

Experiments with Fluxon ratchet

Edward Goldobin

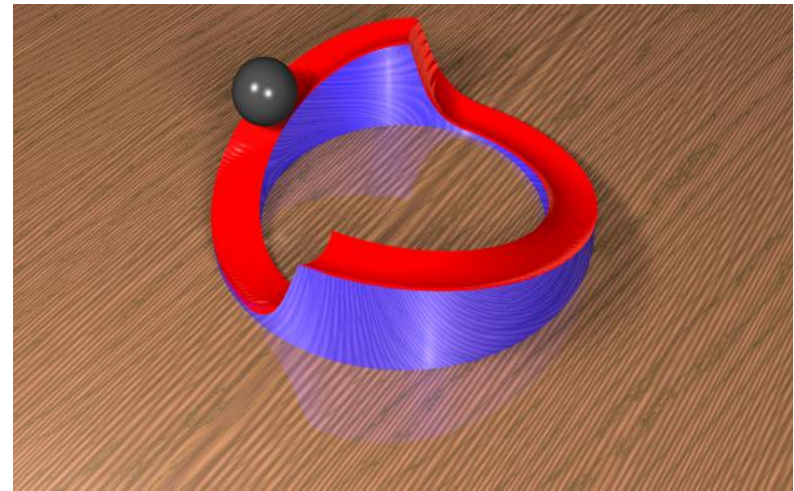
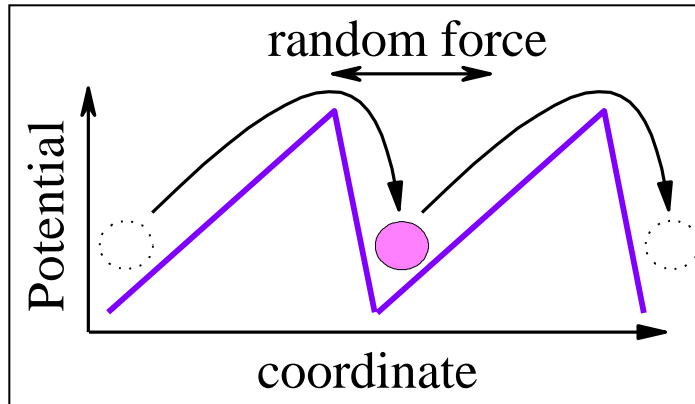
Physikalisches Institut — Experimentalphysik II, University of Tübingen, Germany



Dubna Winter School 05.02.2011

What is a Ratchet?

Point-like Brownian particle in asymmetric periodic potential



Aim: to rectify fluctuations

Reimann, Phys. Rep. **361**, 57 (2002).

P. Hänggi, F. Marchesoni, and F. Nori, Ann. Phys. (N.Y.) **14**, 51 (2005).

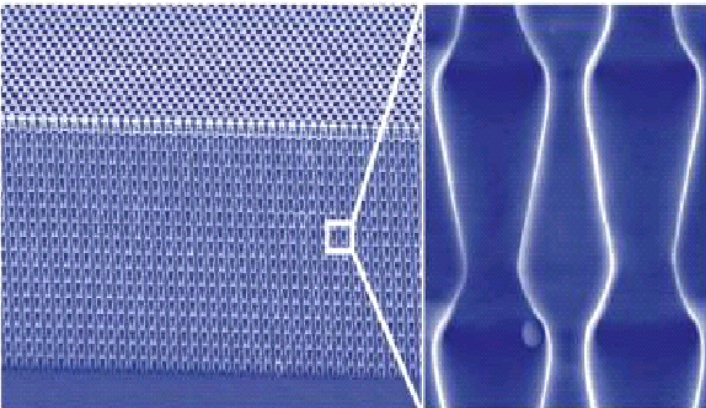
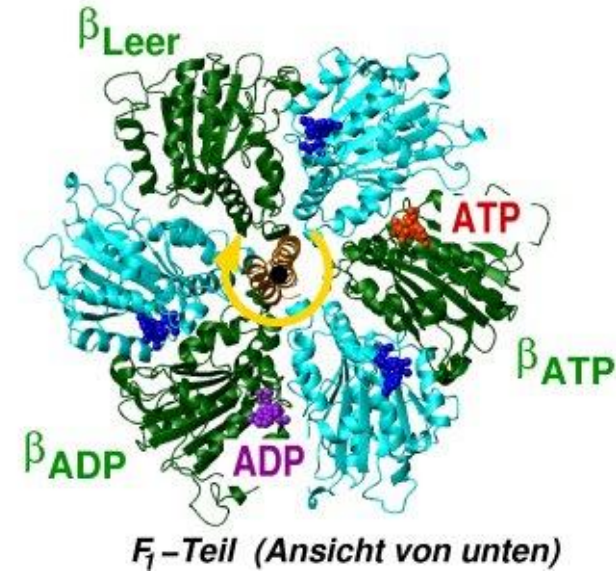
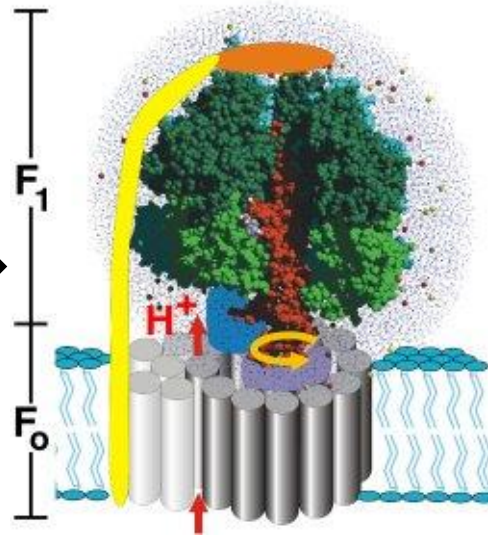
P. Hänggi and F. Marchesoni, Rev. Mod. Phys. **81**, 387 (2009).

Ratchet effect



← Classical ratchet system

Transport
in biological
systems →



← Separation of macromolecules

What is a Josephson Vortex Ratchet (JVR)?

JVR Principle (Pert. theory approx.)

Particle \rightarrow fluxon

Potential $\rightarrow w(x), h(x), j(x)$

Periodic \rightarrow annular LJJ (n periods)

Advantages of "Josephson"

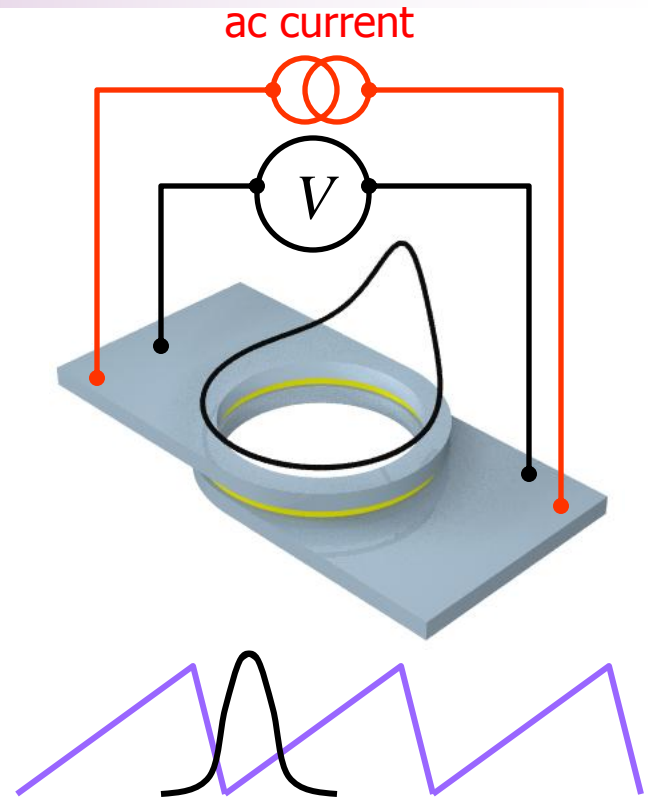
Directed motion \rightarrow voltage


Band \rightarrow up to 100GHz

overdamped \leftrightarrow underdamped
quantum ratchet


Peculiarities

Particle \rightarrow relativistic soliton



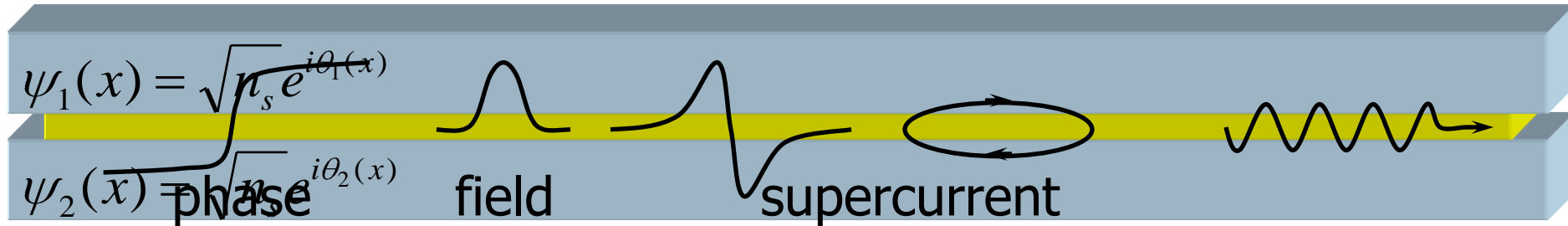
 Carapella, PRB **63**, 54515 (2001)

 Carapella et al., PRL **87**, 77002 (2001)

 E. Goldobin, et al. Phys. Rev. E **63**, 031111 (2001)

 M. Beck, et al. Phys. Rev. Lett. **95**, 090603 (2005)

Long Josephson junctions



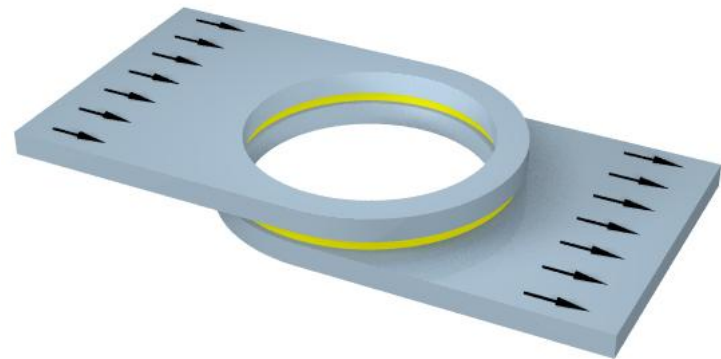
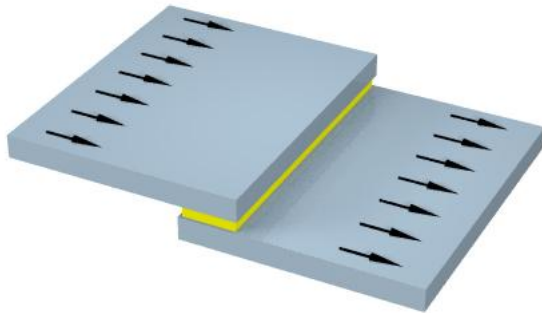
1D model:

$$\phi_{xx} - \phi_{tt} - \sin(\phi) = \alpha\phi_t - \gamma(x)$$

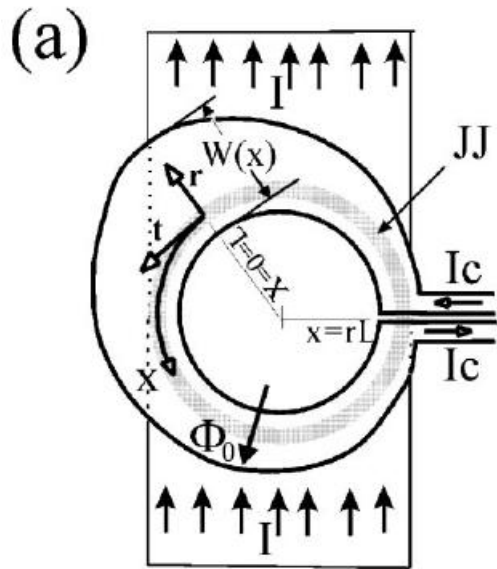
$$\lambda_J \sim 1 \dots 30 \mu\text{m}$$

$$\omega_p \sim 50 \dots 4000 \text{ GHz}$$

$$\lambda_J \cdot \omega_p \equiv \bar{c}_0 \approx \frac{c_{\text{light}}}{20 \dots 40}$$



Different approaches to construct Asymmetric Periodic Potential



Carapella et al., PRL **87**, 77002 (2001)

Carapella, PRB **63**, 54515 (2001)

Internal modes:

$$\phi_{,tt} - \phi_{,xx} + \frac{\partial V(\phi)}{\partial \phi} = - \frac{\partial V_{\text{ext}}}{\partial \phi} - \beta \phi_{,t} + \eta(x,t)$$

Willis et al., PRE **69**, 056612 (2004)

PRE **71**, 016604 (2005)

Morales-Molina, et al., PRL **91**, 234102 (2003)

Constantini, et al., PRL **87**, 114102 (2001)

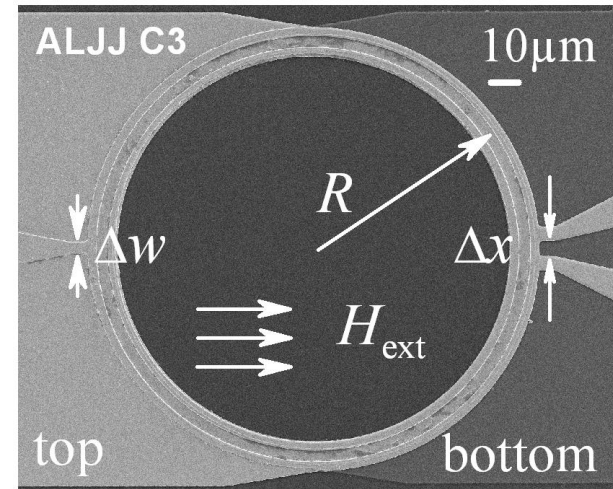
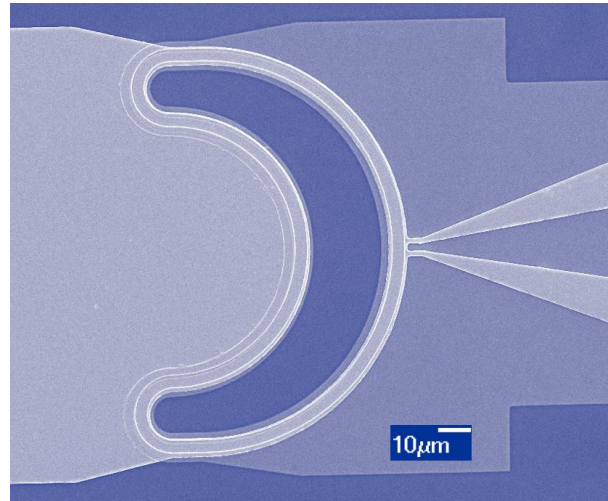
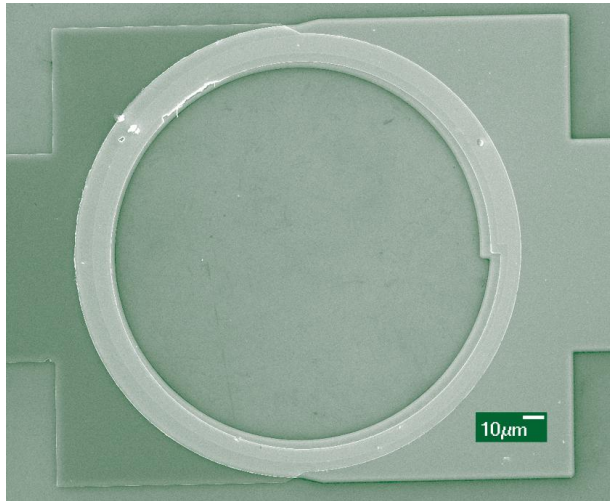
Asymmetric substrate pot.+symm. drive: move soliton in some dir.

Salerno et al., PRE **65**, 25602 (2002);

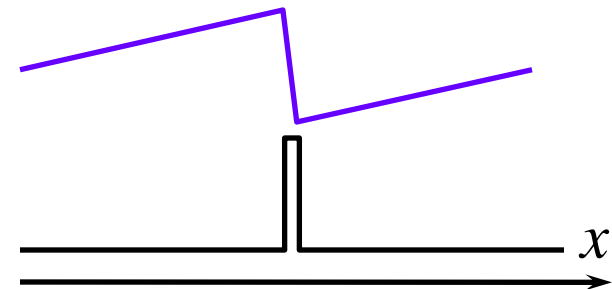
Constantini et al., PRE **65**, 51103 (2002)

Asymmetric potential

$$\phi_{xx} - \phi_{tt} - \sin(\phi) = \alpha\phi_t - \gamma(x) + h_x(x) - \frac{w_x(x)}{w(x)} [\phi_x - h(x)]$$



| $U(x)$ | Potential: | $w(x)$ | $h(x)$ | $j(x)$ | $\phi h(x_0)$ |
|---------------|------------|--------|--------|--------|---------------|
| tunability? | no | no | yes | yes | |
| topo. limits? | no | no | yes | no | |
| I/H source? | no | no | yes | yes | |



JJ width modulation – potential for fluxon

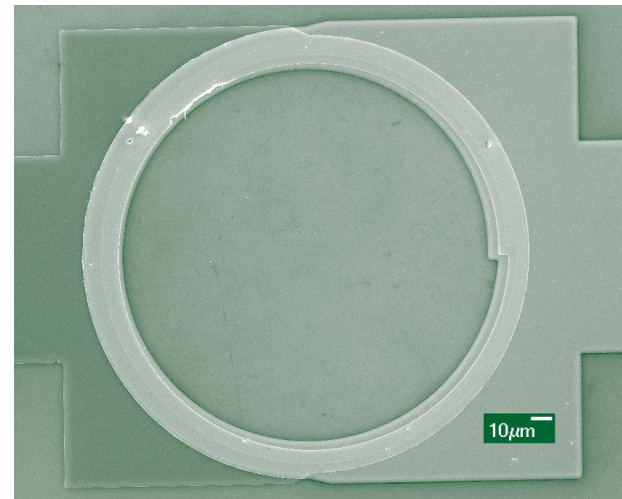
$$\phi_{xx} - \phi_{tt} - \sin(\phi) = \alpha\phi_t - \gamma(t) - \frac{w_x(x)}{w(x)}\phi_x$$

$$H = \int_0^L \left[\frac{\phi_t^2}{2} + \frac{\phi_x^2}{2} + (1 - \cos \phi) \right] w(x) dx$$

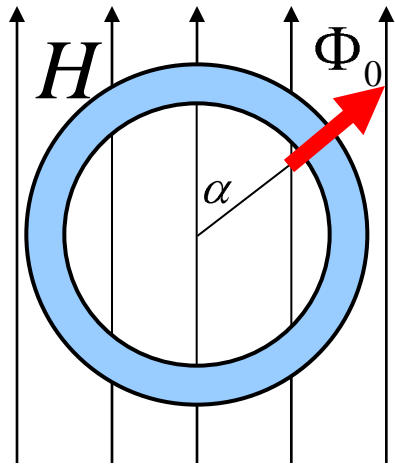
$$\phi(x, t) = 4 \arctan \exp \left(\frac{x - ut - x_0}{\sqrt{1 - u^2}} \right)$$

$$U(x_0) = \int_0^L \frac{4w(x)}{\cosh^2(x - x_0)} \approx 8w(x_0)$$

Non-relativistic approx,
“slow” modulation



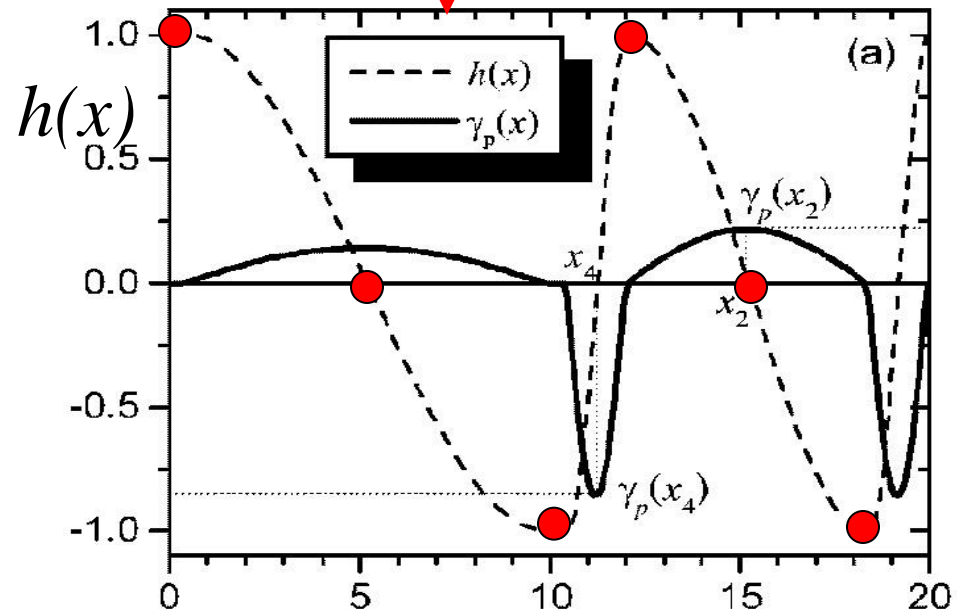
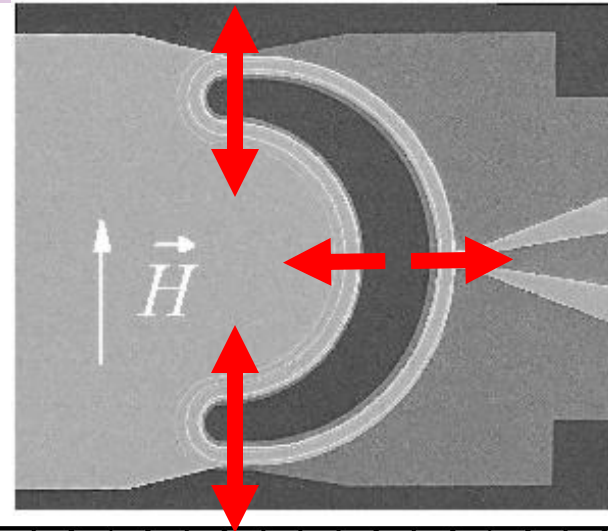
Asymmetry due to non-uniform field



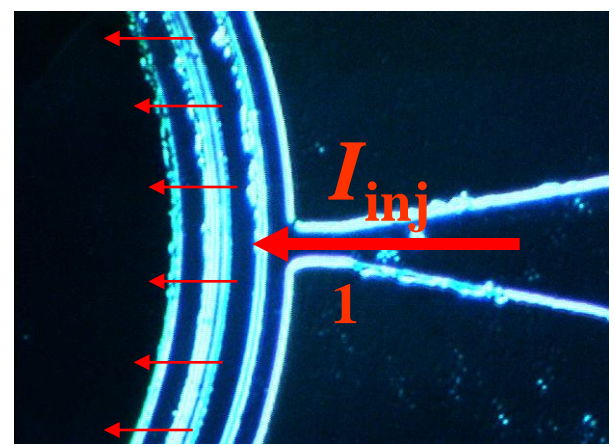
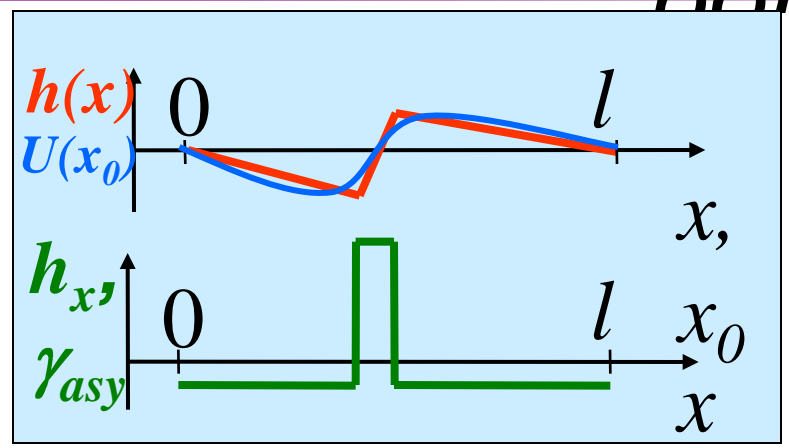
$$H(x) = H \cdot \cos \alpha$$

$$U(x_0) = -2\pi w h(x_0)$$

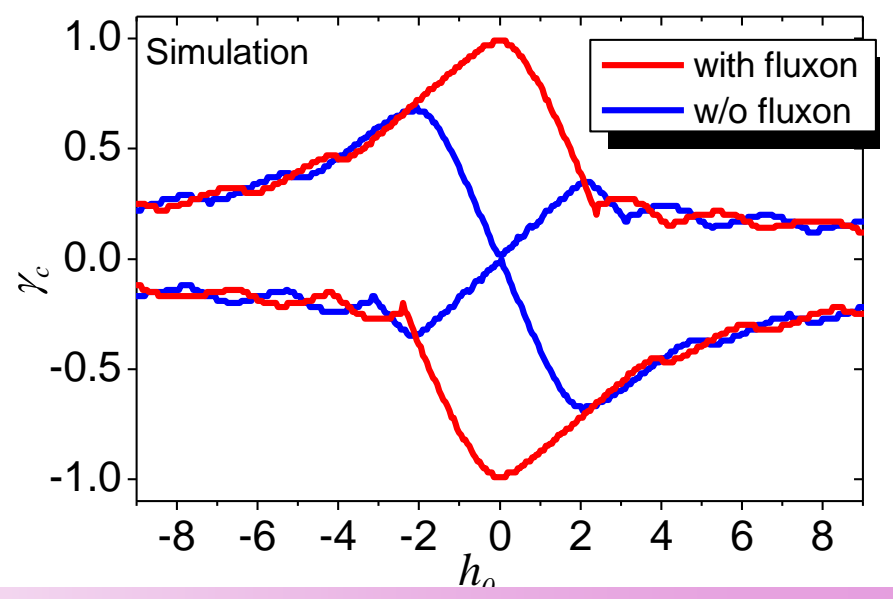
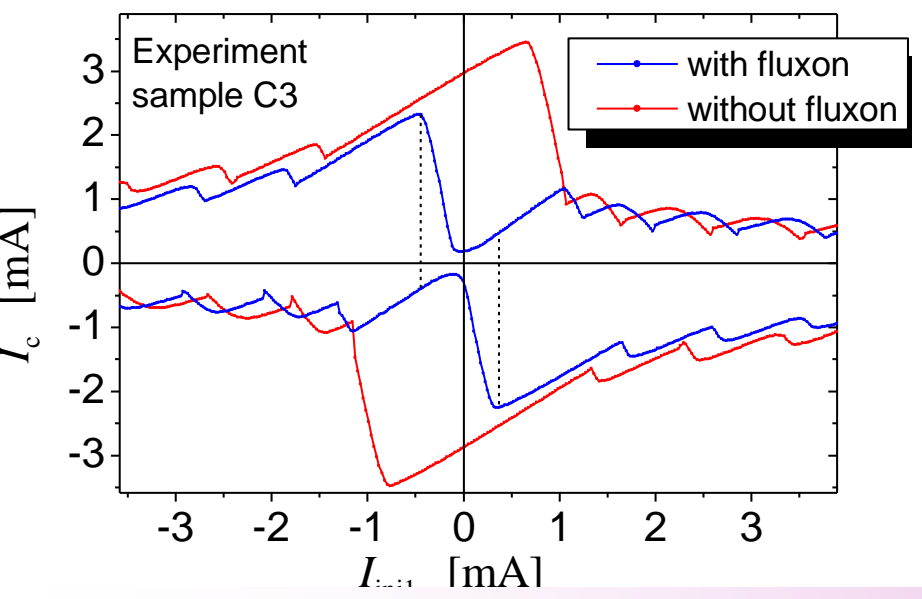
- Tunable potential
- Topological limitations



How an injector can create strongly asymmetric saw-tooth potential?



$$\frac{m_-}{m_+} = 6,6$$



System under investigation: Relativistic Deterministic JVR

For theorists

$$\begin{aligned} \phi_{xx} - \phi_{tt} - \sin \phi &= \\ &= \alpha \phi_t + h_x(x) - \gamma_{ac} \sin(\omega t) \end{aligned}$$

α – damping ($\alpha \ll 1$)



$h(x)$ – potential profile (saw-tooth)

γ_{ac} – amplitude of ac drive

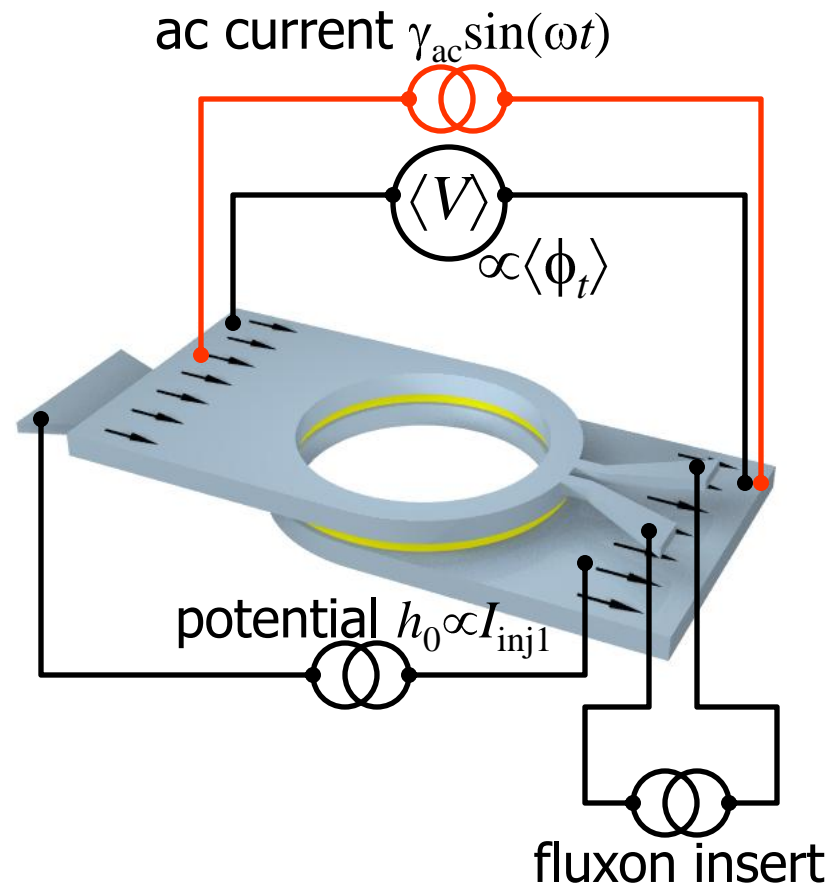
γ_{dc} – dc force (usually =0)

$$\phi(x, t) = 4 \arctan \exp \left[\frac{x - x_0(t)}{\sqrt{1 - \dot{x}^2}} \right]$$

One can derive eq. for $x_0(t)$ in the PT limit, assuming rigid fluxon shape.

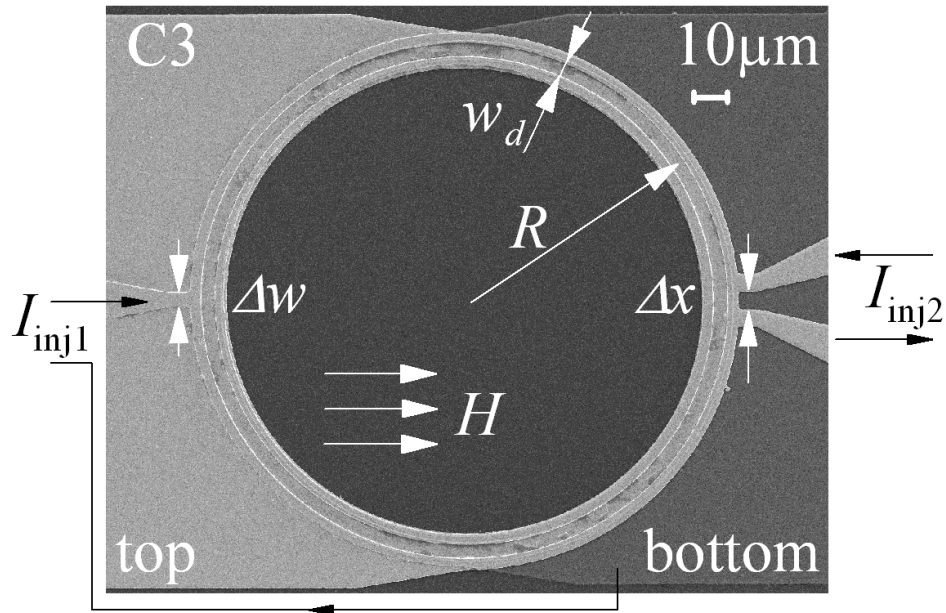
-  E. Goldobin et al, PRE **63**, 031111 (2001)
-  G. Carapella et al., PRB **63**, 054515 (2001).

For experimentalists

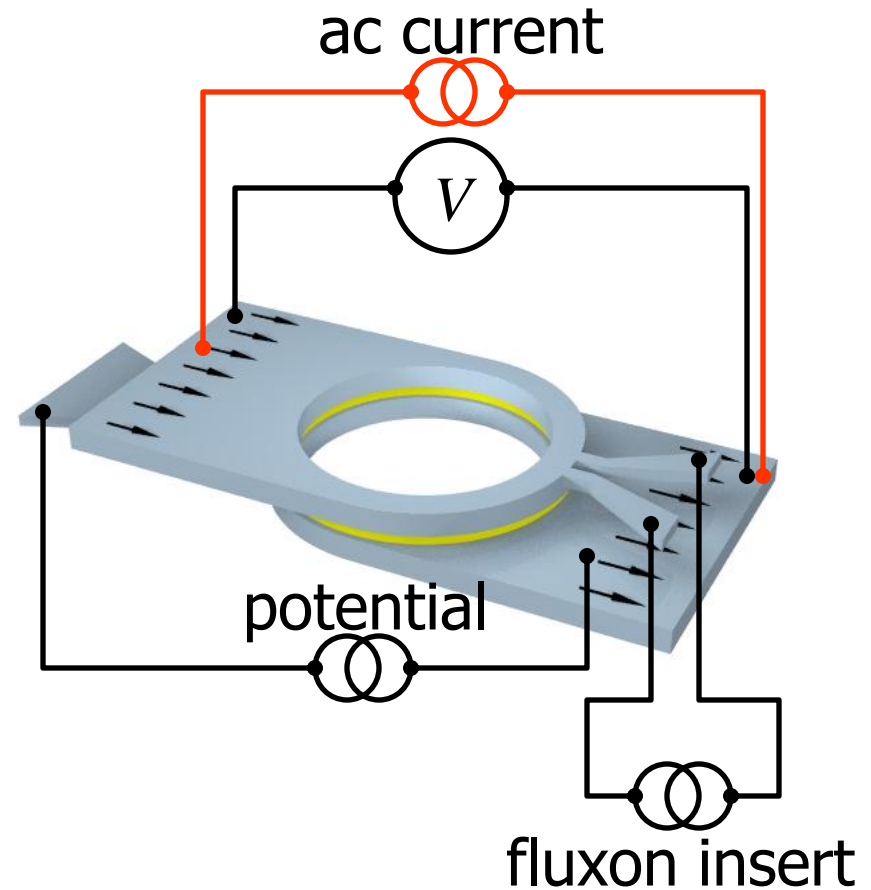


-  M. Beck, et al. PRL **95**, 090603 (2005)

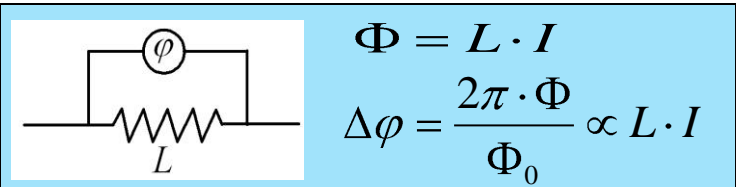
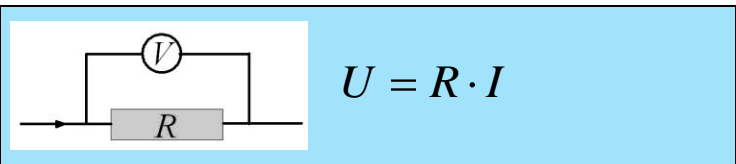
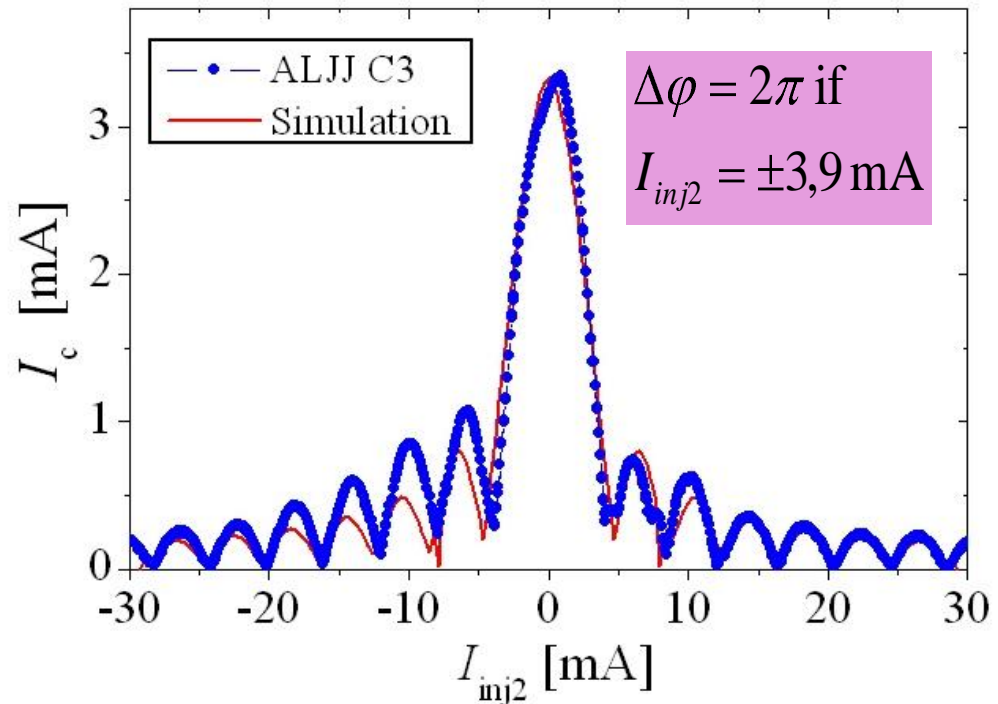
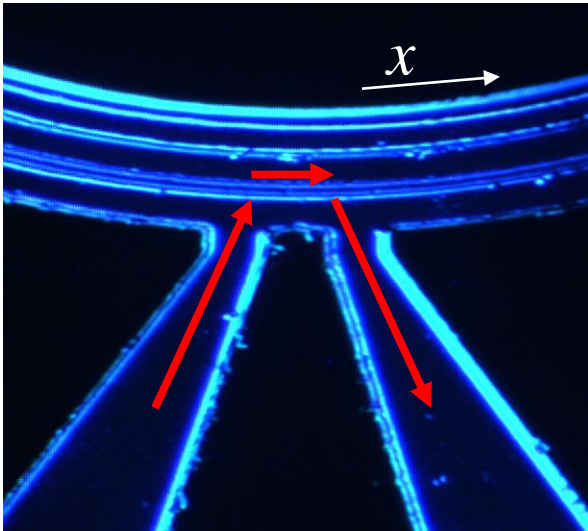
Geometry



$$R = 70\mu\text{m}, \quad \Delta x = 5\mu\text{m}, \quad \Delta w = 5\mu\text{m}$$



Insertion of fluxon



Insertion of a fluxon ($\Delta\varphi=2\pi$):
 A. V. Ustinov, APL **80**, 3153 (2002)

B. Malomed, et al. PRB **69**, 64502 (2004)

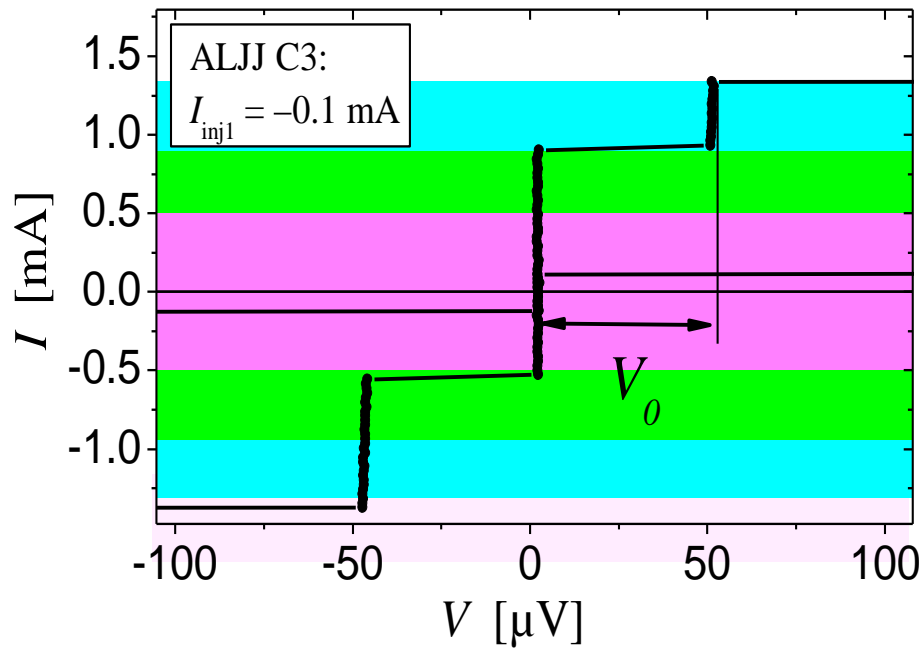
E. Goldobin, et al. PRB **67**, 224515 (2003); PRL **92**, 57005(2005)

Quasi-static drive (100Hz)

Apply : $I(t) = I_{ac} \cdot \cos(\omega t)$

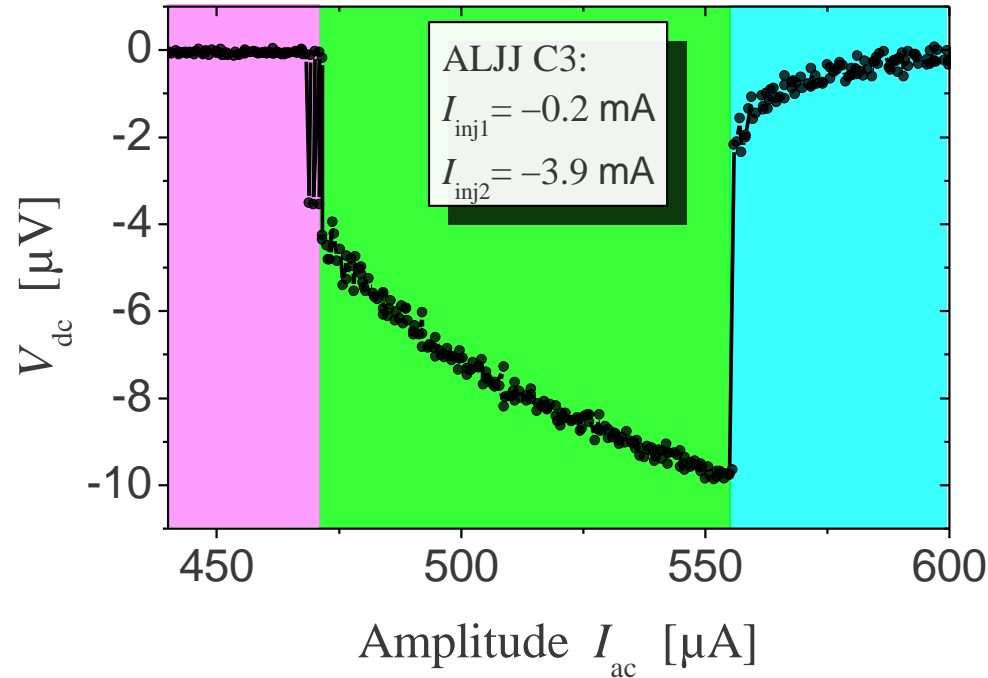
Measure : $V_{dc}(I_{ac}) = \langle V \rangle(I_{ac})$

I-V Characteristic

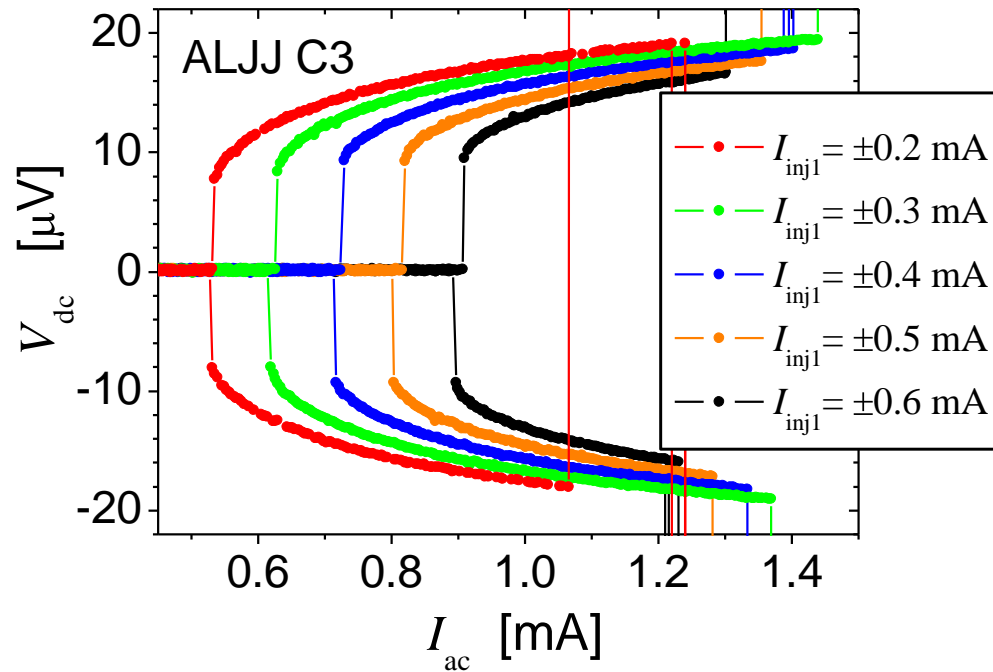
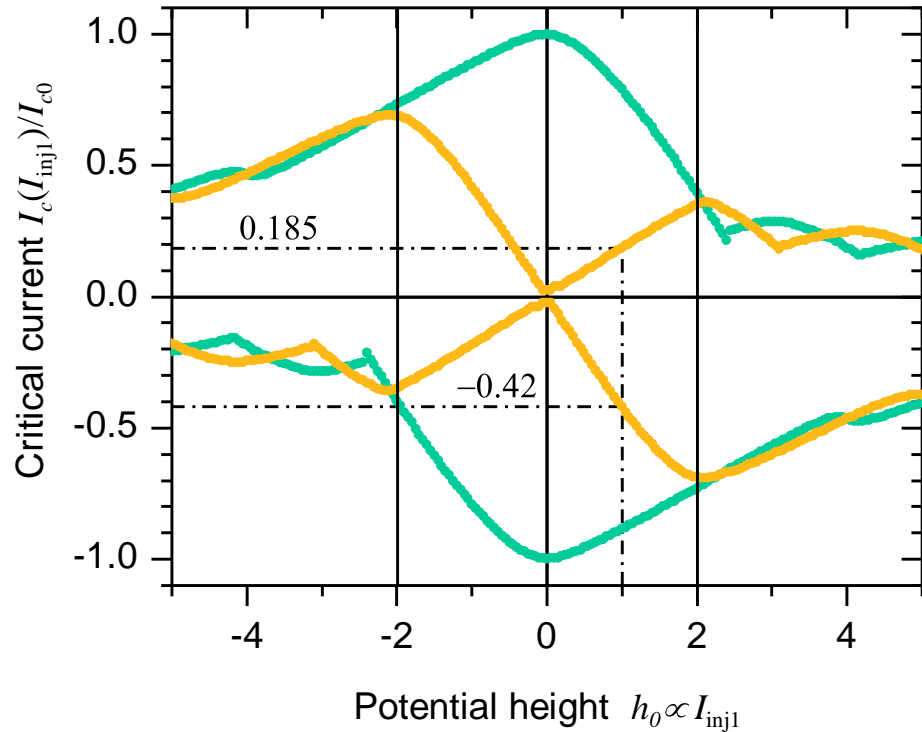


Asymmetric fluxon steps!

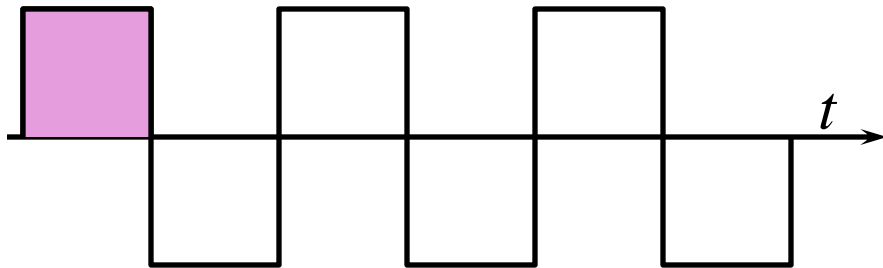
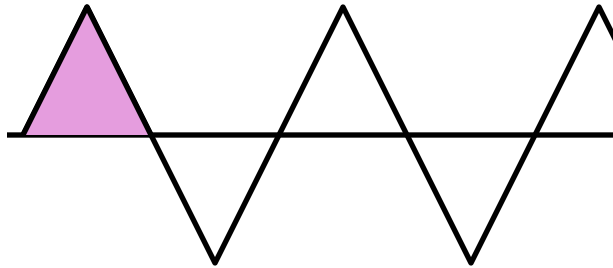
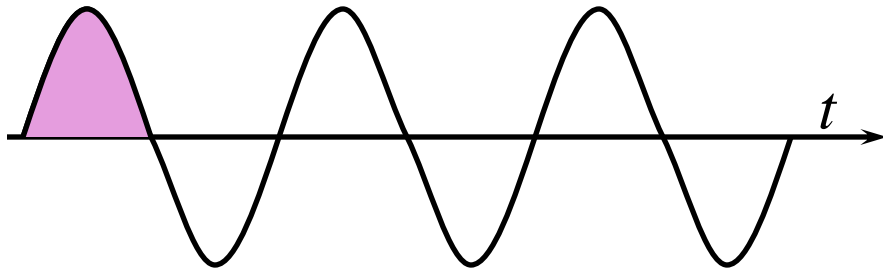
Rectification curve $V_{dc}(I_{ac})$



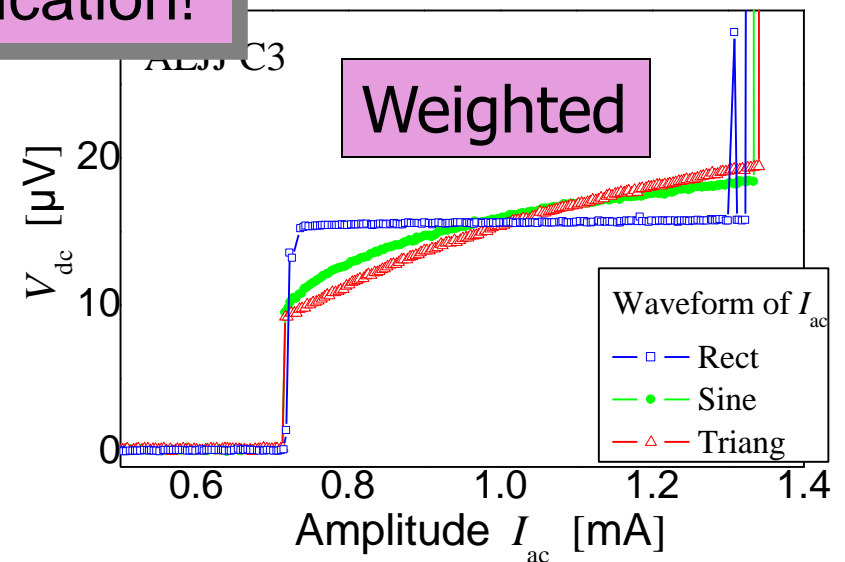
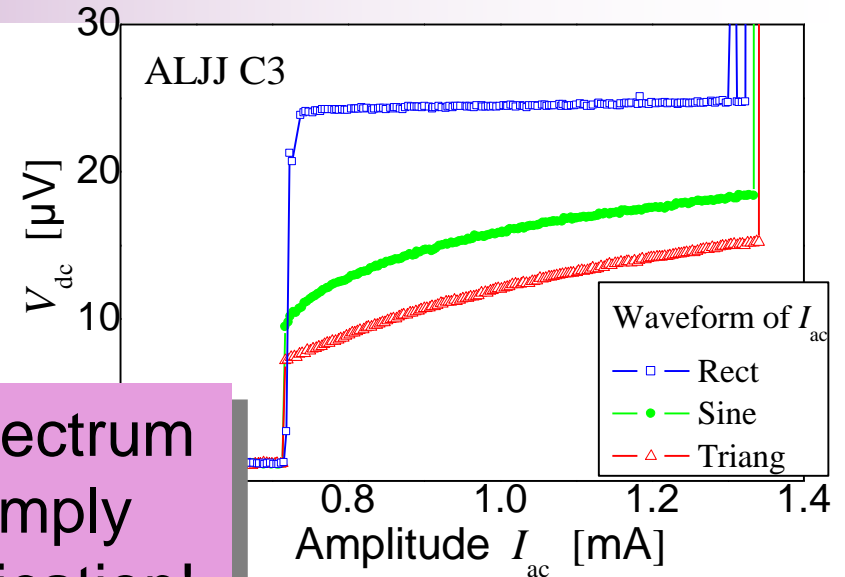
Ratchet effect vs. potential height



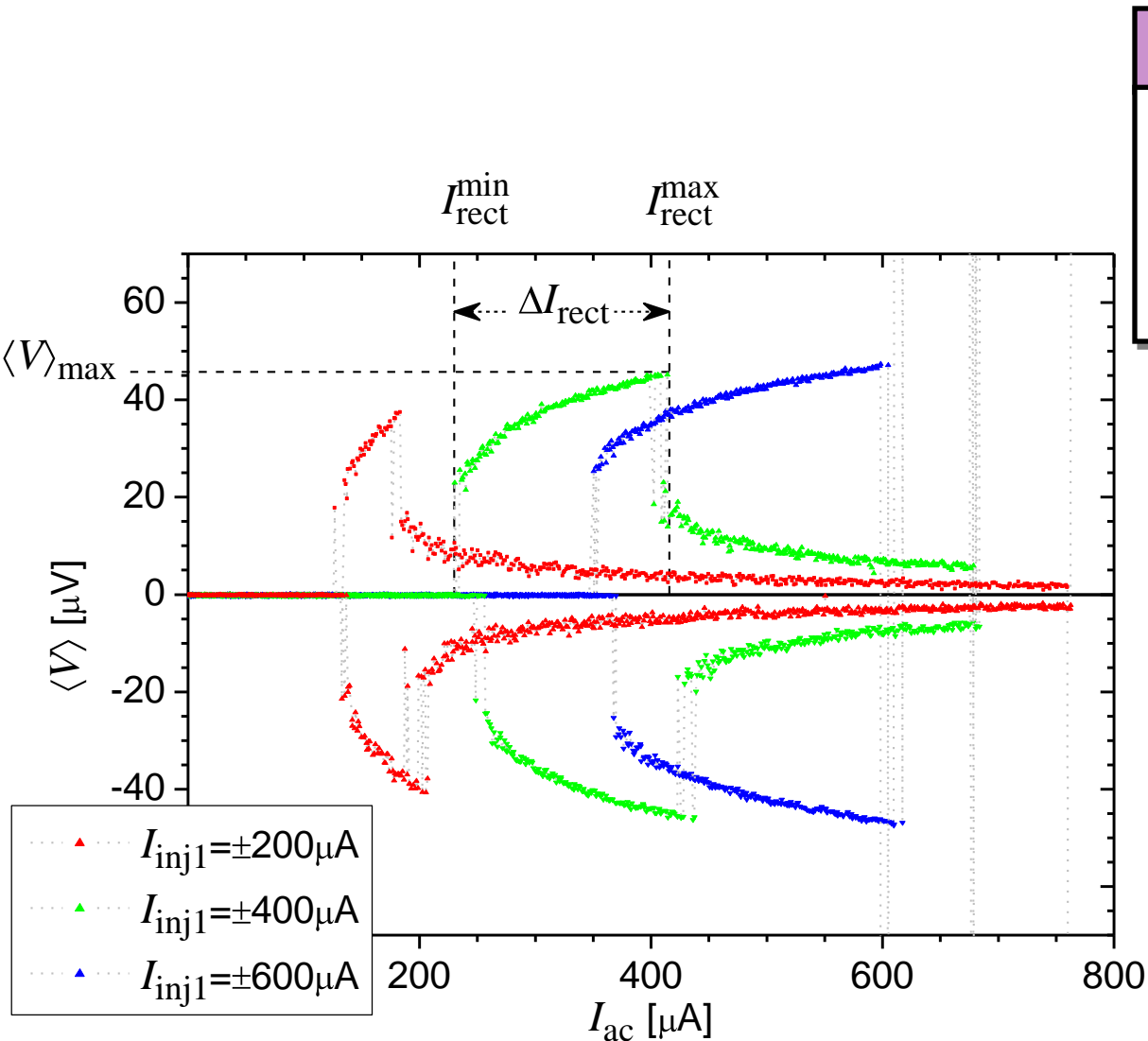
Ratchet effect vs. driver shape



Compact spectrum
does not imply
better rectification!



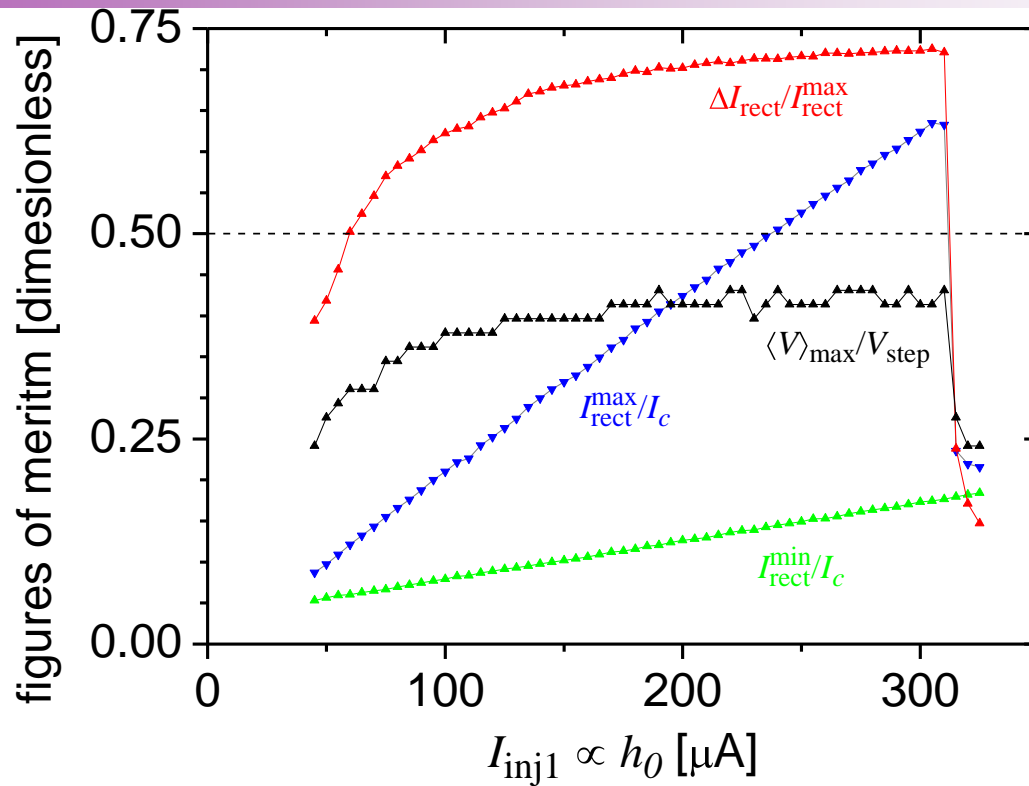
Figures of merit



Figures of merit

- min. I_{rect} (aim: $\downarrow, h_0 \downarrow$)
- max. I_{rect} (aim: $\uparrow, h_0 \uparrow$)
- ΔI_{rect} (aim: $\uparrow, h_0 \uparrow$)
- $\langle V \rangle_{max}$ (aim: $\uparrow, h_0 \uparrow$)

Adiabatic JVR: optimization



Figures of merit

- min. I_{rect} (aim: $\downarrow, h_0 \downarrow$)
- max. I_{rect} (aim: $\uparrow, h_0 \uparrow$)
- ΔI_{rect} (aim: $\uparrow, h_0 \uparrow$)
- $\langle V \rangle_{\text{max}}$ (aim: $\uparrow, h_0 \uparrow$)

Result

- One cannot optimize all figures at ones.
- Increasing h_0 we optimize 3 out of 4: most importantly ΔI_{rect} and $\langle V \rangle_{\text{max}}$.

Loaded quasi-static JVR

How large should be the dc counter force to stop the ratchet?

In adiabatic regime the result can be obtained from the IVC: $V(I) = V(I_{ac} \sin(\omega t) + I_{dc})$

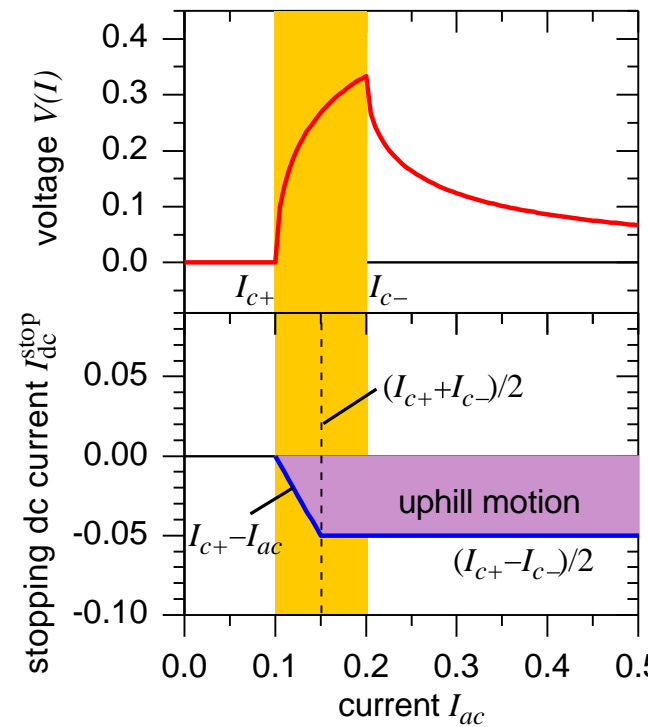
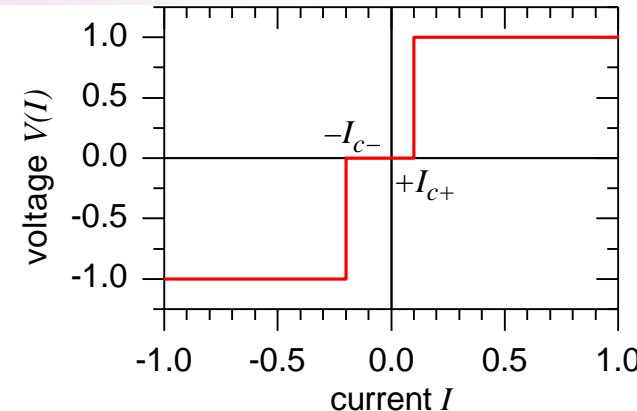
Averaging over one period we get

$$\langle V \rangle(I_{ac}, I_{dc})$$

Solving $\langle V \rangle(I_{ac}, I_{dc}) = 0$ for I_{dc} we get $I_{dc}^{\text{stop}}(I_{ac})$

Results for a simple step-like IVC model (I_{c+}, I_{c-})

- stopping force grows linearly
- it saturates in the middle of rect. window
- saturation value is equal to rect. window size
- make window as large as possible (large h_0)
- operate ratchet in the 2nd half of rect. window!



High frequency drive

– non Adiabatic JVR (nA-JVR)

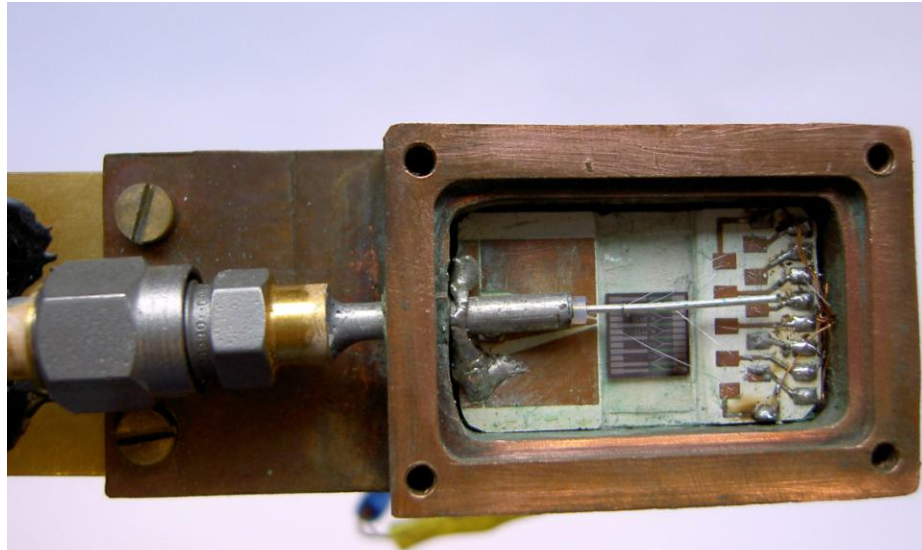
→ What means **high** frequency?

- Period of the drive is comparable with the time it takes for fluxon to move around the annulus
- periodic dynamics → integer number of turns
- discrete average velocity = discrete average voltage steps

→ Underdamped system: inertial effects, chaos

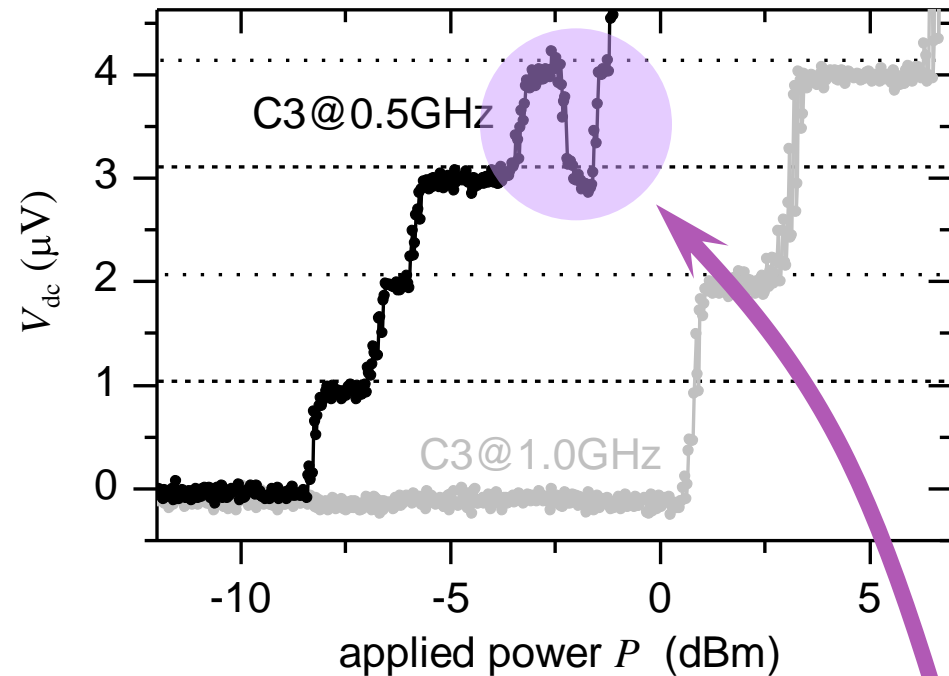
Setup:

- Resonator copper box, the first mode @ 6.0 GHz
- Sender: rf-antenna
- Receiver: JJ-Electrodes
- $I_{ac} \leftrightarrow$ power P

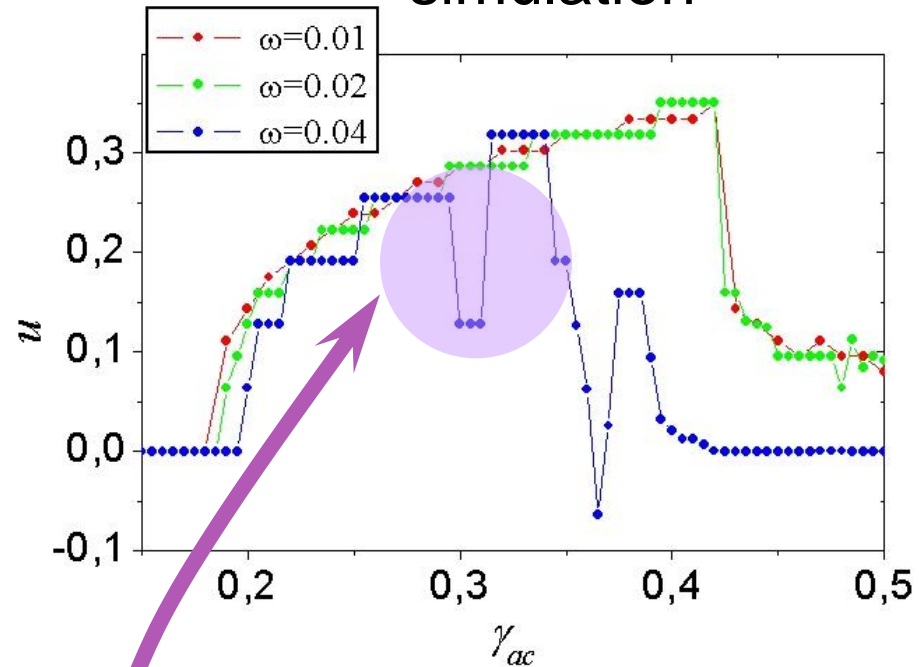


nA-JVR: Quantized rectification

experiment



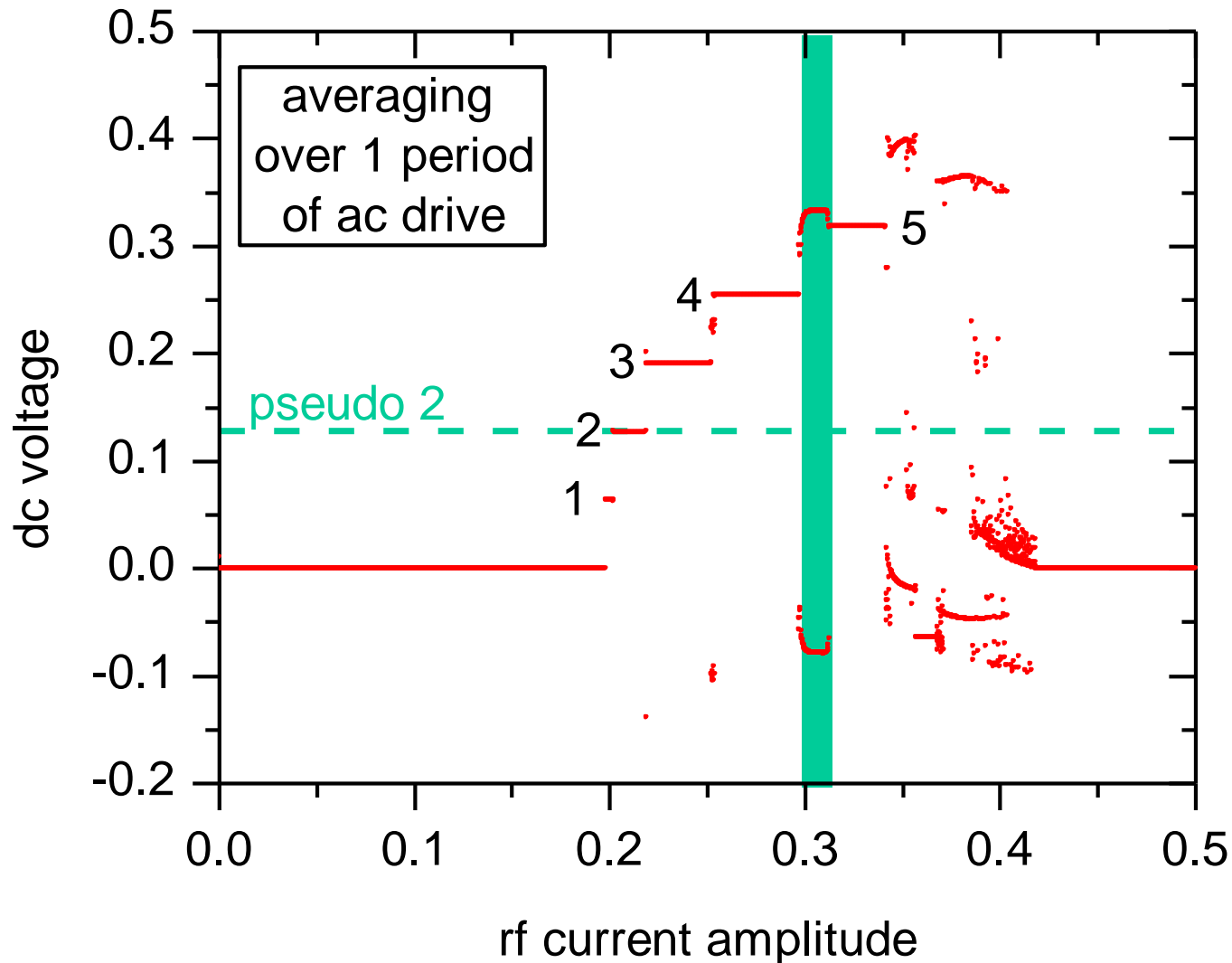
simulation



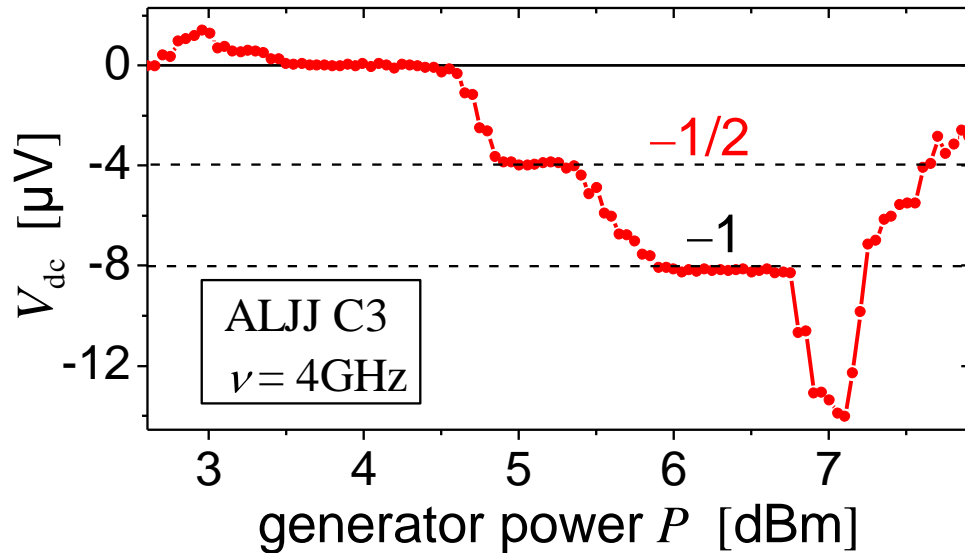
$$V_{dc}^n = n\Phi_0 \frac{\omega}{2\pi} \approx n \left[\frac{2\mu\text{V}}{\text{GHz}} \right]$$

Non-monotonous steps!

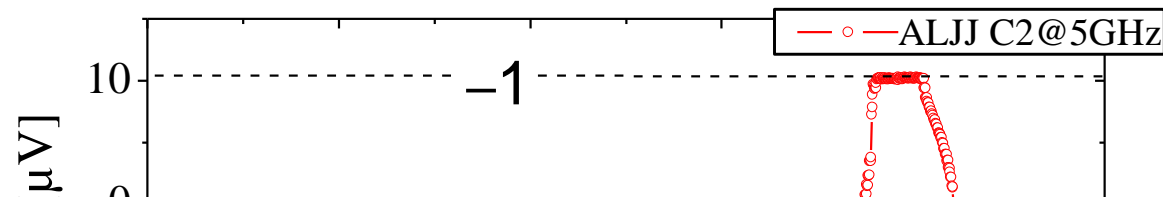
nA-JVR: Period-2 dynamics




nA-JVR: half integer steps & current reversal



One turn takes
two driving periods!



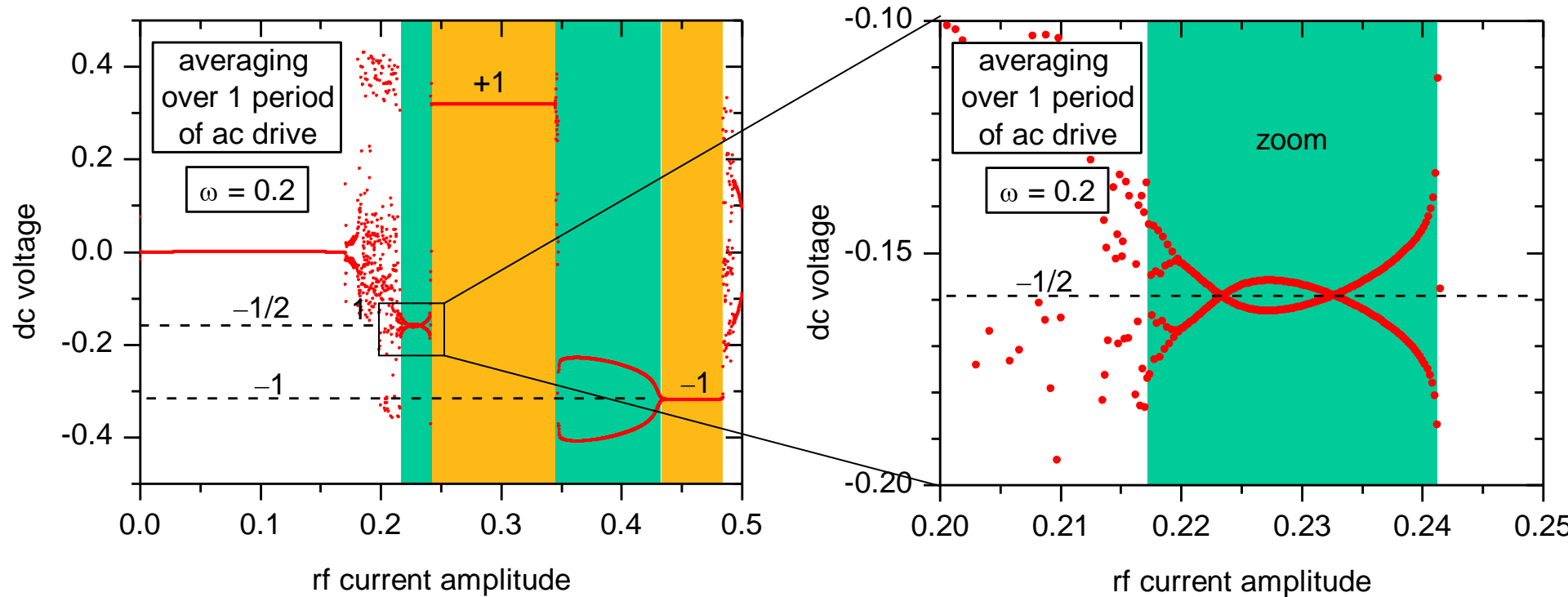
Particle's (fluxon's)
current reversal

 J. Mateos, et al. PRL **84**, 258 (2000)
Physica A **325**, 92 (2003)

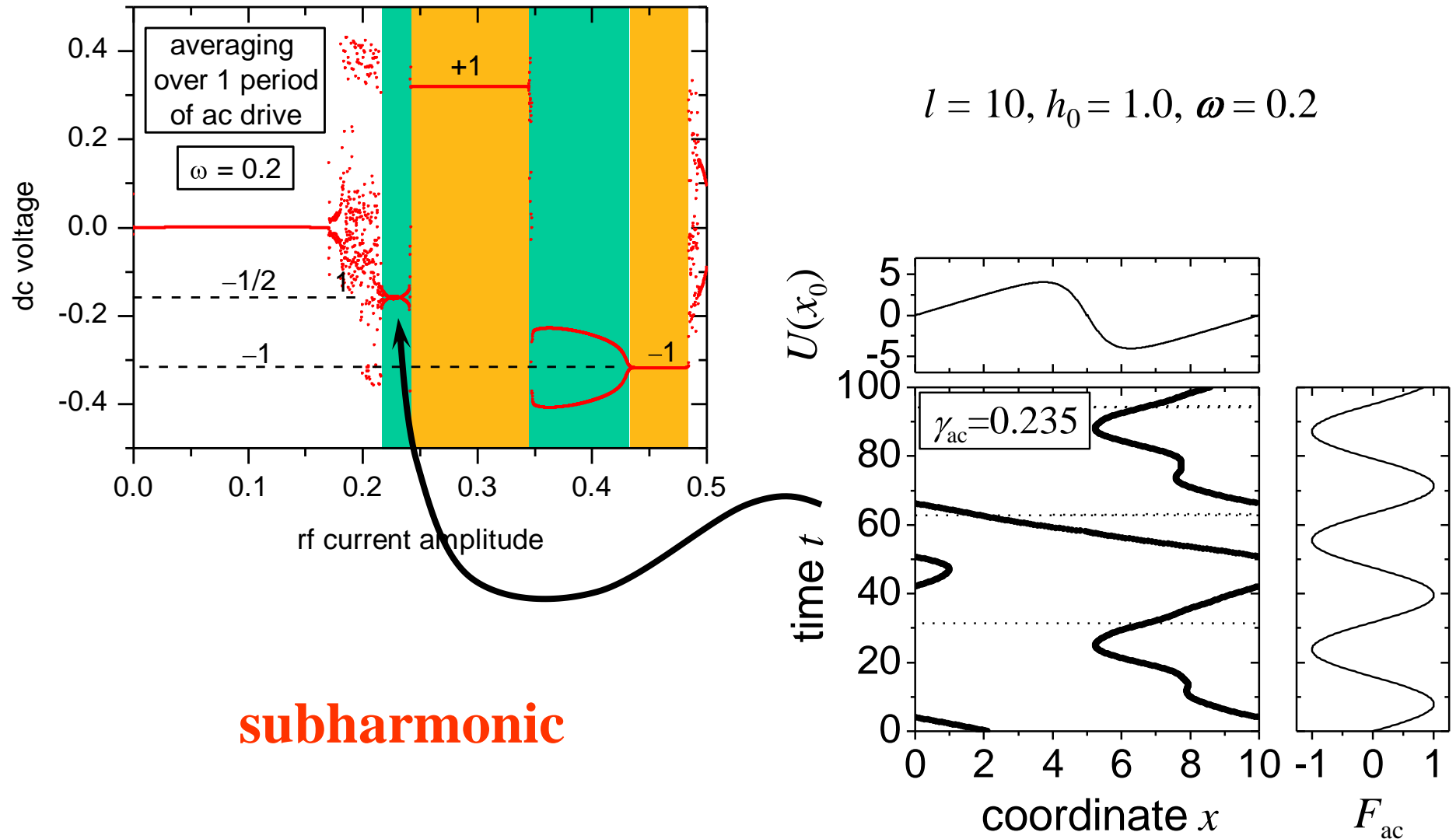
 M. Barbi, et al. Phys. Rev. E **62**, 1988 (2000)

 P. Jung, et al. Phys. Rev. Lett. **76**, 3436 (1996)

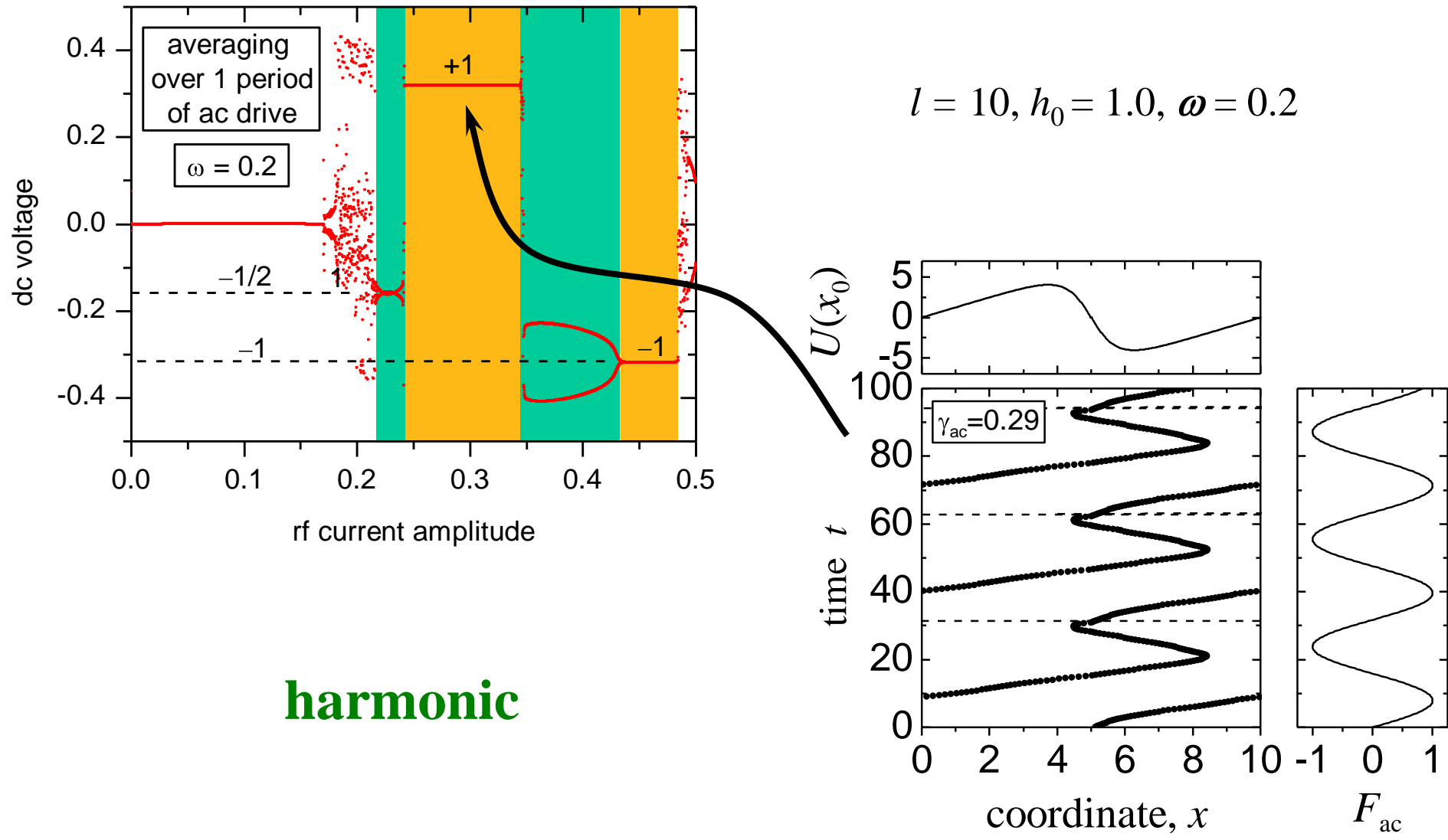
nA -JVR: Half integer steps & current reversal



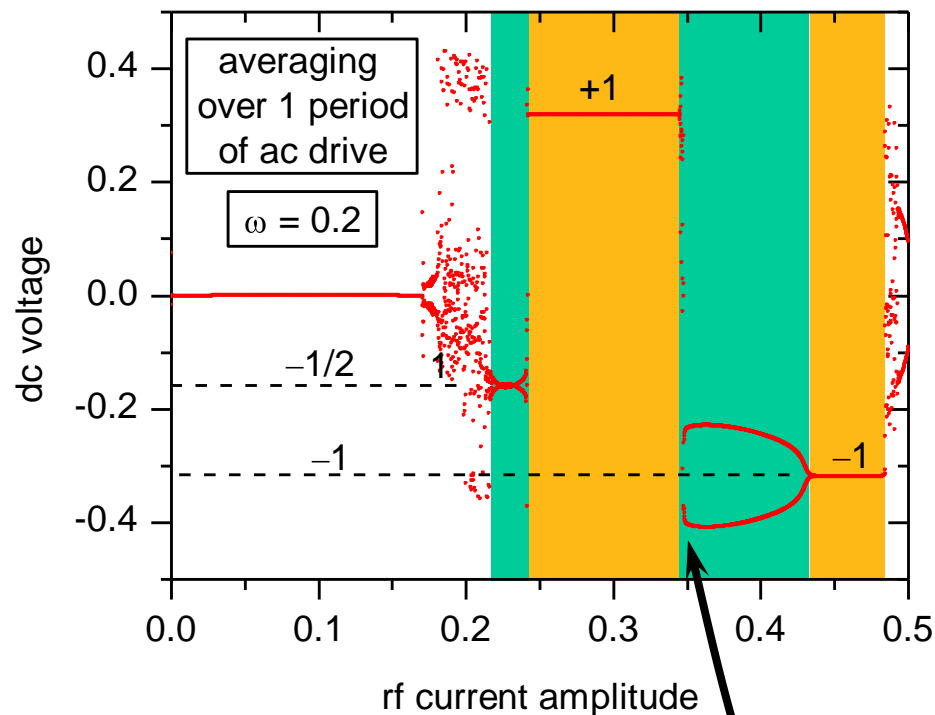
nA-JVR: fluxon trajectories & period 2 dynamics



nA-JVR: Period-1 dynamics: forward step



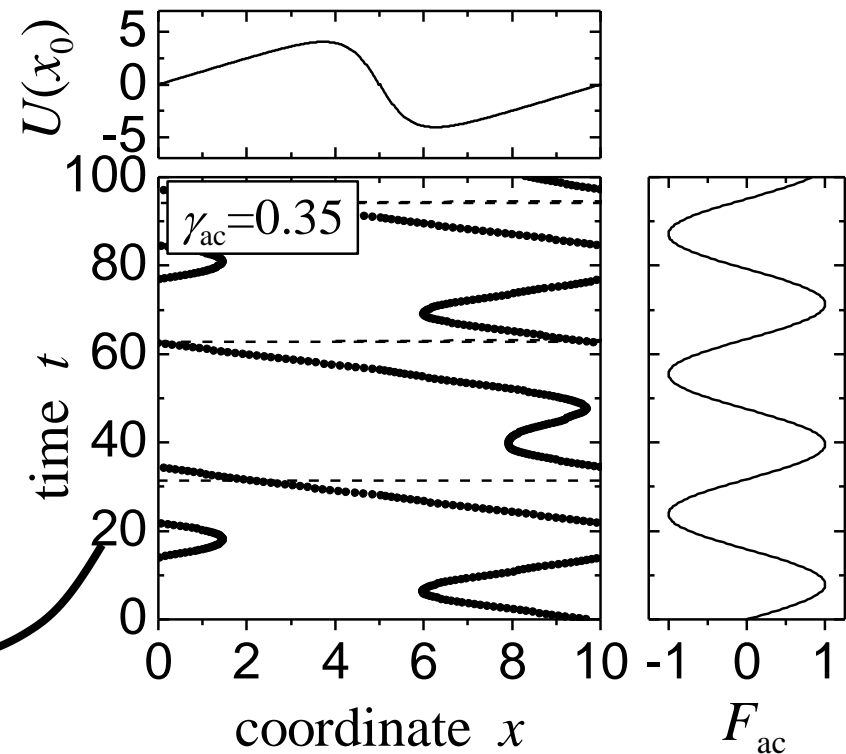
nA-JVR: Period 1&2 dynamics, current (voltage) reversal



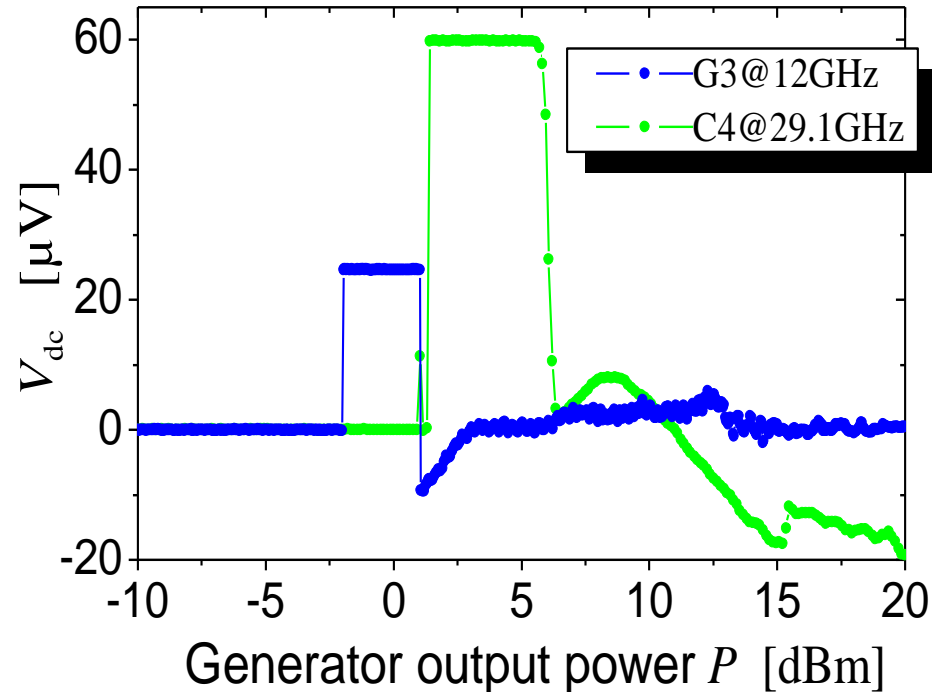
Subharmonic

current reversal: period 1&2

$$l = 10, h_0 = 1.0, \omega = 0.2$$



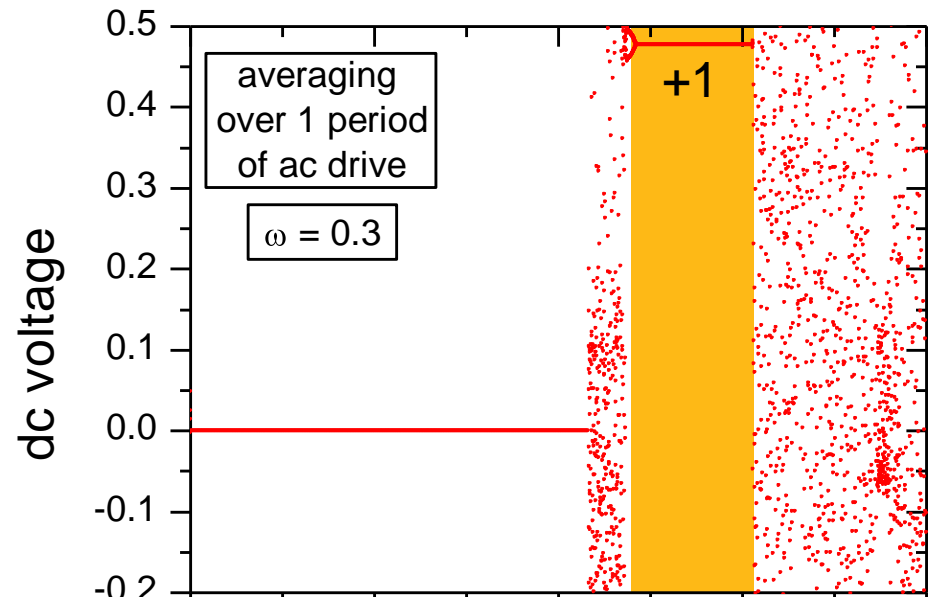
nA-JVR: Peak velocities



Average velocity :

G3 : $\langle u \rangle = 0.91 \bar{c}_0$

Simulation
Max. Average velocity
0.47 (?)



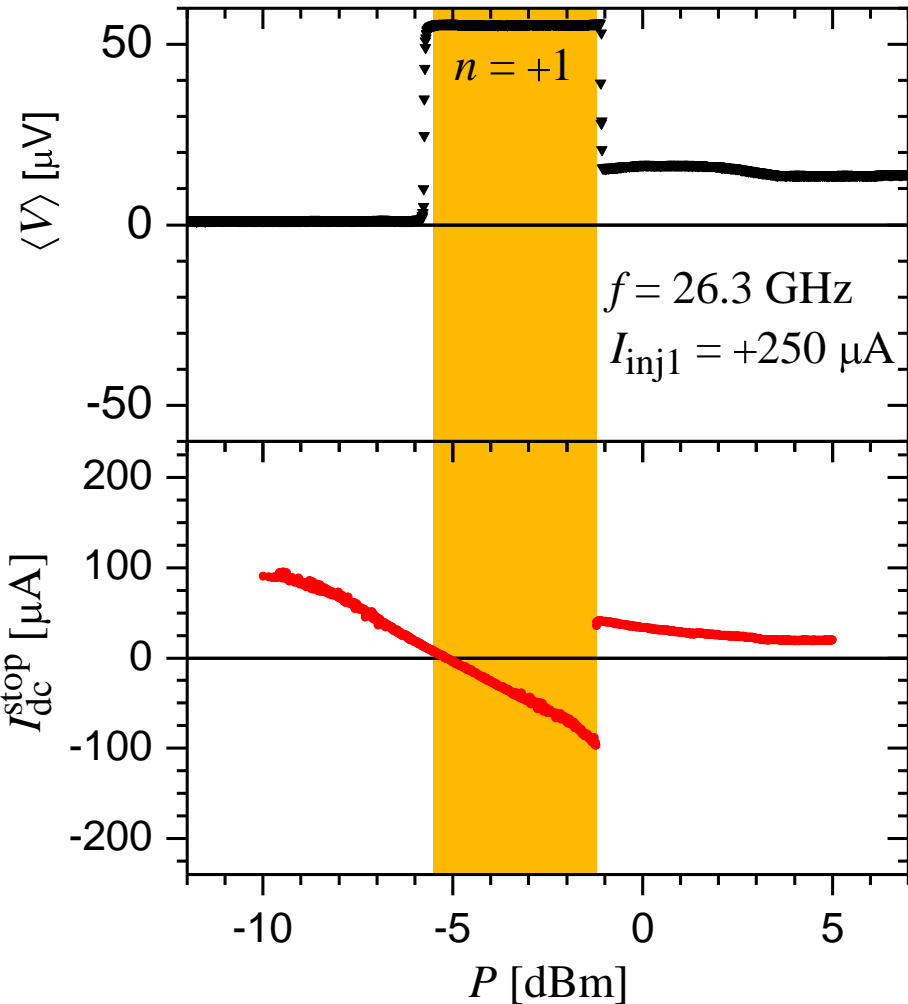
$\langle u \rangle = 0.22$: A. Ustinov, et al. PRL **93**, 87001 (2004)

$\langle u \rangle = 0.33$: G. Carapella, et al. Physica C **382**, 337 (2002)

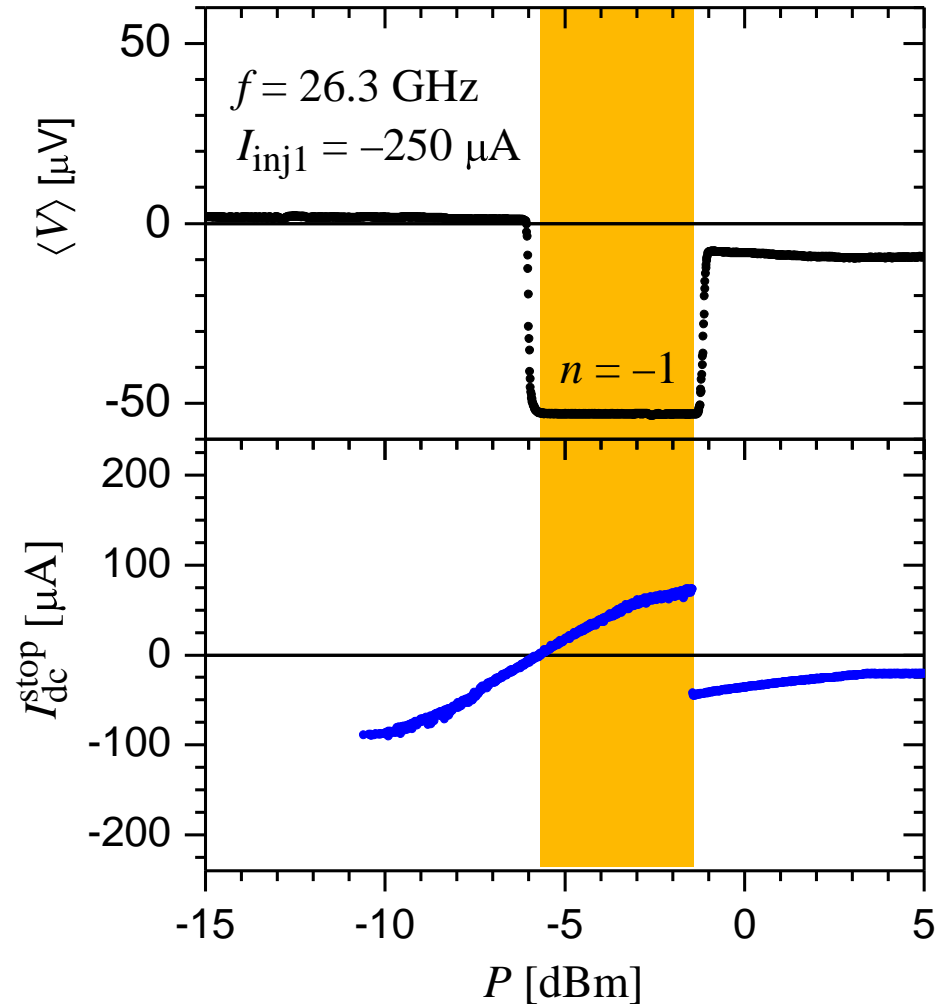
$\langle u \rangle = 0.17$ @ $\alpha = 1$: F. Mertens, et al. PRE **74**, 66602 (2006)

nA-JVR loaded by dc current

Positive potential

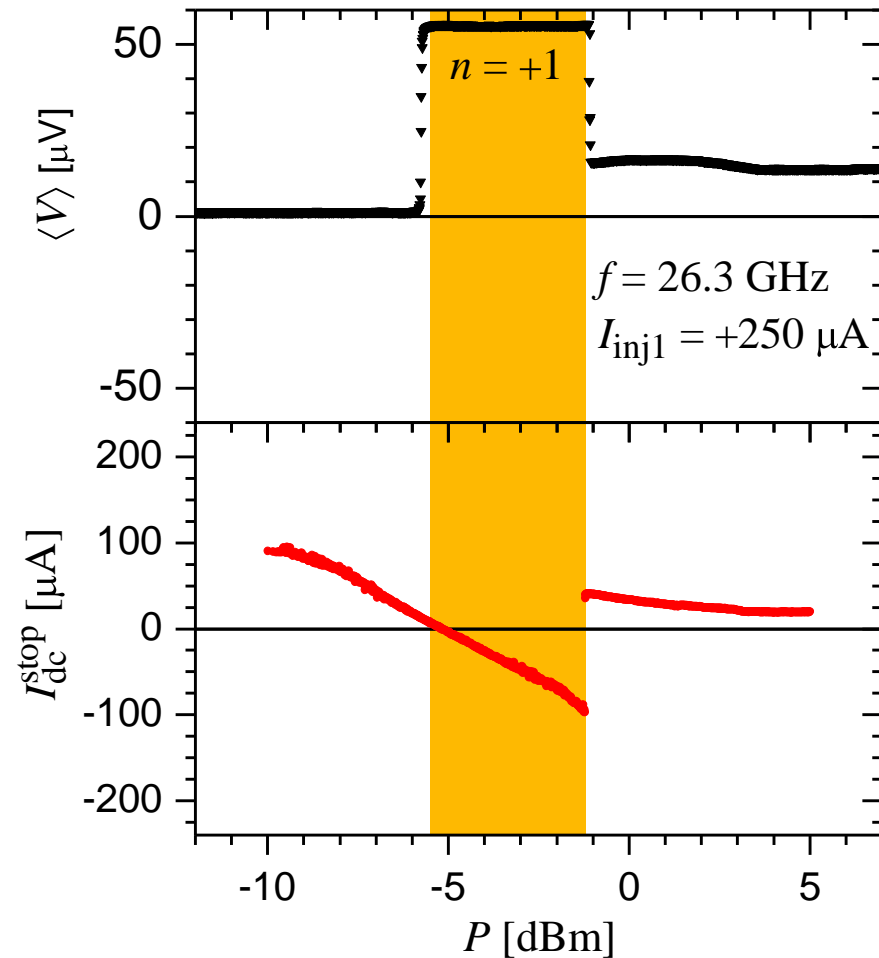


Negative potential

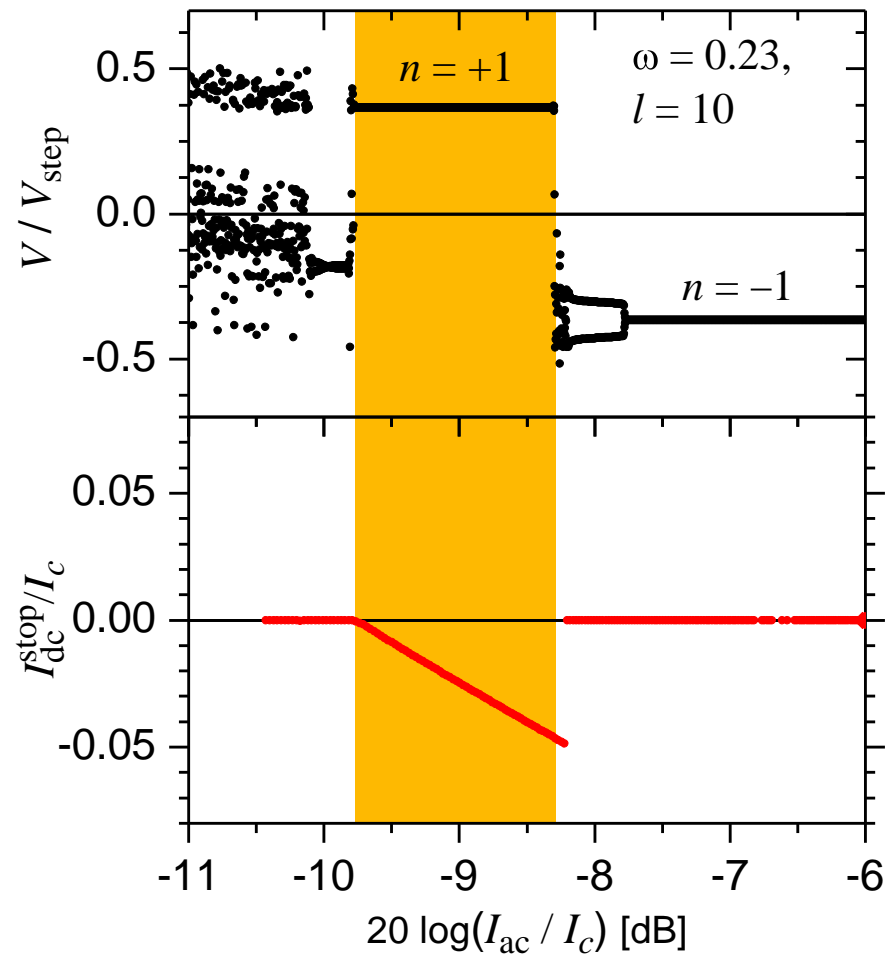


nA-JVR: stopping force (current) experiment vs. simulation

Experiment



Simulation

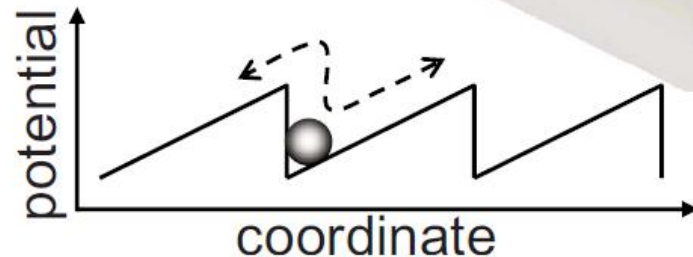
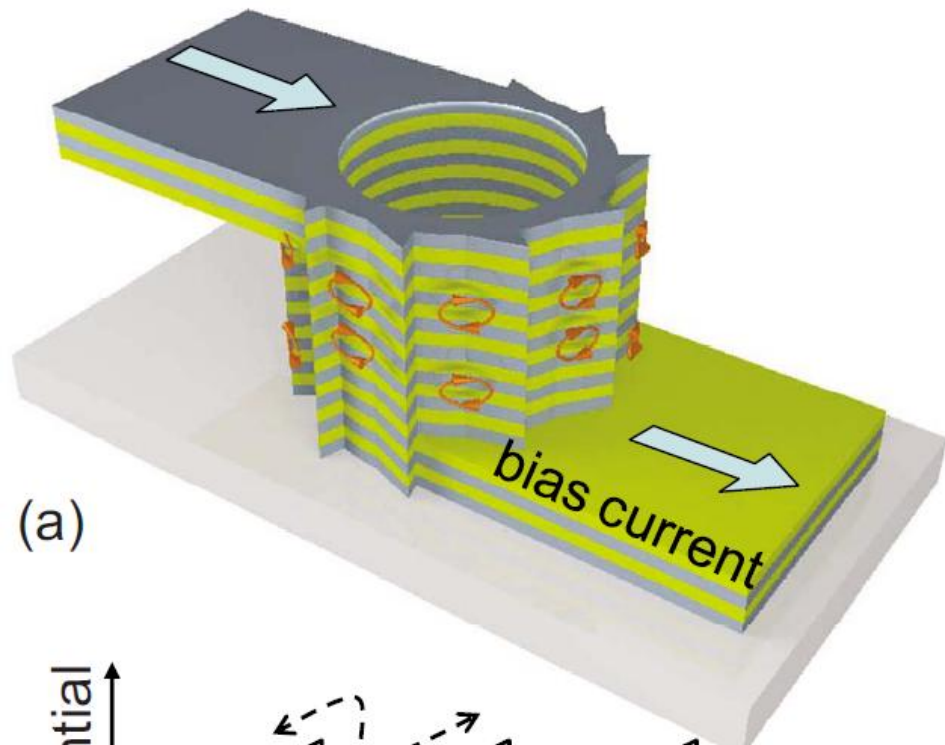


JVR made of Intrinsic JJ stack

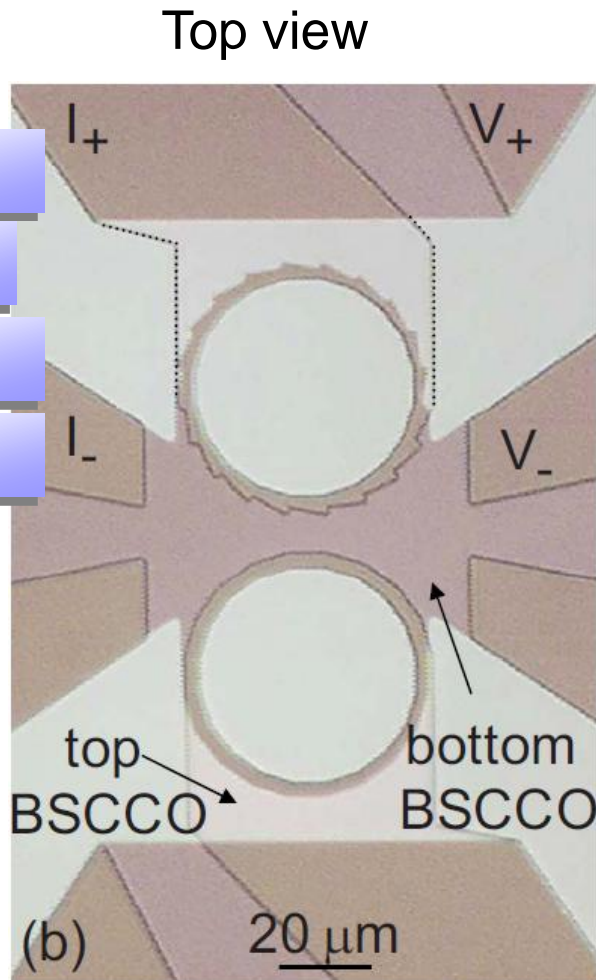
- To increase both f_{\max} and V_{\max} we need to increase v_{\max}
- To increase V_{\max} one can connect several ratchets in series
- Idea: use **intrinsic Josephson junctions** (IJJ)
- $f_{\text{pl}} \sim 150 \text{ GHz}$
- However, IJJ are not independent (coupled).

$$c_q = \frac{c_s}{\sqrt{1 - 2s \cos(\pi q / (N - 1))}}, \quad q = 1 \dots N.$$

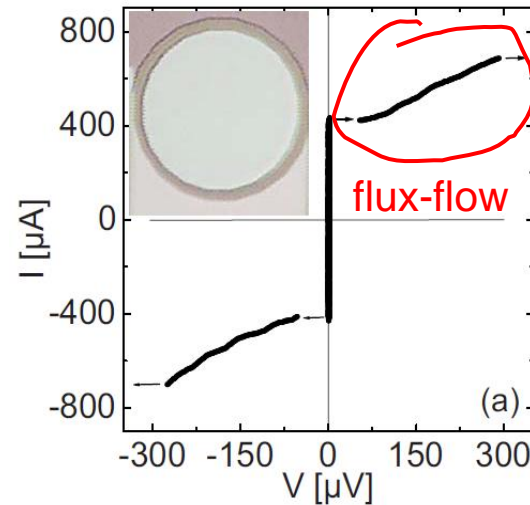
In-phase mode: $c_1 \rightarrow c_0 / \sqrt{\epsilon} \gg c_s$.



Samples, I - V characteristic



$$V_{\text{max}}^{\text{FF}} = m \Phi_0 v_{\text{max}} / L$$

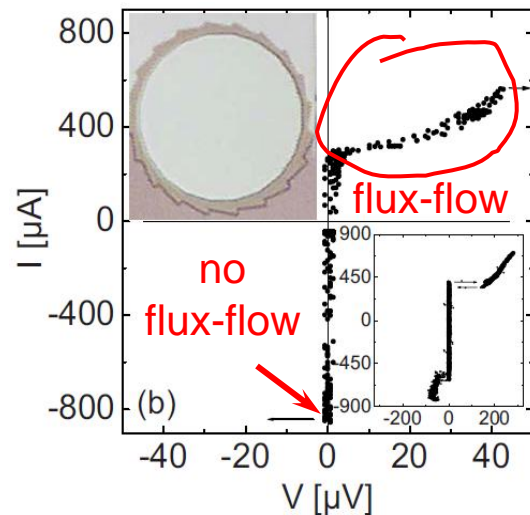


$$T = 4.2 \text{ K}$$

$$V_{\text{max}} = 300 \mu\text{V}$$

$$m v_{\text{max}} = 2.3 \times 10^7 \text{ m/s}$$

$$m_{\text{max}} = 100 \text{ fluxons}$$



$$T = 8.0 \text{ K}$$

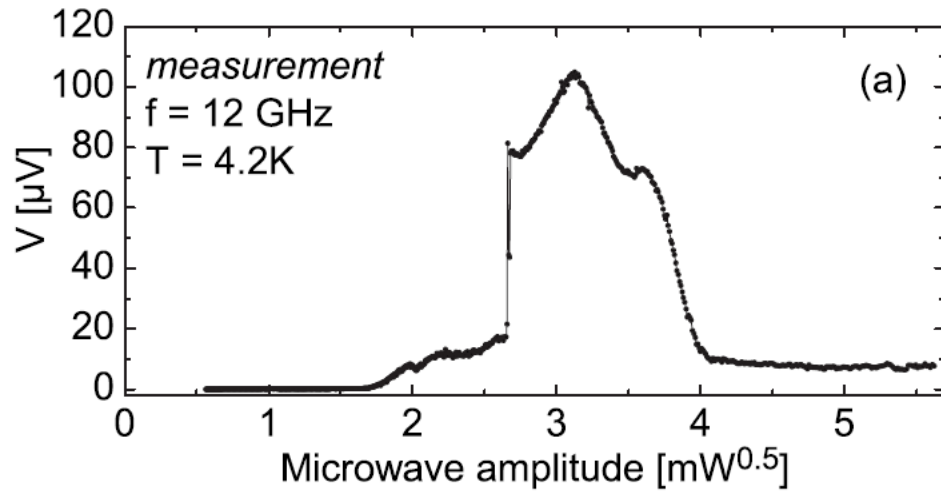
$$V_{\text{max}} = 40 \mu\text{V}$$

$$m v_{\text{max}} = 1.5 \times 10^7 \text{ m/s}$$

$$m_{\text{max}} = 15$$

Rectification

Experiment



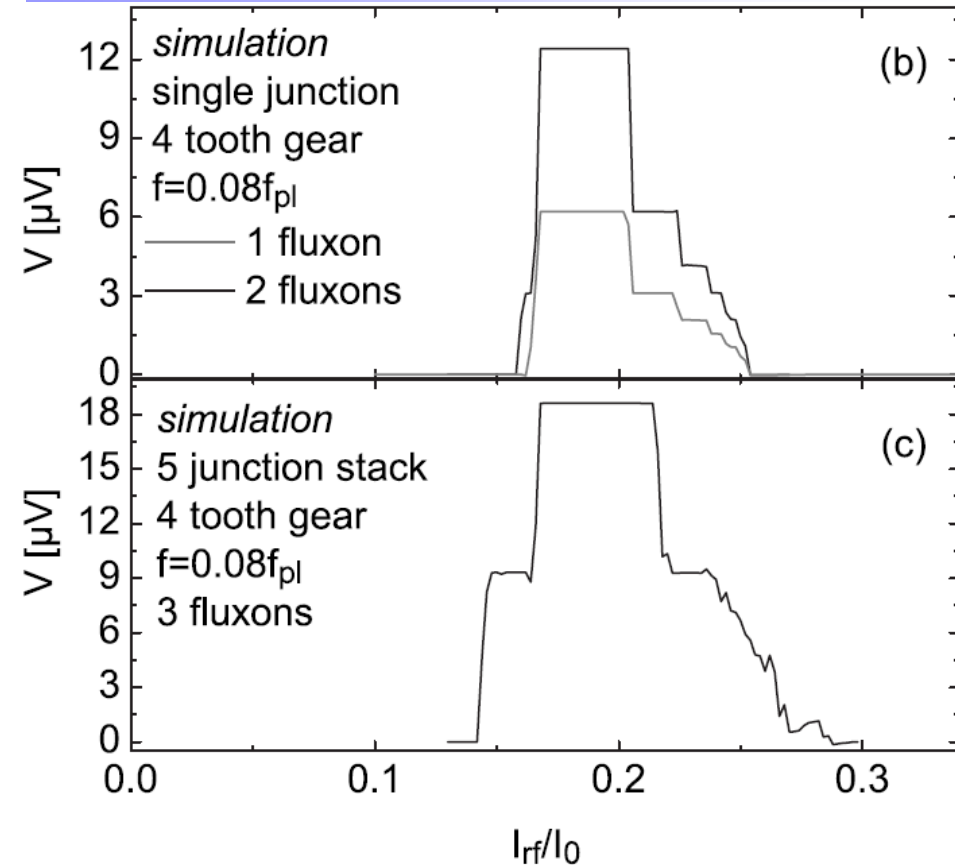
$\langle V \rangle_{max} = 100 \mu\text{V}$ (record!)

$m \sim 40$ fluxons

random fluxon lattice

Why no (Shapiro) steps?

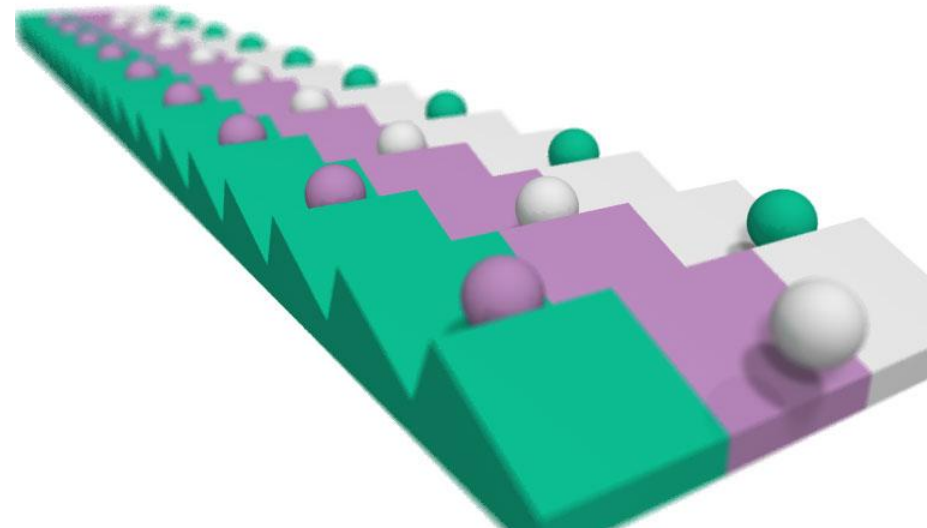
Simulation



For large m and N the steps are smeared!

Summary & Outlook

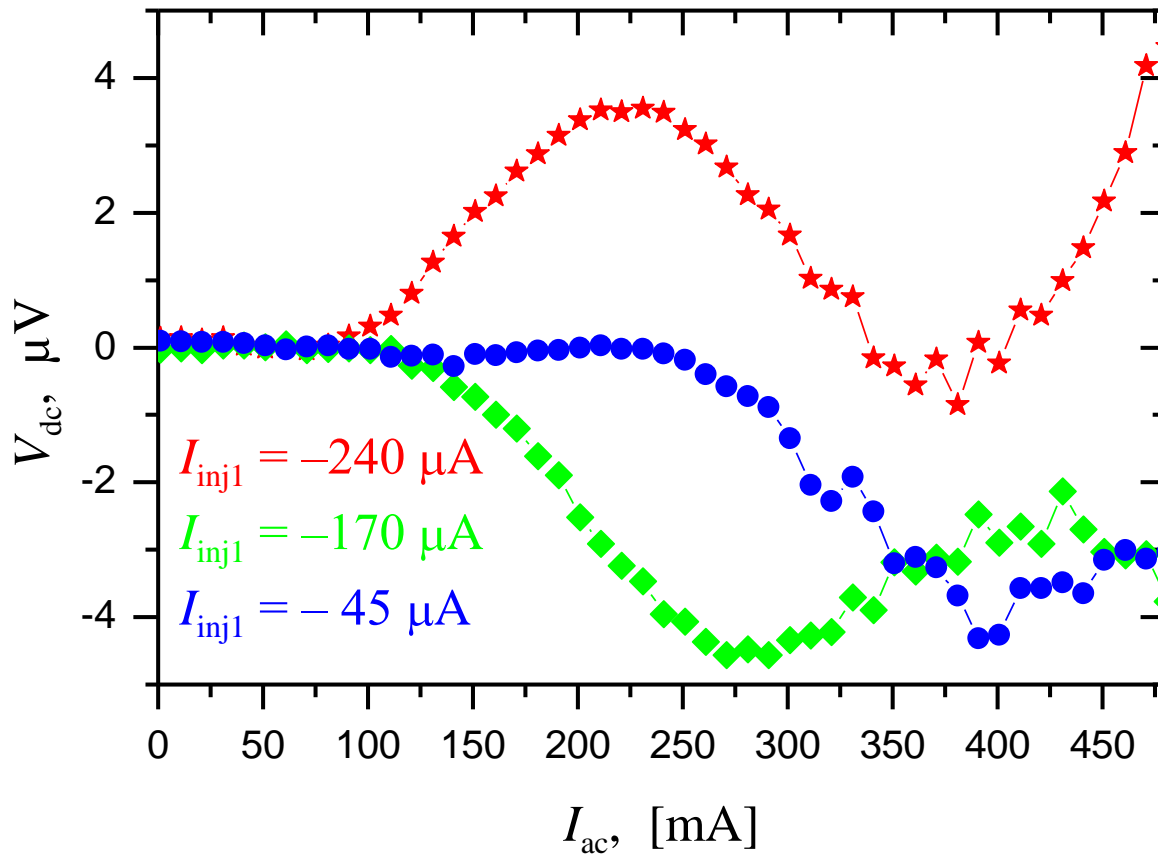
- Fluxon (Josephson Vortex) ratchet
 - principles of operation (asymmetric periodic potential, particle=fluxon)
 - advantages
- Quasi-static (adiabatic) drive:
 - rectification curves vs. potential height, driver shape
 - figures of merit
 - stopping force
- Non-adiabatic drive:
 - quantized rectification
 - period 2 dynamics:
 - subharmonic steps,
 - pseudo-integer steps
 - current (voltage) reversal
 - loaded ratchet, stopping force
- Intrinsic Josephson vortex ratchet



*Thanks for your
attention!*

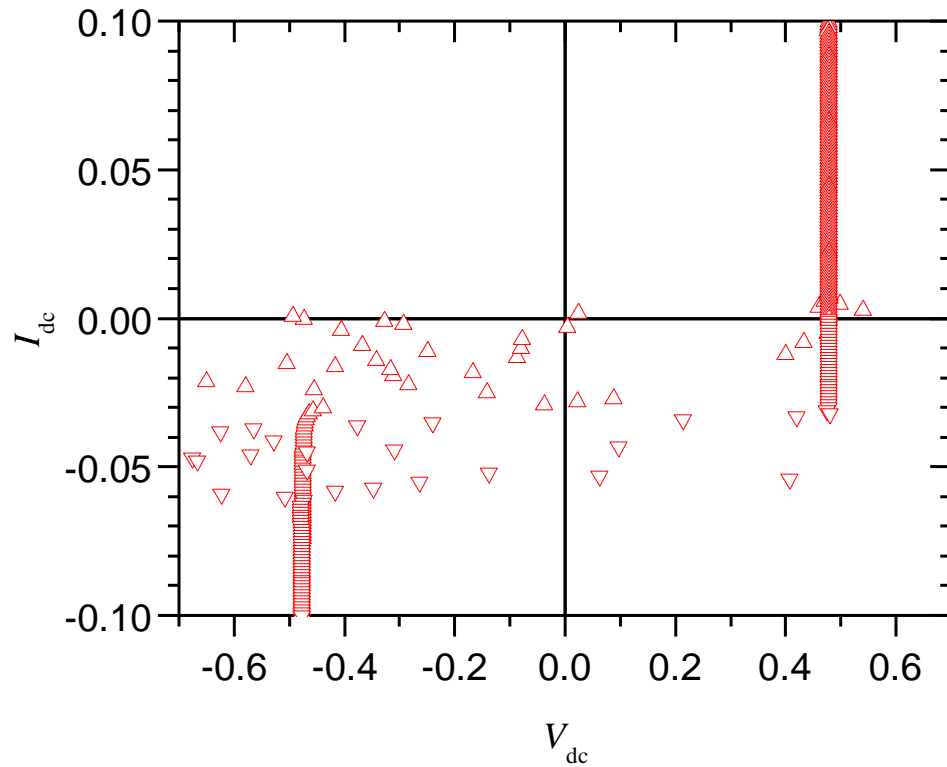
Questions?

Can I rectify a white noise?

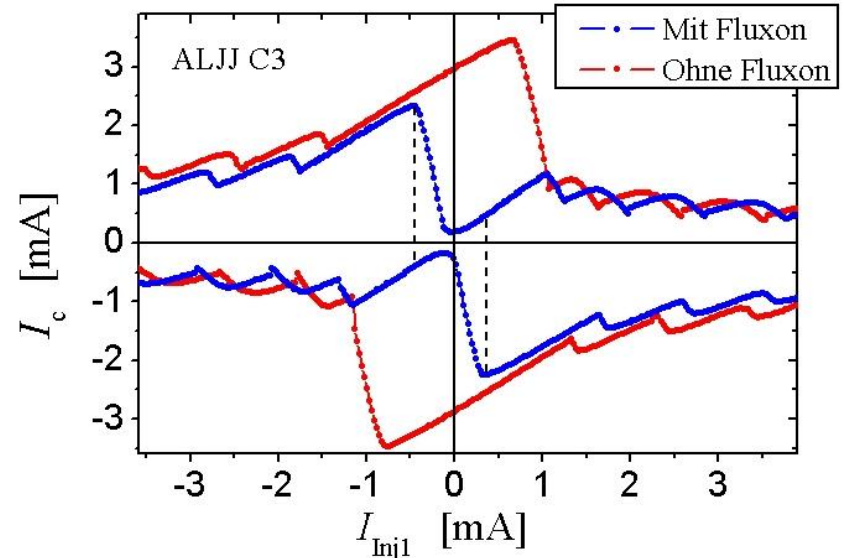
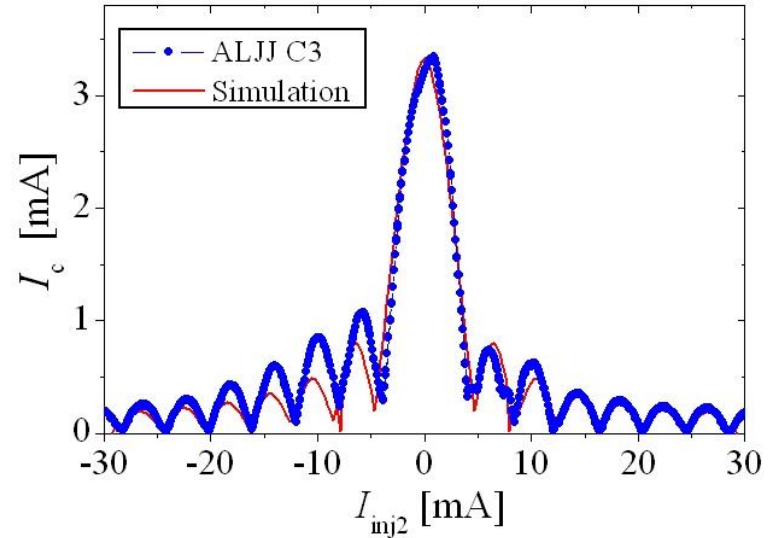
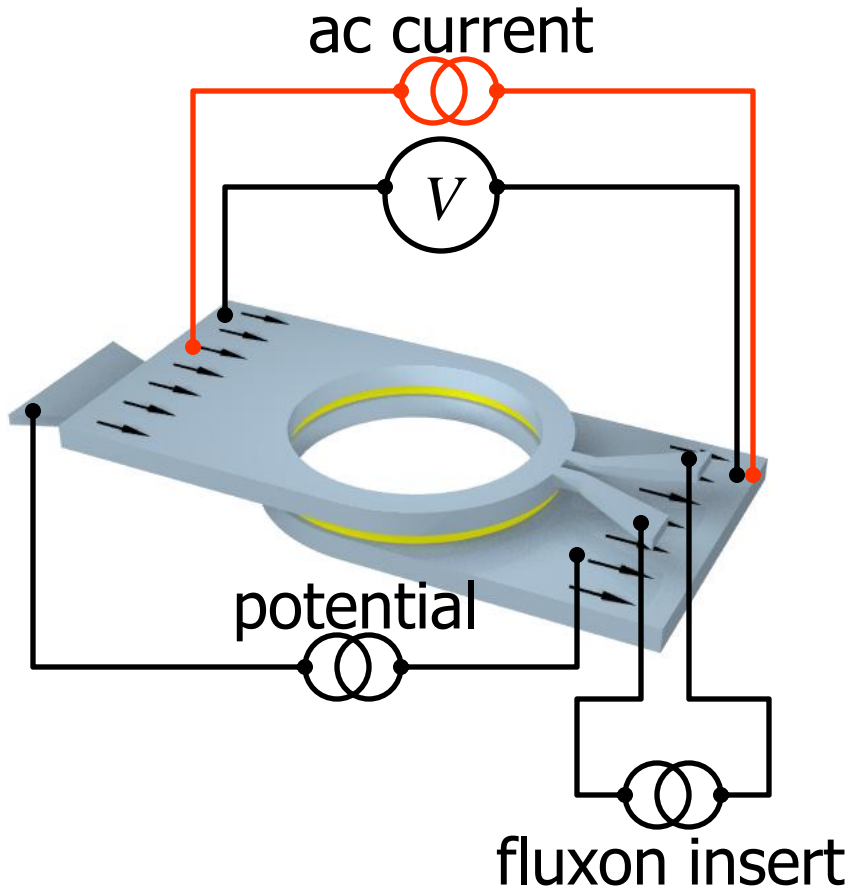


“white” noise cut off at 50MHz

Ratchet or just synchronization?



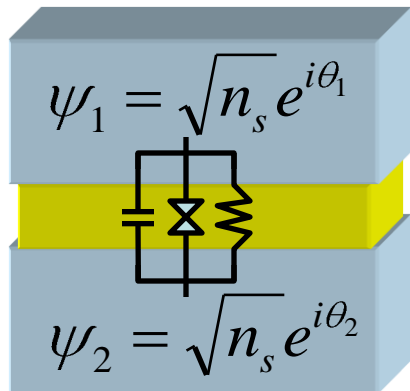
How do we measure?



Apply : $I(t) = I_{ac} \cdot \cos(\omega t)$

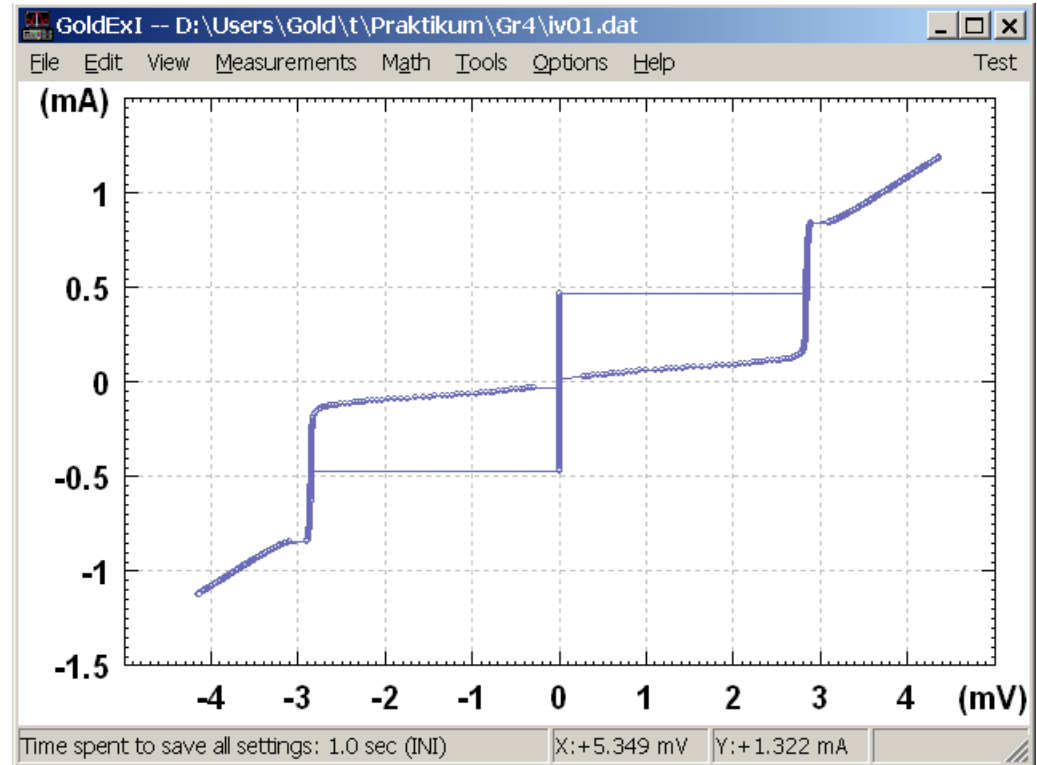
Measure : $V_{dc}(I_{ac}) = \langle V \rangle(I_{ac})$

Josephson junction



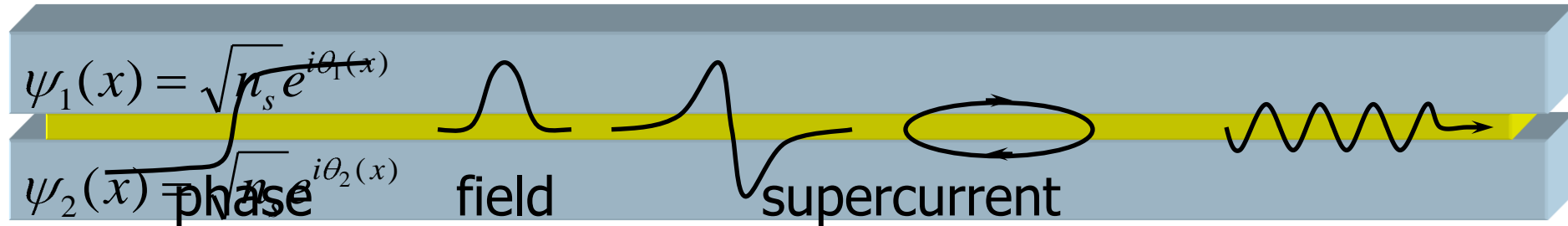
$$I_s = I_c \sin(\theta_2 - \theta_1) = I_c \sin(\phi)$$

$I_c \propto$ overlap of ψ



 B. D. Josephson, Phys. Lett. 1, 251 (1962).

Long Josephson junctions



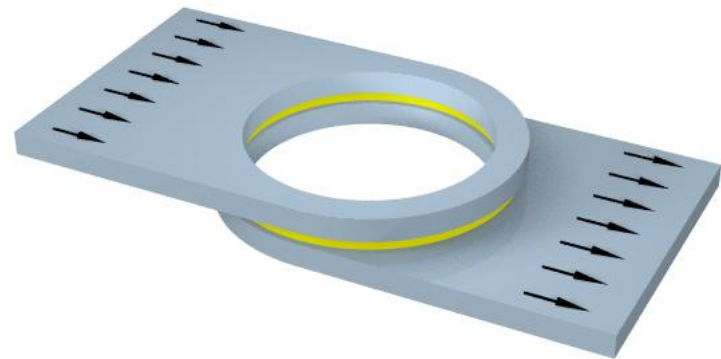
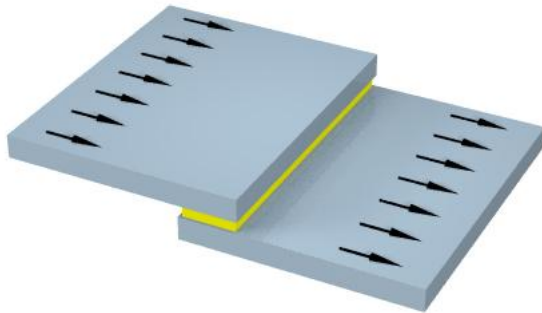
1D model:

$$\phi_{xx} - \phi_{tt} - \sin(\phi) = \alpha\phi_t - \gamma(x)$$

$$\lambda_J \sim 1 \dots 30 \mu\text{m}$$

$$\omega_p \sim 50 \dots 4000 \text{ GHz}$$

$$\lambda_J \cdot \omega_p \equiv \bar{c}_0 \approx \frac{c_{\text{light}}}{20 \dots 40}$$



Content

→ Introduction

- What is (long) Josephson junction?
- What is a ratchet?
- Why to build ratchets using long Josephson junctions?

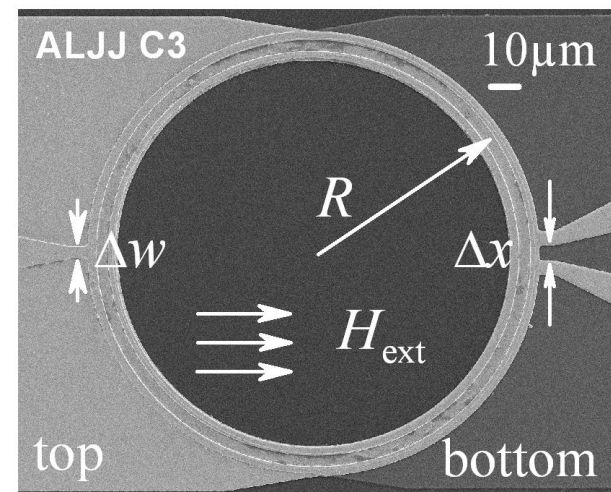
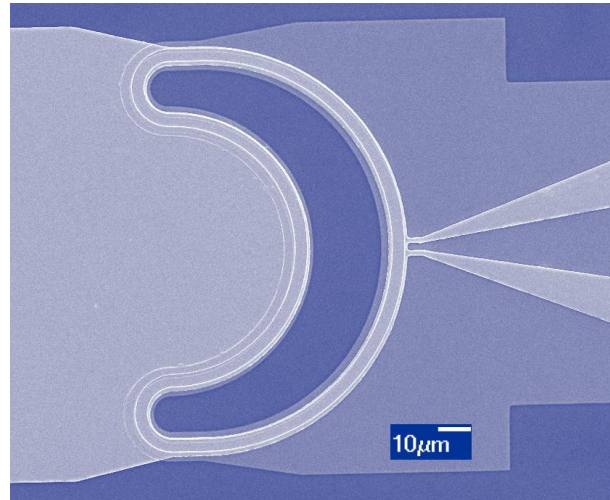
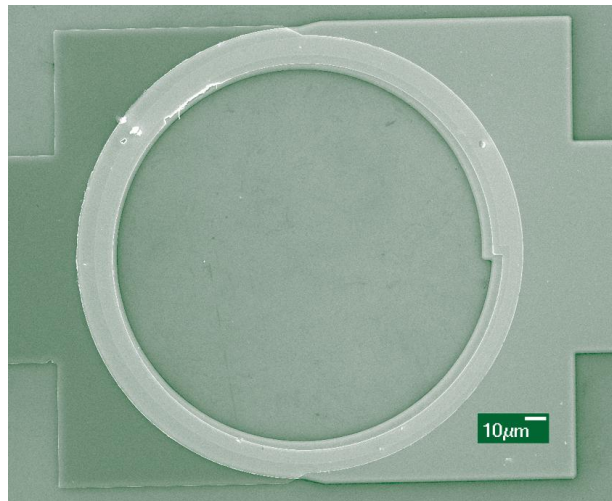
→ How to create an efficient asymmetric periodic potential?

→ Results (experiment and simulations)

- Everything under control: 1 injector, 2 injectors..(calibration)
- Quasi-statically (adiabatically) driven ratchet
- Non-adiabatic effects

→ Conclusions & outlook

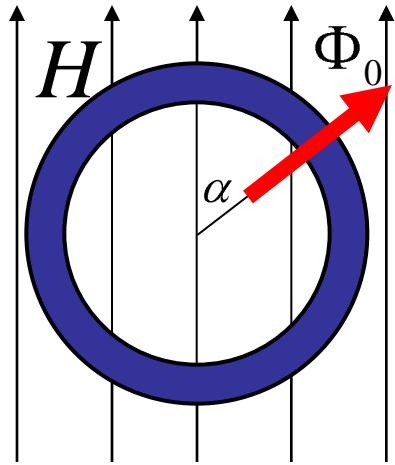
Asymmetric potential



| Potential: | $w(x)$ | $h(x)$ | $j(x)$ |
|-----------------------|--------|--------|--------|
| tunability? | no | yes | yes |
| topological. limits? | no | yes | no |
| current/field source? | no | yes | yes |



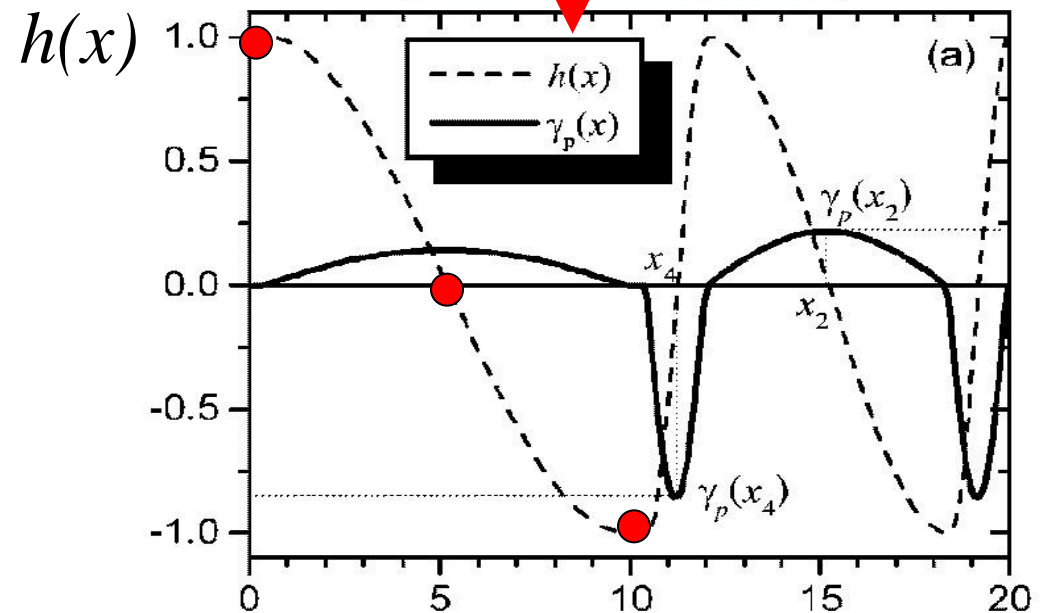
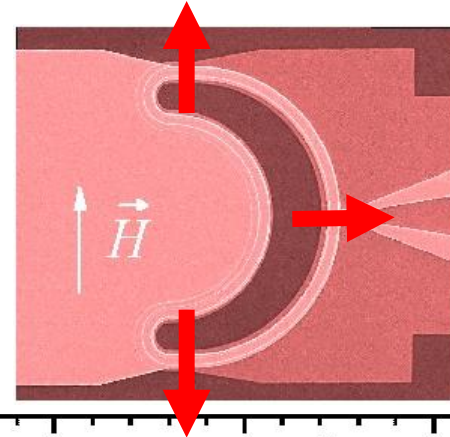
Asymmetry due to non-uniform field



$$H(x) = H \cdot \cos \alpha$$

$$U(x_0) = -2\pi\omega h(x_0)$$

- Tunable potential
- Topological limitations

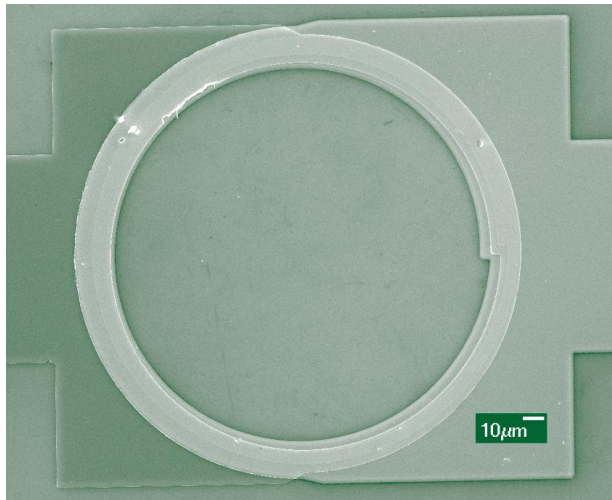


Asymmetric potential

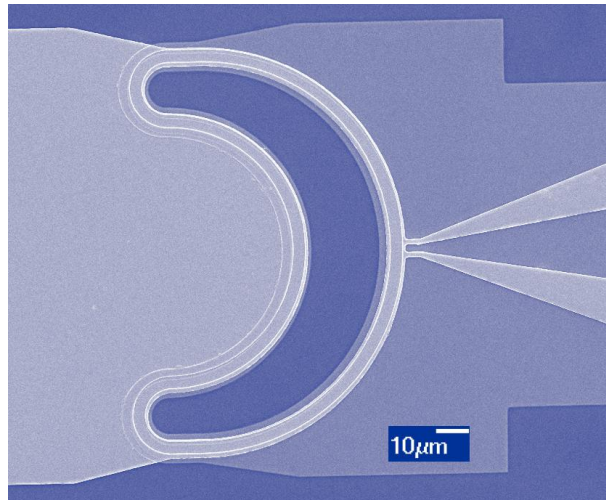
$$\phi_{xx} - \phi_{tt} - \sin(\phi) = \alpha\phi_t - \gamma(x) + h_x(x) - \frac{w_x(x)}{w(x)} [\phi_x - h(x)]$$

$$H = \int_0^L w(x) \left[\frac{\phi_t^2}{2} + \frac{(\phi_x - h)^2}{2} + (1 + \cos \phi) \right] dx$$

$$U(x_0) = \int_{-\infty}^{+\infty} \frac{4w(x)}{\cosh^2(x-x_0)} - \frac{2w(x)h(x)}{\cosh(x-x_0)} dx$$



$$U(x_0) \approx 8w(x_0)$$



$$U(x_0) = -2\pi wh(x_0)$$

