

# The Thermal-Statistical Model for Particle Production I.

J. Cleymans

23 - 25 July 2008 / JINR, Dubna



# Outline

South Africa and the University of Cape Town

Statistical Model

Strangeness

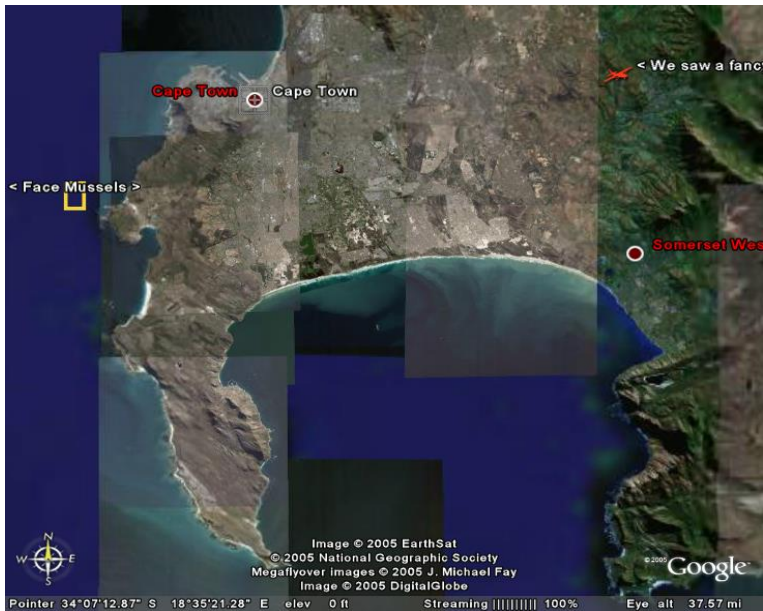
$E_T/N_{ch}$  vs.  $E/N$

Dependence on the Size of System.

The Horn in the  $K^+/\pi^+$  Ratio



# Cape Peninsula



# South Africa



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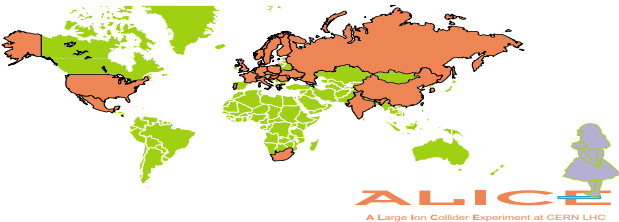
## UCT



## ALICE Collaboration



**The ALICE Collaboration**  
27 countries, 83 Institutes and 1000 members

**America**

- Wisconsin, Wiscconsin Physics Institute

**China**

- Beijing, Institute of Atomic Energy
- Wuhan, Wuhan Normal University, Institute of Particle Physics

**Croatia**

- Zagreb, University of Zagreb, Petruskovalova Matematicka Fakultet
- Zagreb, Ruđer Bosković Institute

**Czech Republic**

- Praha, Academy of Sciences of the Czech Republic, Institute of Physics
- Praha, Academy of Sciences of the Czech Republic, Nuclear Physics Institute

**Denmark**

- Copenhagen, Niels Bohr Institute, University of Copenhagen

**Ireland**

- Cork, University of York, Department of Physics and Heraldic Institute of Physics

**Italy**

- Arezzo, Università della Valle (Clemente Panzani II), Istituto Nazionale di Fisica Nucleare, Istituto dei Fisici per l'Infineon (INFN) - CNRS
- Ferrara, Laboratorio di Fisica Corpuscolare (LPC)
- Firenze, Scuola Normale Superiore, Istituto Nazionale di Fisica Nucleare, Laboratoire de Physique Subatomique et des Techniciens Associés (LNSAT/INAF)

**Japan**

- Chiba, Université de Chiba, Institut National de Physique Nucléaire et de Physique des Particules (INPNC-CNRS), Institut de Physique des Particules (INPNC-CNRS)

**Norway**

- Trondheim, Universitetet i Luleå, Fakultet for Naturvitenskap og Ingeniørvesen, Institutt for Fysikk, Institut National de Physique Nucléaire et de Physique des Particules (INPNC-CNRS), Institut de Physique des Particules (INPNC-CNRS), Institut de Physique Nucléaire et de Physique des Particules (INPNC-CNRS), Institut de Physique Nucléaire et de Physique des Particules (INPNC-CNRS)

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- Darmstadt, Technische Universität Darmstadt, Institut für Kernphysik
- Heidelberg, Institut für Experimentelle Kernphysik, Institut für Experimentelle Kernphysik, Institut für Experimentelle Kernphysik, Institut für Experimentelle Kernphysik, Institut für Experimentelle Kernphysik
- München, Fakultät für Physik, Ludwig-Maximilians-Universität München, Institut für Experimentelle Kernphysik, Institut für Experimentelle Kernphysik, Institut für Experimentelle Kernphysik, Institut für Experimentelle Kernphysik

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**India**

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- Bhubaneswar, Institute of Physics
- Calcutta, Variable Energy Cyclotron Centre (VECC)
- Chandigarh, Punjab University, High Energy Physics group
- Japan, University of Hyderabad, Physics Department
- Agra, University of Agra, Physics Department

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- Alessandria, Università del Piemonte Orientale, Istituto Nazionale di Fisica Nucleare (INFN)
- Bari, Politecnico di Bari
- Bari, Università degli Studi di Bari and INFN
- Bologna, Università degli Studi di Bologna and INFN
- Cagliari, Università degli Studi di Cagliari and INFN
- Catania, Università di Catania and INFN
- Legnano, Laboratorio Nazionali di Legnano, INFN
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- Salerno, Università degli Studi di Salerno and INFN
- Trieste, Università degli Studi di Trieste and INFN

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- Mexico, D.F., Universidad Nacional Autónoma de México, Instituto de Física and Instituto de Ciencias Nucleares
- Morelia, Michoacán, Universidad Michoacana de San Nicolás de Hidalgo, and Center of Mathematics
- Puebla, Universidad Autónoma de Puebla

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**Norway**

- Bergen, University of Bergen, Department of Physics
- Bergen, Bergen University College Faculty of Engineering
- Oslo, University of Oslo, Department of Physics

**Poland**

- Cracow, H. Niewodniczański Institute of Nuclear Physics, High Energy Physics Department
- Warszawa, Sokołowski Institute for Nuclear Studies
- Warszawa, Warsaw University of Technology

**Portugal****Rep. of Korea**

- Gangneung, Gangneung National University
- Pohang, Pohang Accelerator Laboratory

**Romania**

- Bucharest, National Institute for Physics and Nuclear Engineering

**Russia**

- Moscow, Moscow Region, Joint Institute for Nuclear Research (JINR)
- Gubkin, Gubkin Pedagogical Nuclear Physics Institute (GPNPI)
- Moscow, Academy of Sciences, Institute for Nuclear Research (INR)
- Moscow, Institute for Theoretical and Experimental Physics (ITEP)
- Moscow, Moscow Research Center Skobeltsyn Institute of Nuclear Physics (SINP)
- Novosibirsk, Skobeltsyn Institute for Nuclear Physics (SINP)
- Protvino, Institute for High Energy Physics (IHEP)
- Sarov, Russian Federal Nuclear Center (VNIIEP)
- St. Petersburg, Saint Petersburg Institute for Physics, Member Institute for Mathematics and Modern Science Association

**Slovenia**

- Ljubljana, Comenius University Faculty of Mathematics, Physics and Informatics, Institute of Experimental Physics, Slovene Academy of Sciences
- Ljubljana, Faculty of Sciences, P.J. Štefan University

**South Africa**

- Rondebosch, University of Cape Town (UCT)

**Sweden**

- Lund, University of Lund, Division of Cosmic and Subatomic Physics

**Switzerland**

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- Lausanne, Ecole Polytechnique Fédérale de Lausanne (EPFL)

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- Kiev, Ukrainian Academy of Sciences, National Scientific Center
- Kiev, Institute of Physics and Technology
- Cherkassy, Scientific and Technological Research Institute of Instrument Engineering (ISTEAC)
- Kiev, Ukrainian Academy of Sciences, Bogolyubov Institute for Theoretical Physics, Department of High Energy Density Physics

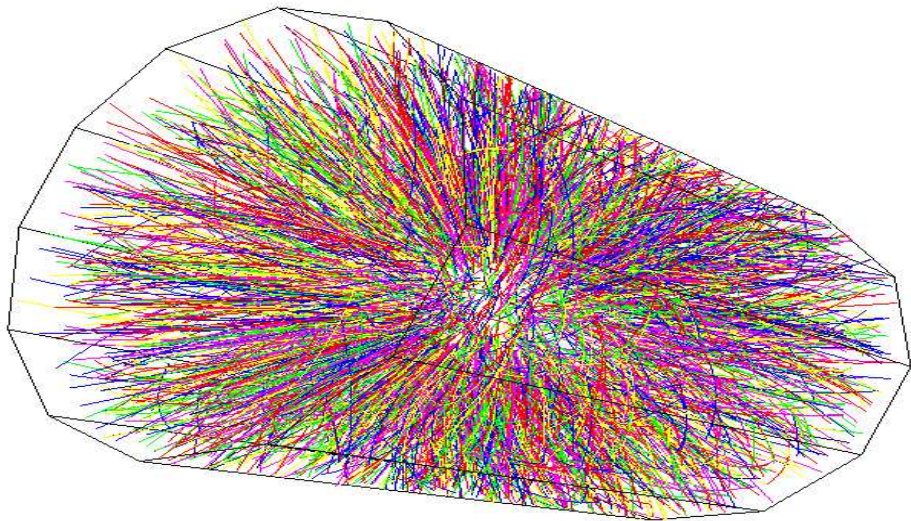
**United Kingdom**

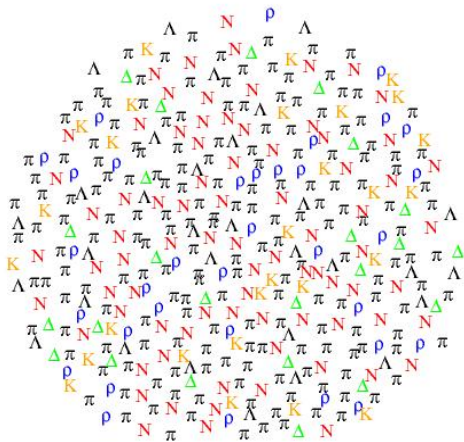
- Birmingham, University of Birmingham, School of Physics and Space Research

**United States of America**

- Columbia, Ohio, Ohio State University, Department of Physics
- Columbus, Ohio, Ohio State University, Department of Physics
- Oak Ridge, TN, Oak Ridge National Laboratory (ORNL), Instrumentation and Control, Division







J.C. and H. Satz, Zeitschrift fuer Physik C57, 135 (1993)

# Thermal Equilibrium

In thermal equilibrium

$$Z = \text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}$$

$$\langle N \rangle = \frac{\text{Tr} N e^{-\frac{H}{T} + \frac{\mu N}{T}}}{\text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}}$$

$$\langle E \rangle = \frac{\text{Tr} E e^{-\frac{H}{T} + \frac{\mu N}{T}}}{\text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}}$$





# Full Hydrodynamic Flow

Bjorken scaling + Transverse expansion

After integration over  $m_T$

$$\frac{dN_i/dy}{dN_j/dy} = \frac{N_i^0}{N_j^0}$$

where  $N_i^0$  is the particle yield  
as calculated in a fireball **AT REST!**

**Effects of hydrodynamic flow cancel out in ratios.**



# Thermal Equilibrium

## Particle Number

$$\begin{aligned}\langle N \rangle &= \frac{\text{Tr} N e^{-\frac{H}{T} + \frac{\mu N}{T}}}{\text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}}} \\ &= \frac{T}{Z} \frac{\partial}{\partial \mu} \text{Tr} e^{-\frac{H}{T} + \frac{\mu N}{T}} \\ &= T \frac{1}{Z} \frac{\partial Z}{\partial \mu} \\ &= T \frac{\partial}{\partial \mu} \ln Z\end{aligned}$$



# Thermal Equilibrium

## Average Energy

$$\begin{aligned}\langle E \rangle &= \frac{\text{Tr } H e^{\frac{-H}{T} + \frac{\mu N}{T}}}{\text{Tr } e^{\frac{-H}{T} + \frac{\mu N}{T}}} \\ &= -\frac{1}{Z} \frac{\partial Z}{\partial \beta} + \mu \langle N \rangle \\ &= T^2 \frac{\partial}{\partial T} \ln Z + \mu \langle N \rangle\end{aligned}$$



# Thermal Equilibrium

$$\begin{aligned} N_i &= g_i V \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E}{T}\right) e^{\frac{\mu_i}{T}} \\ &= g_i V \frac{4\pi}{(2\pi)^3} \int p^2 dp \exp\left(-\frac{\sqrt{p^2 + m_i^2}}{T}\right) e^{\frac{\mu_i}{T}} \\ &= g_i V \frac{4\pi}{(2\pi)^3} T^3 \int x^2 dx \exp\left(-\sqrt{x^2 + m_i^2/T^2}\right) e^{\frac{\mu_i}{T}} \\ &= g_i V \frac{1}{2\pi^2} T m_i^2 K_2\left(\frac{m_i}{T}\right) e^{\frac{\mu_i}{T}} \end{aligned}$$



# Thermal Equilibrium

$$n_i = g_i \frac{1}{2\pi^2} T m_i^2 K_2 \left( \frac{m_i}{T} \right) e^{\frac{\mu_i}{T}}$$

$$\epsilon_i = g_i \frac{1}{2\pi^2} T m_i^3 \left[ K_1 \left( \frac{m_i}{T} \right) + 3 \frac{T}{m} K_2 \left( \frac{m_i}{T} \right) \right] e^{\frac{\mu_i}{T}}$$

$$s_i = g_i \frac{1}{2\pi^2} m_i^3 \left[ K_1 \left( \frac{m_i}{T} \right) + \frac{4T}{m} K_2 \left( \frac{m_i}{T} \right) - \frac{\mu_i}{m} K_2 \left( \frac{m_i}{T} \right) \right] e^{\frac{\mu_i}{T}}$$

$$P_i = g_i \frac{1}{2\pi^2} T^2 m_i^2 K_2 \left( \frac{m_i}{T} \right) e^{\frac{\mu_i}{T}}$$



# Chemical Equilibrium

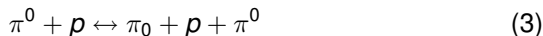
In equilibrium

$$E_1 + E_2 + \dots = E_3 + E_4 + E_5 + \dots \quad (1)$$

for the chemical potentials

$$\mu_1 + \mu_2 + \dots = \mu_3 + \mu_4 + \mu_5 + \dots \quad (2)$$

As an example



leads to

$$\mu_{\pi^0} + \mu_p = \mu_{\pi^0} + \mu_p + \mu_{\pi^0} \quad (4)$$

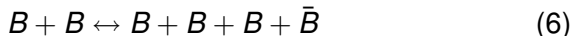
which leads to

$$\mu_{\pi^0} = 0 \quad (5)$$



## Chemical Equilibrium

In equilibrium



$$dE = -pdV + TdS + \mu_B dN_B + \mu_{\bar{B}} dN_{\bar{B}}$$

Due to baryon number conservation one has

$$N_B - N_{\bar{B}} = \text{constant}$$

and

$$dN_B = dN_{\bar{B}}$$

The energy is a minimum for

$$dE = (\mu_B + \mu_{\bar{B}})dN_B = 0 \quad (7)$$

$$\mu_B = -\mu_{\bar{B}} \quad (8)$$



# Chemical Equilibrium

In equilibrium

$$N_B = g V \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E}{T} + \frac{\mu_B}{T}\right)$$

$$N_{\bar{B}} = g V \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E}{T} - \frac{\mu_B}{T}\right)$$

$$N_B = N_{\bar{B}} \rightarrow \mu_B = 0$$

$$N_B \geq N_{\bar{B}} \rightarrow \mu_B \geq 0$$

$$N_B \leq N_{\bar{B}} \rightarrow \mu_B \leq 0$$





	Chemical Equilibrium	No Chem. Equil.
$\pi$	$\exp\left[-\frac{E_\pi}{T}\right]$	$\exp\left[-\frac{E_\pi}{T} + \frac{\mu_\pi}{T}\right]$
$N$	$\exp\left[-\frac{E_N}{T} + \frac{\mu_B}{T}\right]$	$\exp\left[-\frac{E_N}{T} + \frac{\mu_N}{T}\right]$
$\bar{N}$	$\exp\left[-\frac{E_N}{T} - \frac{\mu_B}{T}\right]$	$\exp\left[-\frac{E_N}{T} + \frac{\mu_{\bar{N}}}{T}\right]$
$\Lambda$	$\exp\left[-\frac{E_\Lambda}{T} + \frac{\mu_B}{T} - \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_\Lambda}{T} + \frac{\mu_\Lambda}{T}\right]$
$\bar{\Lambda}$	$\exp\left[-\frac{E_\Lambda}{T} - \frac{\mu_B}{T} + \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_\Lambda}{T} + \frac{\mu_{\bar{\Lambda}}}{T}\right]$
$K$	$\exp\left[-\frac{E_K}{T} + \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_K}{T} + \frac{\mu_K}{T}\right]$
$\bar{K}$	$\exp\left[-\frac{E_K}{T} - \frac{\mu_S}{T}\right]$	$\exp\left[-\frac{E_K}{T} + \frac{\mu_{\bar{K}}}{T}\right]$



The number of particles of type  $i$  is determined by:

$$N_i = V g_i \int \frac{d^3 p}{(2\pi)^3} \exp\left(-\frac{E_i}{T} + \frac{\mu_i}{T}\right)$$

For bosons:

$$N_i = V g_i \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\exp\left(\frac{E_i}{T} - \frac{\mu_i}{T}\right) - 1}$$

For fermions:

$$N_i = V g_i \int \frac{d^3 p}{(2\pi)^3} \frac{1}{\exp\left(\frac{E_i}{T} - \frac{\mu_i}{T}\right) + 1}$$



# Chemical Equilibrium

Only conserved quantum numbers matter for chemical equilibrium: In equilibrium

$$\mu_i = B_i\mu_B + Q_i\mu_Q + S_i\mu_S + C_i\mu_C + \dots \quad (9)$$



$g_i$	$m_i$	stat	$S_i$	$B_i$	$Q_i$	Particle $i$
1	0.140	-1	0	0	1.	$\pi^+$
1	0.135	-1	0	0	0.	$\pi^0$
1	0.140	-1	0	0	-1.	$\pi^-$
1	0.547	-1	0	0	0.	$\eta$
3	0.770	-1	0	0	1.	$\rho^+$
3	0.770	-1	0	0	0.	$\rho^0$
3	0.770	-1	0	0	-1.	$\rho^-$
3	0.782	-1	0	0	0.	$\omega$
1	0.958	-1	0	0	0.	$\eta'$
1	0.980	-1	0	0	0.	$f_0$
1	0.982	-1	0	0	1.	$a_0^+$
1	0.982	-1	0	0	0.	$a_0^0$
1	0.982	-1	0	0	-1.	$a_0^-$
3	1.019	-1	0	0	0.	$\phi$
3	1.170	-1	0	0	0.	
3	1.230	-1	0	0	1.	
3	1.230	-1	0	0	0.	
3	1.230	-1	0	0	-1.	
3	1.229	-1	0	0	1.	
3	1.229	-1	0	0	0.	
3	1.229	-1	0	0	-1.	
5	1.275	-1	0	0	0.	
3	1.282	-1	0	0	0.	
1	1.297	-1	0	0	0.	
1	1.300	-1	0	0	1.	
1	1.300	-1	0	0	0.	



# The Role of Resonances

## Example: $\rho$ 's

$$\rho \rightarrow \pi^+ \pi^-$$

Final, observed, number of  $\pi^+$  is given by

$$N_{\pi^+} = N_{\pi^+}(\text{thermal}) + N_{\pi^+}(\text{resonance decays})$$

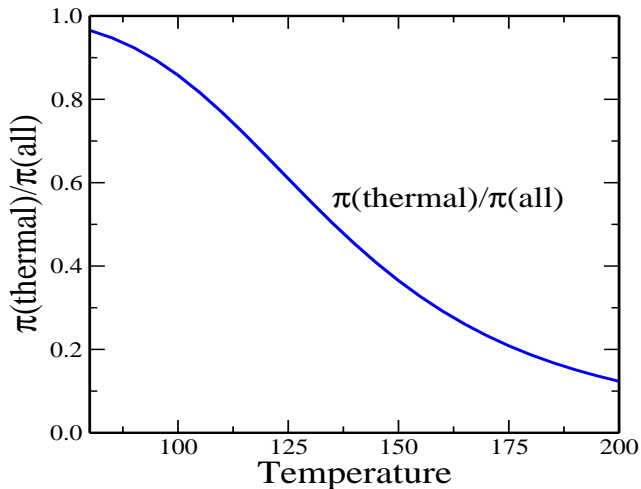
depending on the temperature, over 80% of observed pions are due to resonance decays



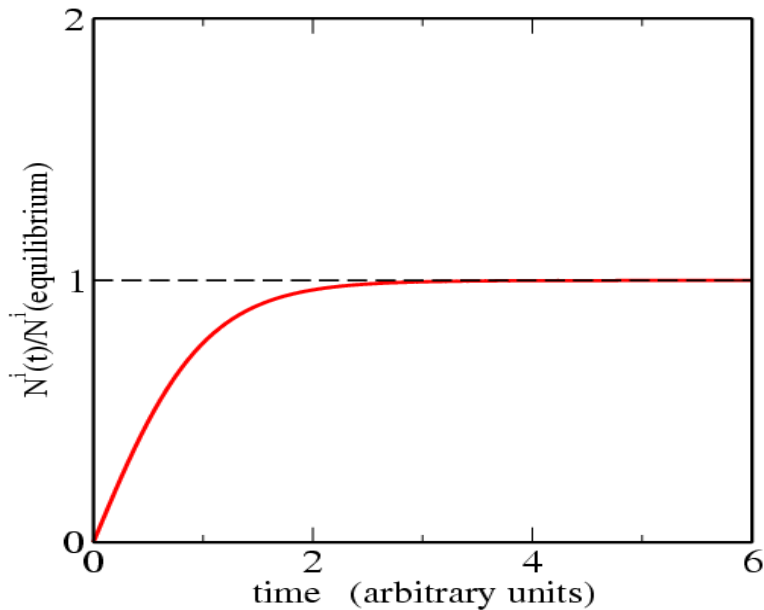
$g_i$	$m_i$	stat	$S_i$	$B_i$	$Q_i$	BR $\rightarrow \pi^+$	Particle $i$
1	0.140	-1	0	0	1.	1.000	$\pi^+$
1	0.135	-1	0	0	0.	0.000	$\pi^0$
1	0.140	-1	0	0	-1.	0.000	$\pi^-$
1	0.547	-1	0	0	0.	0.285	$\eta$
3	0.770	-1	0	0	1.	1.000	$\rho^+$
3	0.770	-1	0	0	0.	1.000	$\rho^0$
3	0.770	-1	0	0	-1.	0.000	$\rho^-$
3	0.782	-1	0	0	0.	0.910	$\omega$
1	0.958	-1	0	0	0.	0.965	$\eta'$
1	0.980	-1	0	0	0.	0.521	$f_0$
1	0.982	-1	0	0	1.	1.285	$a_0^+$
1	0.982	-1	0	0	0.	0.285	$a_0^0$
1	0.982	-1	0	0	-1.	0.285	$a_0^-$
3	1.019	-1	0	0	0.	0.155	$\phi$
3	1.170	-1	0	0	0.	1.000	$h_1$
3	1.230	-1	0	0	1.	1.500	
3	1.230	-1	0	0	0.	0.50	
3	1.230	-1	0	0	-1.	0.50	
3	1.229	-1	0	0	1.	1.91	
3	1.229	-1	0	0	0.	0.91	
3	1.229	-1	0	0	-1.	0.91	
5	1.275	-1	0	0	0.	0.69	
3	1.282	-1	0	0	0.	1.00	
1	1.297	-1	0	0	0.	1.11	
1	1.300	-1	0	0	1.	2.00	
1	1.300	-1	0	0	0.	1.50	



## Importance of Resonances.



## Strangeness saturation?





## Strangeness saturation?

$$N_i = \boxed{\gamma_s^{|S|}} V g_i \int \frac{d^3p}{(2\pi)^3} \exp\left(-\frac{E_i}{T} + \frac{\mu_i}{T}\right)$$

with

$\gamma_s < 1$  strangeness under-saturation

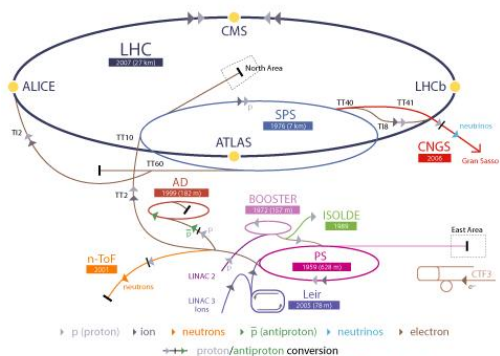
$\gamma_s = 1$  strangeness in chemical equilibrium

$\gamma_s > 1$  strangeness over-saturation



## SPS

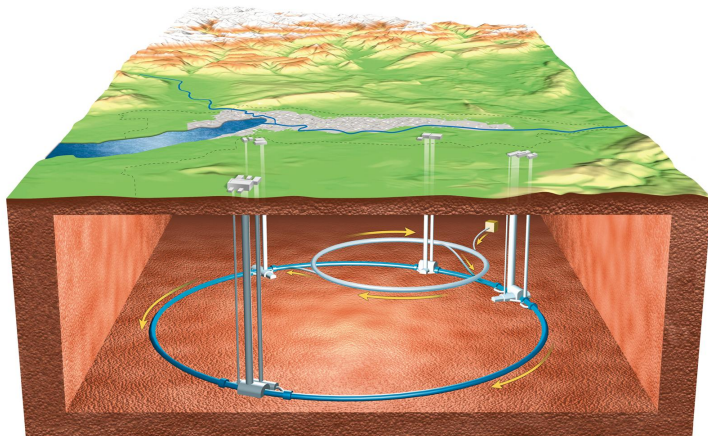
## CERN Accelerator Complex



LHC Large Hadron Collider SPS Super Proton Synchrotron PS Proton Synchrotron  
 AD Antiproton Decelerator CTF3 Clic Test Facility  
 CNGS Cern Neutrinos to Gran Sasso ISOLDE Isotope Separator OnLine DEvice  
 LEIR Low Energy Ion Ring LINAC LINear ACcelerator n-ToF Neutrons Time Of Flight



# SPS



## SPS data.

	Measurement
Pb–Pb 158A GeV	
$(\pi^+ + \pi^-)/2.$	$600 \pm 30$
$K^+$	$95 \pm 10$
$K^-$	$50 \pm 5$
$K_S^0$	$60 \pm 12$
$p$	$140 \pm 12$
$\bar{p}$	$10 \pm 1.7$
$\phi$	$7.6 \pm 1.1$
$\Xi^-$	$4.42 \pm 0.31$
$\Xi^-$	$0.74 \pm 0.04$
$\bar{\Lambda}/\Lambda$	$0.2 \pm 0.04$



## SPS data.

SPS: Chemical Freeze-Out Parameters:

$$T = 156.0 \pm 2.4 \text{ MeV}$$

$$\mu_B = 239 \pm 12 \text{ MeV}$$

$$\gamma_s = 0.862 \pm 0.036$$

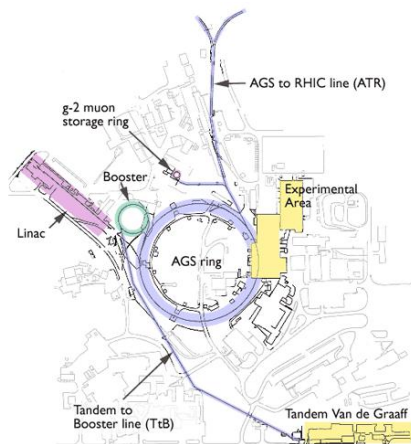
F. Becattini, J.C., A. Keränen, E. Suhonen and K. Redlich  
Physical Review C64 (2001) 024901.



# AGS



# AGS



# AGS data.

	Measurement
<b>Au–Au 11.6A GeV</b>	
<b>Participants</b>	$363 \pm 10$
$K^+$	$23.7 \pm 2.9$
$K^-$	$3.76 \pm 0.47$
$\pi^+$	$133.7 \pm 9.9$
$\Lambda$	$20.34 \pm 2.74$
$p/\pi^+$	$1.234 \pm 0.126$
$\bar{p}$	$>0.0185 \pm 0.0018$





## AGS data.

AGS: Chemical Freeze-Out Parameters:

$$T = 130.6 \pm 5.5 \text{ MeV}$$

$$\mu_B = 594 \pm 26 \text{ MeV}$$

$$\gamma_s = 0.883 \pm 0.124$$

F. Becattini, J.C., A. Keränen, E. Suhonen and K. Redlich  
Physical Review C64 (2001) 024901.

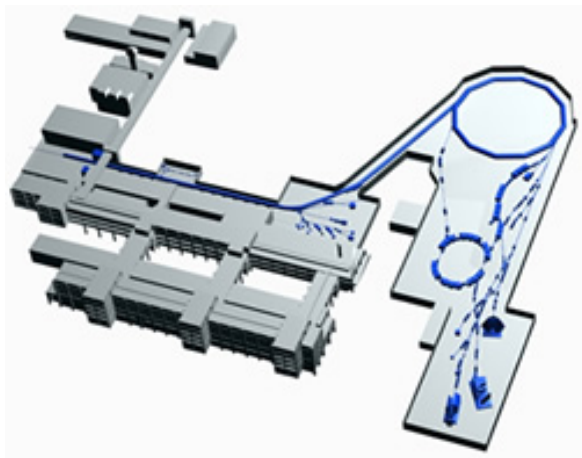


# SIS data.

	Measurement
<b>Au–Au 1.7A GeV</b>	
$\pi^+/\text{p}$	$0.052 \pm 0.013$
$\text{K}^+/\pi^+$	$0.003 \pm 0.00075$
$\pi^-/\pi^+$	$2.05 \pm 0.51$
$\eta/\pi^0$	$0.018 \pm 0.007$



# GSI



## SIS data.

SIS: Chemical Freeze-Out Parameters:

$$T = 49.7 \pm 1.1 \text{ MeV}$$

$$\mu_B = 818 \pm 15 \text{ MeV}$$

$$\gamma_s = 1 \text{ (fixed)}$$

J. C., H. Oeschler and K. Redlich)  
Physical Review C59, (1999) 1663.



## RHIC data.

J. C., B. Kämpfer, M. Kaneta, S. Wheaton, N. Xu, Phys. Rev. C71, 0409071 (2005)

Ratio	Experiment	Central	Mid-Central	Peripheral
$\pi_{(2)}^-/\pi_{(2)}^+$	BRAHMS	0.990±0.100		
	PHENIX	0.960±0.177	0.920±0.170	0.933±0.172
	PHOBOS	1.000±0.022		
	STAR	1.000±0.073	1.000±0.073	1.000 ± 0.073
$K_{(2)}^+/K_{(2)}^-$	PHENIX	1.152±0.240	1.292±0.268	1.322±0.284
	PHOBOS	1.099±0.111		
	STAR	1.109±0.022	1.105±0.036	1.120±0.040
$\bar{p}_{(1)}/p_{(1)}$	PHENIX	0.680±0.149	0.671±0.142	0.717±0.157
$\bar{p}_{(2)}/p_{(2)}$	BRAHMS	0.650±0.092		
	PHOBOS	0.600±0.072		
	STAR	0.714±0.050	0.724±0.050	0.764±0.053
$\bar{\Lambda}_{(1)}/\Lambda_{(1)}$	PHENIX	0.750±0.180	0.798±0.197	0.795±0.197
$\bar{\Lambda}_{(2)}/\Lambda_{(2)}$	STAR	0.719±0.090	0.739±0.092	0.744±0.100
$\Xi_{(2)}^+/\Xi_{(2)}^-$	STAR	0.840±0.053	0.822±0.114	0.815±0.096
$\bar{\Omega}^+/\Omega^-$	STAR	1.062±0.410		
$K_{(2)}^-/\pi_{(2)}^-$	PHENIX	0.151±0.030	0.134±0.027	0.116±0.023
	STAR	0.151±0.022	0.147±0.022	0.130±0.019
$K_S^0/\pi_{(2)}^-$	STAR	0.134±0.022	0.131±0.022	0.108±0.018
$\bar{p}_{(1)}/\pi_{(2)}^-$	PHENIX	0.049±0.010	0.047±0.010	0.045±0.009
$\bar{p}_{(2)}/\pi_{(2)}^-$	STAR	0.069±0.019	0.067±0.019	0.067±0.019
$\Lambda_{(1)}/\pi_{(2)}^-$	STAR	0.043±0.008	0.043±0.008	0.039±0.007
$\Lambda_{(2)}/\pi_{(2)}^-$	PHENIX	0.072±0.017	0.068±0.016	0.074±0.017
$< K^{*0} > / \pi_{(2)}^-$	STAR	0.039±0.011		
$\phi/\pi_{(2)}^-$	STAR	0.022±0.003	0.021±0.004	0.022±0.004
$\Xi_{(2)}^-/\pi_{(2)}^-$	STAR	0.0093±0.0012	0.0072±0.0011	0.0060±0.0008



## RHIC data.

RHIC: Chemical Freeze-Out Parameters:

$$T = 169 \pm 4.2 \text{ MeV}$$

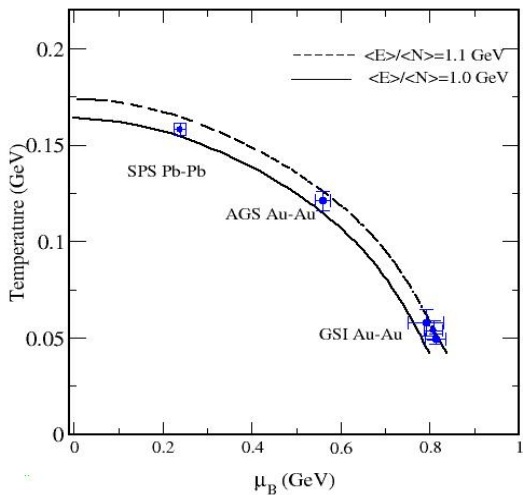
$$\mu_B = 39.6 \pm 6 \text{ MeV}$$

$$\gamma_s = 0.9 \pm 0.1$$

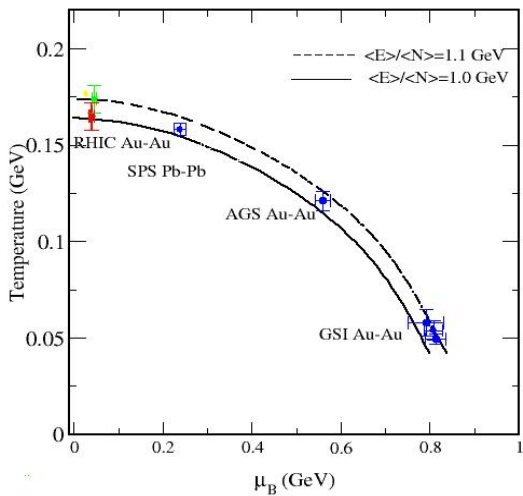
J. C., B. Kämpfer, M. Kaneta, S. Wheaton, N. Xu  
Phys. Rev. C71, 0409071 (2005)



## E/N in 1999

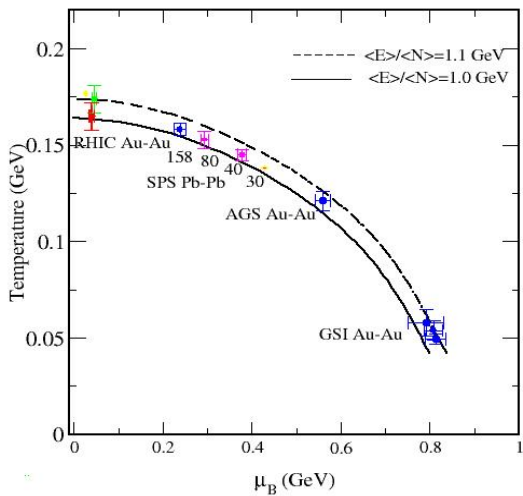


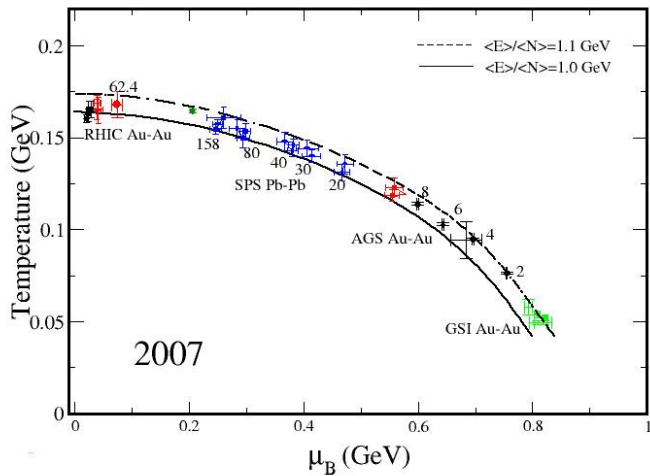
## E/N in 2000



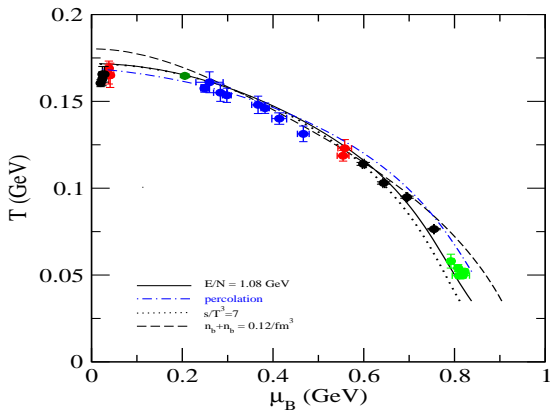


## E/N in 2005





# Chemical Freeze-Out: Status in 2005

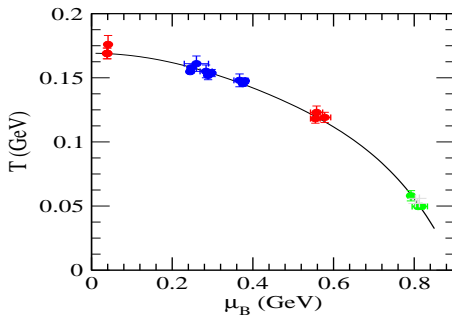


V. Magas and H. Satz, Eur. Phys. J. **C32** 115 (2003).

P. Braun-Munzinger and J. Stachel, J. Phys. G:Nucl. Part. Phys. **28** 1971 (2002).



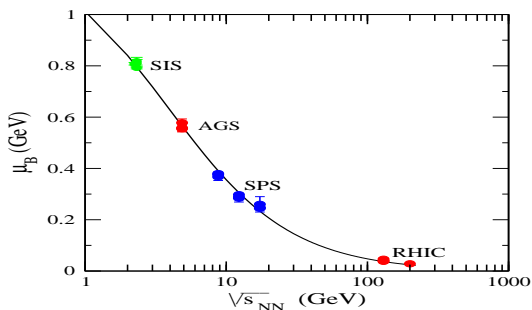
# Chemical Freeze-Out: Status in 2005



$$T(\mu_B) = 0.169 - 0.189\mu_B^2 + 0.165\mu_B^4 - 0.229\mu_B^6.$$

J. C., H. Oeschler, K. Redlich, S. Wheaton



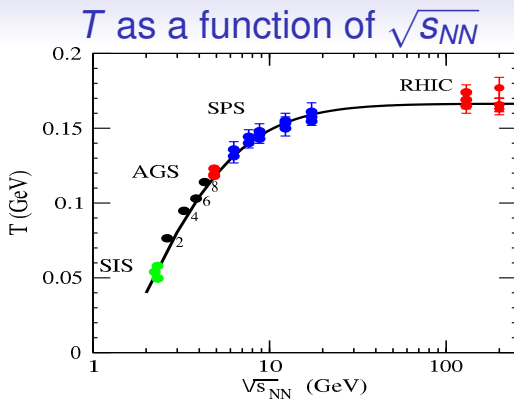
$\mu_B$  as a function of  $\sqrt{s_{NN}}$ 

$$\mu_B(\sqrt{s}) = \frac{1.273 \text{ GeV}}{1 + 0.258 \text{ GeV}^{-1} \sqrt{s}}$$

This predicts at LHC  $\mu_B \approx 1 \text{ MeV}$ .

J. C., H. Oeschler, K. Redlich, S. Wheaton

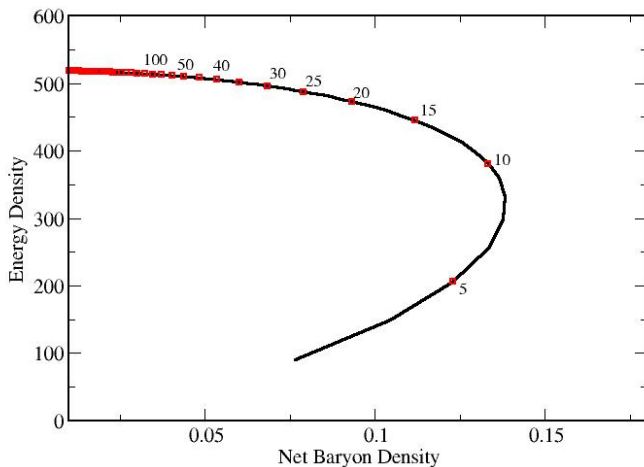




This predicts at LHC  $T \approx 170$  MeV.

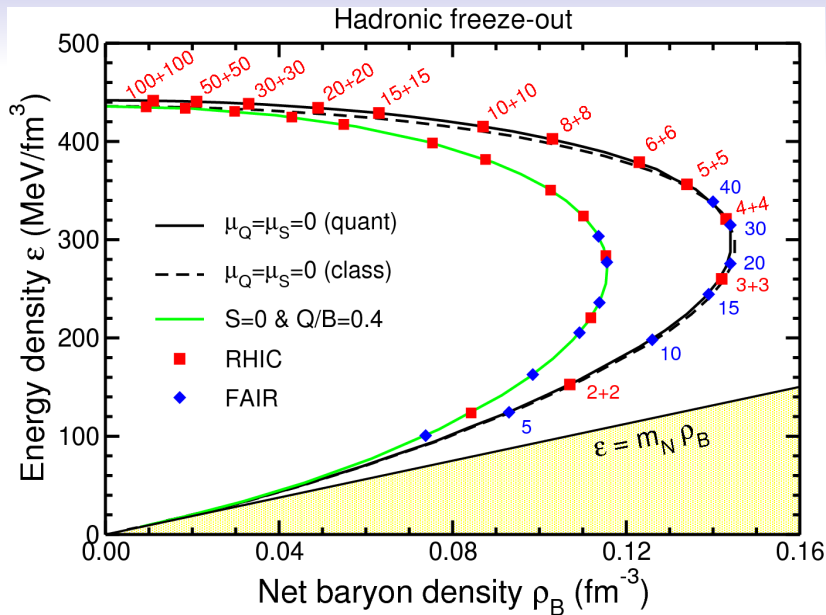
J. C., H. Oeschler, K. Redlich, S. Wheaton





J. Randrup and J.C., Phys. Rev. C74 (2006) 047901





J. Randrup and J.C., Phys. Rev. C74 (2006) 047901





## Will it be possible to determine directly $E/N$ ?

$E$ : energy of primordial hadrons

$N$ : number of primordial hadrons



$$\begin{aligned}\langle E_T \rangle &= \langle E \sin \theta \rangle \\ &= \frac{\pi}{4} \langle E \rangle\end{aligned}$$



Low energy limit

$$\lim \frac{E_T}{N_{ch}} = \frac{\frac{\pi}{4} m_N}{0.4} \approx 1.8 \text{ GeV}$$

High energy limit

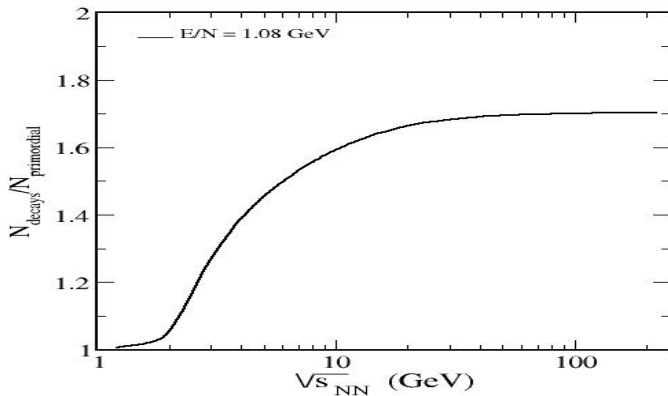
$$\lim \frac{E_T}{N_{ch}} = \frac{\frac{\pi}{4} \langle M \rangle}{2/3} \approx 0.9 \text{ GeV}$$



However

$E_T$  : subtract  $m_N$  for baryons  
add  $m_N$  for antibaryons.

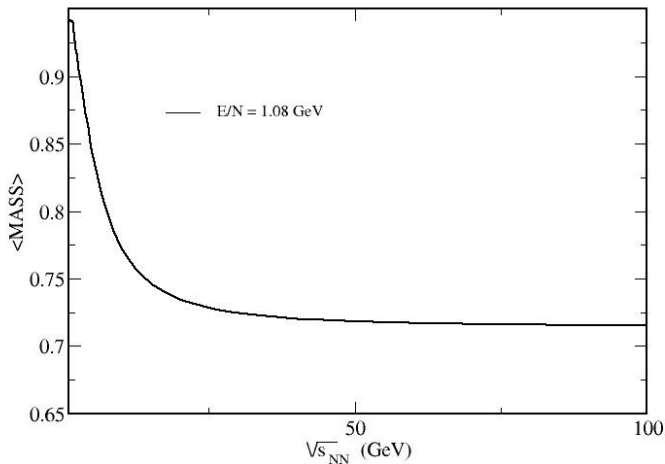
## Primordial vs Final State Hadrons



J.C., R. Sahoo, D.K. Srivastava, S. Wheaton



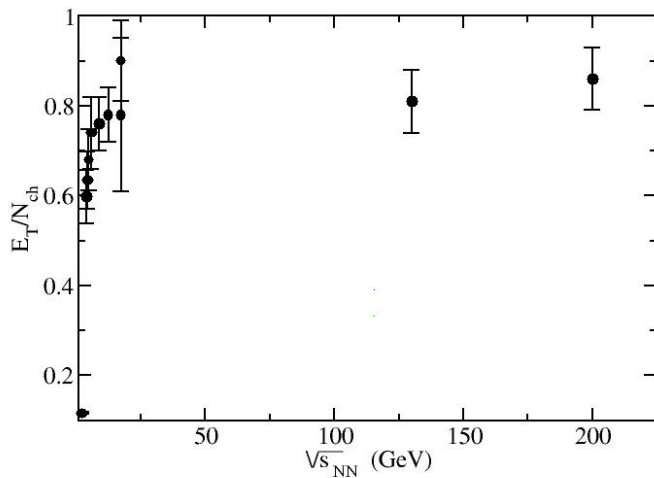
## Average Mass in Fireball



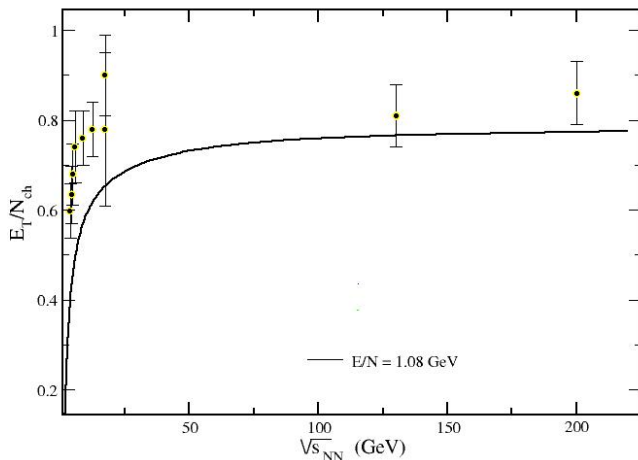
J.C., R. Sahoo, D.K. Srivastava, S. Wheaton



## Transverse Energy per Charge



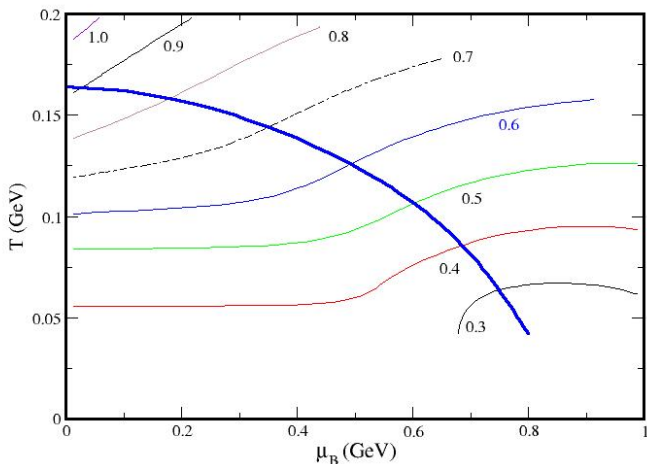
## Transverse Energy per Charged Hadron



J.C., R. Sahoo, D.K. Srivastava, S. Wheaton





Lines of constant  $E_T/N_{ch}$ 

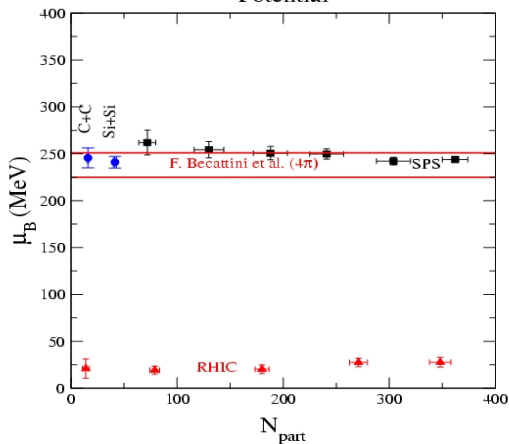
J.C., R. Sahoo, D.K. Srivastava, S. Wheaton



$E_T/N_{ch}$  mainly follows  $T$  and is determined by  $E/N$ ,



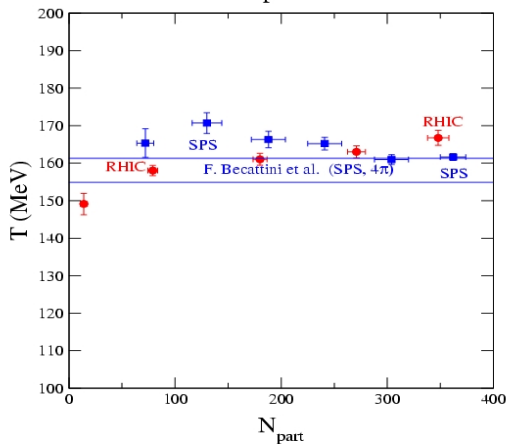
## Centrality Dependence of the Baryon Chemical Potential



J. C., B. Kämpfer, P. Steinberg and S. Wheaton, Journal of Physics G30 S595-S598 (2004).

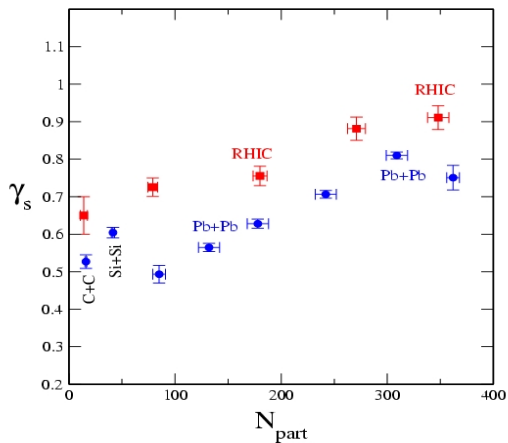


## Centrality Dependence of the Chemical Freeze-out Temperature



J. C., B. Kämpfer, P. Steinberg and S. Wheaton, Journal of Physics G30 S595-S598 (2004).





J. C., B. Kämpfer, P. Steinberg and S. Wheaton, Journal of Physics G30 S595-S598 (2004).

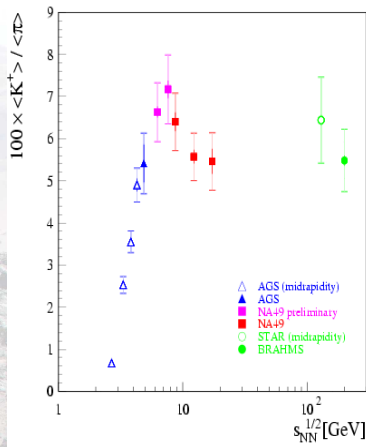


The NA49 Collaboration has recently performed a series of measurements of Pb-Pb collisions at 20, 30, 40, 80 and 158 AGeV beam energies . When these results are combined with measurements at lower beam energies from the AGS they reveal an unusually sharp variation with beam energy in the  $\Lambda/\langle\pi\rangle$ , with  $\langle\pi\rangle \equiv 3/2(\pi^+ + \pi^-)$ , and  $K^+/\pi^+$  ratios. Such a strong variation with energy does not occur in pp collisions and therefore indicates a major difference in heavy-ion collisions. This transition has been referred as the “horn”.



# The Elephant in the Room

Friese  
Dinkelaker  
Blume  
Speltz



Difficult to avoid, Hard to Model

→ But no unambiguous corroborating evidence



## Strangeness in Heavy Ion Collisions vs Strangeness in pp - collisions

Use the Wroblewski factor

$$\lambda_s = \frac{2 \langle s\bar{s} \rangle}{\langle u\bar{u} \rangle + \langle d\bar{d} \rangle}$$

This is determined by the number of **newly** created quark – anti-quark pairs and **before** strong decays, i.e. before  $\rho$ 's and  $\Delta$ 's decay.

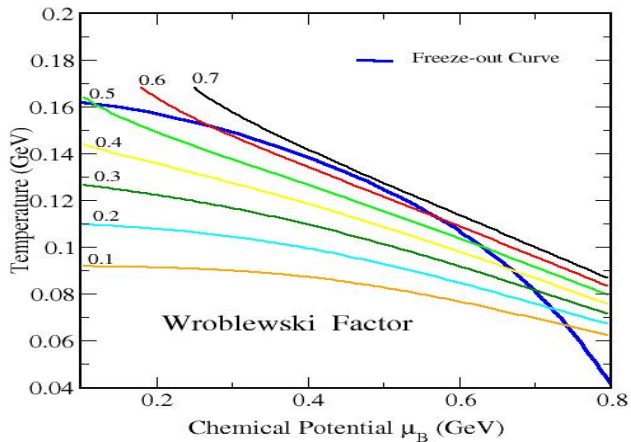
Limiting values :

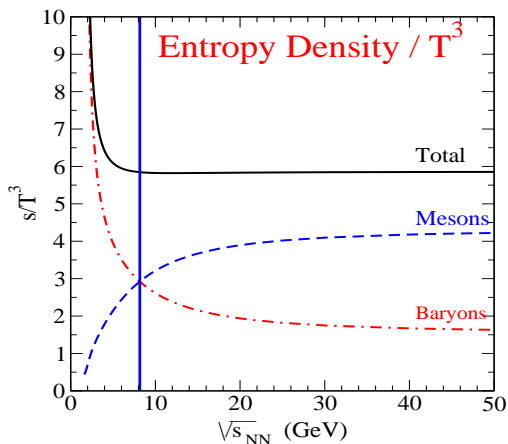
$\lambda_s = 1$  all quark pairs are equally abundant, SU(3) symmetry.

$\lambda_s = 0$  no strange quark pairs.



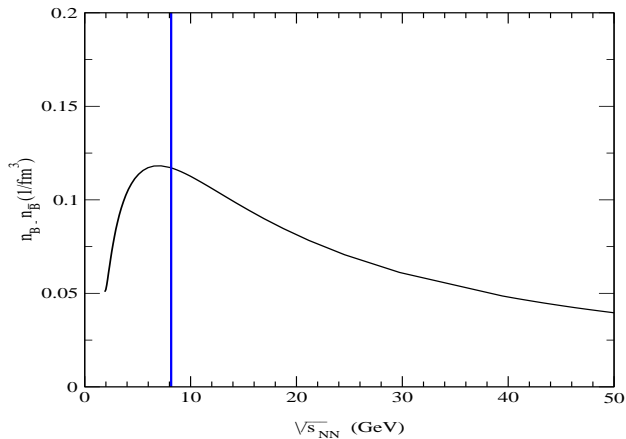






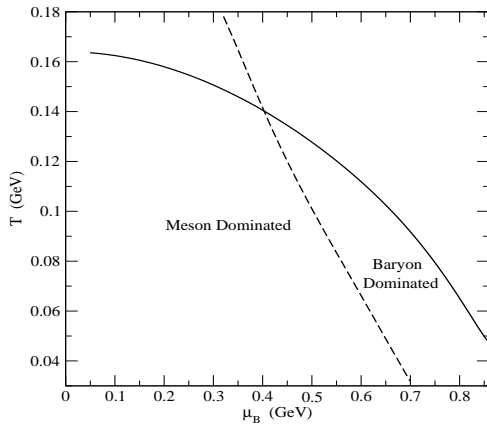
J. C., H. Oeschler, K. Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54.





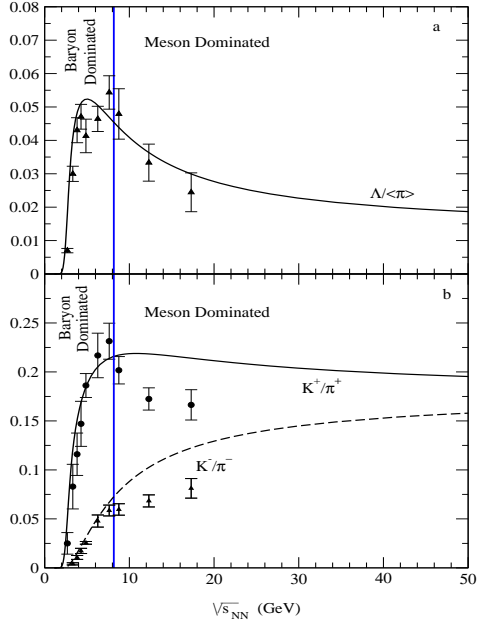
J. C., H. Oeschler, K. Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54.





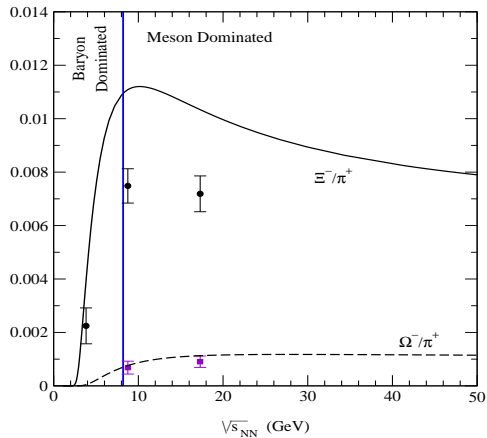
J. C., H. Oeschler, K. Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54.





J. C., H. Oeschler, K. Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54.





J. C., H. Oeschler, K. Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54.



## Maxima in Particle Ratios predicted by the Thermal Model.

Ratio	Maximum at $\sqrt{s_{NN}}$ (GeV)	Maximum Value
$\Lambda / \langle \pi \rangle$	5.1	0.052
$\Xi^- / \pi^+$	10.2	0.011
$K^+ / \pi^+$	10.8	0.22
$\Omega^- / \pi^+$	27	0.0012

J. C., H. Oeschler, K. Redlich and S. Wheaton, Physics Letters B615 (2005) 50-54.

