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Importance of Fluctuations of cross section in muon-catalysed $t\text{-}t$ fusion

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$\mu d-t$ fusion

- Recent Experimental Observation: T dependence of λ_c and ω^{eff}

cf. N. Kawamura, et al., Phys.Rev.Lett. 90, 043401(2003)

“Discovery of Temperature-Dependent Phenomena of Muon-Catalyzed Fusion in Solid Deuterium and Tritium Mixtures”

- due to the resonant molecular formation in



- Importance of the resonant state of $dt\mu$ mesomolecules



$\mu t\text{-}t$ fusion



cf. T. Matsuzaki, Phys.Lett.B 557, 176(2003)

"Evidence for strong n- α correlations in the t+t reaction proved by the neutron energy distribution of muon catalysed t-t fusion"

- No shallow bound state in $tt\mu$
- λ_c can be estimated using "in-flight" fusion model?
⇒ No T dependence



- the μ cycling rate and the reaction rate

$$\lambda = \rho_{LH} < \sigma v >$$

$$\rho_{LH} = 4.25 \times 10^{22} cm^{-2}$$

as a function of T

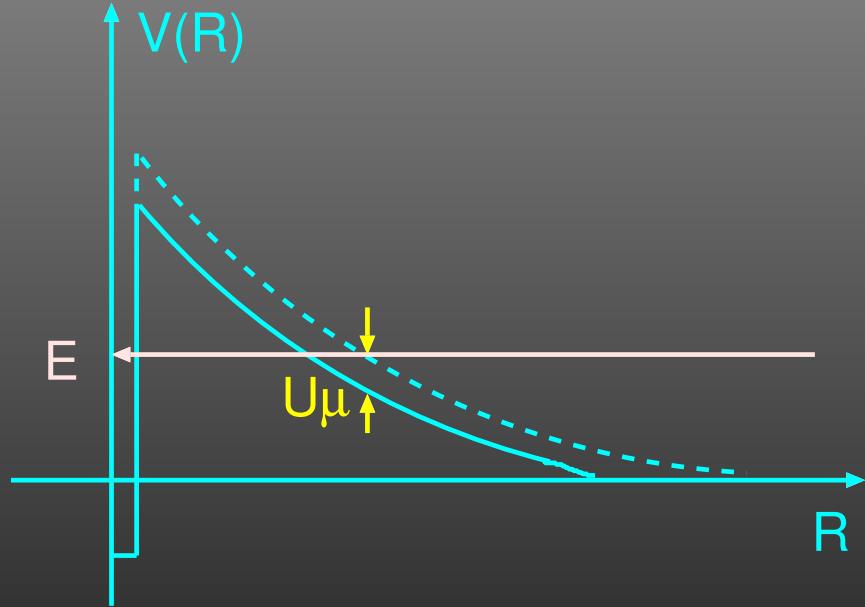
- Large Enhancement of the cross section by μ^-

$$\sigma(E) = f_\mu \sigma_0(E)$$

- Fluctuation of $f_\mu \Leftarrow$ Chaotic dynamics of the 3-body system



Screening Potential



U_μ : Screening Potential

$$f_\mu \equiv \frac{\sigma(E)}{\sigma_0(E)} = \frac{\sigma_0(E + U_\mu)}{\sigma_0(E)}$$

$$\sim \exp\left\{ \pi\eta(E) \frac{U_\mu}{E} \right\}$$

$$U_\mu \sim \frac{E}{\pi\eta(E)} \log f$$



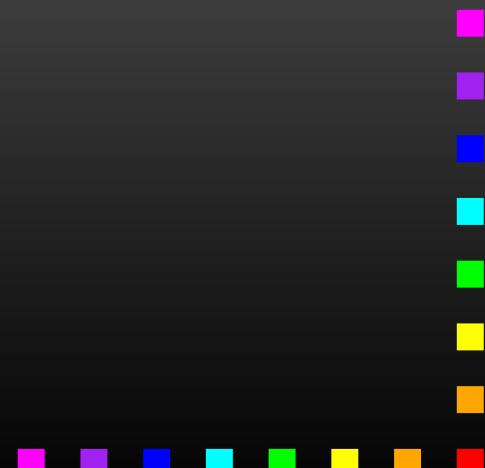
Constrained Molecular Dynamics (CoMD)

S.Kimura and A.Bonasera, Phys. Rev. A 72, 014703 (2005)

Lagrange multiplier method for constraints

$$\mathcal{L} = \sum_i \frac{\mathbf{p}_i^2}{2m_i} - \sum_{i,j(\neq i)} U(\mathbf{r}_{ij}) + \sum_{i,j(\neq i)} \lambda_i \left(\frac{\mathbf{r}_{ij}\mathbf{p}_{ij}}{\xi\hbar} - 1 \right)$$

$$\mathbf{r}_{ij} = |\mathbf{r}_i - \mathbf{r}_j|; \quad \mathbf{p}_{ij} = |\mathbf{p}_i - \mathbf{p}_j|$$



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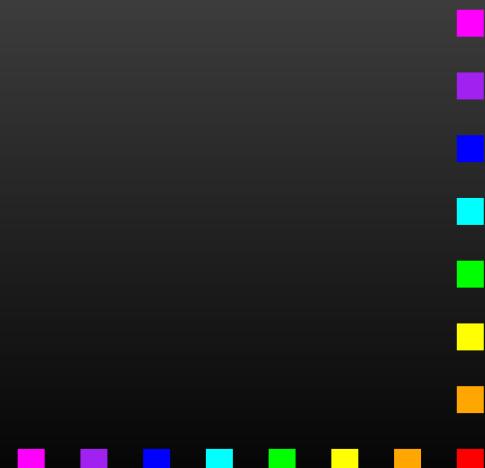
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Variational calculus leads Hamilton Equation with Constraint:

$$\frac{d\mathbf{r}_i}{dt} = \frac{\mathbf{p}_i}{m_i} + \frac{\lambda_i \mathbf{r}_{ij}}{\xi\hbar} \frac{\partial \mathbf{p}_{ij}}{\partial \mathbf{p}_i}$$

$$\frac{d\mathbf{p}_i}{dt} = -\nabla_{\mathbf{r}} U(\mathbf{r}_i) - \frac{\lambda_i \mathbf{p}_{ij}}{\xi\hbar} \frac{\partial \mathbf{r}_{ij}}{\partial \mathbf{r}_i}$$

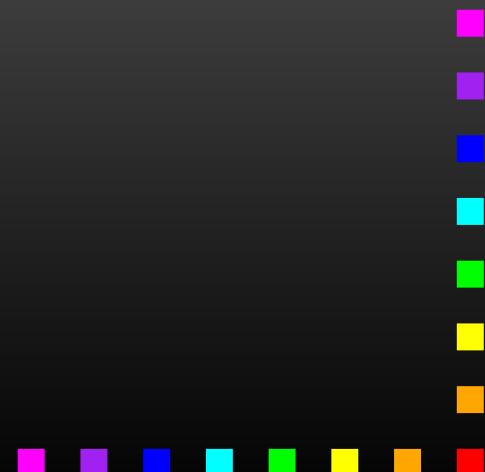


Tunneling process

$$\frac{d\mathbf{r}_i}{dt} = \frac{\mathbf{p}_i}{m_i}; \quad \frac{d\mathbf{p}_i}{dt} = -\nabla_{\mathbf{r}} U(\mathbf{r}_i)$$

Collective coordinates and momenta

$$\mathbf{R}^{\text{coll}} \equiv \mathbf{r}_P - \mathbf{r}_T; \quad \mathbf{P}^{\text{coll}} \equiv \mathbf{p}_P - \mathbf{p}_T; \quad \mathbf{F}_P^{\text{coll}} \equiv \dot{\mathbf{P}}^{\text{coll}}$$



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$$\frac{d\mathbf{r}_{T(P)}^{\mathcal{S}}}{d\tau} = \frac{\mathbf{p}_{T(P)}^{\mathcal{S}}}{m_{T(P)}}; \quad \frac{d\mathbf{p}_{T(P)}^{\mathcal{S}}}{d\tau} = -\nabla_{\mathbf{r}} U(\mathbf{r}_{T(P)}^{\mathcal{S}}) - 2\mathbf{F}_{T(P)}^{\text{coll}}$$

Tunneling penetrability: $\Pi(E) = (1 + \exp(2\mathcal{A}(E)/\hbar))^{-1}$

$$\mathcal{A}(E) = \int_{\mathbf{r}_b}^{\mathbf{r}_a} \mathbf{P}^{\text{coll}} d\mathbf{R}^{\text{coll}}$$

without muon $\Rightarrow \Pi_0(E)$

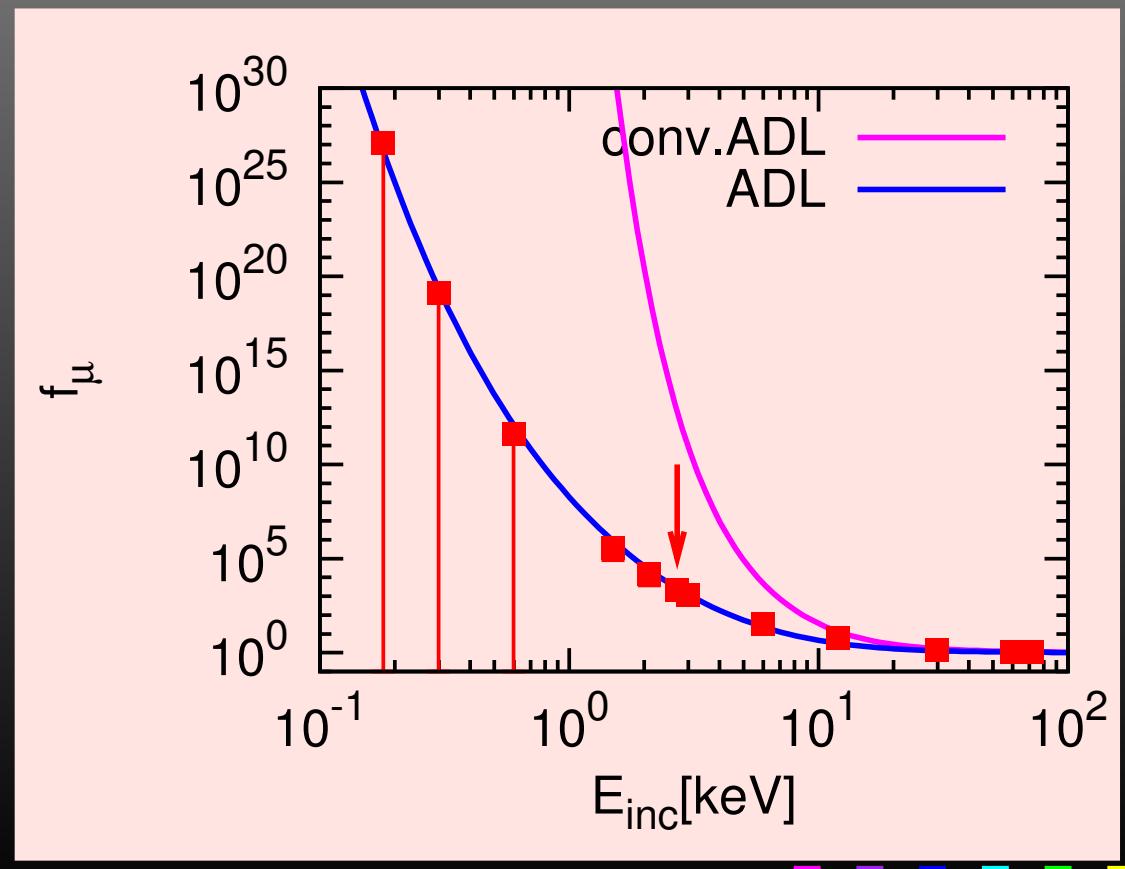
Enhancement factor: $f_\mu = \Pi(E)/\Pi_0(E)$



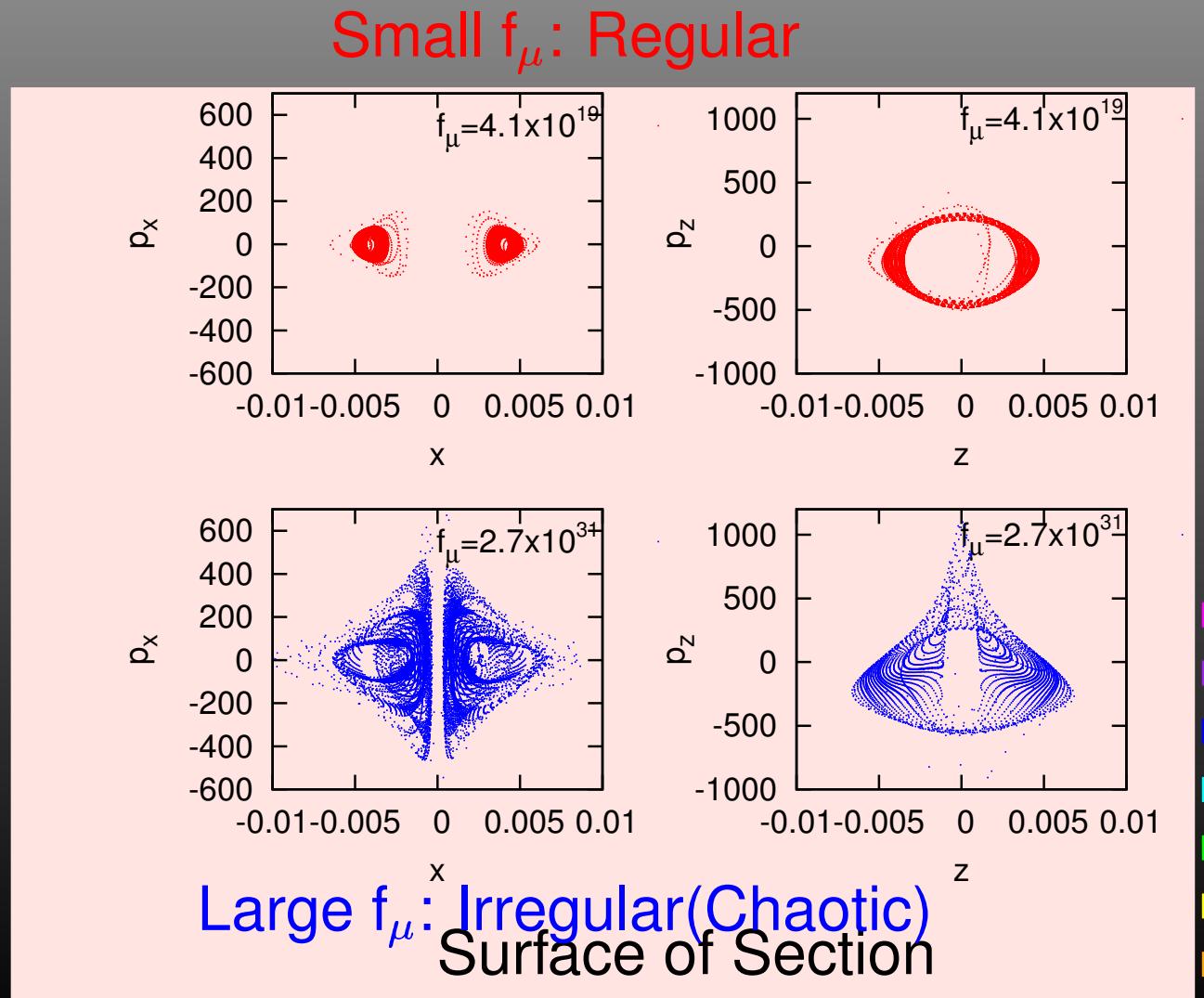
Enhancement factor

and

Variance $\Sigma = \sqrt{\bar{f}_\mu^2 - (\bar{f}_\mu)^2}$

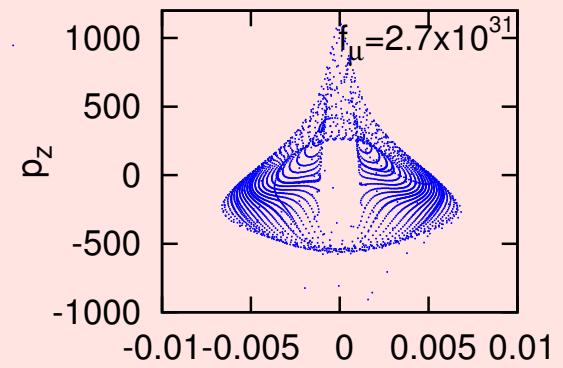
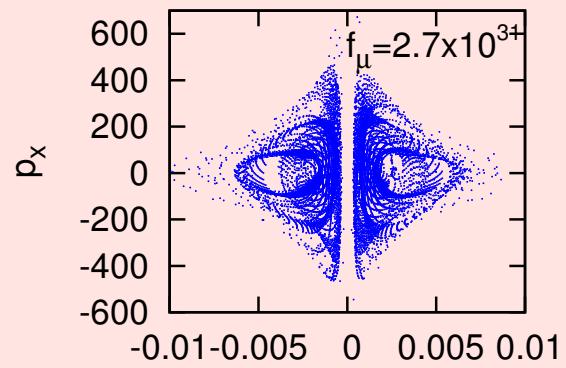
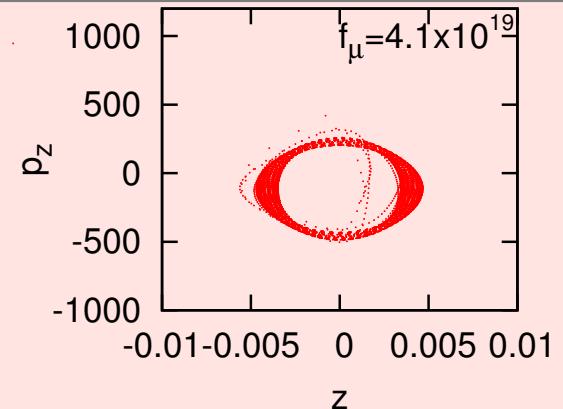
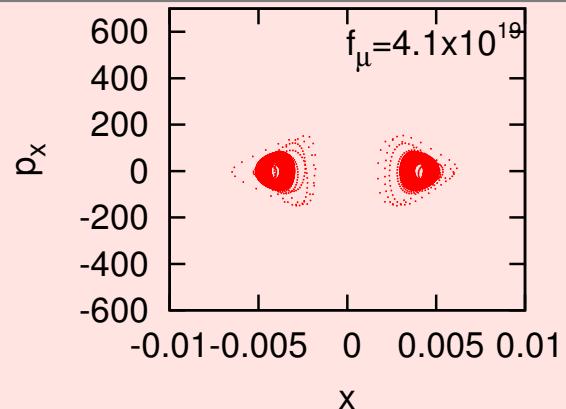


Muonic Motion and Enhancement factor as an order parameter



Muonic Motion and Enhancement factor as an order parameter

Small f_μ : Regular (\Leftarrow Long TR)

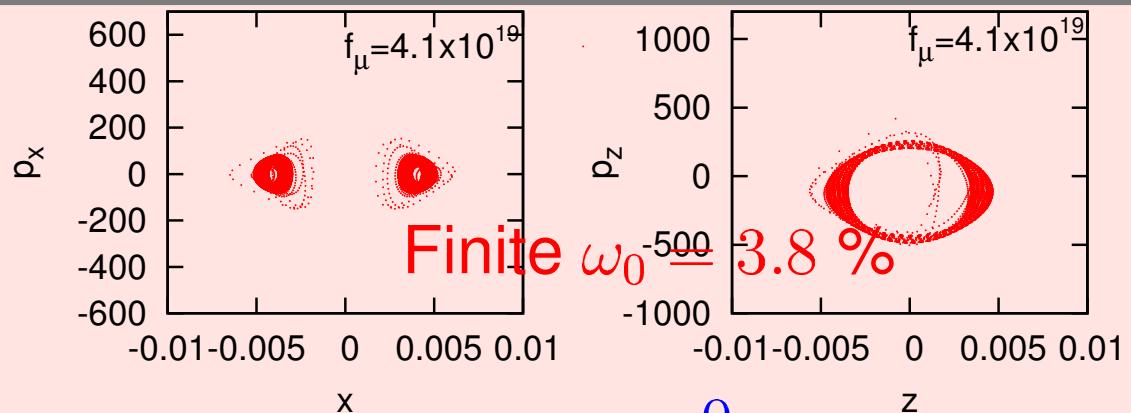


Large f_μ : Irregular(Chaotic) (\Leftarrow Short TR)
Surface of Section



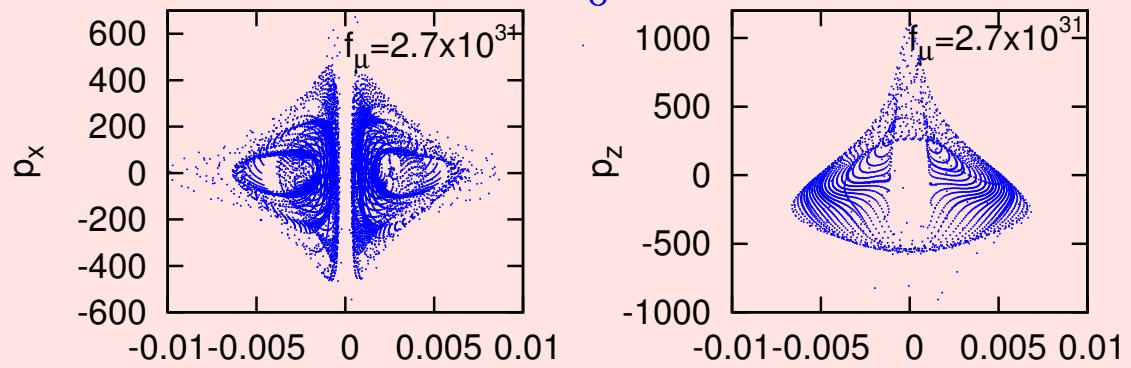
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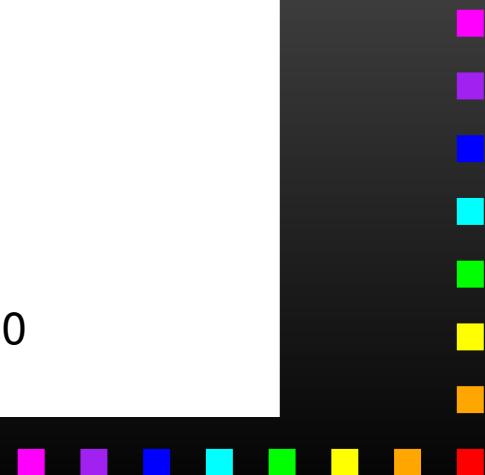
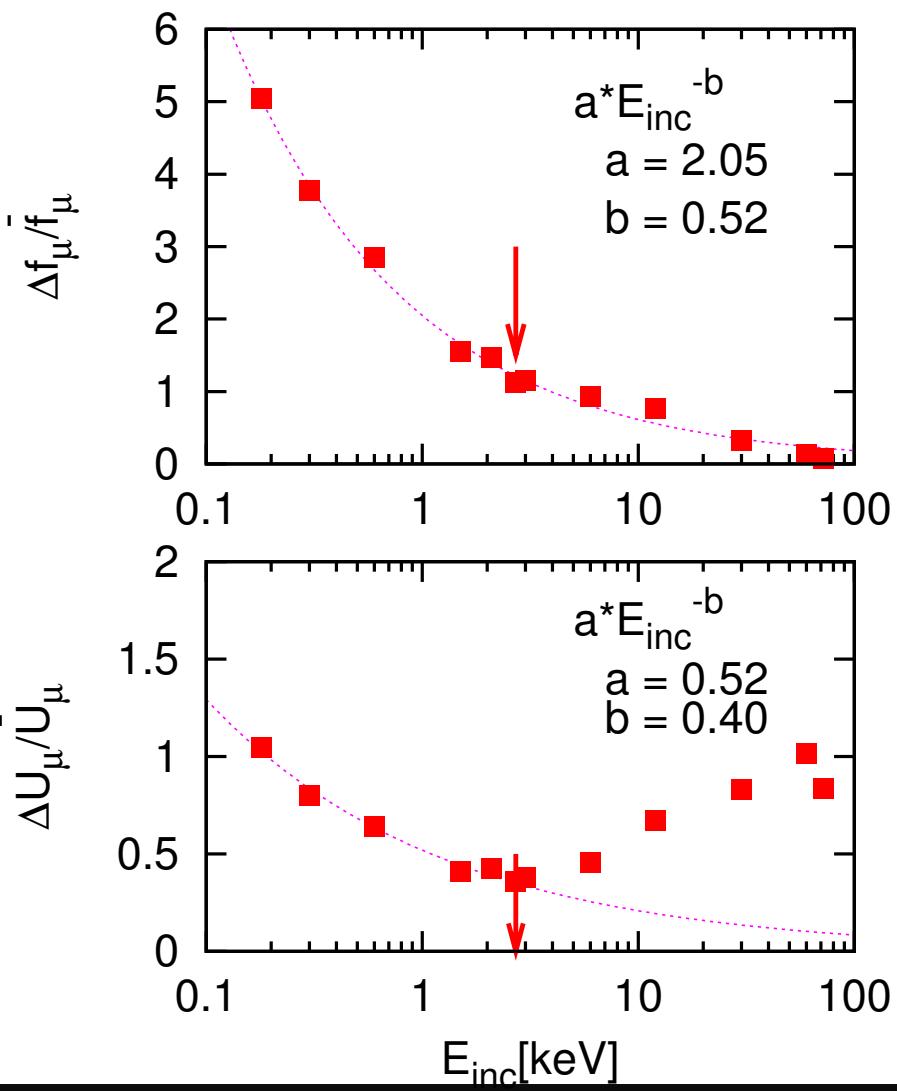
Finite $\omega_0 = 3.8\%$

$\omega_0 = 0$



Large f_μ : Irregular(Chaotic) (\Leftarrow Short TR)
Surface of Section





Reaction Rate

$$\begin{aligned}\lambda &= \rho \langle \sigma v \rangle \\ &= \rho \int \sigma_0(E + U_\mu) v \Psi(E, T) dE\end{aligned}$$



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$$\begin{aligned}\sigma_0(E) &= \frac{S(E)}{E} e^{-2\pi\eta(E)} \\ S(E) &= 0.20 - 0.32E + 0.476E^2 [\text{MeVb}] \\ &\quad (\text{for the } t + t \text{ reaction})\end{aligned}$$

S. Winkler et al., J. Phys. G 18(1991) L147



Reaction Rate

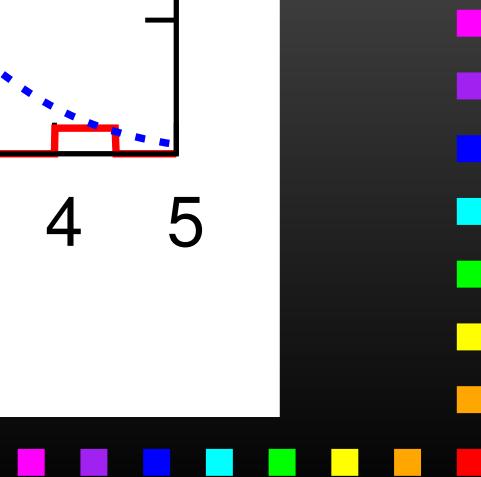
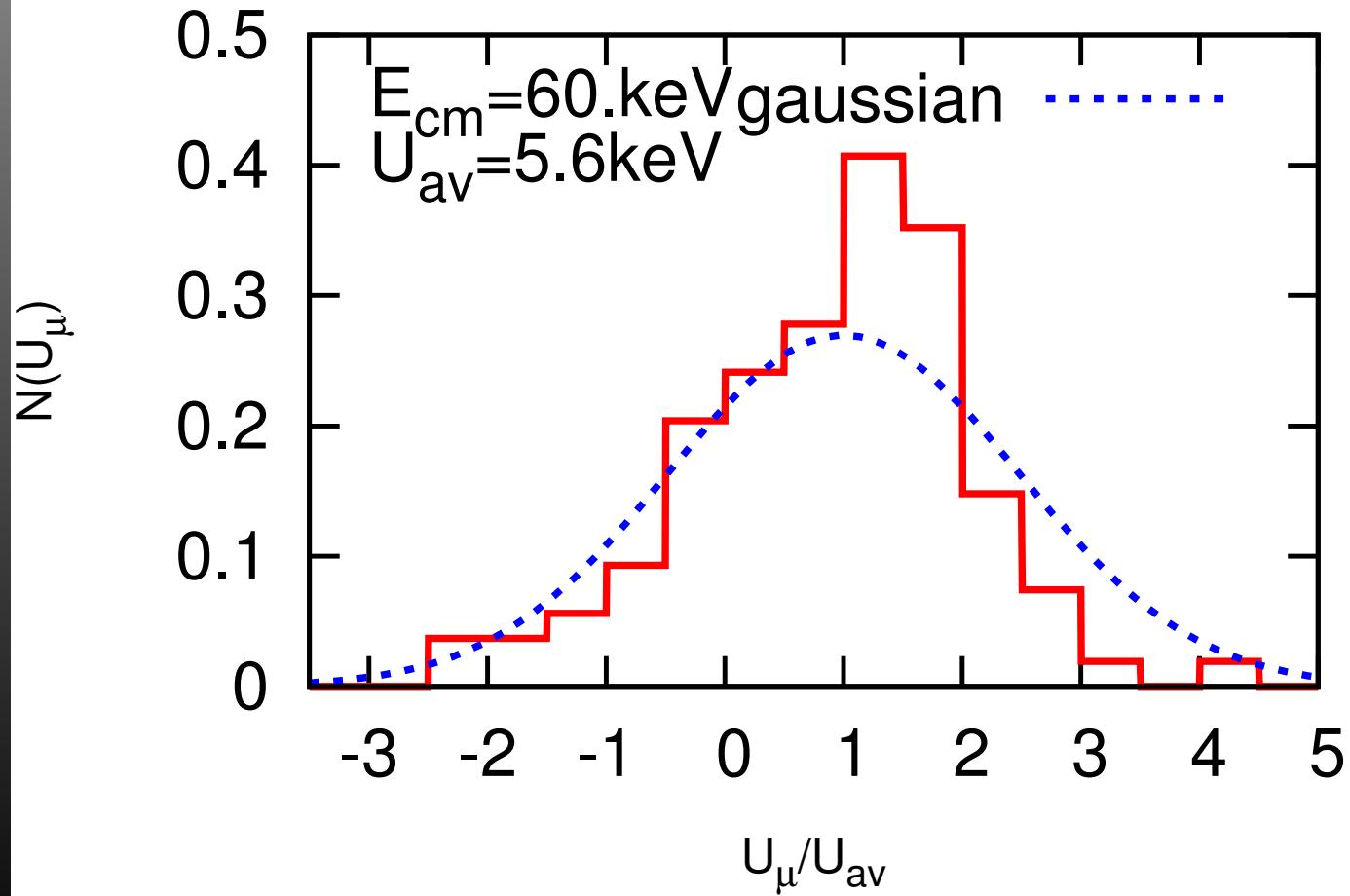
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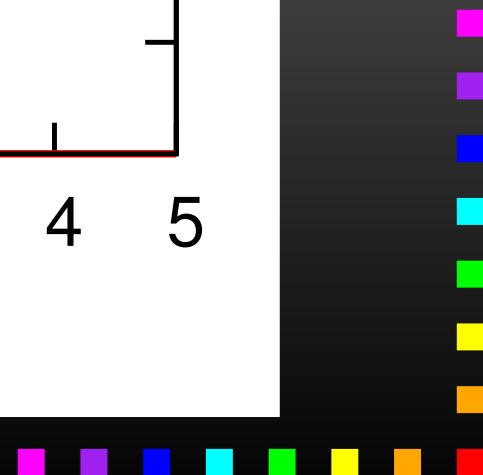
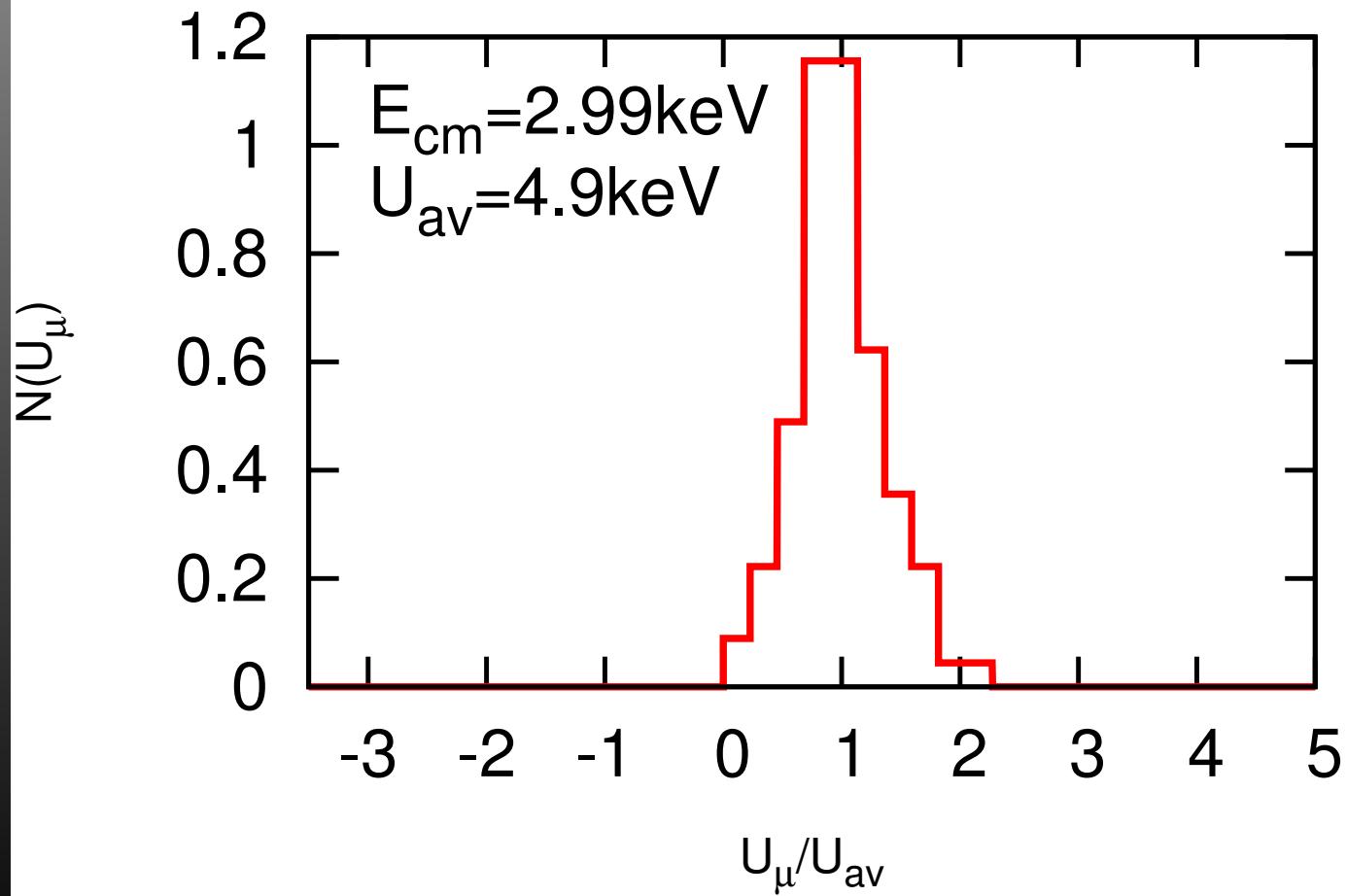
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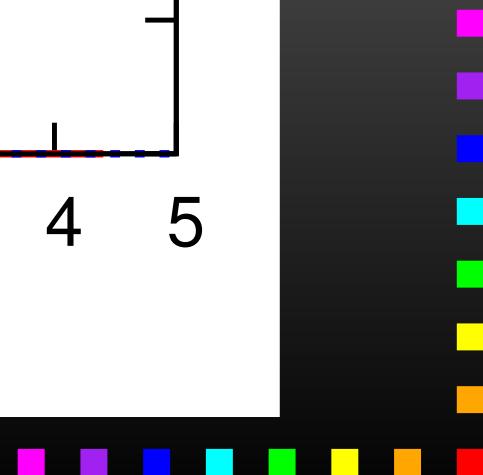
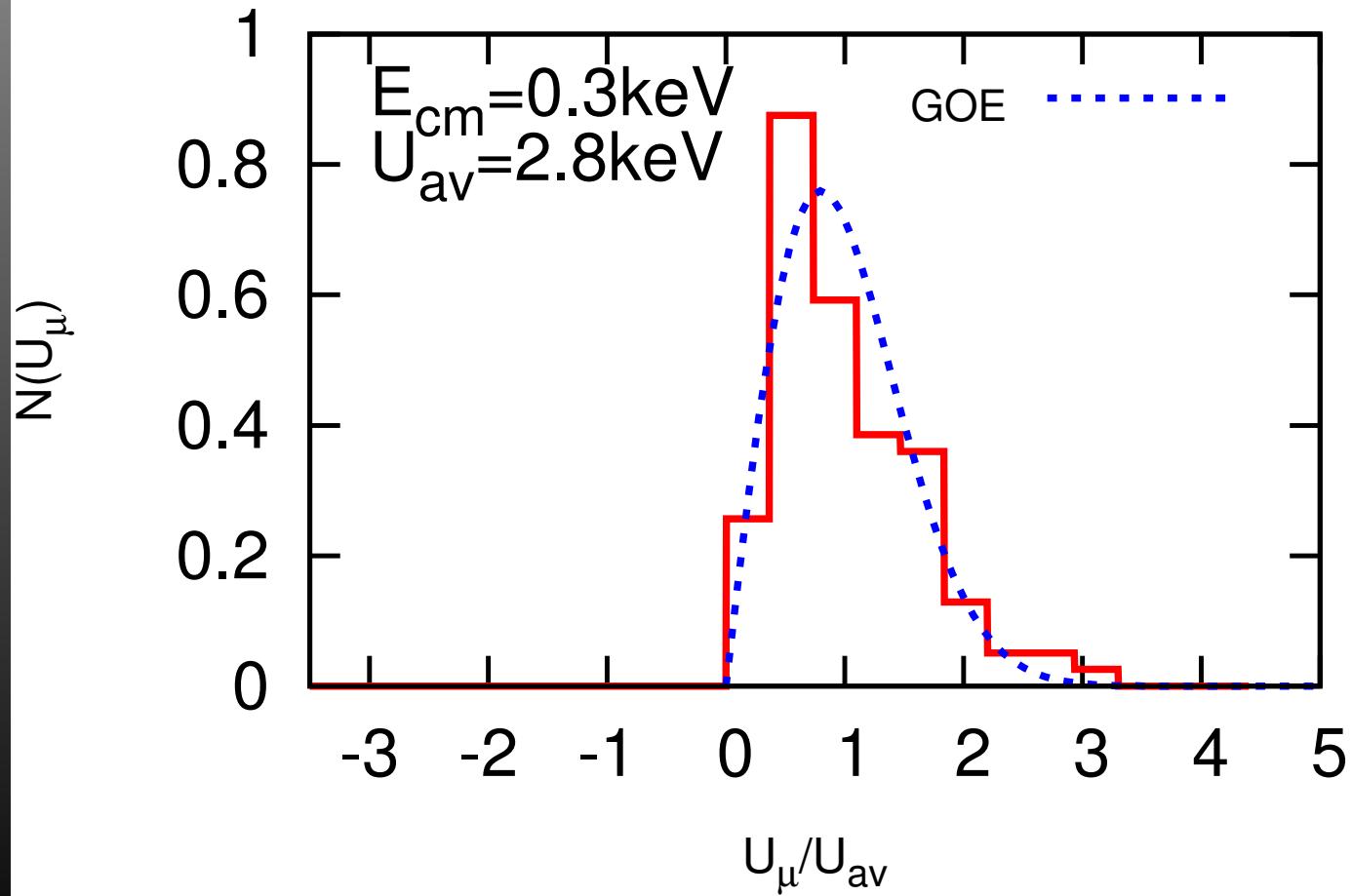
Distributions of U_μ



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Reaction Rate

$$\lambda = \rho \int \sigma_0(E + U_\mu) v \Psi(E, T) N(U_\mu) dE dU_\mu$$



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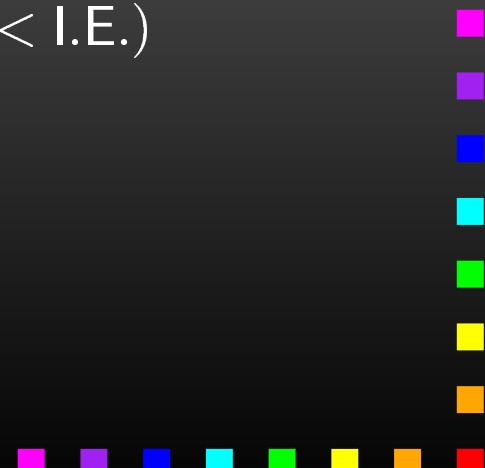
$$\begin{aligned}
 N(U_\mu) &= \frac{1}{\sqrt{2\pi} \Delta U_\mu} \exp\left(-\frac{(U_\mu - \bar{U}_\mu)^2}{2\Delta U_\mu^2}\right) \quad (E > I.E.) \\
 &= \frac{\pi}{2} \left(\frac{U_\mu}{\bar{U}_\mu}\right) \exp\left(-\frac{\pi}{4} \left(\frac{U_\mu}{\bar{U}_\mu}\right)^2\right) \quad (E < I.E.)
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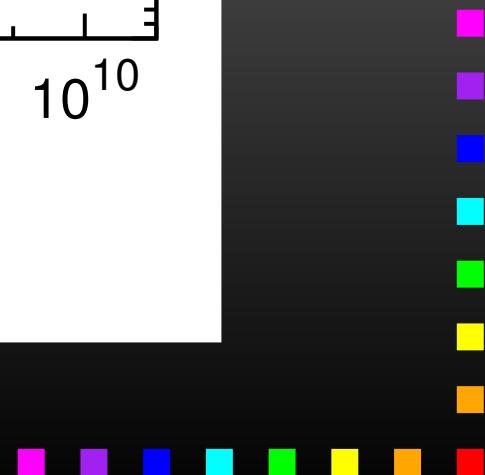
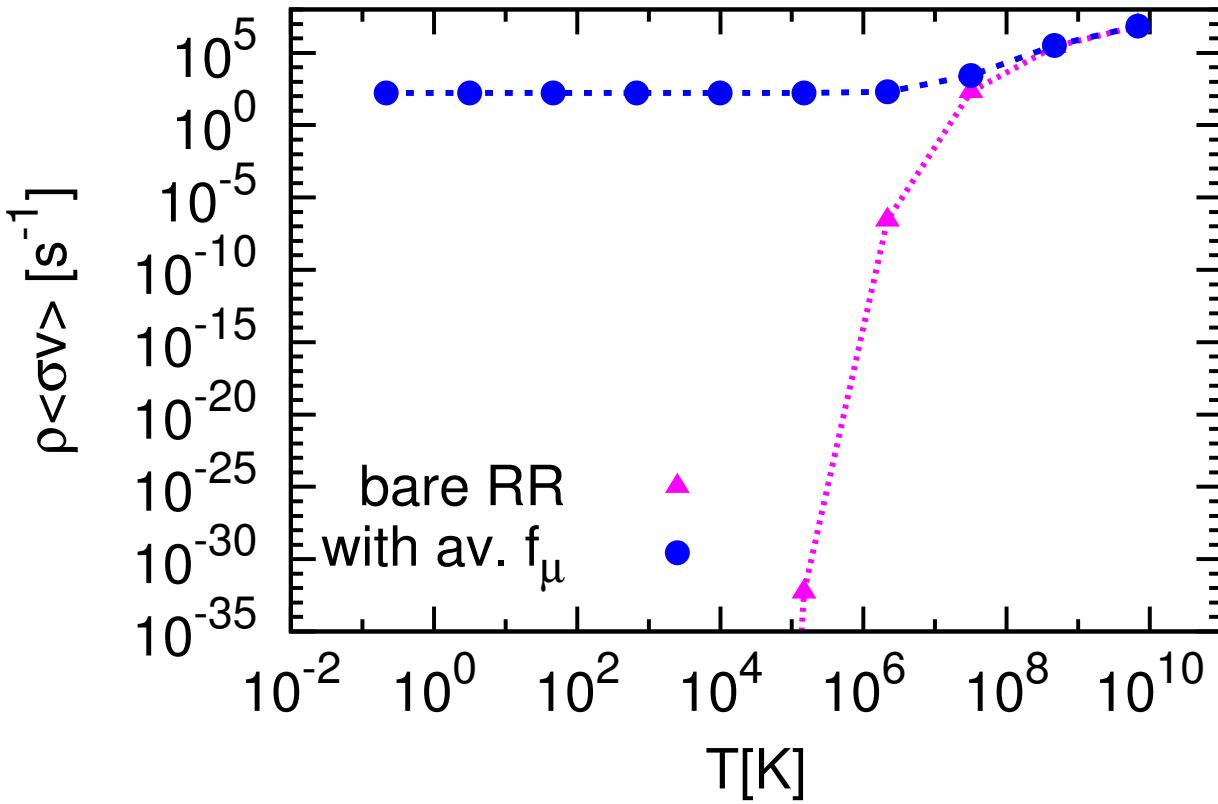
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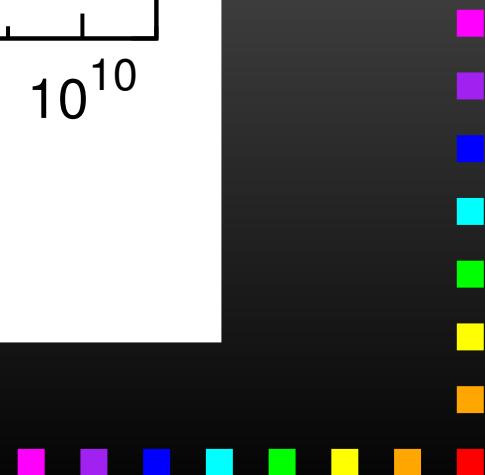
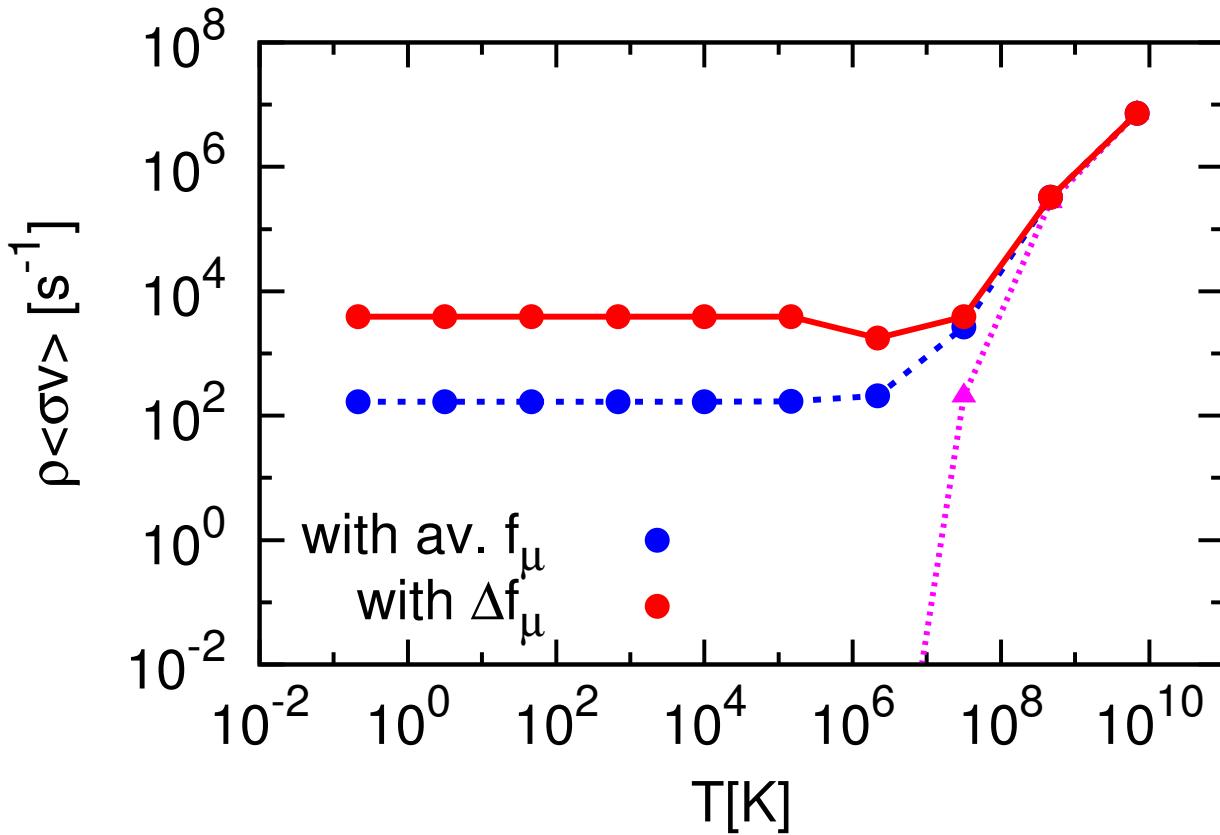
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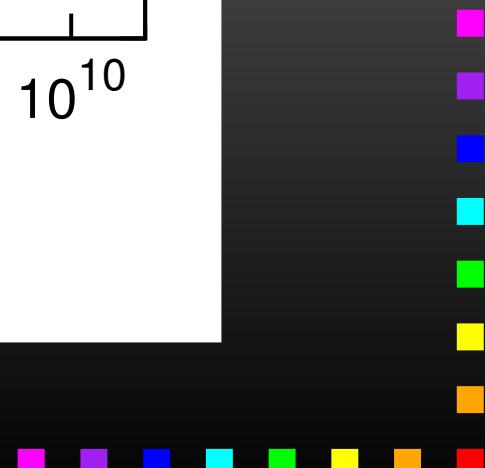
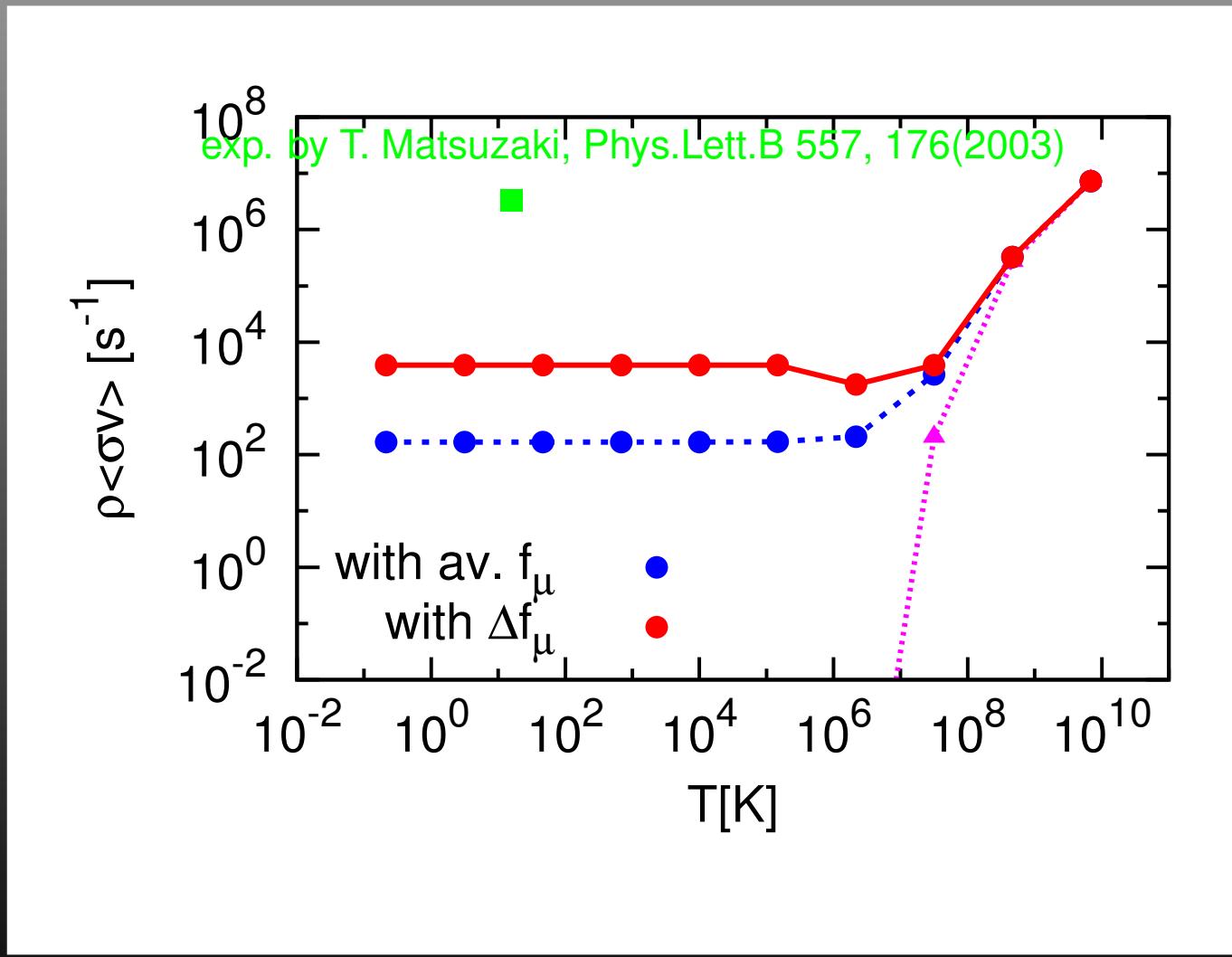
Reaction Rate of the t+t



Reaction Rate of the t+t



Reaction Rate of the t+t



Conclusions

Importance of Fluctuations of cross section in muon-catalysed $t-t$ fusion

- Reaction rate and μ cycling rate
- Numerical simulation by the CoMD
- A characteristic change of the slope of $\Delta U_\mu / \bar{U}_\mu$ at the ionization energy of the μ molecule.
- Fluctuation of σ contributes to enhance the RR
- Obtained RR has no T dependence in the low T region, as expected. However it is smaller than exp. μ cycling rate

