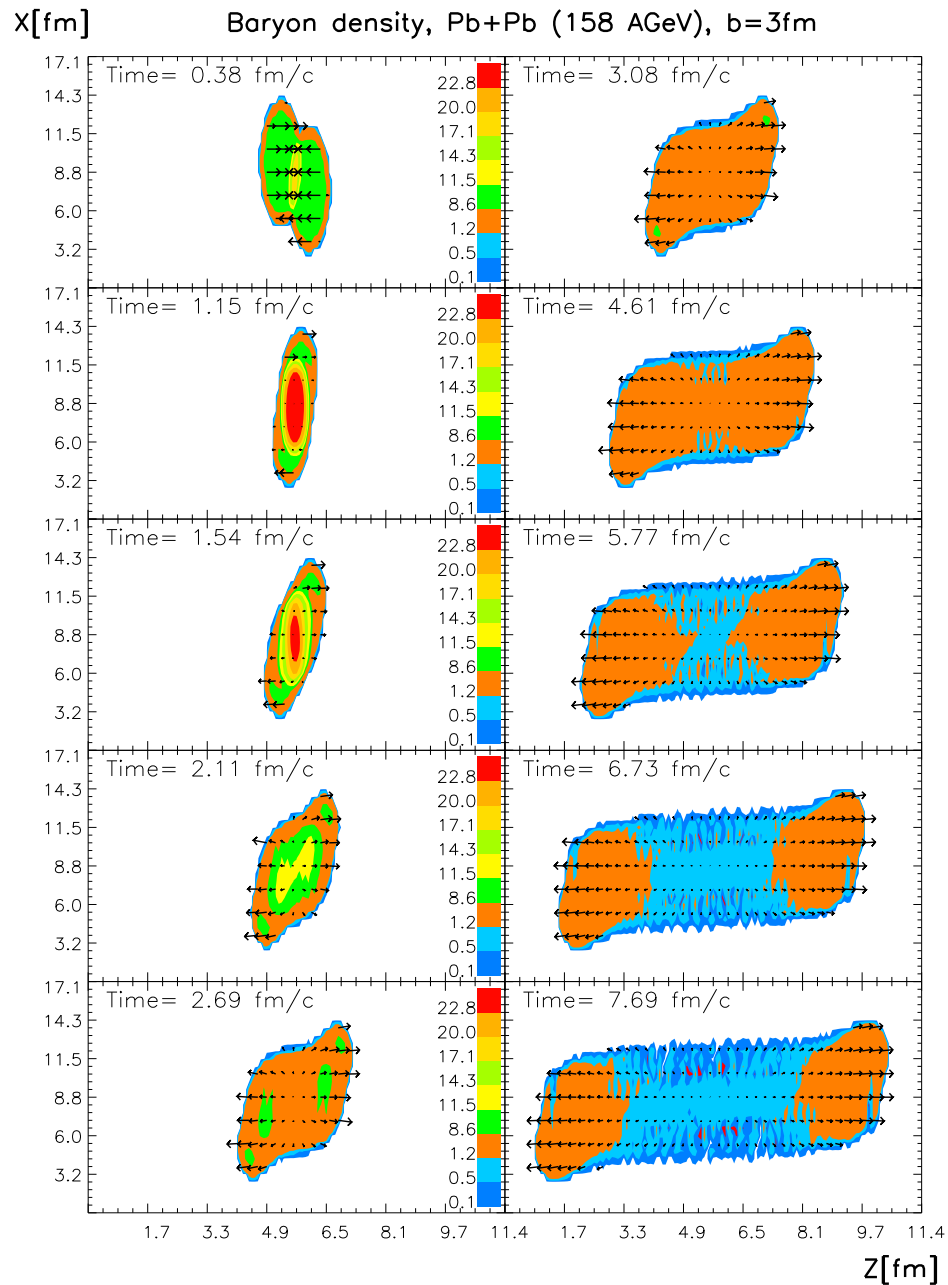
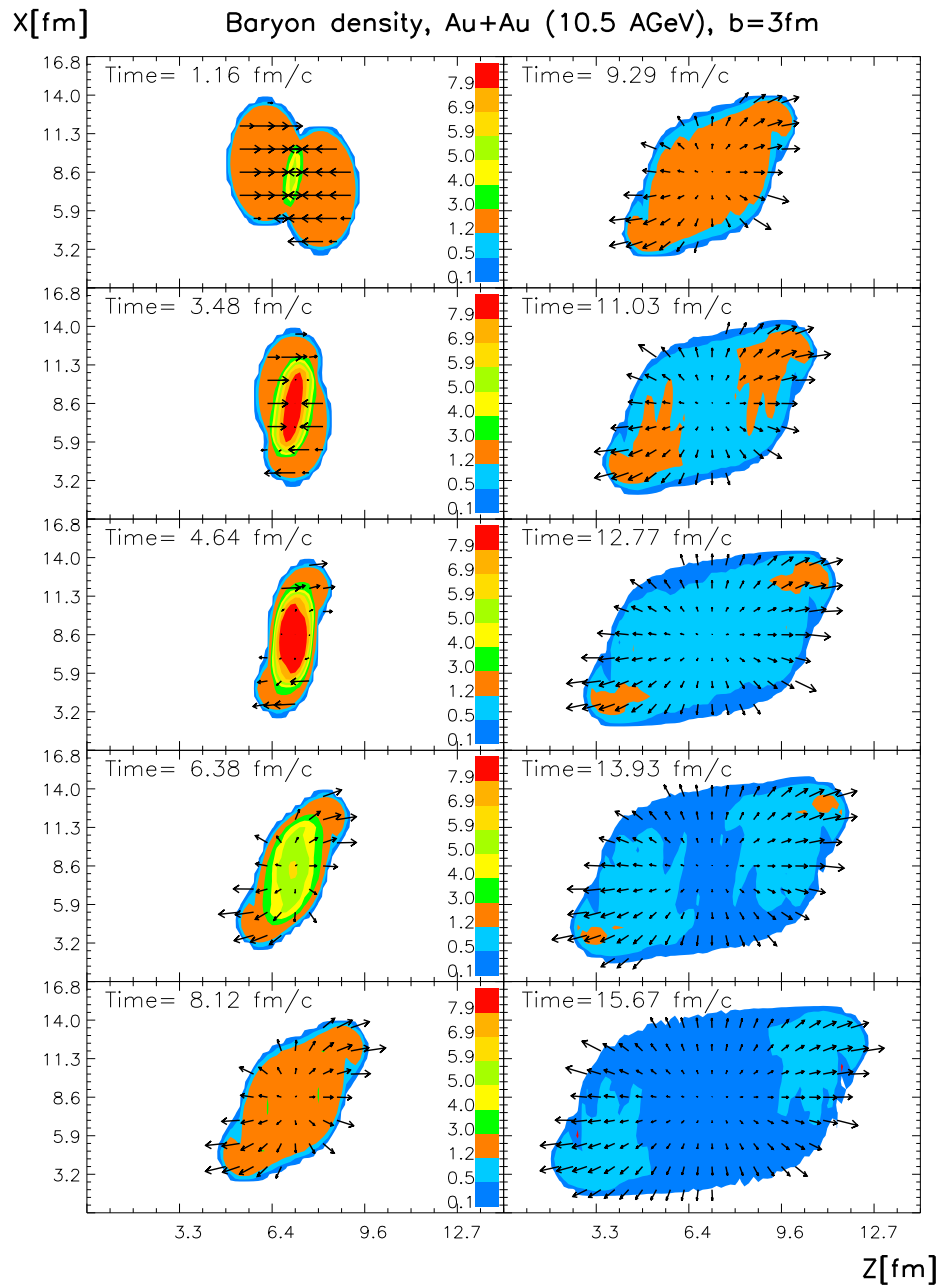
NA49: nucl-ex/0409004

3F-hydro: Hadron gas EoS,  $b=2.2$  fm, grand canonical ensemble, unique freeze-out

# Baryon density evolution

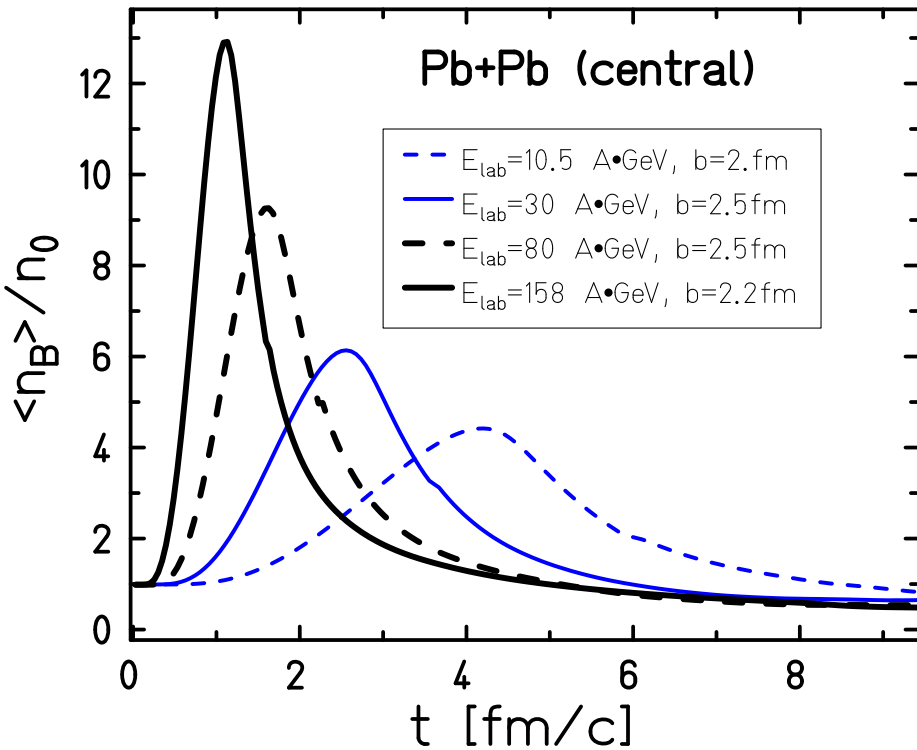
AGS

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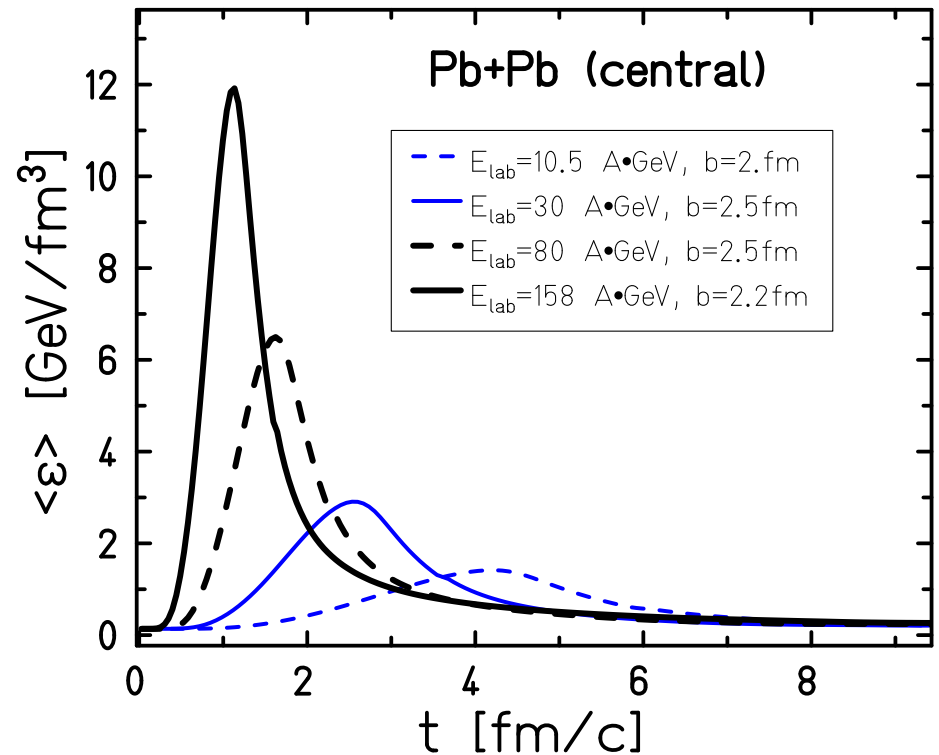


# Temporal evolution of thermodynamic quantities

Baryon density



Energy density



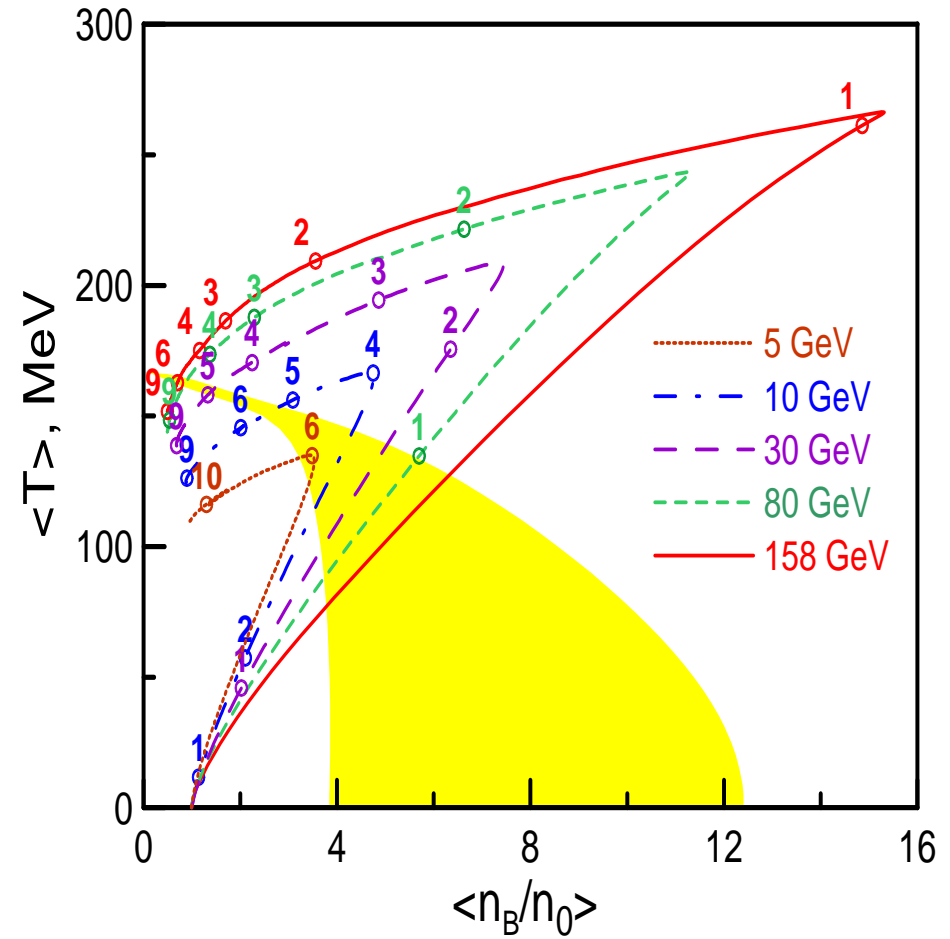
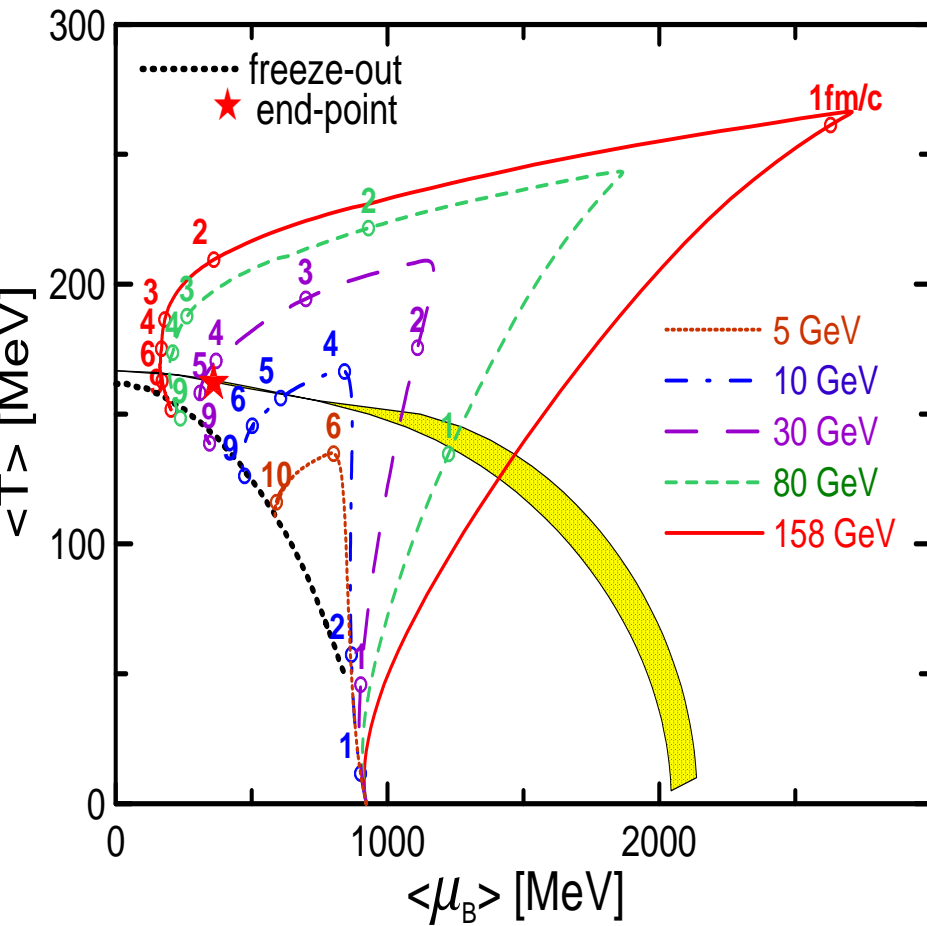
$$\langle n_B \rangle = \int d^3x n_B W(x) / \int d^3x W(x)$$

$$\langle \varepsilon \rangle = \int d^3x \varepsilon W(x) / \int d^3x W(x)$$

where  $W(x) = n_B(x)$

# Phase diagram

central  $Pb + Pb$ , Hadron gas EoS



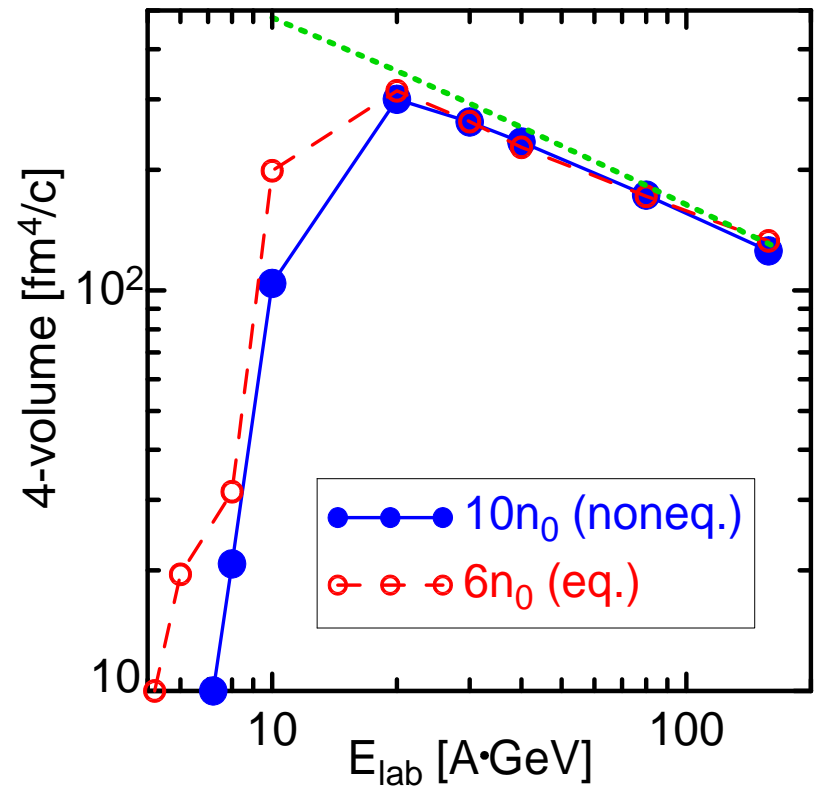
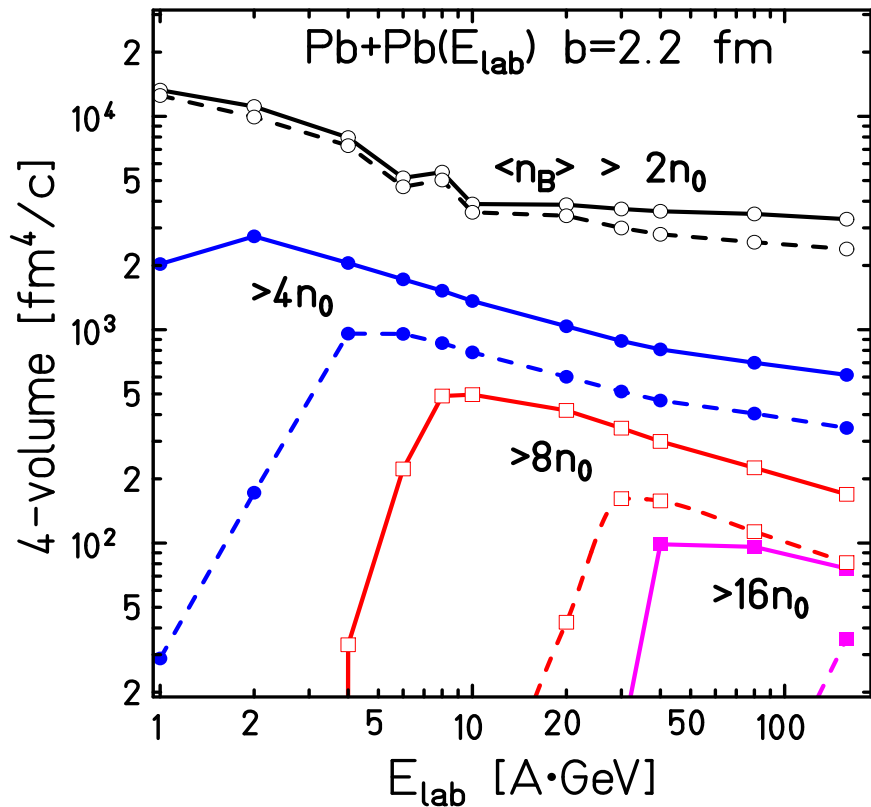
★ **Critical end-point:** Z.Fodor, S.D.Katz, hep-lat/0402006.

Freeze-out curve: J.Cleymans, K.Redlich, Phys. Rev. Lett. **81** (1998) 5284.

# Invariant 4-volume

How long and in which volume a quantity  $Q$  exceeds a  $Q_0$  value

$$V_4(Q) = \int d^4x \Theta(q - Q)$$



1.  $Q = n_B =$  baryonic density (solid lines)

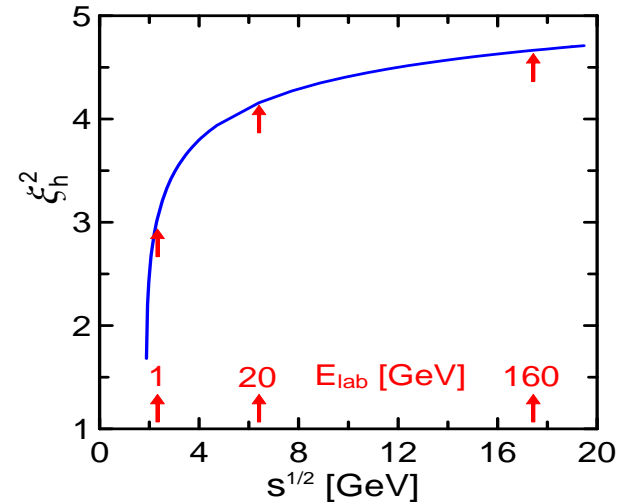
2.  $Q = n_B^{(eq)} =$  baryonic density of thermalized (unified) matter

Lorentz-contracted cylinder:  $V_4(E_{lab}) = \pi R^2 (2R/\gamma_{cm}) \cdot \delta t$  with  $R=4$  fm,  $\delta t=3$  fm/c



# Summary and outlook

- All observables considered here (beside flows) are reasonably reproduced by 3F-hydro with a simple *hadronic EoS*, provided the friction is enhanced as follows



- Is this enhancement reasonable in view of model uncertainties?

(medium effects, multiparticle collisions, poor knowledge of various  $\sigma$ )

At  $T \sim T_c$  masses  $\rightarrow 0$  and the scattering length for  $q - \bar{q}$ , (quasi-)mesons and gluons goes through  $\infty$  ? (at RHIC the enhancement factor  $10 - 10^2$  is needed for partonic  $\sigma$  !;

E.Shuryak and I.Zahed, Phys.Rev. **C70** (2004) 021901; "sticky molasses" : G.E.Brown, C.-H.Lee, M.Rho, hep-ph/0402207)

- Different EoS (with different order of phase transition) should be probed

- Dilepton probes are very promising to disentangle EoS

- $E_{lab} \sim 20$  GeV/nucleon is preferable for production of thermalized matter with  $n_B^{eq} > 6n_0$

# Dilepton production from hydrodynamically expanding fireball

V. Toneev and V. Skokov

The Bogoliubov Laboratory of Theoretical Physics,  
Joint Institute for Nuclear Research, Dubna

June 10, 2005

## 1 Introduction

## 2 Model description

Initial conditions

Expansion dynamics

EoS

Fireball evolution

Evolution in averages

## 3 Hadron observable

## 4 Dileptons

## 5 Outlook

# Introduction

Dilepton  
production  
from hydrody-  
namically  
expanding  
fireball

V. Toneev  
and V.  
Skokov

Outline

Introduction

Model  
description

Initial conditions

Expansion  
dynamics

EoS

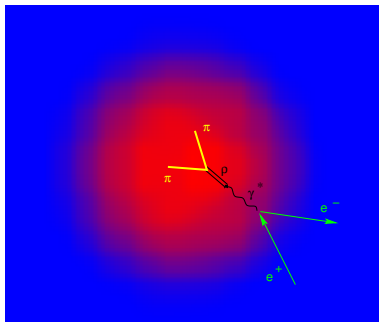
Fireball  
evolution

Evolution in  
averages

Hadron  
observable

Dileptons

Outlook



- in medium effects

$$\frac{d^2 N_{ee}}{dM d\eta} = \frac{M}{\Delta\eta_{e\pm}} \int d\eta \int d^4x \int_0^{2\pi} d\phi \int_0^\infty p_T dp_T \frac{d^8 N_{ee}(T(x), M, \eta, p_T)}{d^4x d^4p} \text{Acc}(M, \eta, p_T)$$

# Introduction

Dilepton production from hydrodynamically expanding fireball

V. Toneev and V. Skokov

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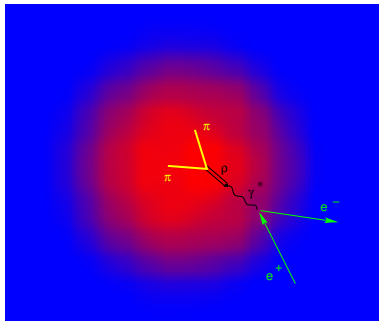
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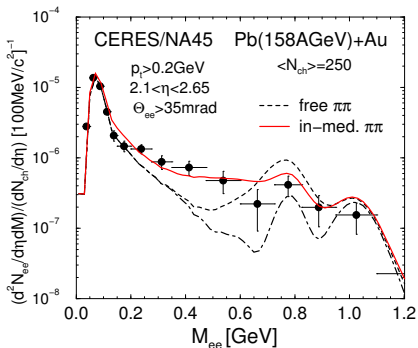


- in medium effects

$$\frac{d^2 N_{ee}}{dM d\eta} = \frac{M}{\Delta\eta_{e\pm}} \int d\eta \int d^4x \int_0^{2\pi} d\phi \int_0^\infty p_T dp_T \frac{d^8 N_{ee}(T(x), M, \eta, p_T)}{d^4x d^4p} \text{Acc}(M, \eta, p_T)$$

R. Rapp and J. Wambach, Eur. Phys. J. A 6 (1999) 415

[arXiv:hep-ph/9907502]



# Initial conditions, entropy creation

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Entropy creation and fireball formation are calculated within the transport code: **QGSM** (quark-gluon string model) that defines an initial state for subsequent hydro stage

$$\frac{df_i(x, p)}{dt} = \sum_j C_{ij}^{gain} - \sum_j C_{ji}^{loss}$$

( $i = N, \Delta, \dots, \text{hyperons}, \pi, K, \eta, \rho, \omega \dots$ )

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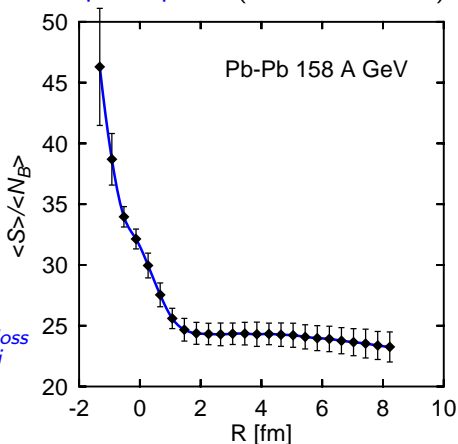
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( $i = N, \Delta, \dots, \text{hyperons}, \pi, K, \eta, \rho, \omega \dots$ )

for participants (central collision)



relative nuclei distance

# Hydrodynamic equations

Isentropic expansion of the formed fireball is described by **3D relativistic hydrodynamics** using calculated energy and baryon densities as well as velocity profile (Flux corrected SHASTA is applied)

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Isentropic expansion of the formed fireball is described by **3D relativistic hydrodynamics** using calculated energy and baryon densities as well as velocity profile (Flux corrected SHASTA is applied)

- **Energy-momentum** conservation

$$\frac{\partial T^{\mu\nu}}{\partial x^\mu} = 0 \quad \text{with} \quad T^{\mu\nu} = (\varepsilon + P) u^\mu u^\nu - P g^{\mu\nu}$$

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# Hydrodynamic equations

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$$\frac{\partial J_B^\mu}{\partial x^\mu} = 0 \quad \text{with} \quad J_B^\mu = n_B u^\mu$$

- Equation of state (EoS)

$$P = P(\varepsilon, n_B).$$

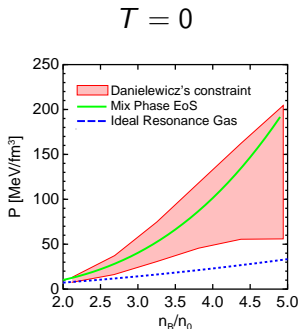
## Statistical quark-hadron mixed phase EoS

(V.D. Toneev et al., J. Phys. G **27** (2001) 827 [nucl-th/0011029])

### Hadronic sector

generalized Zimanyi mean-field  
model, (Nucl. Phys. **A484**, 647 (1988))

- Saturation properties of nuclear matter (Binding energy, pressure, incompressibility at normal density)
- Danielewicz's constraint  
(P. Danielewicz, R. Lacey, and W.G. Lynch,  
Science **298**, 1592 (2002) [nucl-th/0208016])



## Quark-gluon sector

Quasiparticle gas of interacting  
quarks and gluons

- Asymptotic of quark masses  
in HTL approximation (!)

- Comparison with lattice  
QCD results for  $N_f = 2 + 1$

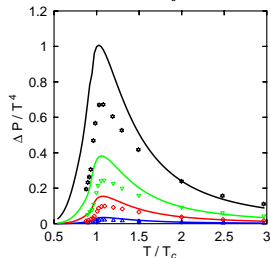
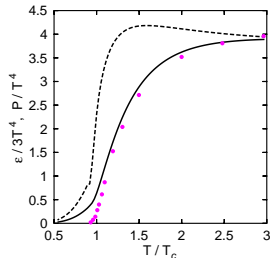
( Z. Fodor and S.D. Katz, JHEP **203**, 14 (2002)

[hep-lat/0106002]; JHEP **404**, 50 (2004)

[hep-lat/0402006])

$$\Delta P = P(T, \mu_B) - P(T, \mu_B = 0)$$

$$\mu_B = 100, 210, 330, 530 \text{ MeV} \quad (\text{from the bottom})$$



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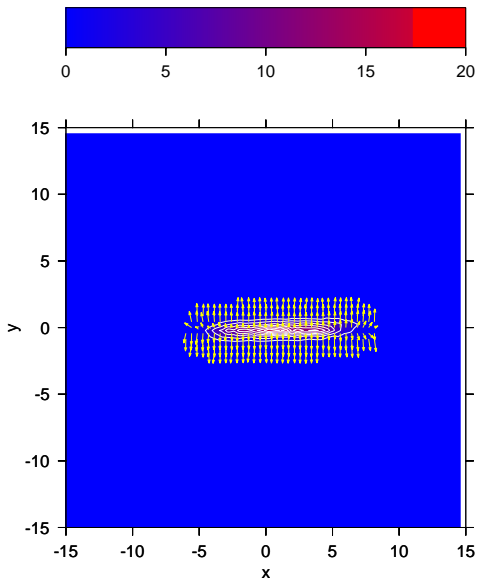
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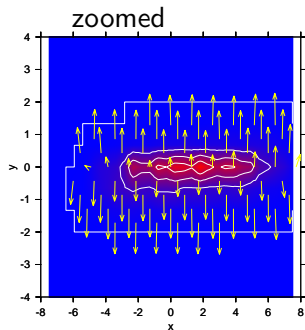
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$t = 0.0$  fm/c;



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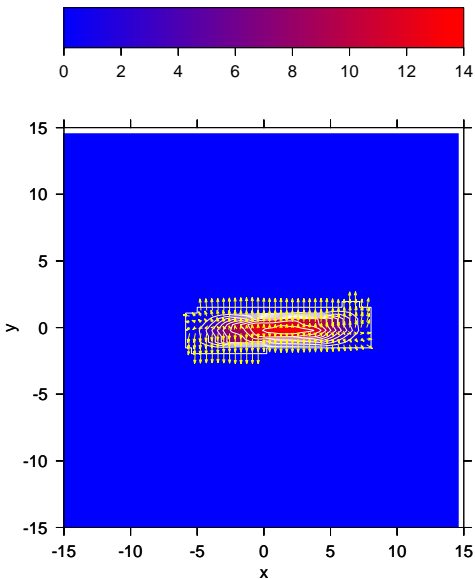
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$t = 0.3 \text{ fm}/c$

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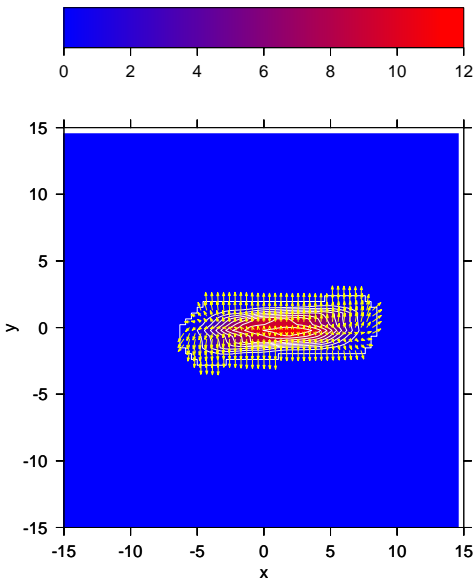
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$t = 0.6 \text{ fm}/c$



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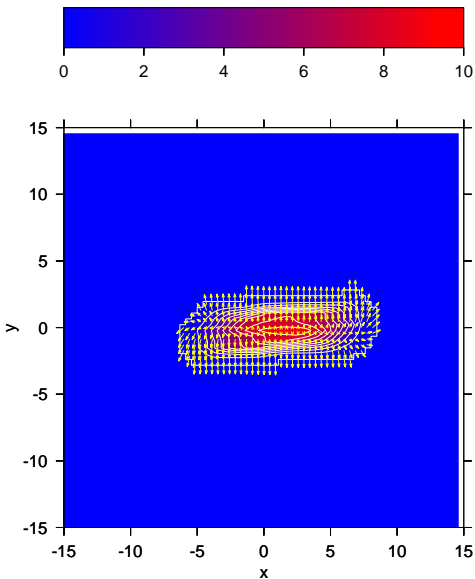
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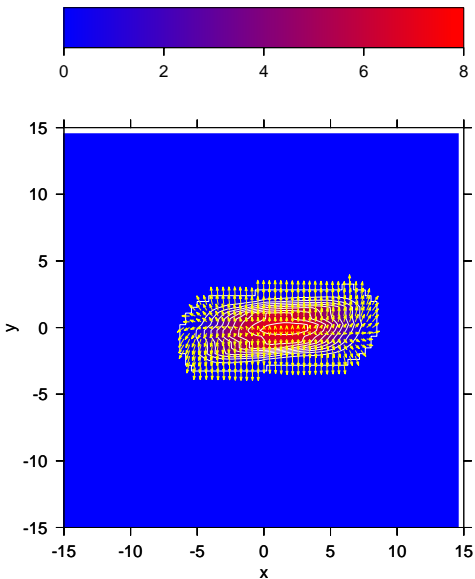
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$t = 1.2 \text{ fm}/c$

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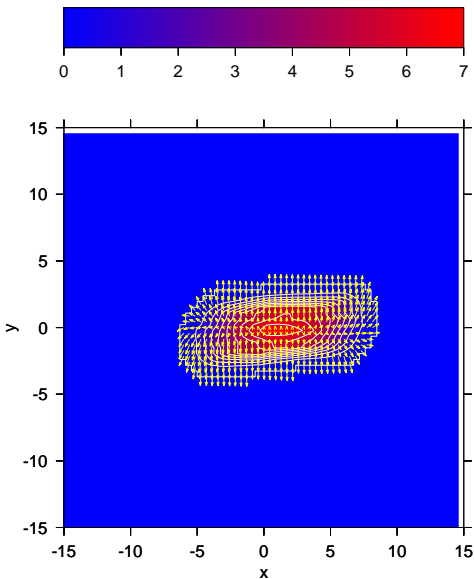
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$t = 1.5 \text{ fm}/c$

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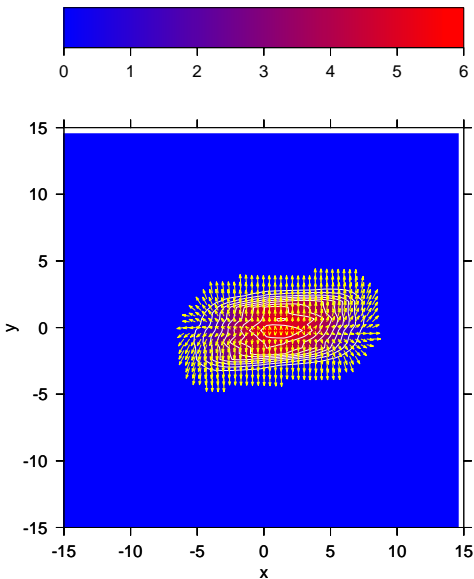
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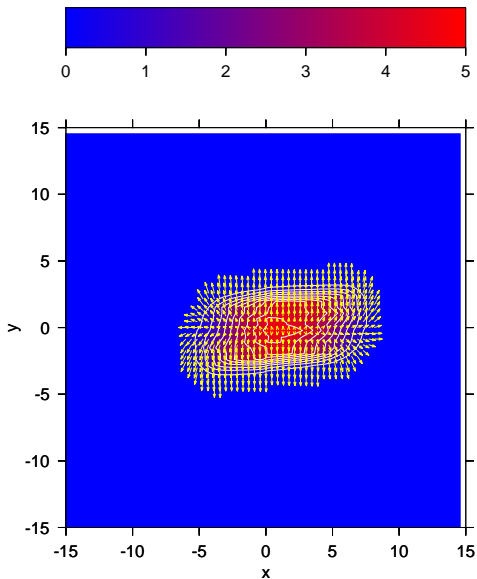
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$t = 2.1 \text{ fm}/c$

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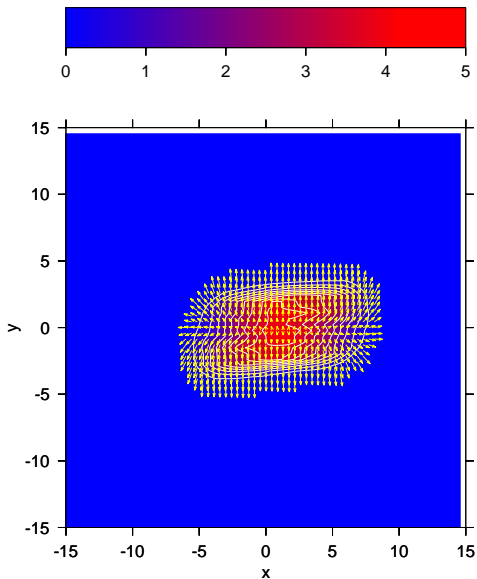
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$t = 2.4 \text{ fm}/c$

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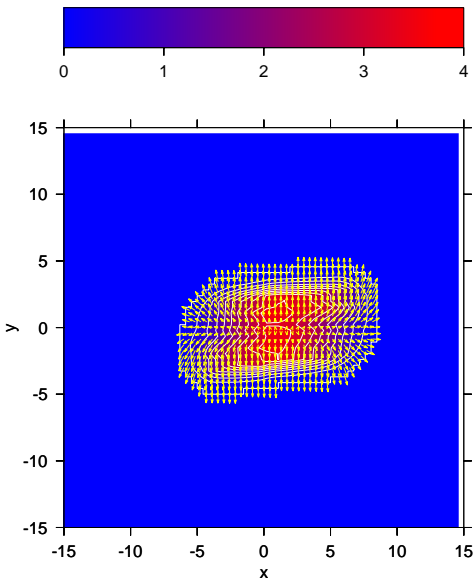
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$t = 2.7 \text{ fm}/c$

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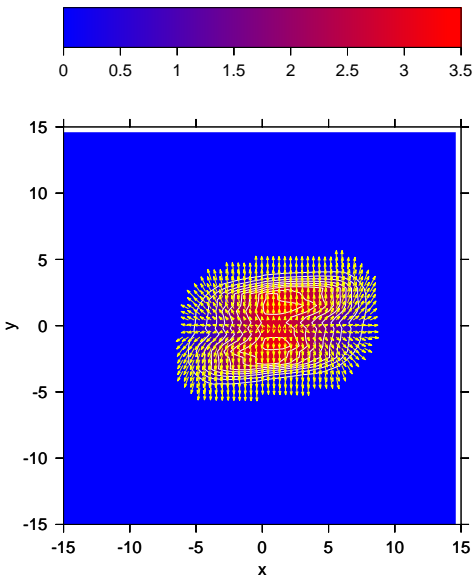
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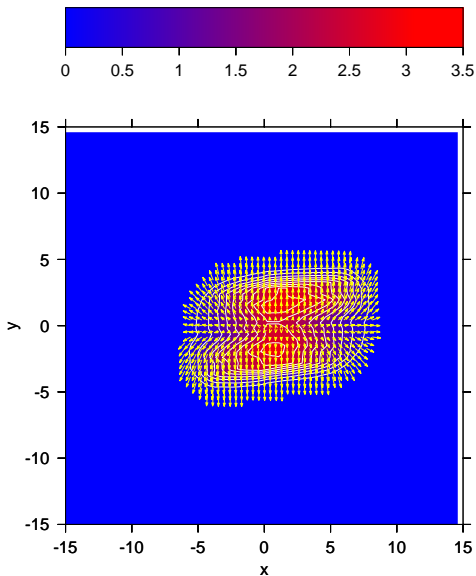
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$t = 3.3 \text{ fm}/c$

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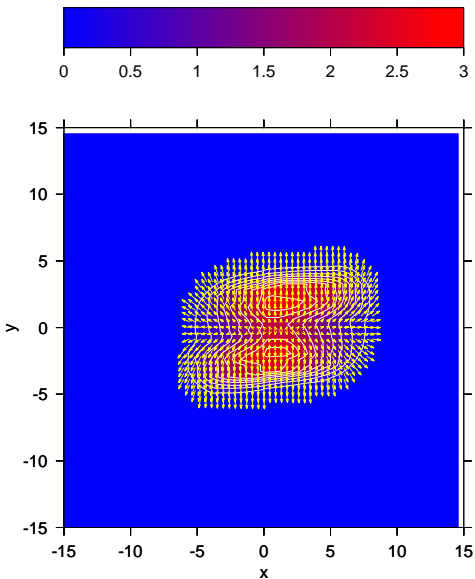
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$t = 3.6 \text{ fm}/c$

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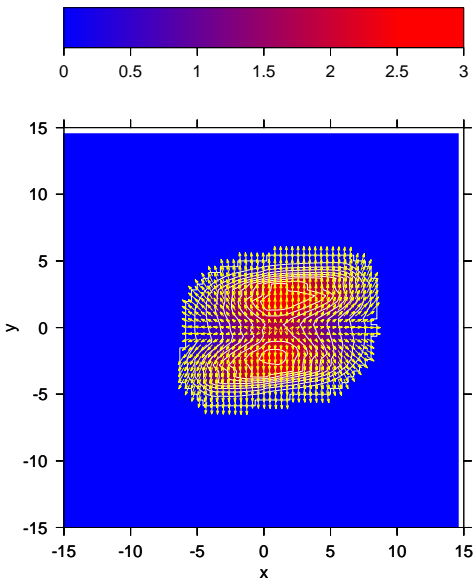
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$t = 3.9 \text{ fm}/c$

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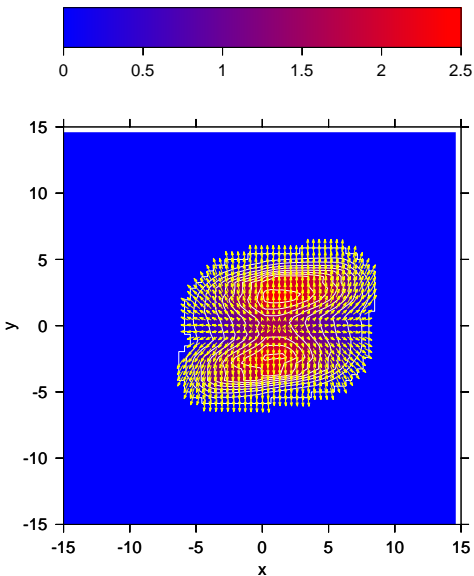
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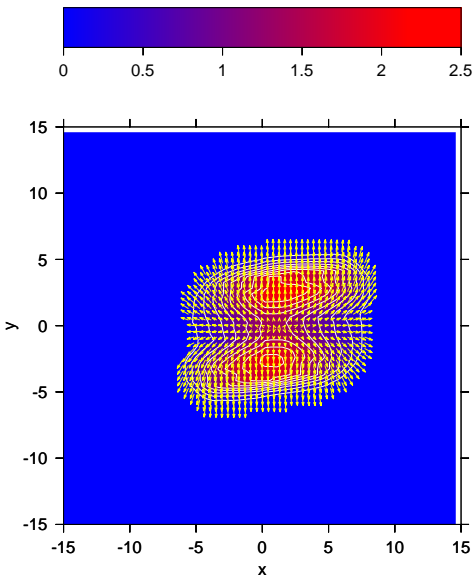
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$t = 4.5 \text{ fm}/c$

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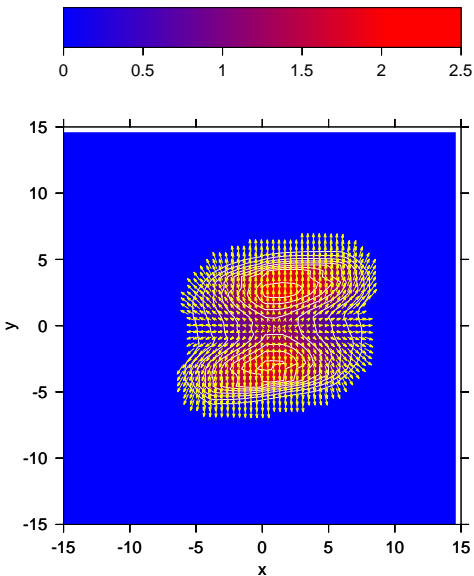
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$t = 4.8 \text{ fm}/c$

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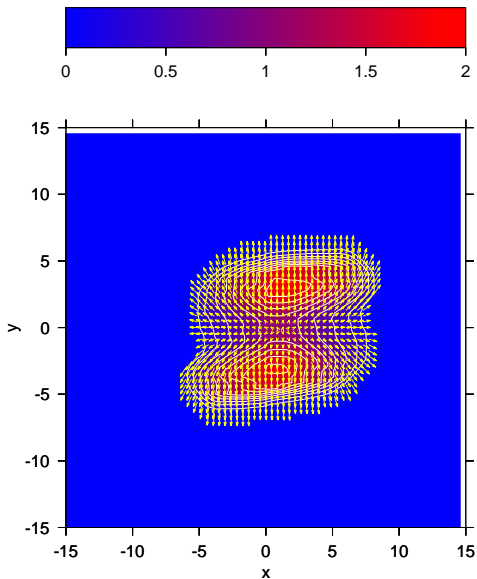
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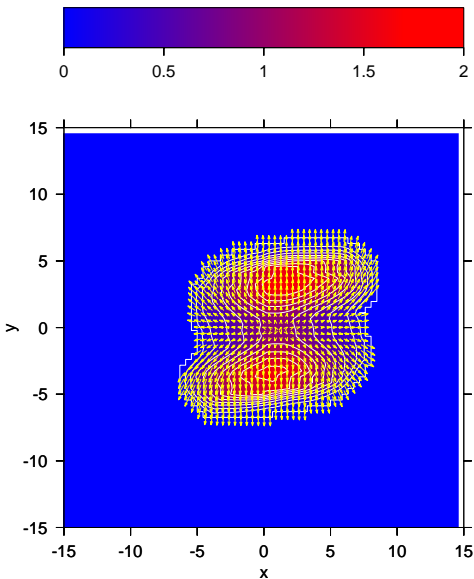
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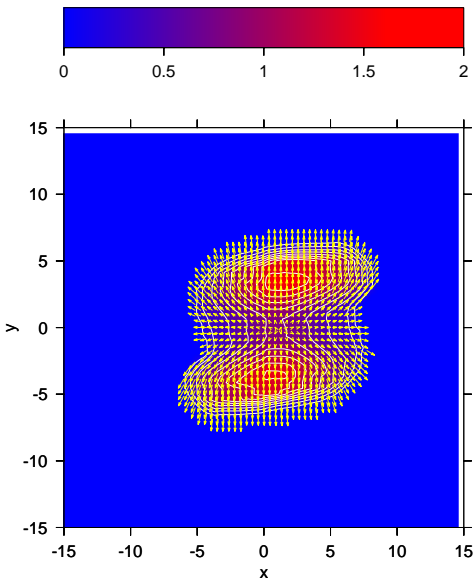
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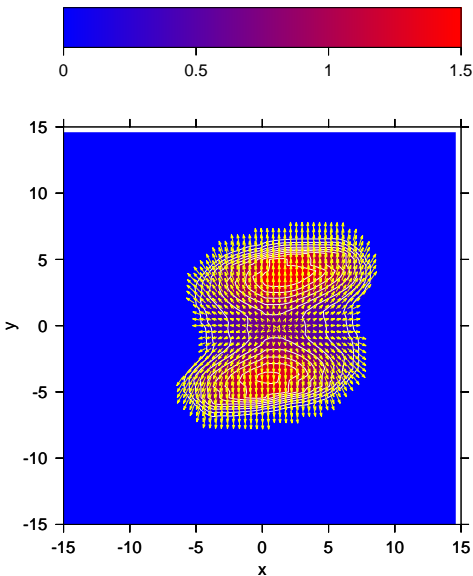
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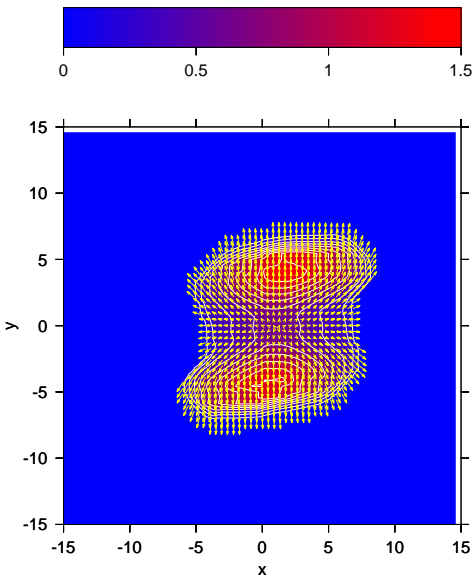
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$t = 6.3 \text{ fm}/c$

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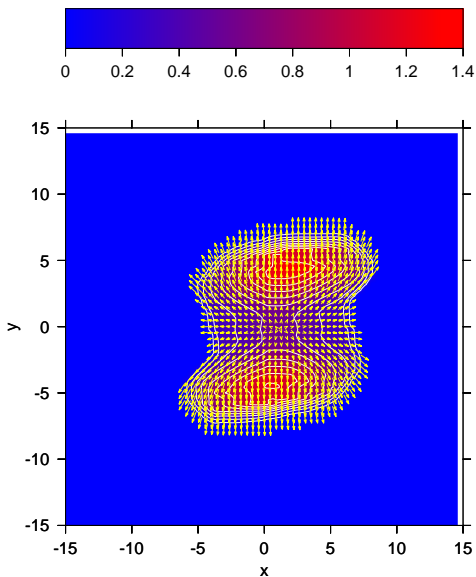
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$t = 6.6 \text{ fm}/c$

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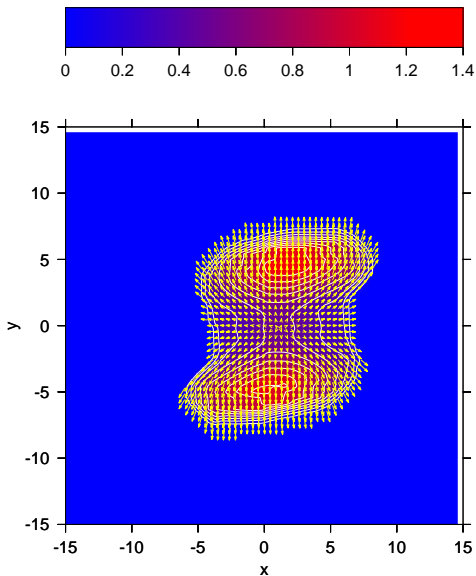
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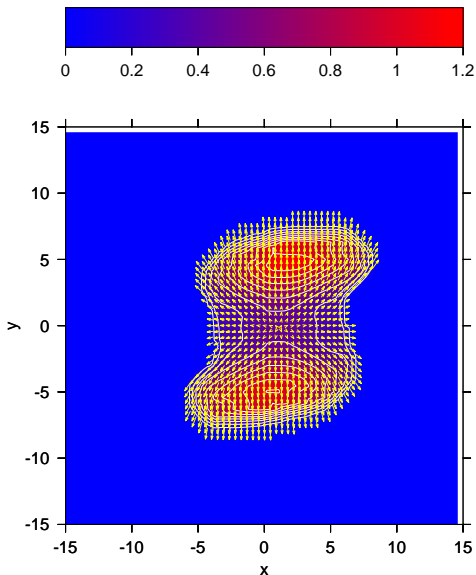
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$t = 7.2 \text{ fm}/c$

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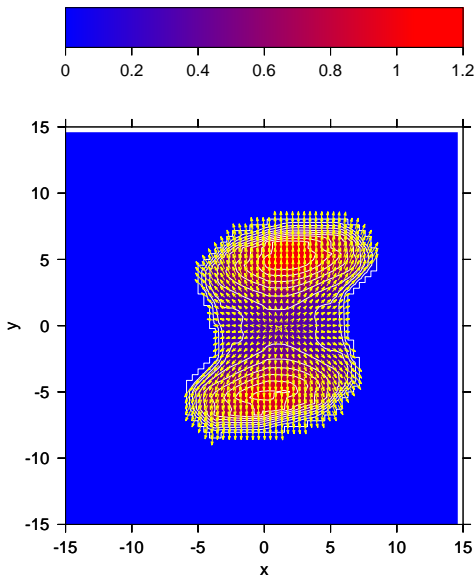
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$t = 7.5 \text{ fm}/c$

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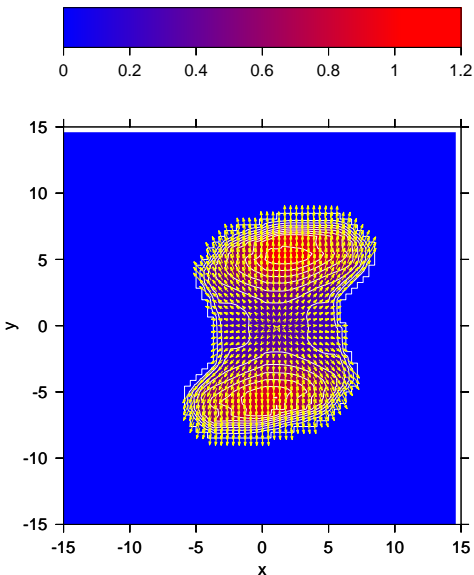
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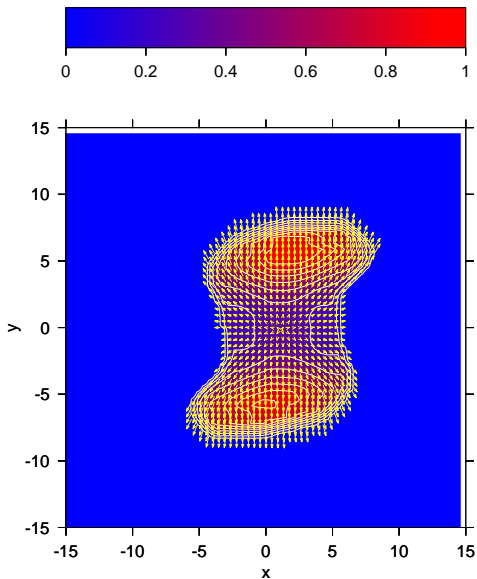
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$t = 8.1 \text{ fm}/c$

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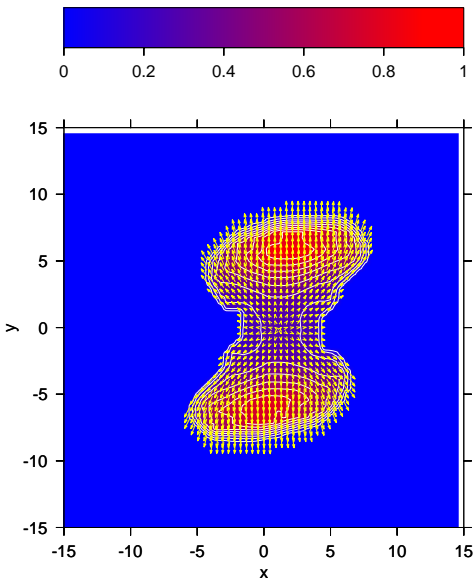
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$t = 8.4 \text{ fm}/c$

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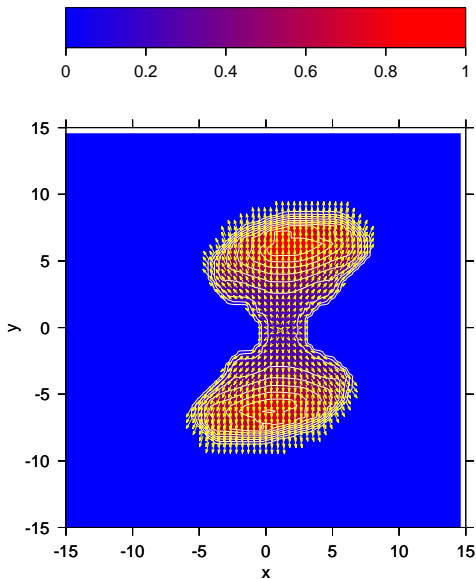
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$t = 8.7 \text{ fm}/c$

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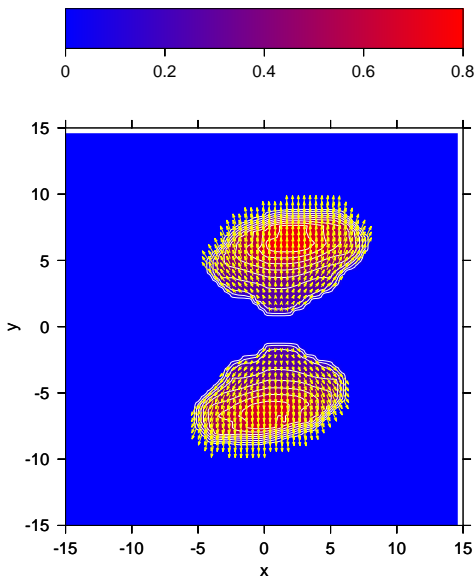
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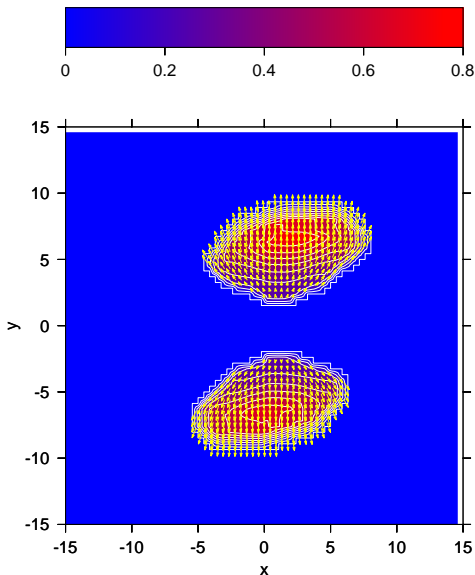
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$t = 9.3 \text{ fm}/c$

# Time evolution of the fireball energy density

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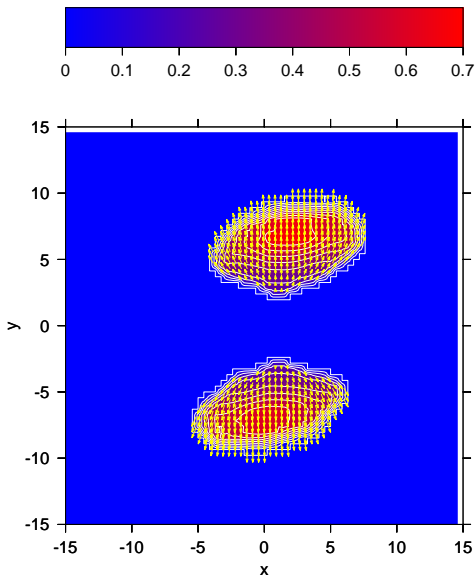
**Fireball evolution**

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$t = 9.6 \text{ fm}/c$

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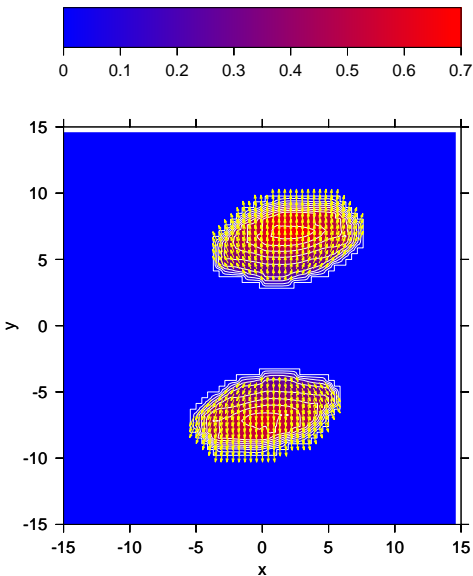
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$t = 9.9 \text{ fm}/c$

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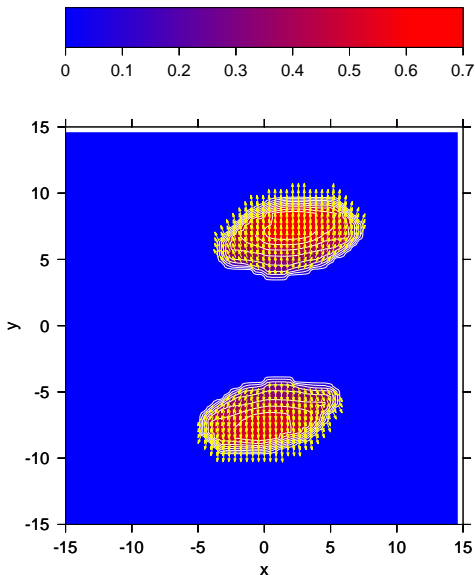
**Fireball evolution**

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$t = 10.2 \text{ fm}/c$



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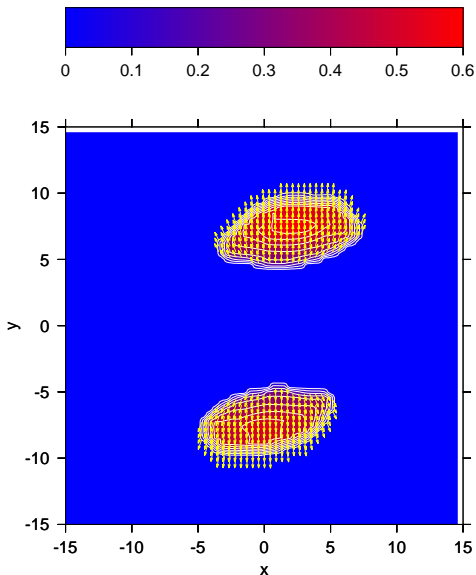
**Fireball evolution**

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$t = 10.5 \text{ fm}/c$

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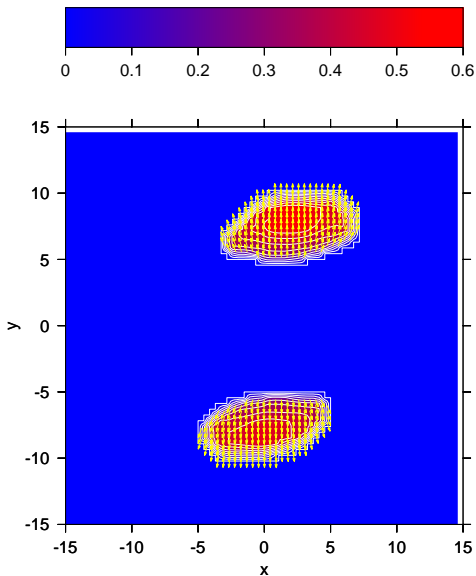
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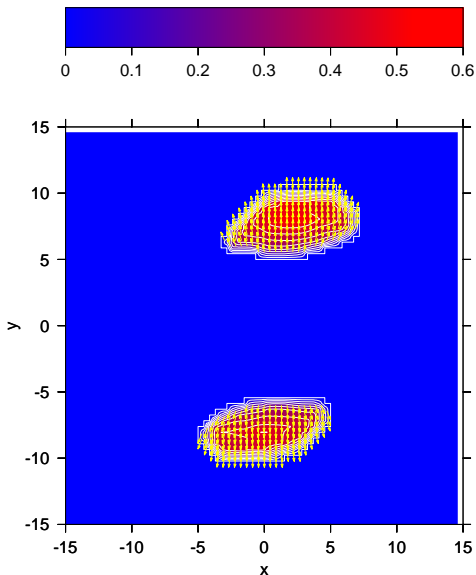
**Fireball evolution**

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$t = 11.1 \text{ fm}/c$

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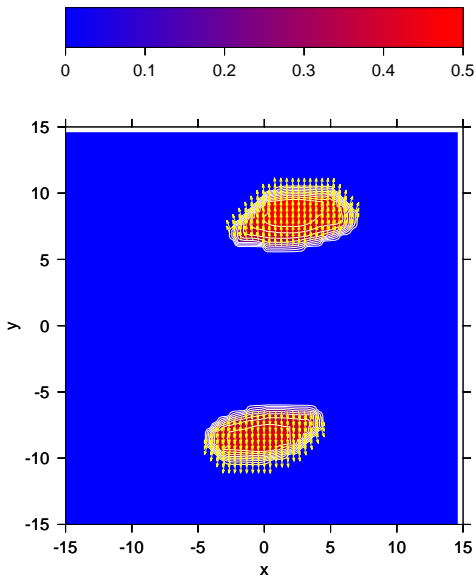
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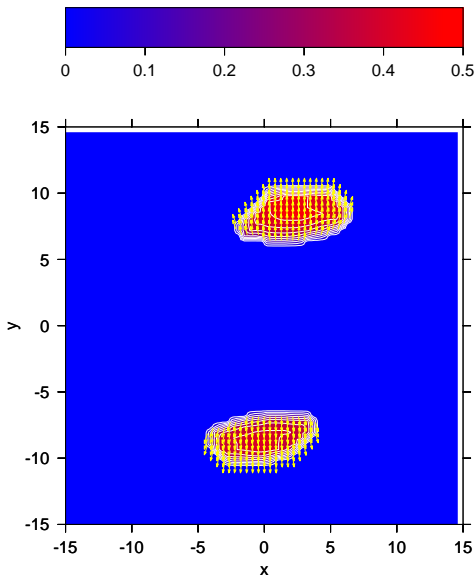
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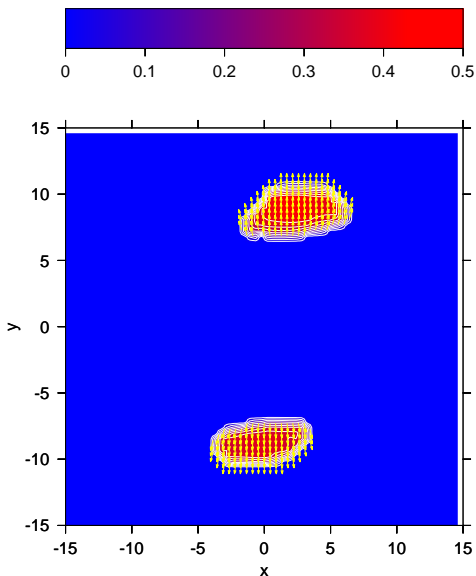
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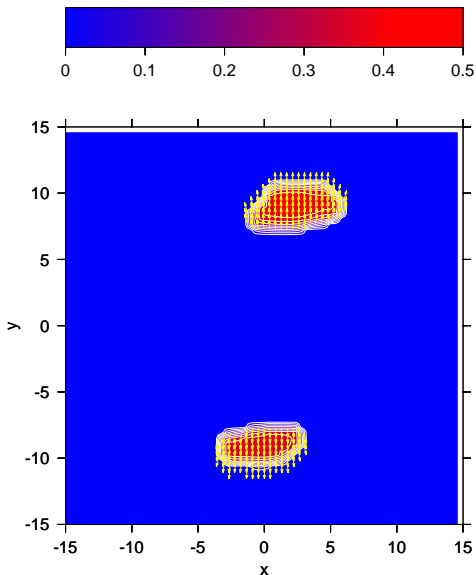
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$t = 12.3 \text{ fm}/c$

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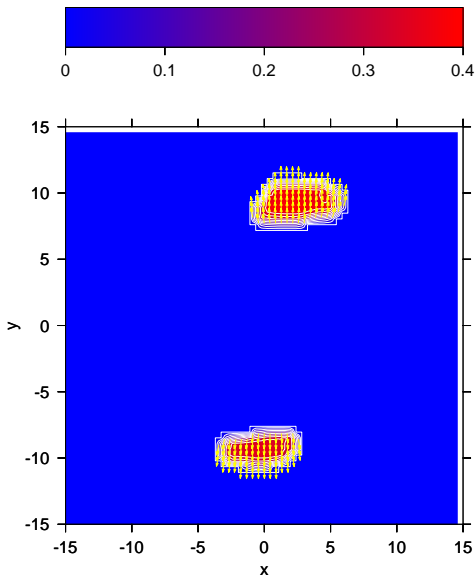
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$t = 12.6 \text{ fm}/c$



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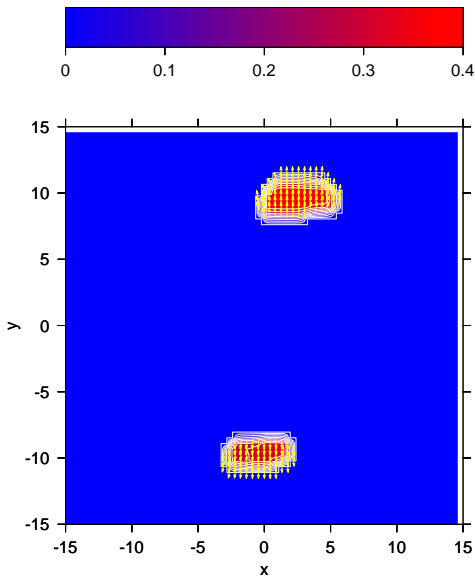
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$t = 12.9 \text{ fm}/c$

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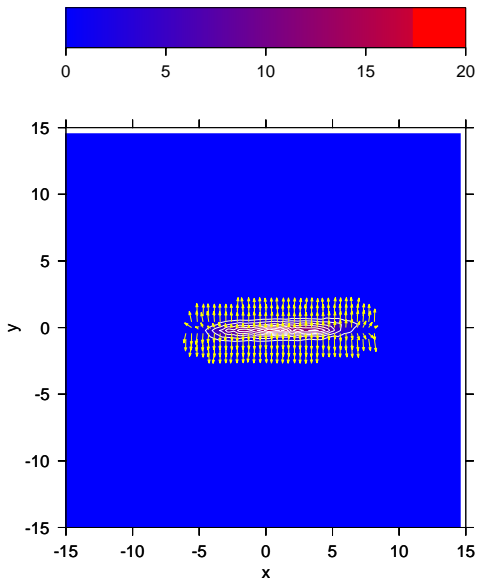
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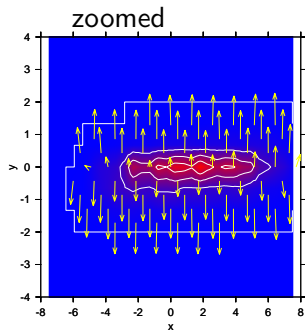
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$t = 0.0$  fm/c;



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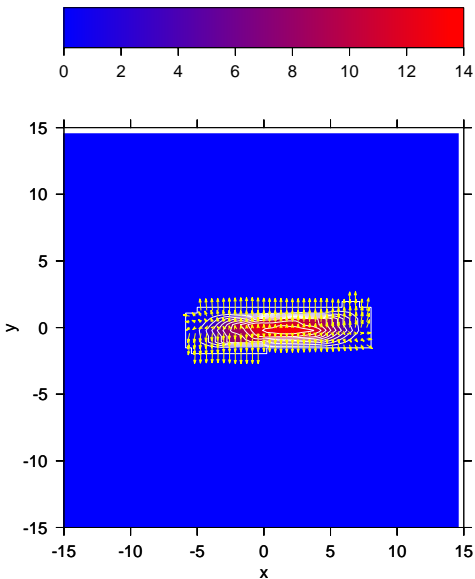
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$t = 0.3 \text{ fm}/c$

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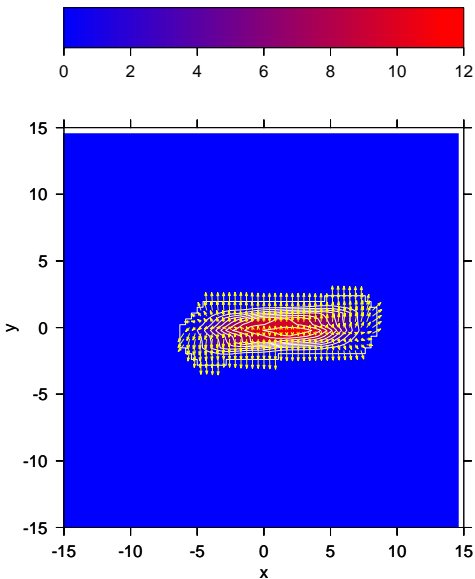
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$t = 0.6 \text{ fm}/c$

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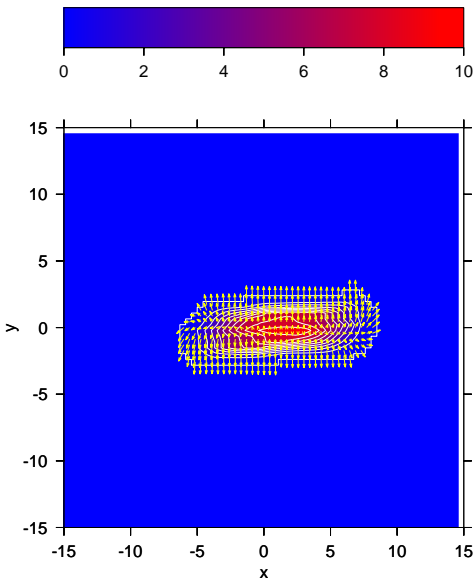
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$t = 0.9 \text{ fm}/c$

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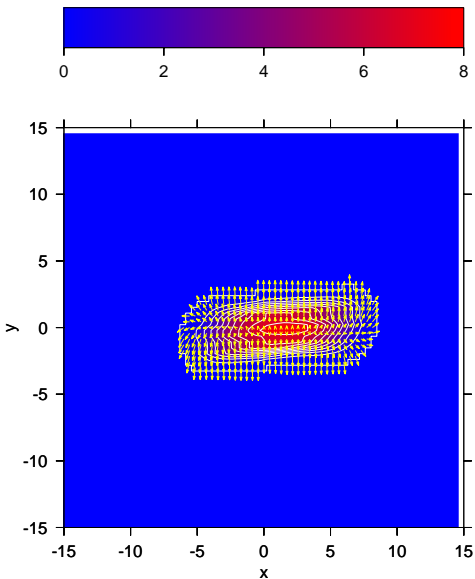
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$t = 1.2 \text{ fm}/c$

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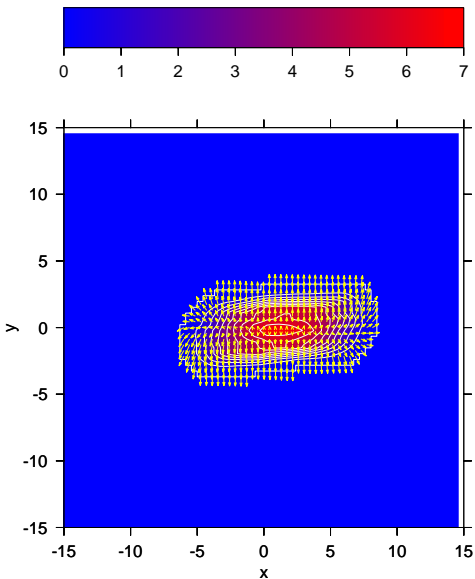
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$t = 1.5 \text{ fm}/c$

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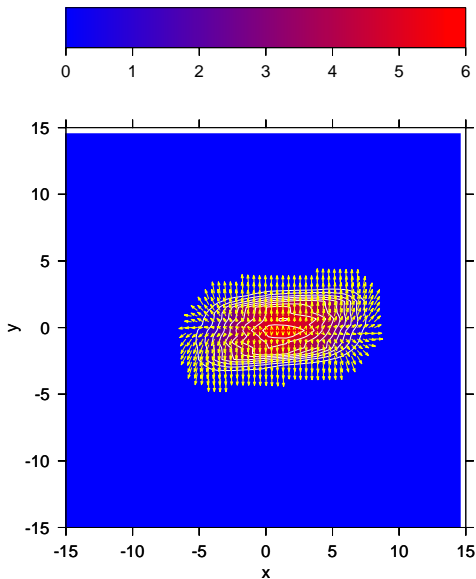
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$t = 1.8 \text{ fm}/c$



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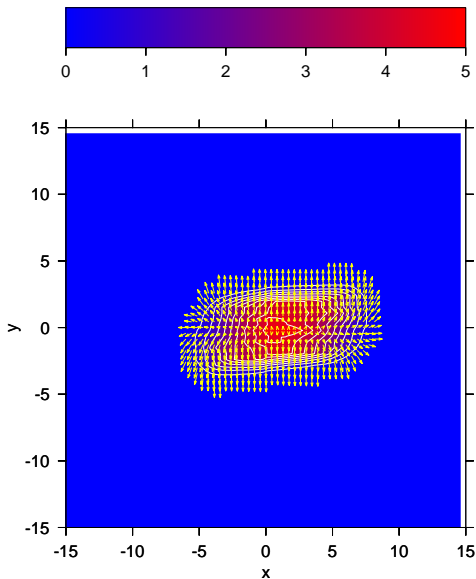
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$t = 2.1 \text{ fm}/c$

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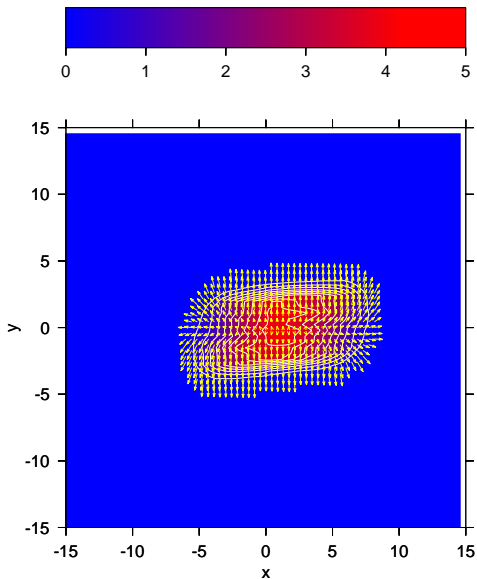
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$t = 2.4 \text{ fm}/c$

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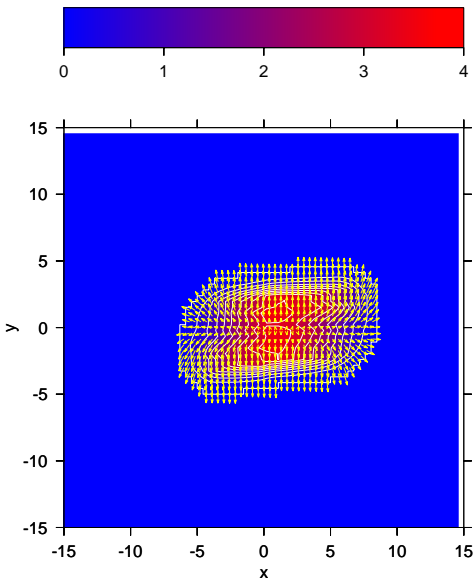
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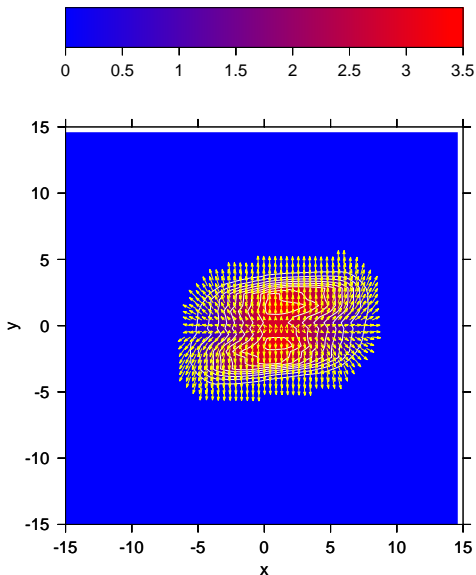
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$t = 3.0 \text{ fm}/c$

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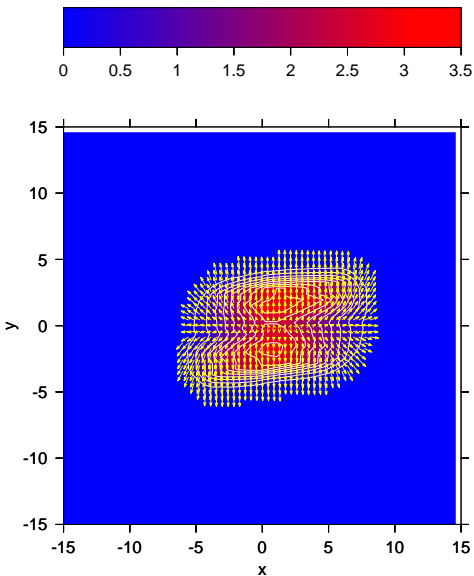
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$t = 3.3 \text{ fm}/c$

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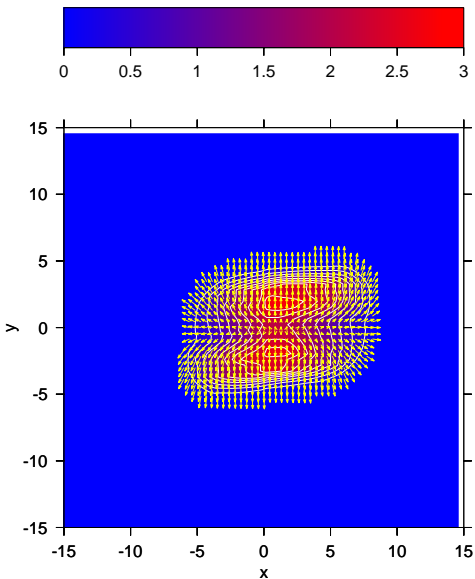
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$t = 3.6 \text{ fm}/c$

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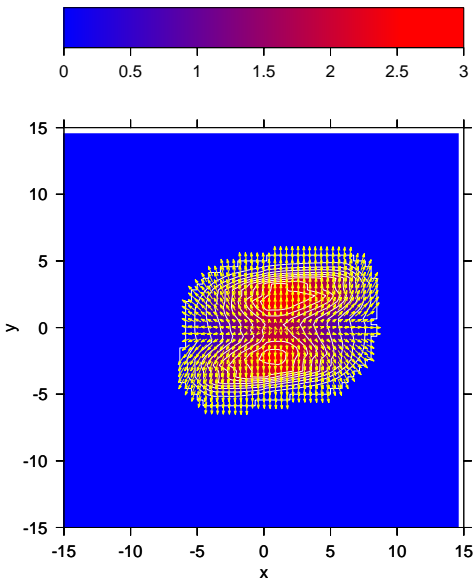
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$t = 3.9 \text{ fm}/c$

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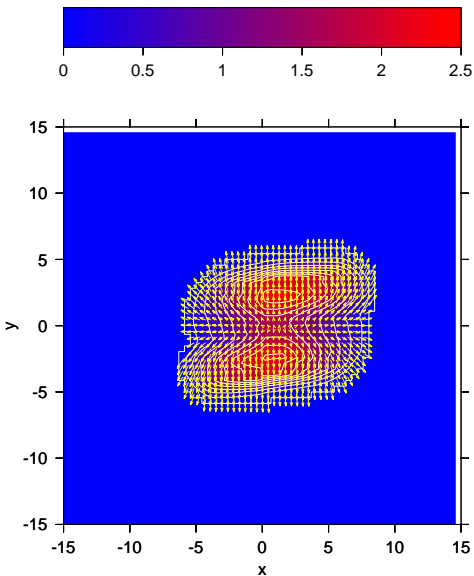
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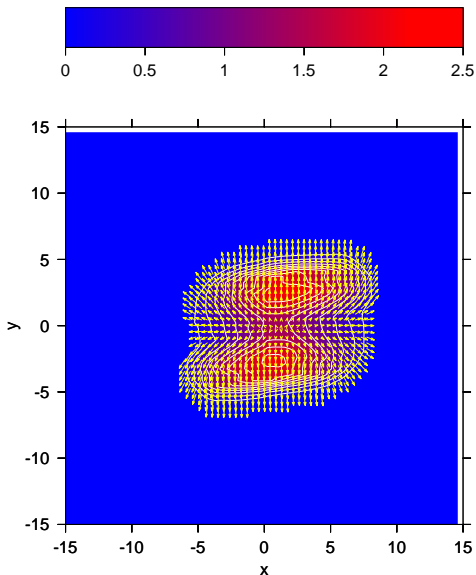
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$t = 4.5 \text{ fm}/c$

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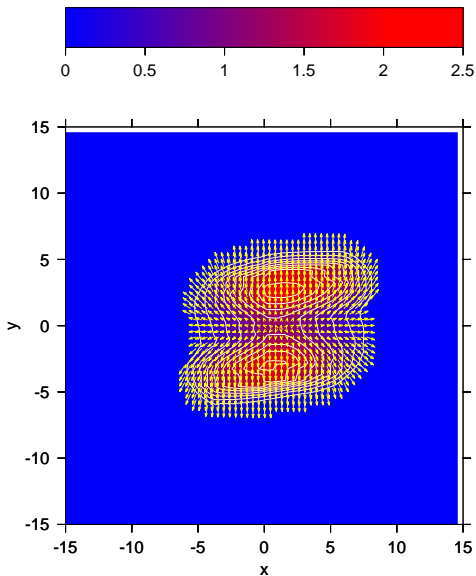
**Fireball evolution**

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$t = 4.8 \text{ fm}/c$

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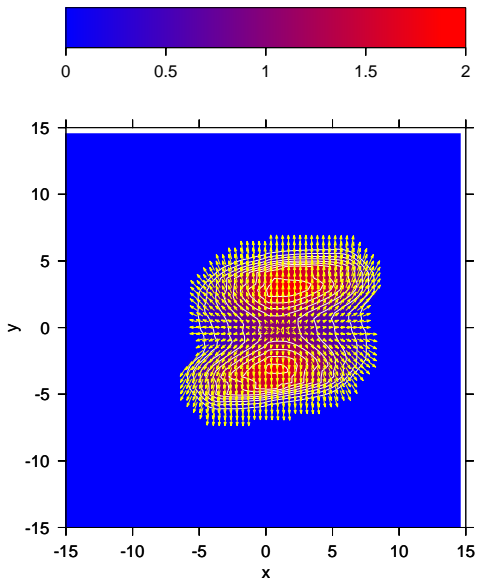
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$t = 5.1 \text{ fm}/c$

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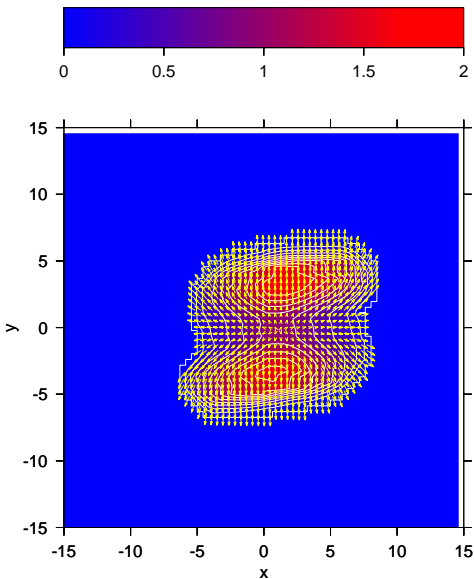
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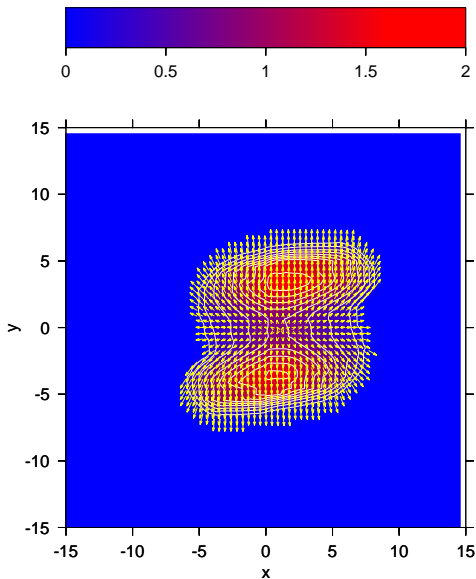
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$t = 5.7 \text{ fm}/c$

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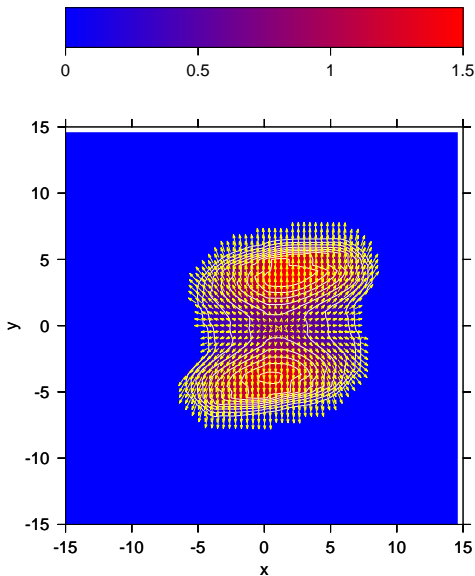
**Fireball evolution**

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$t = 6.0 \text{ fm}/c$

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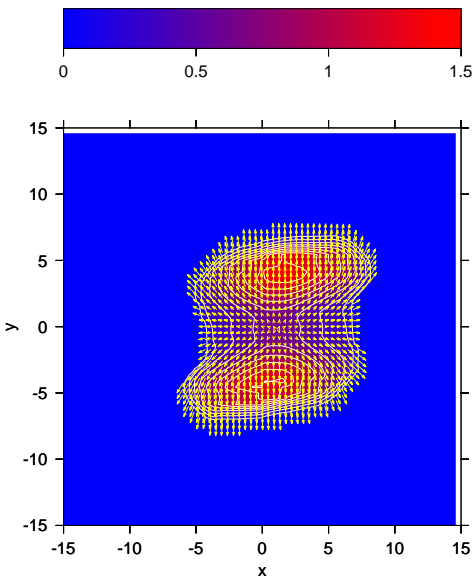
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$t = 6.3 \text{ fm}/c$

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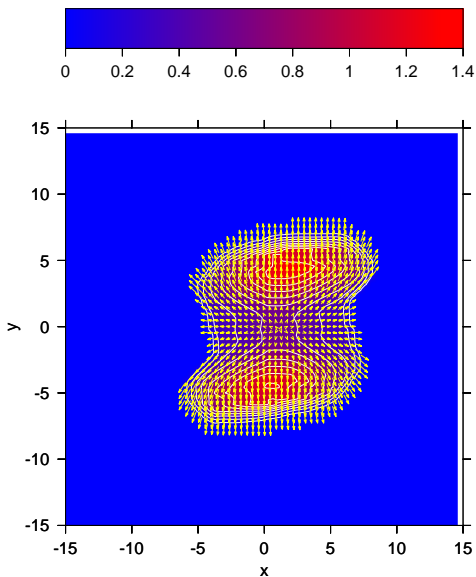
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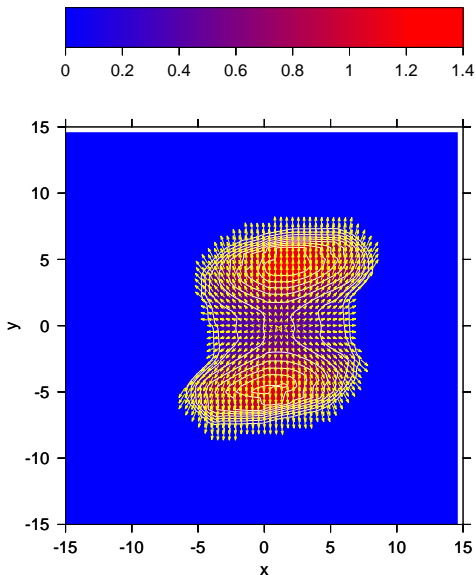
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$t = 6.9 \text{ fm}/c$

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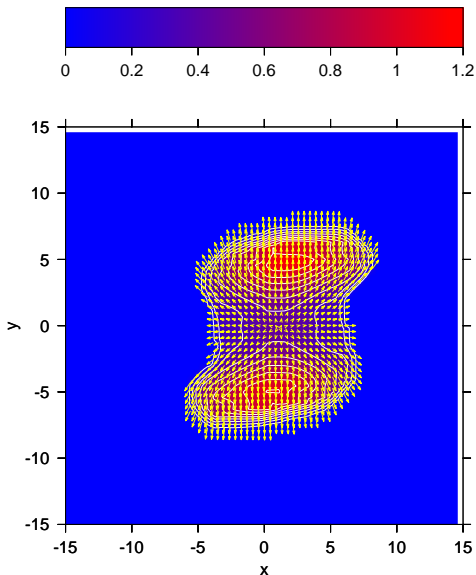
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$t = 7.2 \text{ fm}/c$

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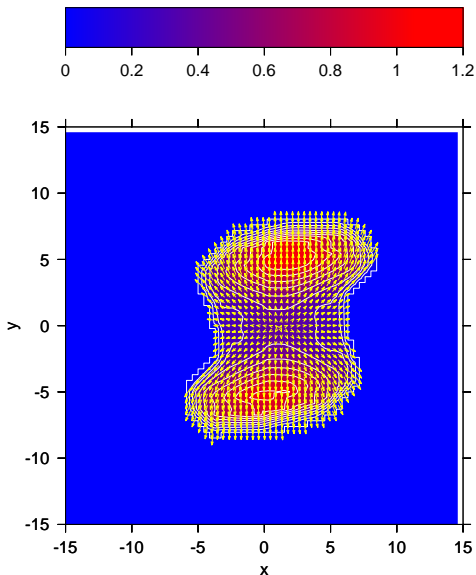
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$t = 7.5 \text{ fm}/c$

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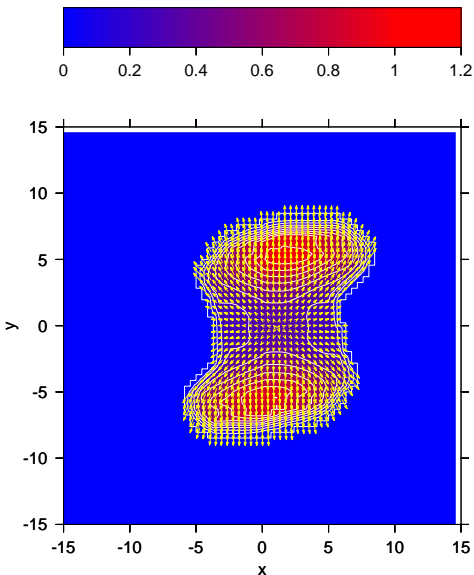
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$t = 7.8 \text{ fm}/c$

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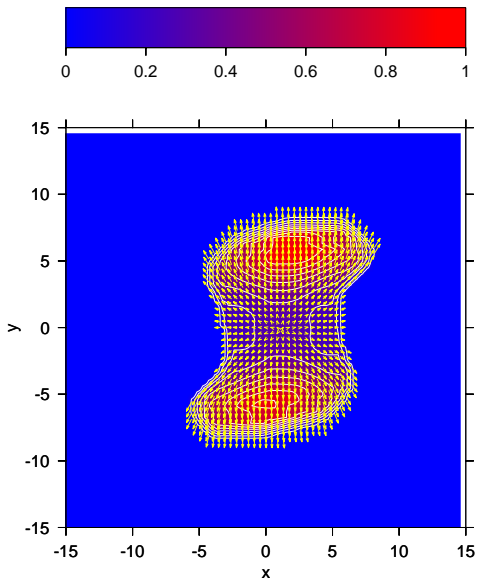
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$t = 8.1 \text{ fm}/c$

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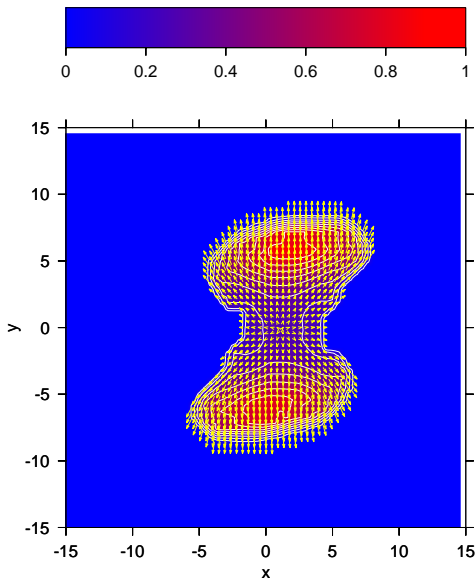
**Fireball evolution**

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$t = 8.4 \text{ fm}/c$

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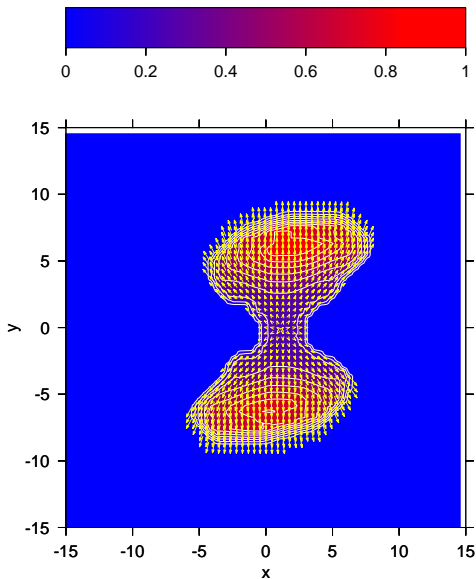
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$t = 8.7 \text{ fm}/c$

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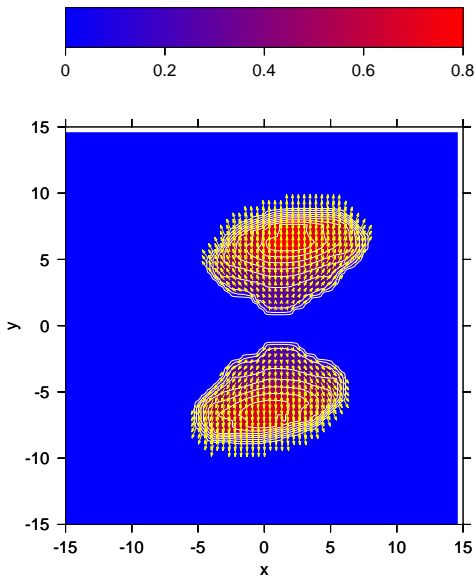
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$t = 9.0 \text{ fm}/c$



# Time evolution of the fireball energy density

Dilepton production from hydrodynamically expanding fireball

V. Toneev and V. Skokov

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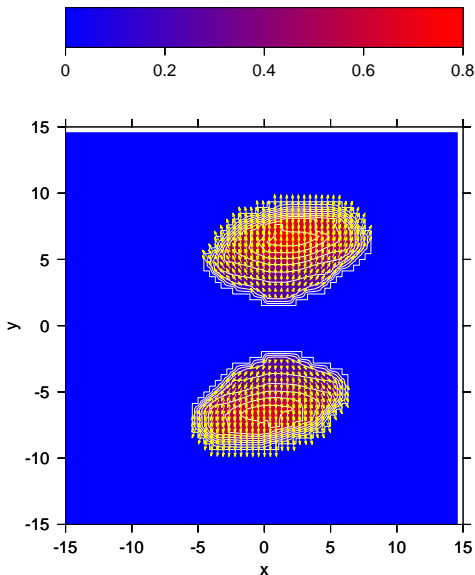
**Fireball evolution**

Evolution in averages

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$t = 9.3 \text{ fm}/c$

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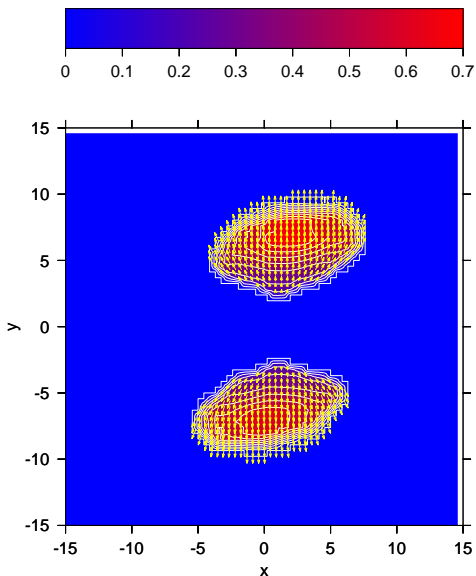
**Fireball evolution**

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$t = 9.6 \text{ fm}/c$

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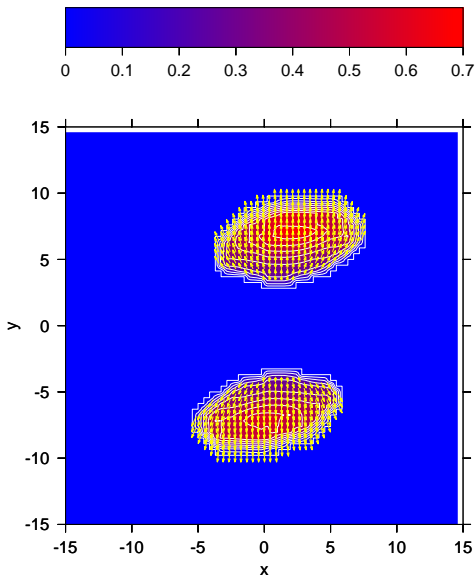
**Fireball evolution**

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$t = 9.9 \text{ fm}/c$

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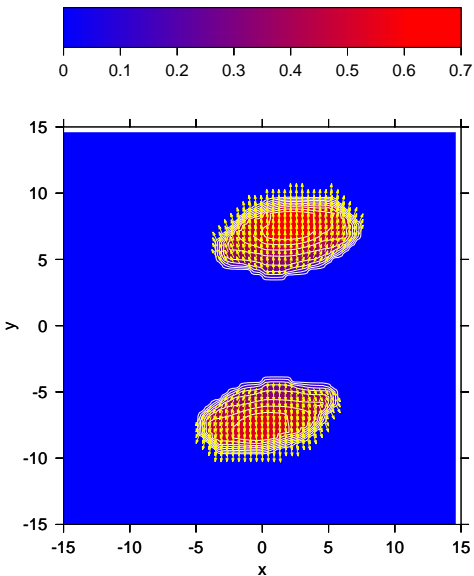
**Fireball evolution**

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$t = 10.2 \text{ fm}/c$

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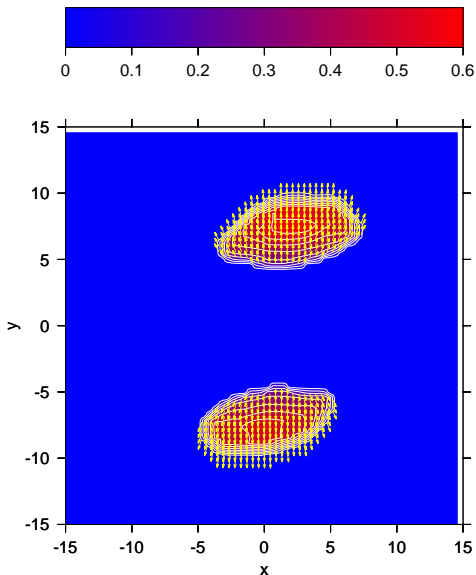
**Fireball evolution**

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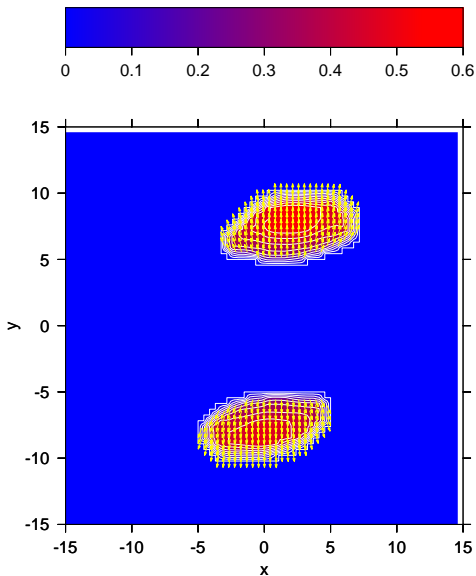
**Fireball evolution**

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$t = 10.8 \text{ fm}/c$

# Time evolution of the fireball energy density

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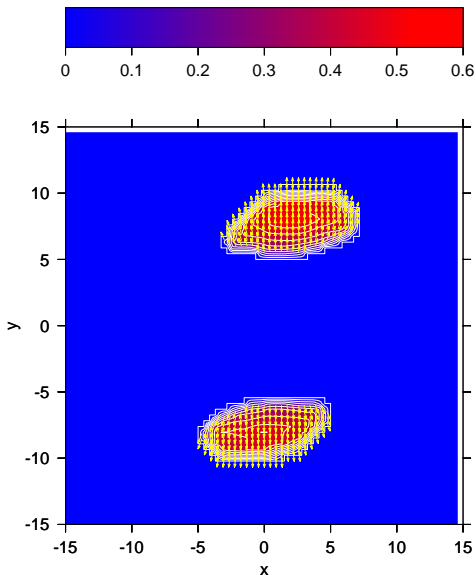
**Fireball evolution**

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$t = 11.1 \text{ fm}/c$

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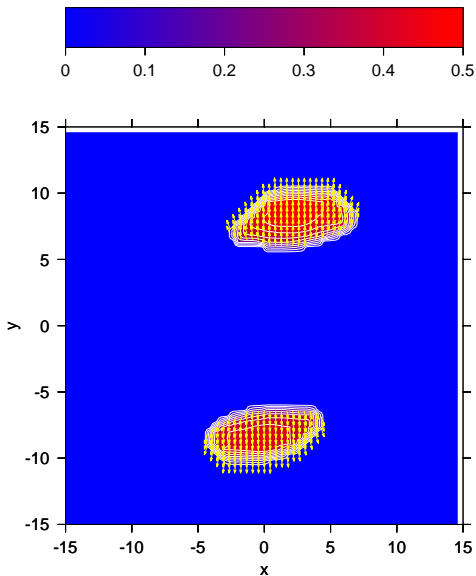
**Fireball evolution**

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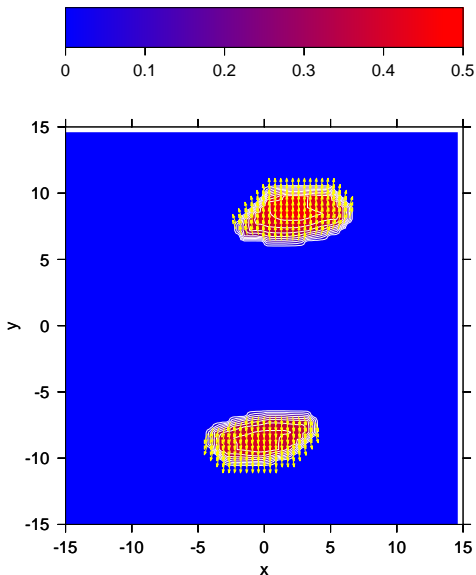
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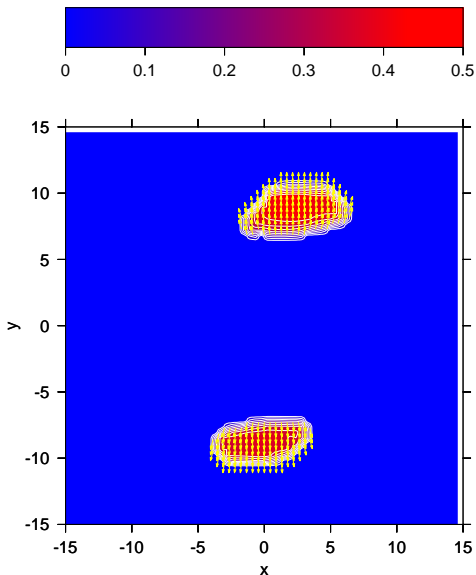
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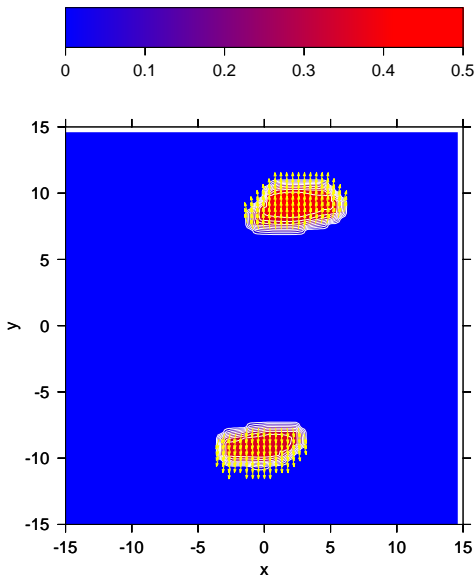
**Fireball evolution**

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$t = 12.3 \text{ fm}/c$

# Time evolution of the fireball energy density

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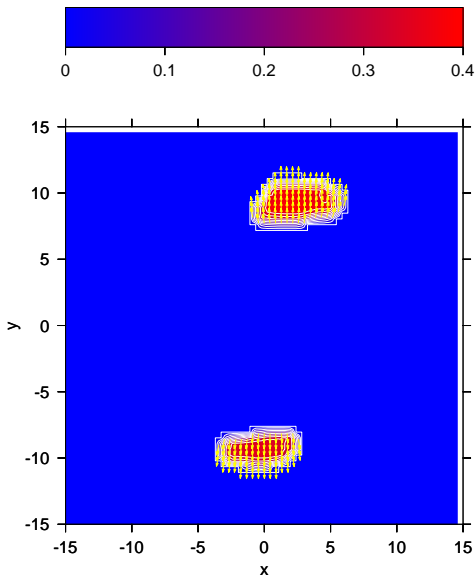
**Fireball evolution**

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$t = 12.6 \text{ fm}/c$

# Time evolution of the fireball energy density

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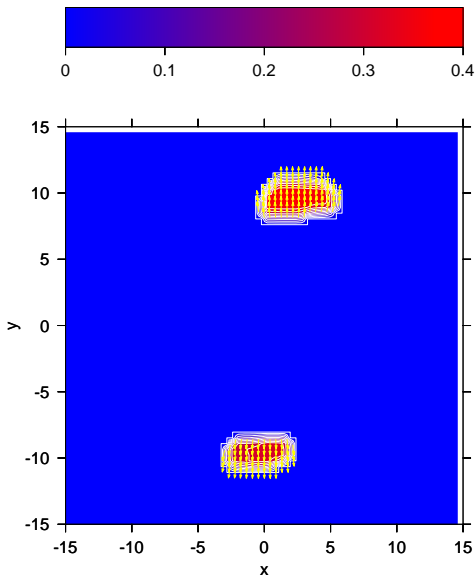
**Fireball evolution**

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$t = 12.9 \text{ fm}/c$

# Average quantities

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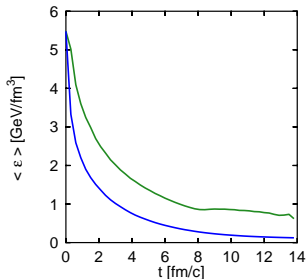
Evolution in averages

Hadron observable

Dileptons

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- Pb + Pb (158 AGeV)*
- Average energy density (mixed phase EoS, mixed phase EoS without freeze-out)



# Average quantities

Dilepton  
production  
from hydrody-  
namically  
expanding  
fireball

V. Toneev  
and V.  
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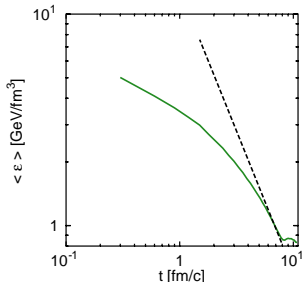
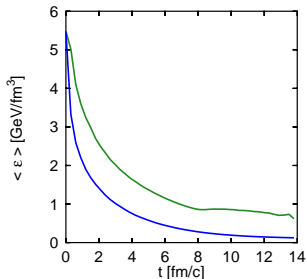
Dileptons

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$Pb + Pb$  (158 AGeV)

- Average energy density  
(mixed phase EoS,  
mixed phase EoS without  
freeze-out)

and comparison with the  
Bjorken scaling regime  
(ultrarelativistic ideal gas  
EoS:  $\varepsilon = \frac{1}{3}P$ , dashed line  
with the slope  $-4/3$ )



# Average quantities

Dilepton  
production  
from hydrody-  
namically  
expanding  
fireball

V. Toneev  
and V.  
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$Pb + Pb$  (158 AGeV)

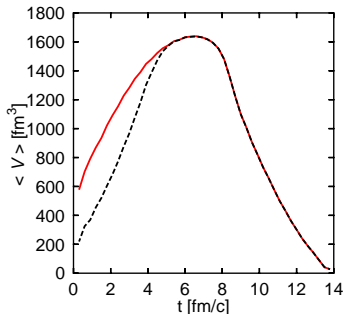
- Evolution of the system volume

There are **two** stages:  
"pure" expansion and  
freeze-out

**Hadron fraction** (dashed)  
is defined by condition

$$N_{quarks}^H > N_{quarks}^{Q+G}$$

expansion ↓

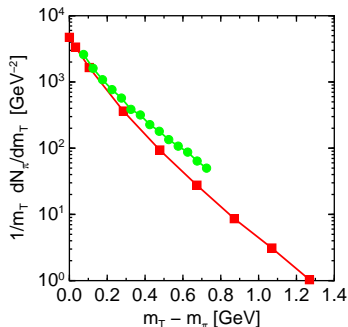
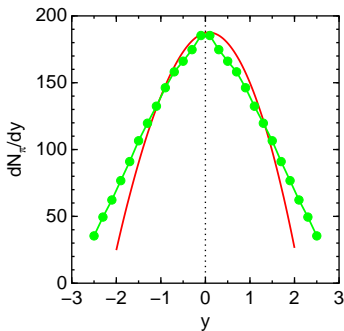


↑ freeze-out



# Comparison with observable

Pion distributions from central ( $\frac{\sigma_{trig}}{\sigma_{tot}} = 7\%$ )  
 $Pb + Pb$  (158 AGeV) collisions



S. V. Afanasiev *et al.* (NA49 Collab.), Phys. Rev. C **66**, 054902 (2002) [nucl-ex/0205002]

# Dileptons; time slices

Dilepton production from hydrodynamically expanding fireball

$$\frac{d^8 N_{ee}}{d^4 x d^4 q} = \frac{\alpha^2}{12\pi^4} \frac{R(q, T, \mu_b)}{e^{q_0/T} - 1}$$

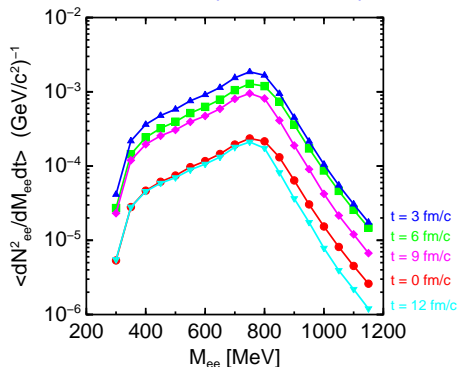
$$q^2 = M^2 = q_0^2 - \vec{q}^2$$



(C. Gale and J. I. Kapusta,

Nucl. Phys. B 357 (1991) 65)

Pb + Au (158 AGeV)



with CERES acceptance

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# Dileptons; comparison with CERES data

Dilepton production from hydrodynamically expanding fireball

V. Toneev and V. Skokov

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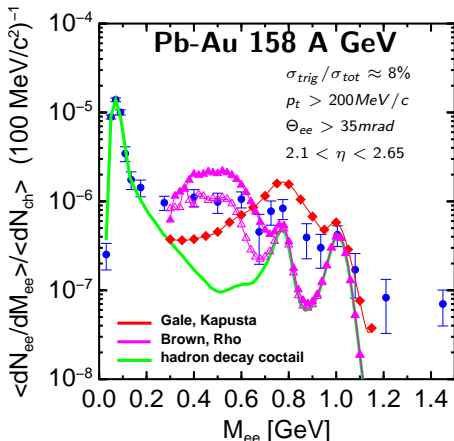
Outlook

CERES/NA45 Collaboration, J. Phys. G 30 (2004) S1007, J. Phys. G 30 (2004)

2027, PRELIMINARY

S.Yurevich (private communication)

G.E. Brown and M. Rho, Phys. Rep. 363 (2002) 85 ( $k = 1/6$ )



$$\frac{d^8 N_{ee}}{d^4 x d^4 q} = \frac{\alpha^2}{48\pi^4} \left(1 - \frac{4m_\pi^2}{M^2}\right)^{3/2} \frac{|F_\pi(M)|^2}{e^{q_0/T} - 1}, \quad F_\pi(M) = \frac{m_\rho^4}{(M^2 - m_\rho^2)^2 - m_\rho^2 \Gamma_\rho^2},$$

$$m_\rho \sim \left(1 - \frac{T^2}{T_c^2}\right)^k \left(1 - 0.18 \frac{n_B}{n_0}\right) m_{\rho 0}, \quad \Gamma_\rho^2 = \Gamma_\rho^2(m_\rho^2)$$

- Influence of full dynamics on the dilepton yield  $\Rightarrow$  field calculations vs. averages, EoS, freeze-out

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- **Phase transition ? Intermediate mass dileptons ?**

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- **Dilepton emission rates at high  $n_B$  ( $\mu_B$ )?**

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- **Other sources  $\Rightarrow$  mixed phase**

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- Other sources  $\Rightarrow$  mixed phase
- **Other observable  $\Rightarrow p_T$  selection,  $M$ -selected  $p_T$  distributions**



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- Other observable  $\Rightarrow p_T$  selection,  $M$ -selected  $p_T$  distributions
- $C + C$  (DLS problem), hydrodynamics ??